

Nucleosynthesis and Galactic Chemical Evolution

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Third Santa Cruz School on Multi-Messenger Astrophysics
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July 16 & 17, 2024



Network for Neutrinos,
Nuclear Astrophysics,
and Symmetries



Outline (7/16 & 7/17)

1. Overview of nucleosynthesis

2. Stellar Evolution & Nucleosynthesis

3. Cosmological Structure Formation

4. Galactic Chemical Evolution

Textbooks

Clayton: Principles of Stellar Evolution & Nucleosynthesis

Arnett: Supernovae & Nucleosynthesis

Pagel: Nucleosynthesis & Chemical Evolution of Galaxies

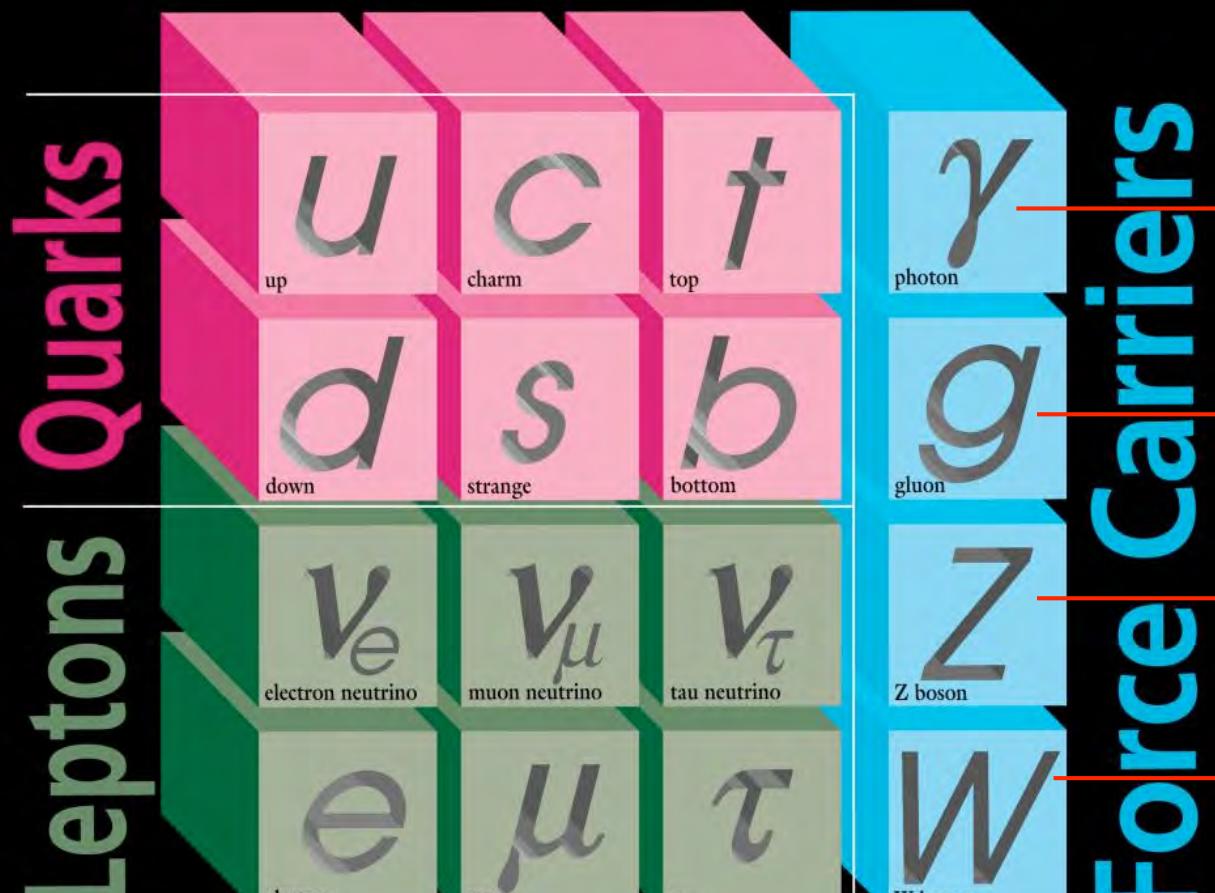
Periodic Table of Elements

1	2	3	4	5	6	7	8	9	10	11	12	13	14	15	16	17	18	
1 H Hydrogen 1.00794	2 He Helium 4.002602																	
3 Li Lithium 6.941	4 Be Beryllium 9.012162																	
11 Na Sodium 22.98976925	12 Mg Magnesium 24.3050																	
19 K Potassium 39.0983	20 Ca Calcium 40.078	21 Sc Scandium 44.955912	22 Ti Titanium 47.867	23 V Vanadium 50.9415	24 Cr Chromium 51.9861	25 Mn Manganese 54.938045	26 Fe Iron 55.845	27 Co Cobalt 58.93186	28 Ni Nickel 58.6934	29 Cu Copper 63.545	30 Zn Zinc 65.38	31 Ga Gallium 69.723	32 Ge Germanium 72.64	33 As Arsenic 74.9216	34 Se Selenium 78.903	35 Br Bromine 79.904	36 Kr Krypton 83.798	
37 Rb Rubidium 85.4675	38 Sr Strontium 87.62	39 Y Yttrium 88.90385	40 Zr Zirconium 91.224	41 Nb Niobium 92.90688	42 Mo Molybdenum 95.96	43 Tc Technetium (97.9072)	44 Ru Ruthenium 101.07	45 Rh Rhodium 102.90580	46 Pd Palladium 106.42	47 Ag Silver 107.8682	48 Cd Cadmium 112.411	49 In Indium 114.810	50 Sn Tin 115.710	51 Sb Antimony 121.760	52 Te Tellurium 127.80	53 I Iodine 126.90447	54 Xe Xenon 131.293	
55 Cs Cesium 132.9054519	56 Ba Barium 137.327	57-71		72 Hf Hafnium 178.49	73 Ta Tantalum 180.94738	74 W Tungsten 183.84	75 Re Rhenium 186.207	76 Os Osmium 190.23	77 Ir Iridium 192.217	78 Pt Platinum 195.064	79 Au Gold 196.995568	80 Hg Mercury 200.58	81 Tl Thallium 204.3833	82 Pb Lead 207.2	83 Bi Bismuth 208.66040	84 Po Polonium (208.6624)	85 At Astatine (218.8611)	86 Rn Radon (222.0176)
87 Fr Francium (223)	88 Ra Radium (226)	89-103		104 Rf Rutherfordium (261)	105 Db Dubnium (262)	106 Sg Seaborgium (265)	107 Bh Bohrium (264)	108 Hs Hassium (277)	109 Mt Meitnerium (268)	110 Ds Darmstadtium (271)	111 Rg Roentgenium (272)	112 Uub Ununbium (286)	113 Uut Ununtrium (284)	114 Uup Ununpentium (285)	115 Uuh Ununhexium (286)	116 Uuo Ununoctium (290)	117 Uus Ununseptium (289)	118 Uuo Ununoctium (294)
For elements with no stable isotopes, the mass number of the isotope with the longest half-life is in parentheses.																		

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57 La Lanthanum 138.90547	58 Ce Cerium 140.116	59 Pr Praseodymium 140.90785	60 Nd Neodymium 144.242	61 Pm Promethium (145)	62 Sm Samarium 150.38	63 Eu Europium 151.904	64 Gd Gadolinium 157.25	65 Tb Terbium 158.92535	66 Dy Dysprosium 162.500	67 Ho Holmium 164.93032	68 Er Erbium 167.259	69 Tm Thulium 168.93421	70 Yb Ytterbium 173.054	71 Lu Lutetium 174.9068			
89 Ac Actinium (227)	90 Th Thorium 232.03800	91 Pa Protactinium 231.03688	92 U Uranium 238.02891	93 Np Neptunium (237)	94 Pu Plutonium (244)	95 Am Americium (243)	96 Cm Curium (247)	97 Bk Berkelium (247)	98 Cf Californium (251)	99 Es Einsteinium (252)	100 Fm Fermium (257)	101 Md Mendelevium (258)	102 No Nobelium (259)	103 Lr Lawrencium (262)			

