



# Simulating and Interpreting the Multimessenger Picture of Neutron Star Mergers



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- collapsed core of a massive star

smallest and densest known class of stellar compact objectsStellar ev



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- collapsed core of a massive star

- smallest and densest known class of stellar compact objects Stellar



- collapsed core of a massive star
- smallest and densest known class of stellar compact objects
- typical size of 12 kilometer and masses between one and two solar masses



# ... how to study them?

- collapsed core of a massive star
- smallest and densest known class of stellar compact objects
- typical size of 12 kilometer and masses between one and two solar masses



### Single Neutron Stars



radio pulsarsthrough X-ray emission

# ... how to study them?

- collapsed core of a massive star
- smallest and densest known class of stellar compact objects
- typical size of 12 kilometer and masses between one and two solar masses



### **Binary Neutron Stars**



- gravitational-wave sources
- electromagnetic transients
- neutrino sources







## **Combining different constraints on the EOS from different research fields**

Combined Equation of SLX +	Science case:
← → C Q A 0.0.0.5000	
An overview of existing and new nuclear and astrophysical constraints on neutron-rich dense matter	the equation of state of arXiv:2402.04172v1
his tool can be used to combine various constraints on the equation of state (EOS) for dense matter. Select the constraints you are interested in. Clicking on the buttons rovide the figures for either EOS-derived quantities or show how the estimate for the canonical neutron star radius changes. Dependencies are taken into account autom y clicking on the images, you can switch between the M-R curve and the corresponding pressure-density relation. ou can also choose weights for the individual inputs, so when the log-likelihoods are added, the weight will be used as a coefficient. We emphasize that the weights are found statistical interpretation.	below will then give you the combined posterior and atically.
Microscopic Theory	
Microscopic Experiments	Accessible here
Astrophysical Limits on the TOV Mass	
Astrophysical M-R Constraints	
Gravitational-Wave and Multimessenger Constraints	
Prior	
Compare Evolution Compare C	bservables
he Numanji Collaboration G Theoretische Astrophysik stitut für Physik und Astronomie niversität Potsdam arl-Liebknecht-Str. 24/25 4476 Potsdam ermany	
chrome	

## **Combining different constraints on the EOS from different research fields**



Science case: Koehn et al. 2024 arXiv:2402.04172v1



## **Combining different constraints on the EOS from different research fields**



Set

Α

 $\chi EFT$ 

pQCD

NICER

J0030 + 0451

В

Set A

HIC

J0952-0607

qLMXBs

Heavy pulsars Black Widow

 $\mathbf{C}$ 

Set B

CREX

PREX-II

<sup>208</sup>Pb dipole



# **The Binary Neutron Star Merger Simulation**



gravitational wave emission

# **The Binary Neutron Star Merger Simulation**



gravitational wave emission

deformation before merger, ejection of material, heavy element production

# **The Binary Neutron Star Merger Simulation**



gravitational wave emission

deformation before merger, ejection of material, heavy element production

black hole formation



## **Theoretical Framework:**

- well-posedness of PDEs
- advantageous properties

$$\begin{split} \partial_t \chi &= \frac{2}{3} \chi \left( \alpha (\ddot{K} + 2\Theta) - D_i \beta^i \right), \\ \partial_t \ddot{\gamma}_{ij} &= -2\alpha \ddot{A}_{ij} + \beta^k \partial_i \ddot{\gamma}_{ij} + 2\ddot{\gamma}_{ki} (\partial_j) \beta^k - \frac{2}{3} \ddot{\gamma}_{ij} \partial_k \beta^k, \\ \partial_t \ddot{K} &= -D^i D_i \alpha + \alpha \left( \ddot{A}_{ij} \dot{A}^{ij} + \frac{1}{3} (\ddot{K} + 2\Theta)^2 \right) \\ &+ 4\pi \alpha (S + E) + \beta^k \partial_k \ddot{K} + \alpha \kappa_i (1 - \kappa_2) \Theta, \\ \partial_t \ddot{A}_{ij} &= \chi \left( -D_i D_i \alpha + \alpha \left( ^{3} R_{ij} - 8\pi S_{ij} \right) \right)^{TT} + \alpha \left( (\ddot{K} + 2\Theta) \dot{A}_{ij} - 2 \ddot{A}^{ij} \ddot{A}_{kj} \right) \\ &+ \beta^k \partial_k \ddot{A}_{ij} + 2 \dot{A}_{ki} (\partial_j) \beta^k - \frac{2}{3} \ddot{A}_{ij} \partial_k \beta^k, \\ \partial_t \ddot{\Gamma}^i &= -2 \ddot{A}^{ik} \partial_k \alpha + 2\alpha \left( \ddot{\Gamma}_{ik} \ddot{A}^{il} - \frac{3}{2} \ddot{A}^{ik} \partial_k \ln(\chi) - \frac{1}{3} \ddot{\gamma}^{ik} \partial_k (\ddot{K} + 2\Theta) - 8\pi \ddot{\gamma}^{ik} S_k \right) \\ &+ \ddot{\gamma}^{ik} \partial_k \partial\beta^j + \frac{1}{3} \dot{\gamma}^{ik} \partial_k \partial\beta^j - 2\alpha \kappa_i (\ddot{\Gamma}^i - \Gamma^i) + \beta^k \partial_k \ddot{\Gamma}^i \\ &- \Gamma^k \partial_k \beta^i + \frac{2}{3} \dot{\Gamma}^i \partial_\beta \beta^k, \\ \partial_i \Theta &= \frac{\alpha}{2} \left( (^{3)} R - \ddot{A}_{ij} \ddot{A}^{ij} + \frac{2}{3} (\ddot{K} + 2\Theta)^2 \right) - \alpha \left( 8\pi E + \kappa_1 (2 + \kappa_2) \Theta \right) \\ &+ \beta^i \partial_i \Theta \end{split}$$



