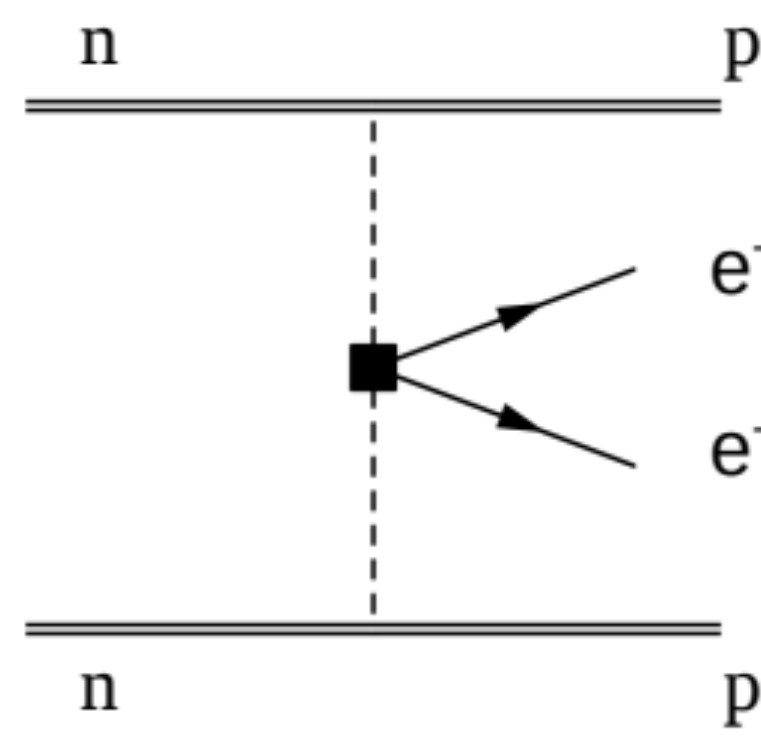


Physics Background

Neutrinoless double beta decay is a hypothetical nuclear transition that has not been discovered yet. Successful observation of this would be a clear indication of BSM.



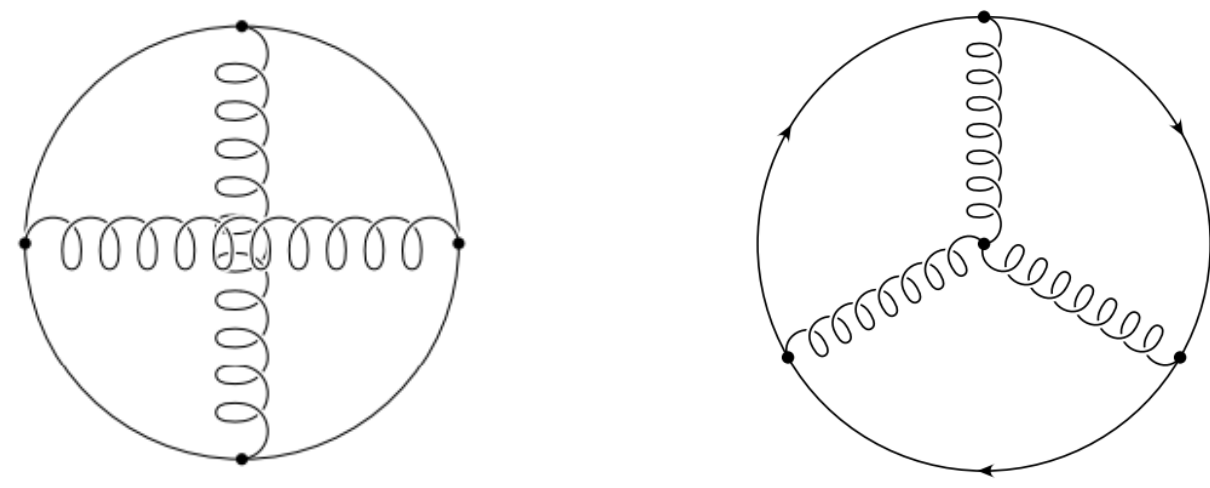
Tools Being Used

EFTs

- encode the relevant low-energy interactions through operators, which are used in many-body nuclear calculations.
- Each operator has an undetermined coefficient, known as the low-energy constants (g LECs), which must be fixed to make accurate predictions.

By extending QCD to a hypothetical limit of many colors, the **large-N** framework imposes a natural hierarchy on the nuclear interactions and decay operators:

- The leading “planar” contributions dominate



vDoBe is an open source code for neutrinoless double β -decay calculations in bounds of couplings and scale of new physics in the SMEFT framework up to dimension 9.