ELEMENTARY PARTICLES



I II III
Three Generations of Matter

Fundamental
Interactions

Electromagnetic

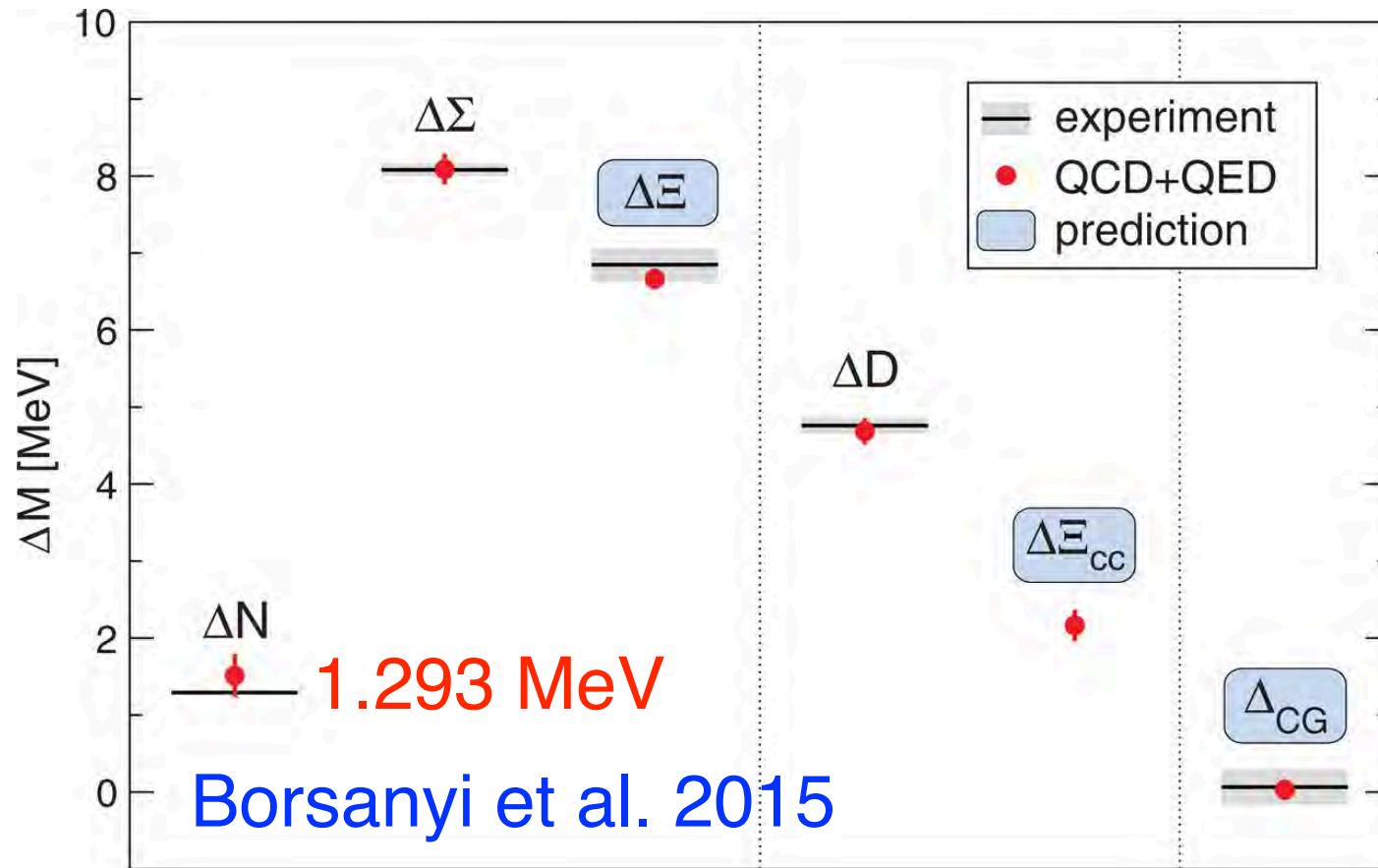
Strong

Weak (NC)

Weak (CC)

Gravitational

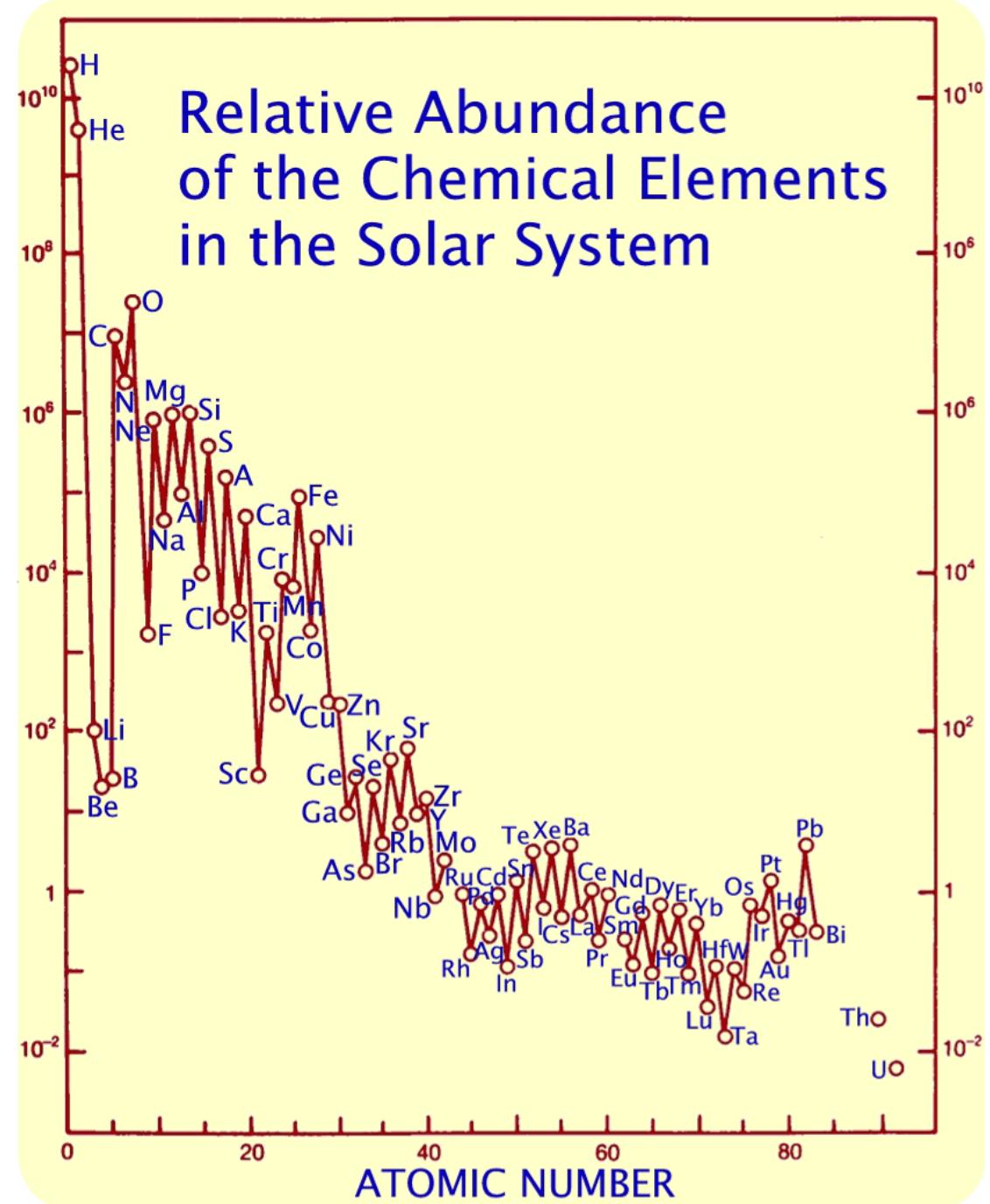
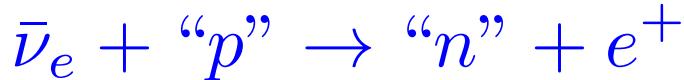
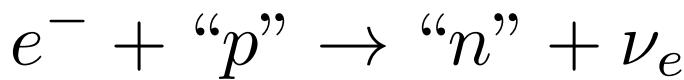
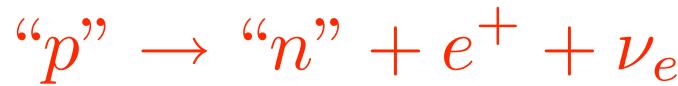
Standard Model of Particle Physics & Life of a Baryon: Big Bang Nucleosynthesis