- well-posedness of PDEs
- advantageous properties

$$\begin{split} \partial_{tX} &= \frac{2}{3} \chi \left( \alpha (\tilde{K} + 2\Theta) - D_{i} \beta^{i} \right), \\ \partial_{t} \tilde{\gamma}_{ij} &= -2\alpha \tilde{A}_{ij} + \beta^{2} \partial_{k} \tilde{\gamma}_{ij} + 2\tilde{\gamma}_{ki} (\partial_{j}) \beta^{k} - \frac{2}{3} \tilde{\gamma}_{ij} \partial_{k} \beta^{k}, \\ \partial_{t} \tilde{K} &= -D^{i} D_{i} \alpha + \alpha \left( \tilde{A}_{ij} \tilde{A}^{ij} + \frac{1}{3} (\tilde{K} + 2\Theta)^{2} \right) \\ &+ 4\pi \alpha (S + E) + \beta^{k} \partial_{k} \tilde{K} + \alpha \kappa_{i} (1 - \kappa_{2}) \Theta, \\ &+ \beta^{k} \partial_{k} \tilde{A}_{ij} + 2 \tilde{A}_{ki} (\partial_{j}) \beta^{k} - \frac{2}{3} \tilde{A}_{ij} \partial_{k} \beta^{k}, \\ \partial_{t} \tilde{A}_{ij} &= \chi \left( -D_{i} D_{i} \alpha + \alpha \left( ^{3} R_{ij} - 8\pi S_{ij} \right) \right)^{TF} + \alpha \left( (\tilde{K} + 2\Theta) \tilde{A}_{ij} - 2 \tilde{A}^{k}_{j} \tilde{A}_{kj} \right) \\ &+ \beta^{k} \partial_{k} \tilde{A}_{ij} + 2 \tilde{A}_{ki} (\partial_{j}) \beta^{k} - \frac{2}{3} \tilde{A}_{ij} \partial_{k} \beta^{k}, \\ \partial_{t} \tilde{I}^{i} &= -2 \tilde{A}^{ik} \partial_{k} \alpha + 2\alpha \left( \tilde{I}_{ki} \tilde{A}^{ki} - \frac{3}{2} \tilde{A}^{ki} \partial_{k} \ln(\chi) - \frac{1}{3} \gamma^{ik} \partial_{k} (\tilde{K} + 2\Theta) - 8\pi \tilde{\gamma}^{ik} S_{k} \right) \\ &+ \tilde{\gamma}^{ki} \partial_{k} \partial_{i} \beta^{i} + \frac{1}{3} \gamma^{ik} \partial_{k} \beta^{j} - 2\alpha \kappa_{i} (\tilde{\Gamma}^{i} - \tilde{\Gamma}^{i}) + \beta^{k} \partial_{k} \tilde{\Gamma}^{i} \\ &- \Gamma^{k} \partial_{k} \beta^{i} + \frac{2}{3} \tilde{I}^{i} \partial_{k} \beta^{k}, \\ \partial_{i} \Theta &= \frac{\alpha}{2} \left( (3R - \tilde{A}_{ij} \tilde{A}_{ij} + \frac{2}{3} (\tilde{K} + 2\Theta)^{2} \right) - \alpha \left( 8\pi E + \kappa_{i} (2 + \kappa_{2}) \Theta \right) \\ &+ \beta^{i} \partial_{i} \Theta \end{split}$$