Our Part of The Problem - LECs from Large N

- Four-quark Operators

$$\begin{aligned} O_1 &= \bar{q}_L^\alpha \gamma_\mu \tau^+ q_L^\alpha \bar{q}_L^\beta \gamma^\mu \tau^+ q_L^\beta, \\ O_2 &= \bar{q}_R^\alpha \tau^+ q_L^\alpha \bar{q}_R^\beta \tau^+ q_L^\beta, \\ O_3 &= \bar{q}_R^\alpha \tau^+ q_L^\beta \bar{q}_R^\beta \tau^+ q_L^\alpha, \\ O_4 &= \bar{q}_L^\alpha \gamma_\mu \tau^+ q_L^\alpha \bar{q}_R^\beta \gamma^\mu \tau^+ q_R^\beta, \\ O_5 &= \bar{q}_L^\alpha \gamma_\mu \tau^+ q_L^\beta \bar{q}_R^\beta \gamma^\mu \tau^+ q_R^\alpha, \end{aligned}$$

Use Fiertz Identity for the rest from

$$O_2$$

$$(1 \pm \gamma^5)_{ij}(1 \pm \gamma^5)_{kl} = \frac{1}{2} \left[(1 \pm \gamma^5)_{il}(1 \pm \gamma^5)_{kj} - \frac{1}{2} [(1 \pm \gamma^5)\sigma^{\mu\nu}]_{il} [(1 \pm \gamma^5)\sigma_{\mu\nu}]_{kj} \right]$$

$$\Rightarrow O_3 = -\frac{1}{2}O_2 + \dots$$

$$[(1 \pm \gamma^5)\gamma^\mu]_{ij}[(1 \mp \gamma^5)\gamma_\mu]_{kl} = 2(1 \pm \gamma^5)_{il}(1 \mp \gamma^5)_{kj}$$

$$\Rightarrow O_5 \equiv -2O_2$$

After large N factorization

- $0\nu\beta\beta$ Lagrangian

$$\mathcal{L}_{\Delta L=2}^{(9)} = \frac{1}{v^5} \sum_i \left[\left(C_{iR}^{(9)} \bar{e}_R C \bar{e}_R^T + C_{iL}^{(9)} \bar{e}_L C \bar{e}_L^T \right) O_i + C_i^{(9)} \bar{e} \gamma_\mu \gamma_5 C \bar{e}^T O_i^\mu \right]$$

- Chiral Lagrangian

$$\begin{aligned} \mathcal{L}_\pi^{\text{scalar}} &= \frac{F_0^4}{4} \left[\frac{5}{3} g_1^{\pi\pi} C_{1L}^{(9)} L_{21}^\mu L_{21\mu} + \left(g_2^{\pi\pi} C_{2L}^{(9)} + g_3^{\pi\pi} C_{3L}^{(9)} \right) \text{Tr} (U \tau^+ U \tau^+) \right. \\ &\quad \left. + \left(g_4^{\pi\pi} C_{4L}^{(9)} + g_5^{\pi\pi} C_{5L}^{(9)} \right) \text{Tr} (U \tau^+ U^\dagger \tau^+) \right] \frac{\bar{e}_L C \bar{e}_L^T}{v^5} + (L \leftrightarrow R) \end{aligned}$$

Using large N Factorization for

$$O_2/O_1/O_4$$

$$\begin{aligned} \langle J_1 J_2 \rangle &= \langle J_1 \rangle \langle J_2 \rangle + \langle J_1 J_2 \rangle_c \\ &= \langle J_1 \rangle \langle J_2 \rangle \left[1 + \frac{\langle J_1 J_2 \rangle_c}{\langle J_1 \rangle \langle J_2 \rangle} \right] \end{aligned}$$

$$\langle J_i \rangle \sim N_c$$

$$\langle J_1 J_2 \rangle_c \sim N_c$$

$$\sim \frac{1}{N_c}$$

QCD & ChiPT

Correlation function matching

$$\frac{1}{Z_{QCD}[0]} \frac{\delta Z_{QCD}}{\delta \chi_{ab}} \bigg|_{\chi_{ab}=0} = \frac{1}{Z_{Chi}[0]} \frac{\delta Z_{Chi}}{\delta \chi_{ab}} \bigg|_{\chi_{ab}=0}$$

$$\langle \bar{q}_L^a q_R^b \rangle \rightarrow -\frac{B_0 F_0^2}{2} U_{ba}^+$$

$$O_2 : \mathcal{L}_{0\nu\beta\beta} \supset \frac{1}{v^5} C_{2L}^{(9)} (\bar{l}) (\bar{u}_R d_L) [\bar{u}_R d_L]$$

$$\Rightarrow \frac{1}{v^5} C_{2L}^{(9)} \langle (\bar{u}_R d_L)_x \rangle \langle [\bar{u}_R d_L]_x \rangle$$

$$\rightarrow \frac{1}{v^5} C_{2L}^{(9)} \frac{F_0^4}{4} \underbrace{(B_0^2)}_{g_2^{\pi\pi}} \text{Tr}(\tau^+ U^\dagger \tau^+ U)$$