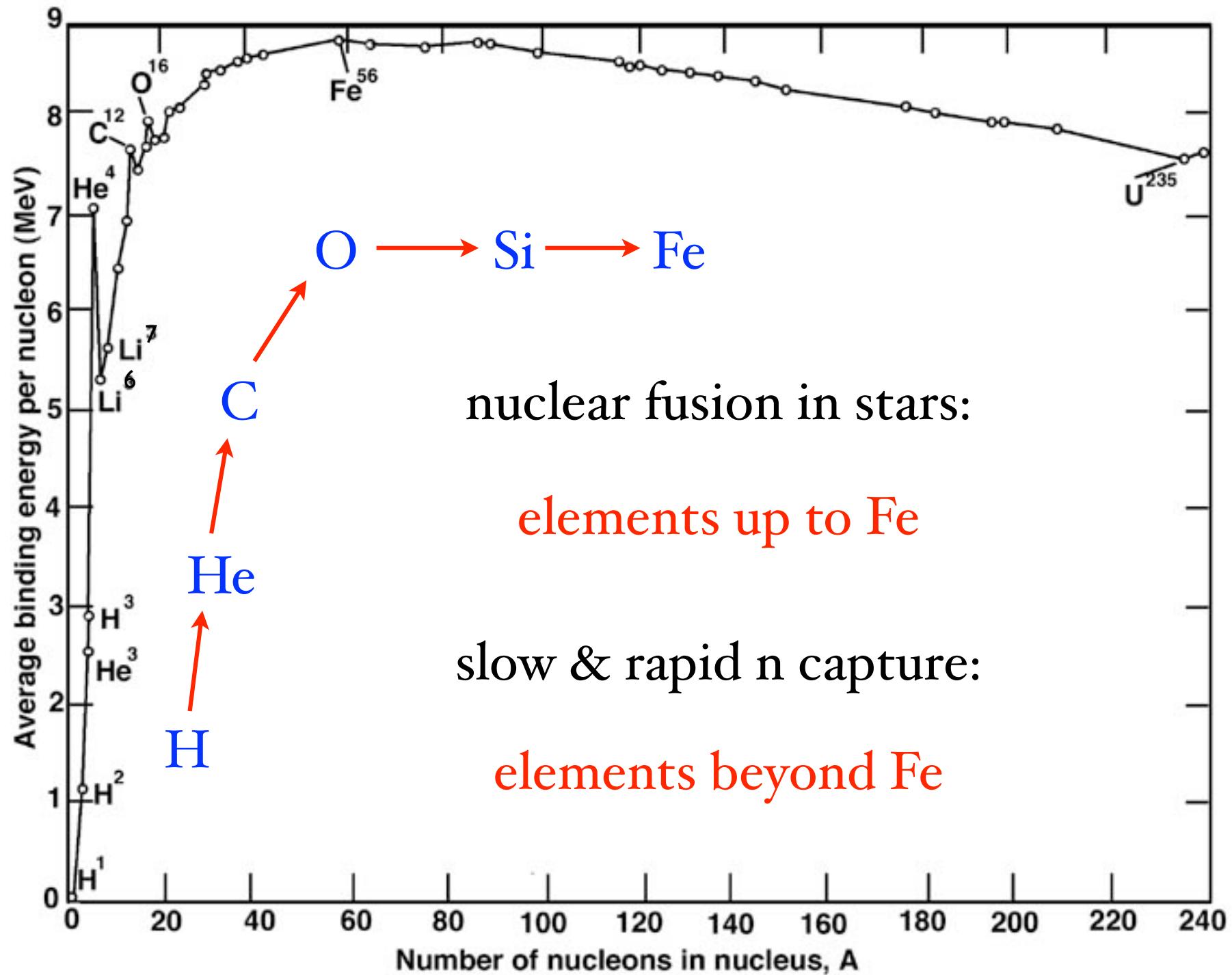


$$\frac{n}{p} = \exp\left(-\frac{M_n - M_p}{T}\right) < 1$$

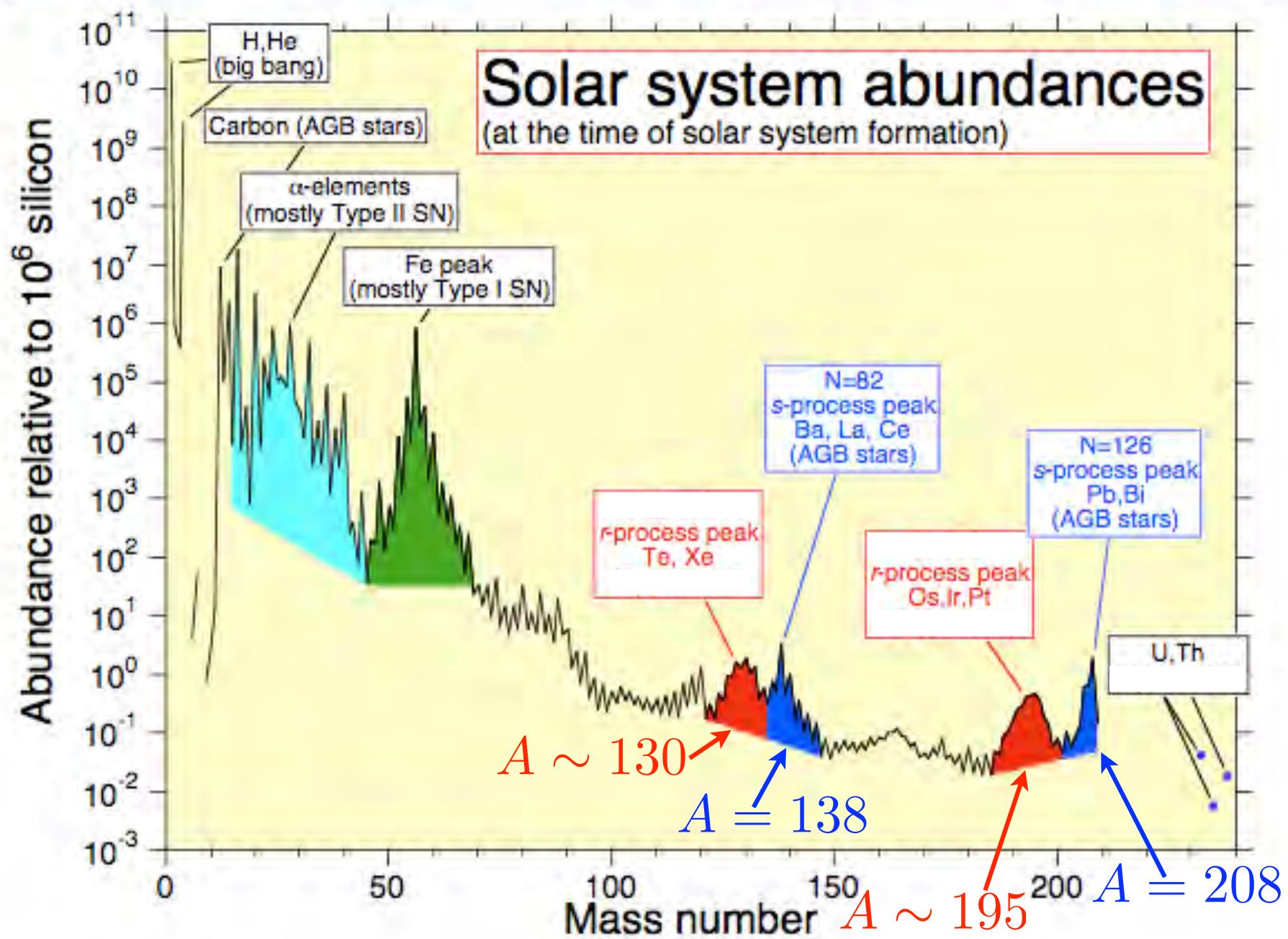
Big Bang:
75% H + 25% He
(by mass)

Sun:
71.1% H + 27.4% He
+ 1.5% “Metals”