### 3+1-decomposition

 $\partial_t \mathbf{u} = \mathbf{A}(\mathbf{u})\mathbf{u} + \mathbf{v}$ 

 $G_{\mu\nu} = 8\pi T_{\mu\nu}$ 

Reformulating as initial value boundary problem



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- advantageous properties



### **Computational Methods:**

- HPC facilities
- parallelizable code
- numerical techniques







## 

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- HPC facilities
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- Microphyics (EOS, Neutrinos)
- Magnetic fields
- Turbulences
- Parameter space





- Microphyics (EOS, Neutrinos)
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- Turbulences
- Parameter space





• Can we test matter above the TOV limit?

$$\begin{split} & \partial_t \Theta = \frac{\alpha}{2} \binom{(3)R - \tilde{A}_{ij}\tilde{A}^{ij} + \frac{2}{3}(\tilde{K} + 2\Theta)^2) - \alpha \left(8\pi E + \kappa_1(2 + \alpha) + \beta^i \partial_i \Theta\right)}{+ \beta^i \partial_i \Theta} \end{split}$$

Ujevic et al., Astrophys.J.Lett. 962 (2024) 1, L3



- Microphyics (EOS, Neutrinos)
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• Can we test matter above the TOV limit?

No!

 $\begin{array}{l} \partial_t \Theta &= \frac{\alpha}{2} \begin{pmatrix} (3)_R - \tilde{\lambda}_{ij} \tilde{A}^{ij} + \frac{2}{3} (\tilde{K} + 2\Theta)^2 \end{pmatrix} - \alpha \left( 8\pi E + \kappa_1 (2 + \kappa_2) + \beta^i \partial_i \Theta \right) \\ &+ \beta^i \partial_i \Theta \end{array}$ 

Ujevic et al., Astrophys.J.Lett. 962 (2024) 1, L3

- Microphyics (EOS, Neutrinos)
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 $\partial_t \Theta = \frac{\alpha}{2} \binom{(3)R - \tilde{A}_{ij}\tilde{A}^{ij} + \frac{2}{3}(\tilde{K} + 2\Theta)^2}{+\beta^i \partial_i \Theta} - \alpha \left(8\pi E + \kappa_1(2 + \kappa_2)e^{-2}\right) + \beta^i \partial_i \Theta$ 

Gieg et al., Universe 8 (2022) 7, 370

- Microphyics (EOS, Neutrinos)
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Schianchi et al., arXiv: 2307.04572 Gieg et al., Universe 8 (2022) 7, 370

Charged Current Processes  $\nu_e + n \leftrightarrow p + e^ \overline{\nu}_e + p \leftrightarrow n + e^+$   $\nu_e + (A, Z) \leftrightarrow (A, Z + 1) + e^-$ Thermal Processes  $e^- + e^+ \leftrightarrow \nu_x + \overline{\nu}_x$   $N + N \leftrightarrow N + N + \nu_x + \overline{\nu}_x$ Elastic Scattering  $\nu + \alpha \rightarrow \nu + \alpha$   $\nu + p \rightarrow \nu + p$   $\nu + n \rightarrow \nu + n$  $\nu + (A, Z) \rightarrow \nu + (A, Z)$ 

 $\partial_t \Theta = \frac{\alpha}{2} \left( {}^{(3)}R - \hat{A}_{ij} \hat{A}^{ij} + \frac{2}{3} (\hat{K} + 2\Theta)^2 \right) - \alpha \left( 8\pi E + \kappa_1 (2 + 2\Theta)^2 \right) + \alpha \left( 8\pi E +$ 

• Inclusion of neutrinos changes matter outflow and remnant's lifetime

• Amount of produced elements and their abundance depend on merger properties and neutrino scheme

- Microphyics (EOS, Neutrinos)
- Magnetic fields
- Turbulences
- Parameter space





	1400		0.00	$\cap$
0.5	1200		7.5	-14
0.4	1000		7.0 🚍	12
0.3 . 2	800		6.5 Ja	10
0.2	600		6.0 (j)=6	
	400		5.0	
0.1	200		-4.5	
0.0	0	1100 1000 100 100 100 100 100 1000 1000	4.0	4

magnetic

fields and

turbulences

PRD 110 (2024) 8, 084046

 $\partial_t \Theta = \frac{\alpha}{2} \binom{(3)R - \tilde{A}_{ij} \tilde{A}^{ij} + \frac{2}{3} (\tilde{K} + 2\Theta)^2}{+\beta^i \partial_i \Theta} - \alpha \left(8\pi E + \kappa_1 (2 + \kappa + 2\Theta)^2 - \alpha (2 + \kappa$ 

t = 0.0 ms $t=8.13~{\rm ms}$  $10^{15}$  $10^{14}$ t = 10.34 mst = 9.61 ms $10^{14}$  $\rho \; [\mathrm{g/cm^3}]$ [B] [G]  $10^{13}$ t = 21.43 mst = 25.12 ms $10^{12}$  $10^{12}$  -