Results From Large N & SM

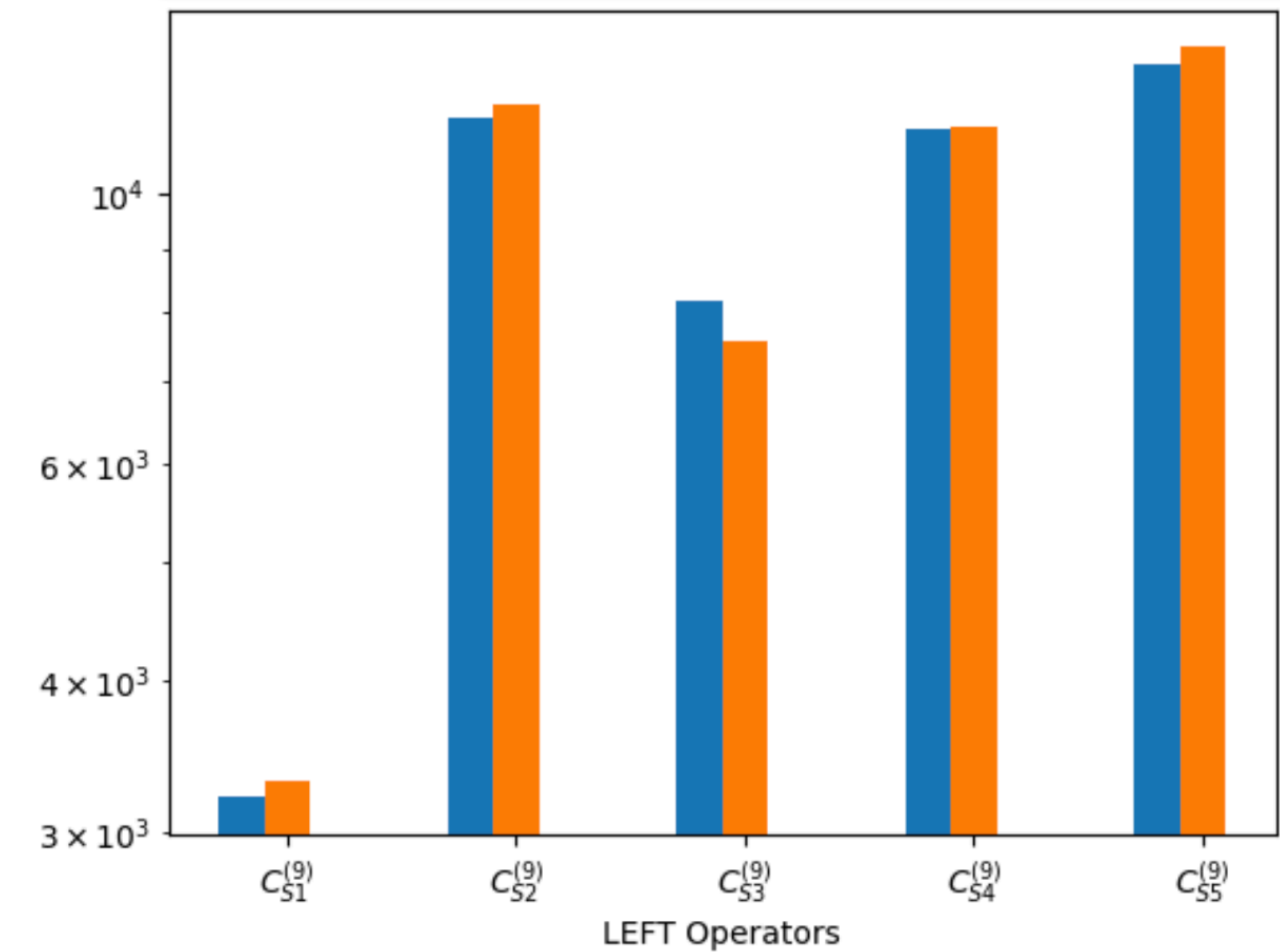
- Blue: Lattice QCD Result

- Orange: Large N EFT Result

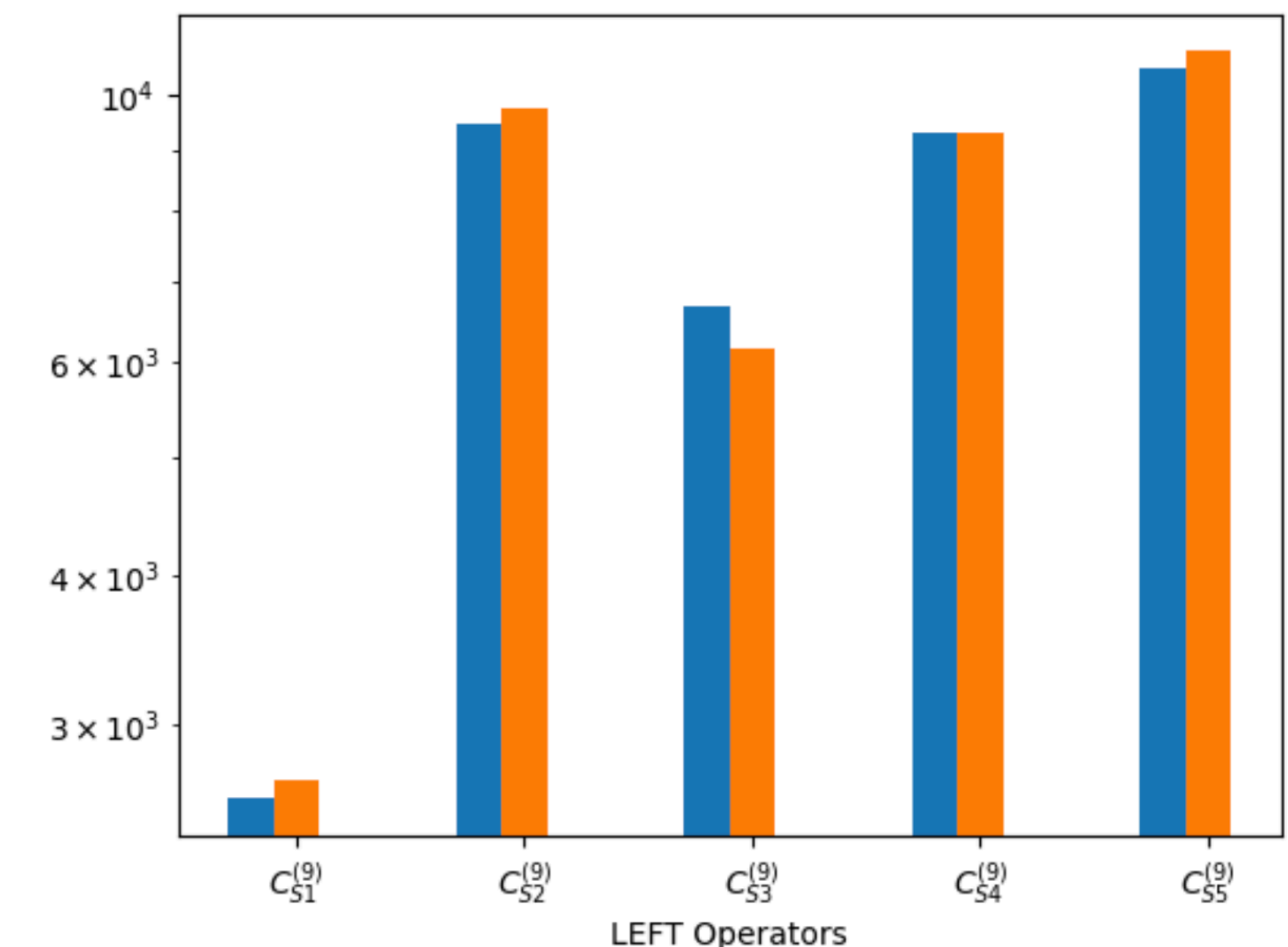
Table of LECs

$g_1^{\pi\pi}$	$g_2^{\pi\pi}$	$g_3^{\pi\pi}$	$g_4^{\pi\pi}$	$g_5^{\pi\pi}$
0.36	2.00	-0.62	-1.90	-8.00
0.60	7.80	-4.00	-1.88	-15.5

Limits On The LEFT Operators (Xe)



Limits On The LEFT Operators (Ge)



Final Remarks / Reference

- Large N factorization works better for vector-vector operators
- Consideration of coulomb for future work
- V. Cirigliano, W. Dekens, J. de Vries, M. L. Graesser & E. Mereghetti, “A neutrinoless double beta decay master formula from effective field theory,” JHEP 12 (2018) 097. / O. Scholer, J. de Vries & L. Gráf, “vDoBe — A Python tool for neutrinoless double beta decay,” JHEP 08 (2023) 043. / V. I. Borodulin, R. N. Rogalyov & S. R. Slabospitskii, “CORE 3.2 (Compendium of Relations, Version 3.2),” arXiv:1702.08246 [hep-ph].