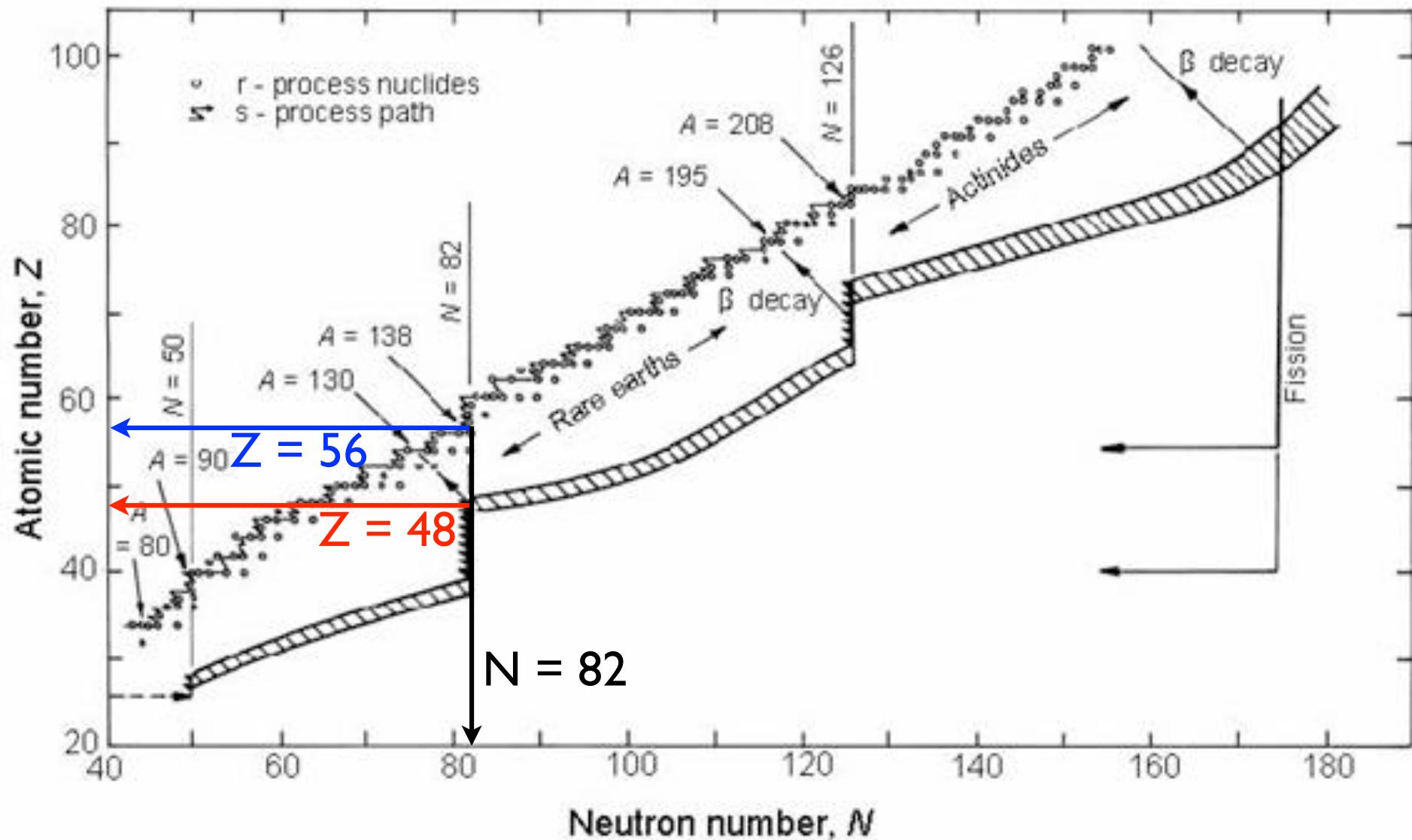




Cosmic Abundances

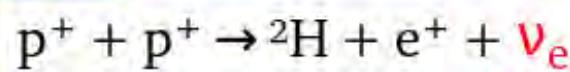


slow (s) and rapid (r) neutron capture processes

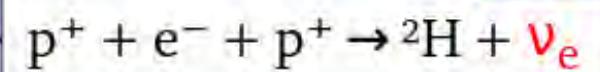


Solar Neutrinos (I)

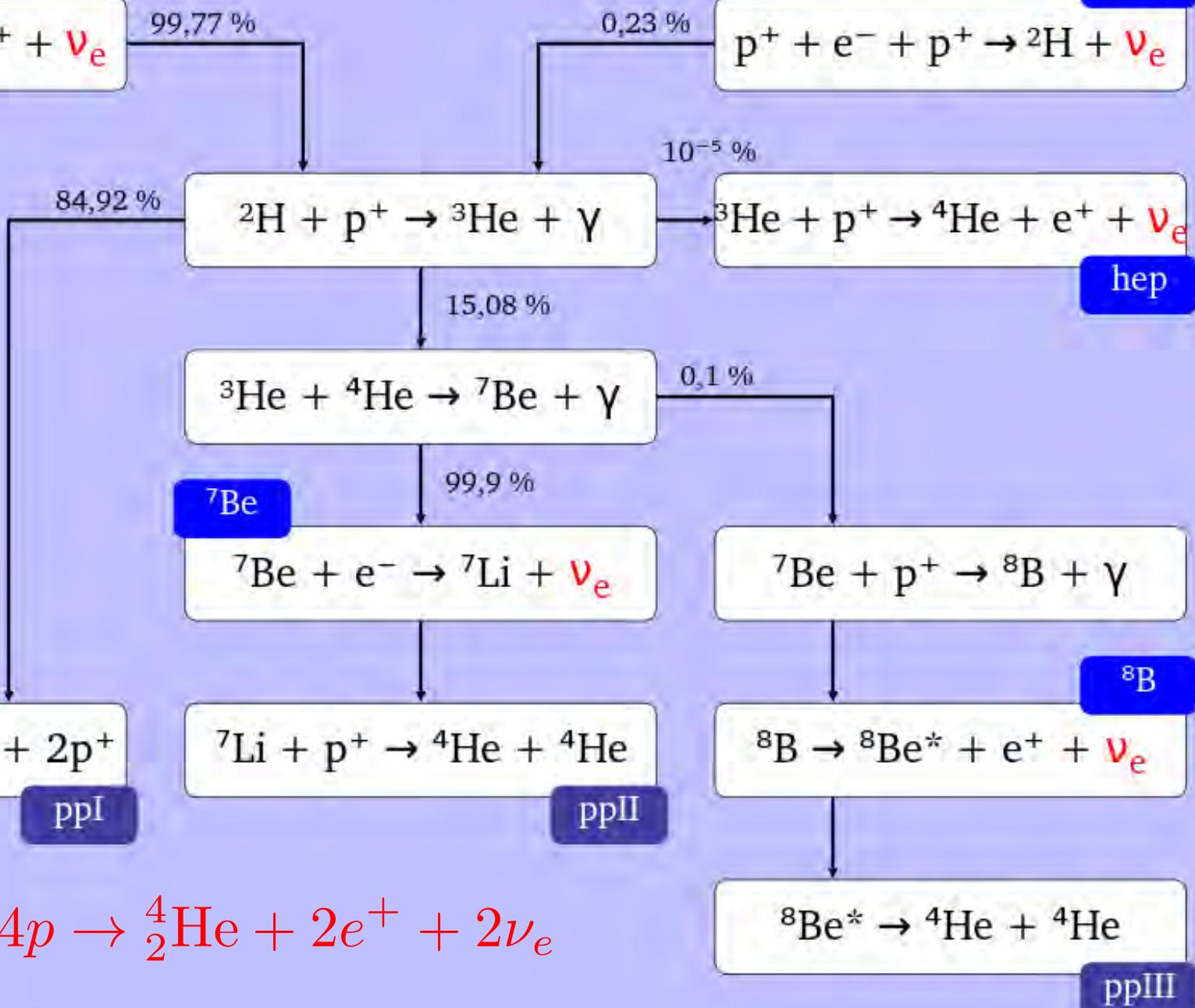
pp



pep

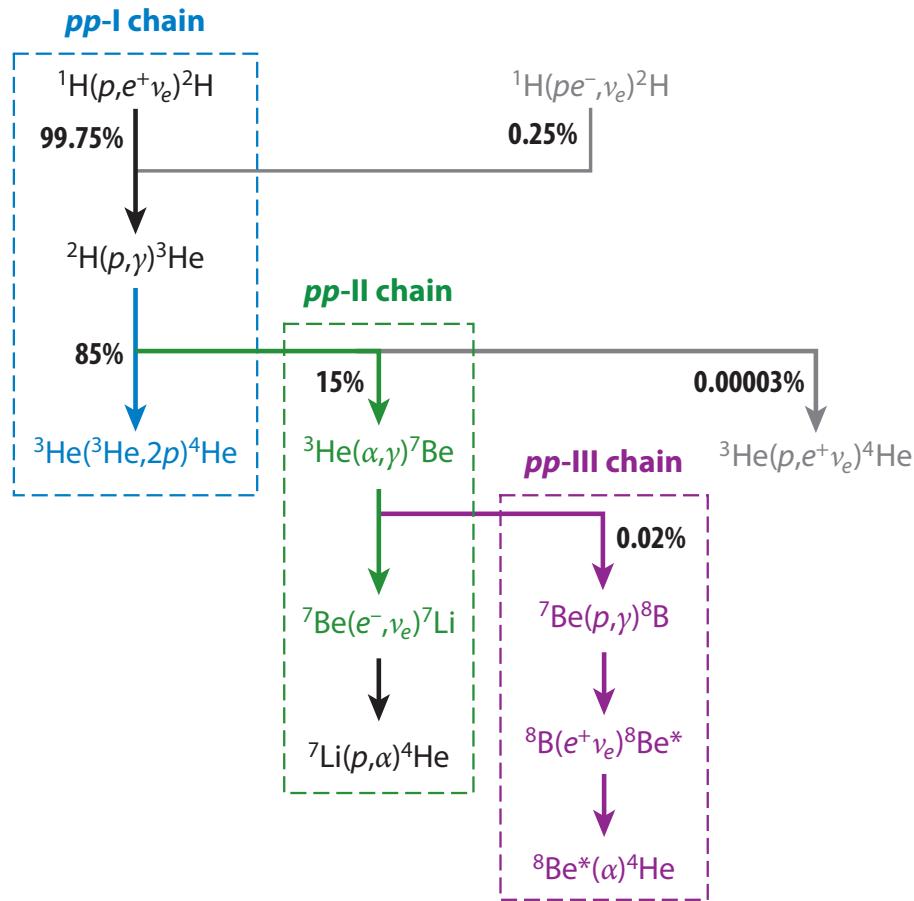


Bethe
(1967)
Davis
(2002)