- Microphyics (EOS, Neutrinos)
- Magnetic fields
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### $\partial_t\Theta = \frac{\alpha}{2} \begin{pmatrix} ^{(3)}R - \bar{\lambda}_{ij}\dot{A}^{ij} + \frac{2}{3}(\dot{K} + 2\Theta)^2 \end{pmatrix} - \alpha \left(8\pi E + \kappa_1(2 + \kappa_2)\right) + \frac{1}{2} \left(\frac{1}{2} - \frac{1}{2}\right) + \frac{1}{2} \left(\frac{1}{2} - \frac{1$ + Bi die

### Markin et al., PRD 108 (2023) 2, 023016

 $\rho$  [g cm - 10<sup>12</sup> Time: 0.0 ms

### **Input Physics:**

- Microphyics (EOS, Neutrinos) •
- Magnetic fields •
- Turbulences •
  - Parameter space



### parameter space coverage



 $M_{\text{NS}} = 1.4~\text{M}_{\odot}$  $M_{BH}=0.5~M_{\odot}$ 







• First simulation of a subsolar mass BH – neutron star merger

- large amount of ejecta
- existing waveform models perform badly when describing



## space coverage

 $\partial_t \Theta = \frac{\alpha}{2} \binom{(3)R - \tilde{A}_{ij}\tilde{A}^{ij} + \frac{2}{3}(\tilde{K} + 2\Theta)^2) - \alpha \left(8\pi E + \kappa_1(2 + \kappa_2) + \beta^i \partial_i \Theta\right)}{+\beta^i \partial_i \Theta}$ 



publicly released more than 590 individual simulations using more than <sup>1</sup>/<sub>2</sub> billion CPUhs

Dietrich et al., CCG 35 (2018) 24, 24LT0 Gonzales et al., QCG 40 2023) 8, 085011

- Microphyics (EOS, Neutrinos)
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  - Parameter space



## **Various models**



Numerical Relativity Simulations



**Post-Newtonian Theory** 

Effective-one-body Formalism



Phenomenological Models

hundreds of millions of templates need to interpret the data





tidal effects lead to an accelerated inspiral

# **Waveform Model Development through NR simulations**

Effective-one-body or Phenomenological Model

confirmation/calibration

Numerical Relativity Simulations



**TD & Hinderer, Phys.Rev.D 95 12, 124006** 



N.Kunert et al., PRD105 (2022) 6, L061301



N.Kunert et al., PRD105 (2022) 6, L061301



N.Kunert et al., PRD105 (2022) 6, L061301

## *GW170817*

 $\Lambda$  determines tidal deformability



## *GW170817*

 $\Lambda$  determines tidal deformability





 $\rightarrow$  no assumption about the type of the compact object

Phys.Rev. X9 (2019) 011001



## *GW170817*

 $\Lambda$  determines tidal deformability



# **EM Signals – Kilonova**

- neutron rich ejecta produce heavy r-process elements
- pseudo-black body radiation from r-process elements
- mergers are major sites for the formation of heavy elements



















## Uncertainties

- 1.) Knowledge about the outflowing material (mass, velocity, geometry, composition)
- 2.) Heating rates depend on the formed elements and ejecta properties
- 3.) Incomplete knowledge about opacities for complicated elements



Cross-code comparisons for numerous geometries and assumptions  $\rightarrow$  estimate on the modelling uncertainty





1.) compute lightcurves for a set (grid) of ejecta properties with a radiative transfer code

2.) interpolate within this grid through Gaussian Process Regression or a Neural Network

3.) link ejecta properties through numerical-relativity predictions to the binary properties





Huth et al., Nature 606 (2022) 276-280



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Article Open Access Published: 08 June 2022

# Constraining neutron-star matter with microscopic and macroscopic collisions

<u>Sabrina Huth</u> ⊡, <u>Peter T. H. Pang</u> ⊡, <u>Ingo Tews</u>, <u>Tim Dietrich</u>, <u>Arnaud Le Fèvre</u>, <u>Achim Schwenk</u>, <u>Wolfgang Trautmann</u>, <u>Kshitij Agarwal</u>, <u>Mattia Bulla</u>, <u>Michael W. Coughlin</u> & <u>Chris Van Den Broeck</u>



Huth et al., Nature 606 (2022) 276-280











# Science Summary and Outlook

 numerical-relativity simulations (microphysics and parameter coverage)





 new nuclear physics and multimessenger astrophysics framework



• constraints on the Hubble constant and supranuclear-dense equation of state

### Gravitational Waves





Electromagnetic Signals





... neutrino detectors, nuclear physics facilities, ....