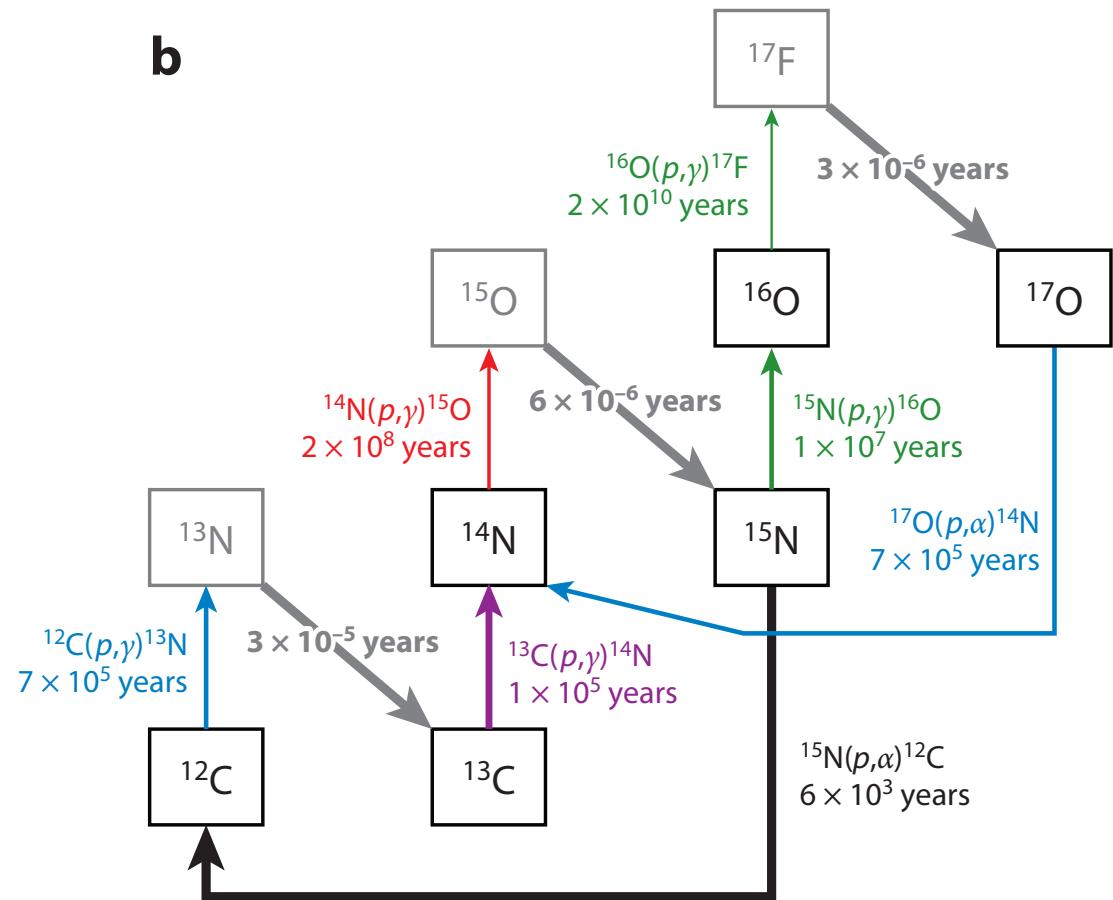


Solar Neutrinos (2)

a

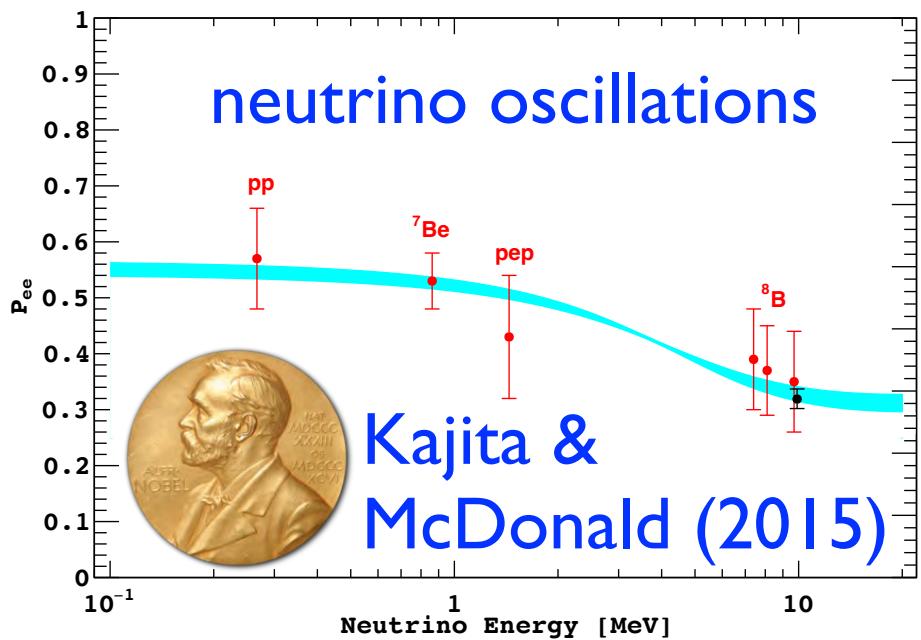
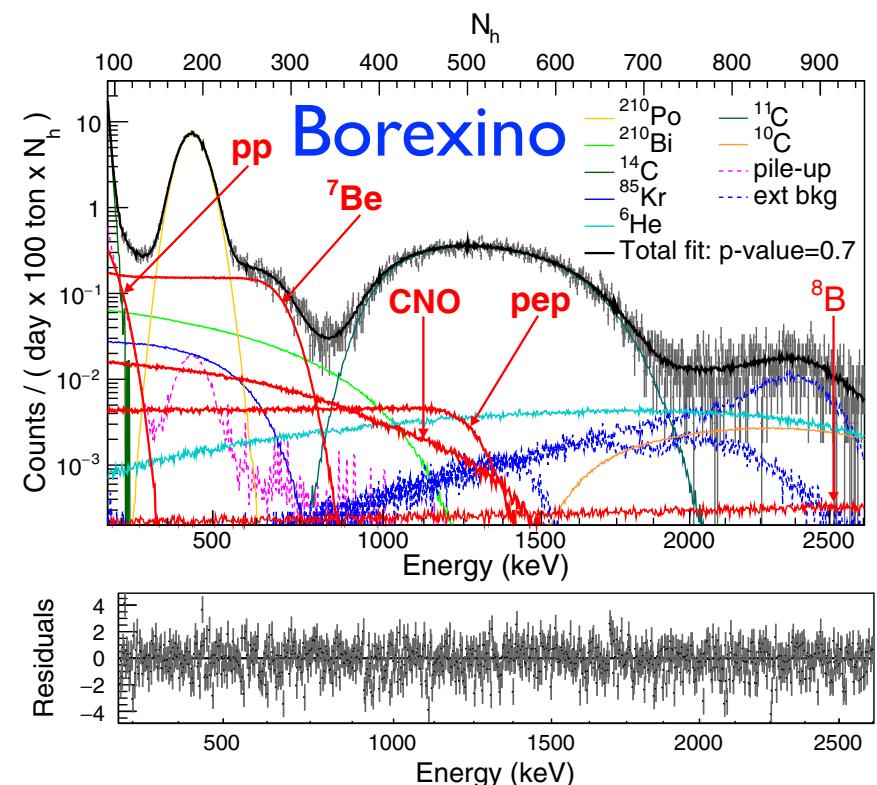
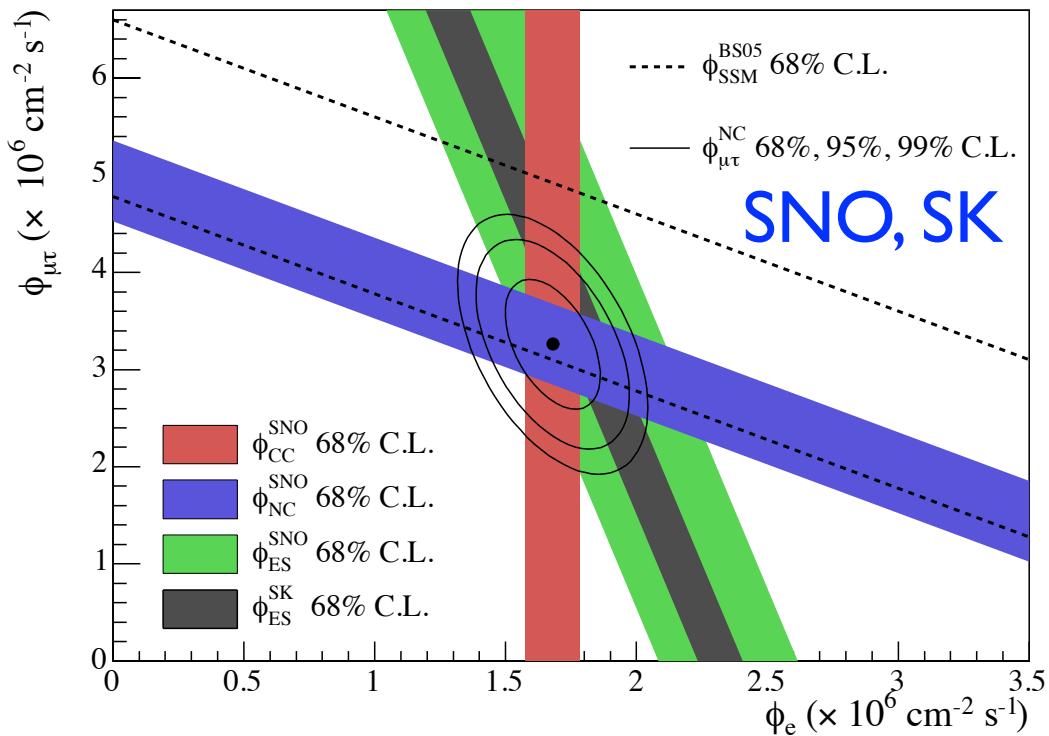
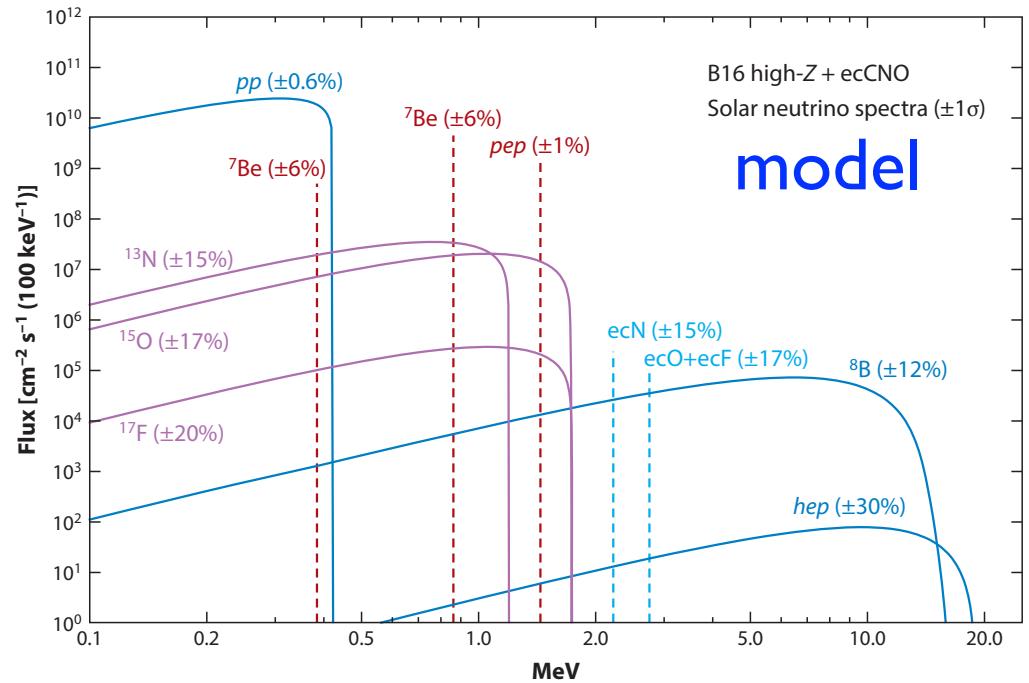


b

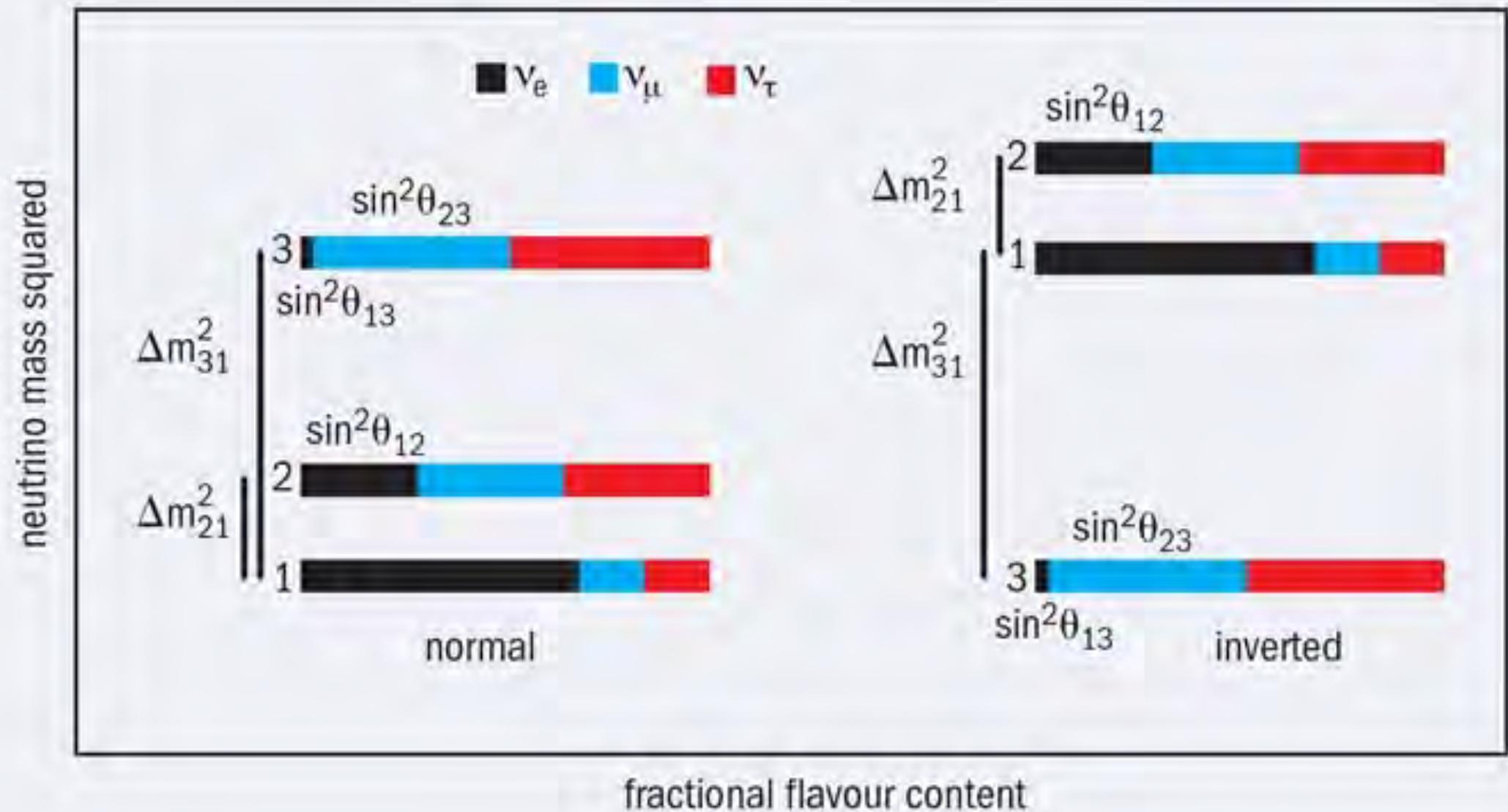


~99% of energy production

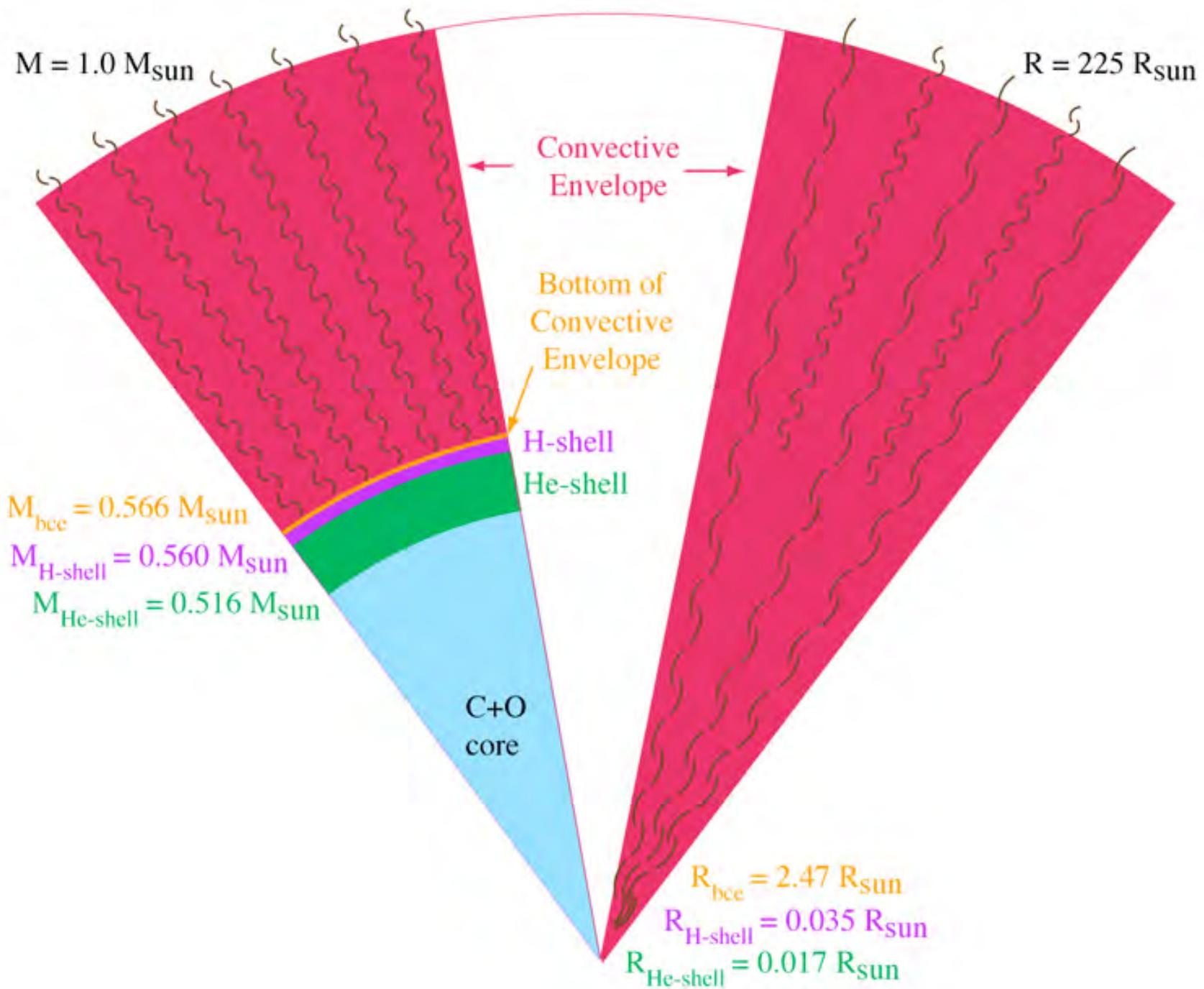
~1% of energy production



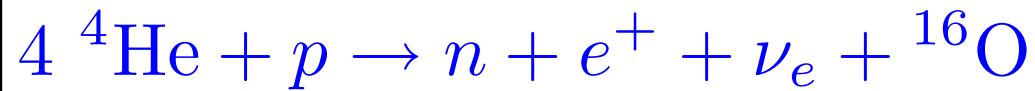
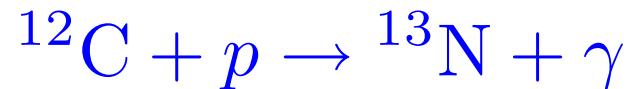
Solar, atmospheric, reactor, & accelerator neutrino experiments



Structure of an Asymptotic Giant Branch star



s-process in low- & intermediate-mass stars

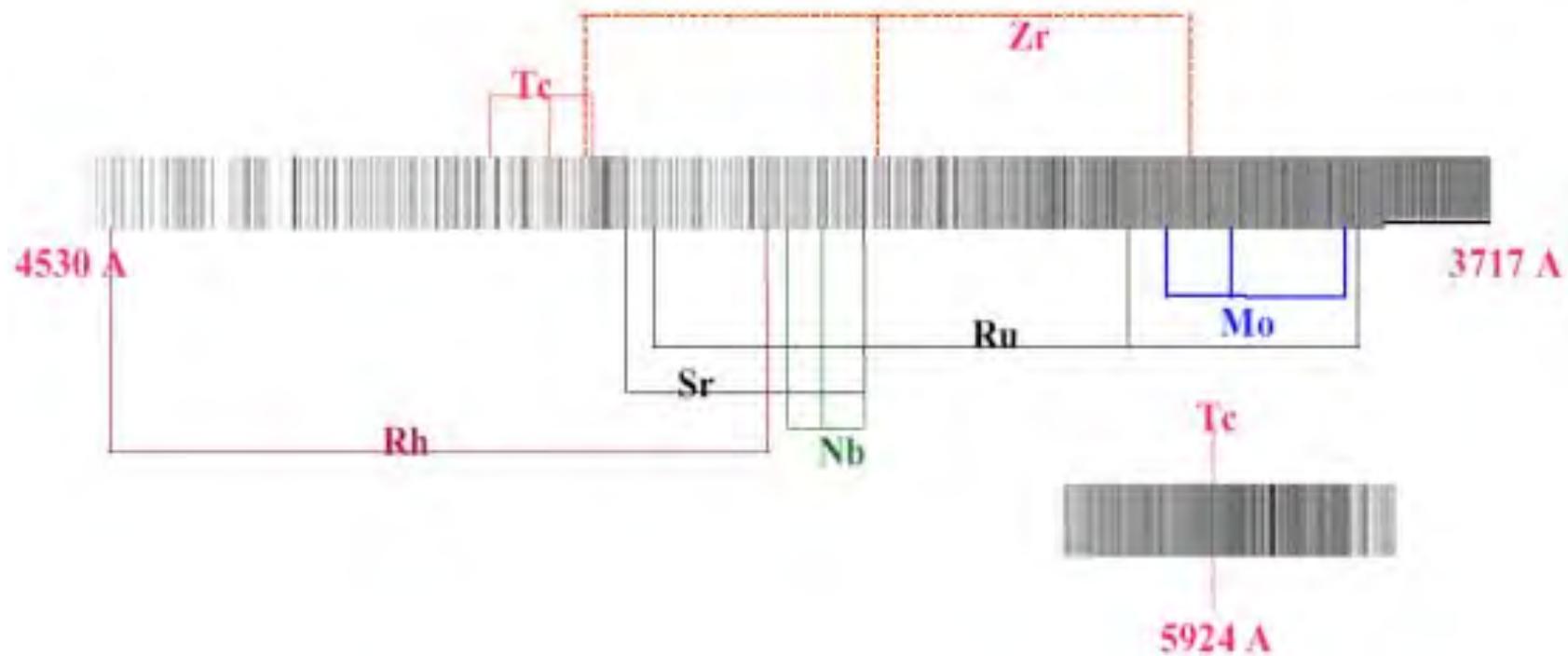


Evidence for Stellar s-Process Nucleosynthesis

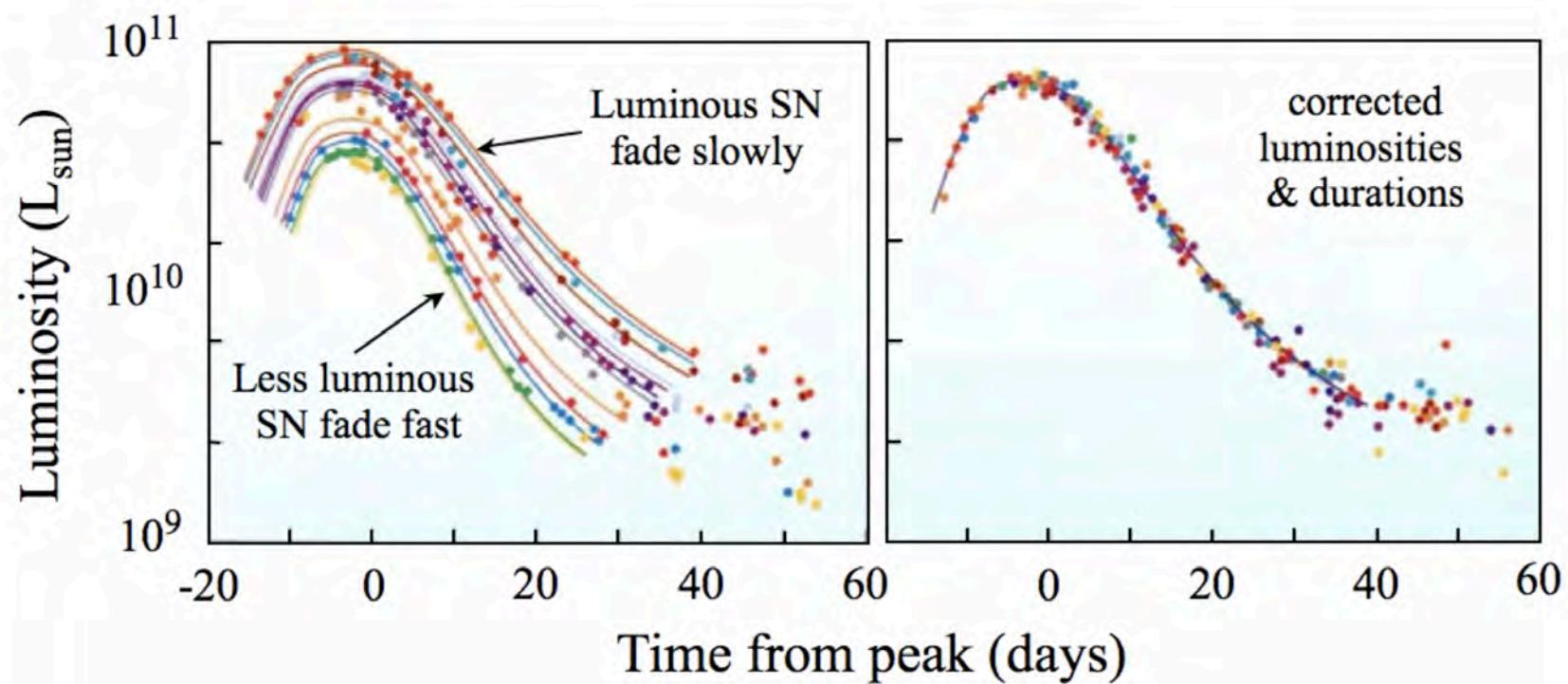
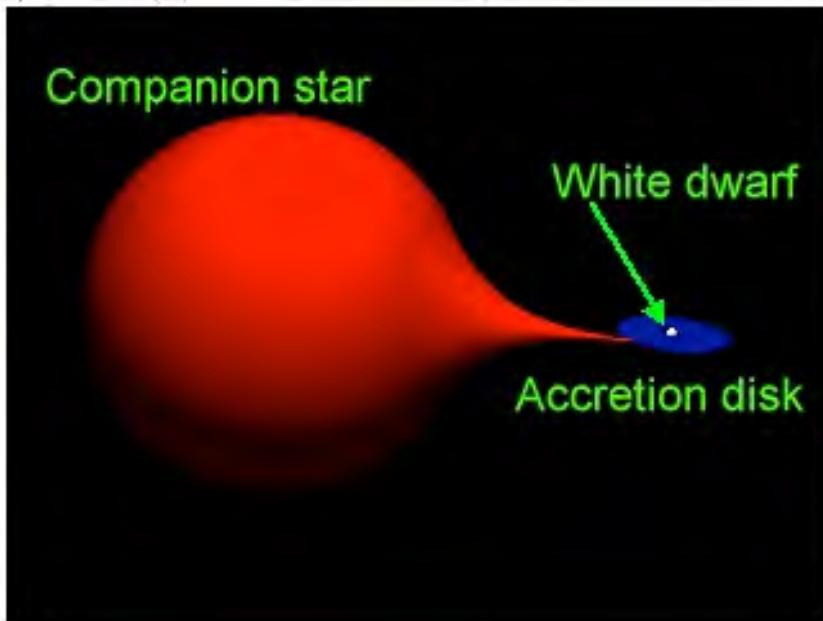
discovery of Tc spectral lines in stars by Merrill in 1952

Tc has no stable isotopes; the half-lives of the longest-lived isotopes are:

2.6×10^6 , 4.2×10^6 , 2.1×10^5 yr for ^{97}Tc , ^{98}Tc , ^{99}Tc



Type Ia Supernovae: Accreting and Merging White Dwarfs



Simple description of stellar structure

$$\frac{dP}{dr} = -\frac{Gm(r)\rho(r)}{r^2}$$

$$\int_0^R \frac{dP}{dr} 4\pi r^3 dr = - \int_0^R \frac{Gm(r)\rho(r)}{r^2} 4\pi r^3 dr$$

$$-3 \int_0^R P(r) 4\pi r^2 dr = - \int_0^R \frac{Gm(r)\rho(r)}{r} 4\pi r^2 dr$$

virial
theorem

$$-3(\gamma - 1)E_K = E_G \Rightarrow E_K = -\frac{E_G}{3(\gamma - 1)}$$

$$P = (\gamma - 1)\mathcal{E}_K = \begin{cases} (2/3)\mathcal{E}_K, & \gamma = 5/3 \text{ (NR)} \\ (1/3)\mathcal{E}_K, & \gamma = 4/3 \text{ (ER)} \end{cases}$$

$$E_{\text{tot}} = E_K + E_G = \frac{3\gamma - 4}{3(\gamma - 1)} E_G$$

Virial theorem for a contracting gas cloud

$$T_c + \frac{\hbar^2}{2m_e d^2} \sim \frac{GMm_p}{R}, \quad \left(\frac{M}{m_p}\right) d^3 \sim R^3$$

$$T_c \sim \frac{GMm_p}{R} - \frac{\hbar^2}{2m_e} \left(\frac{M}{m_p}\right)^{2/3} \frac{1}{R^2} \Rightarrow T_{c,\max} \propto M^{4/3}$$

massive stars reach much higher temperatures during contraction

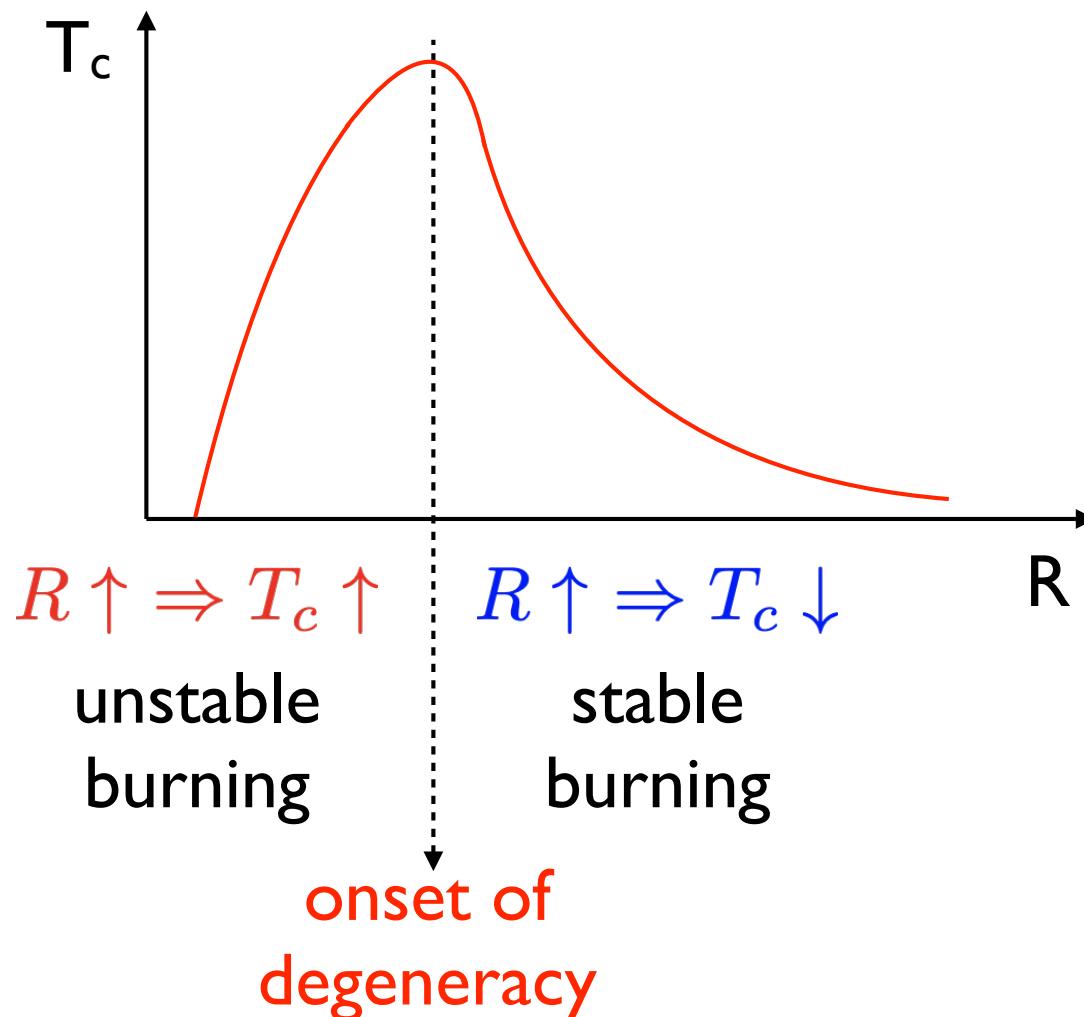
before degeneracy sets in: $T_c \propto \frac{M}{R}$, $\rho_c \propto \frac{M}{R^3} \Rightarrow \frac{T_c^3}{\rho_c} \propto M^2$

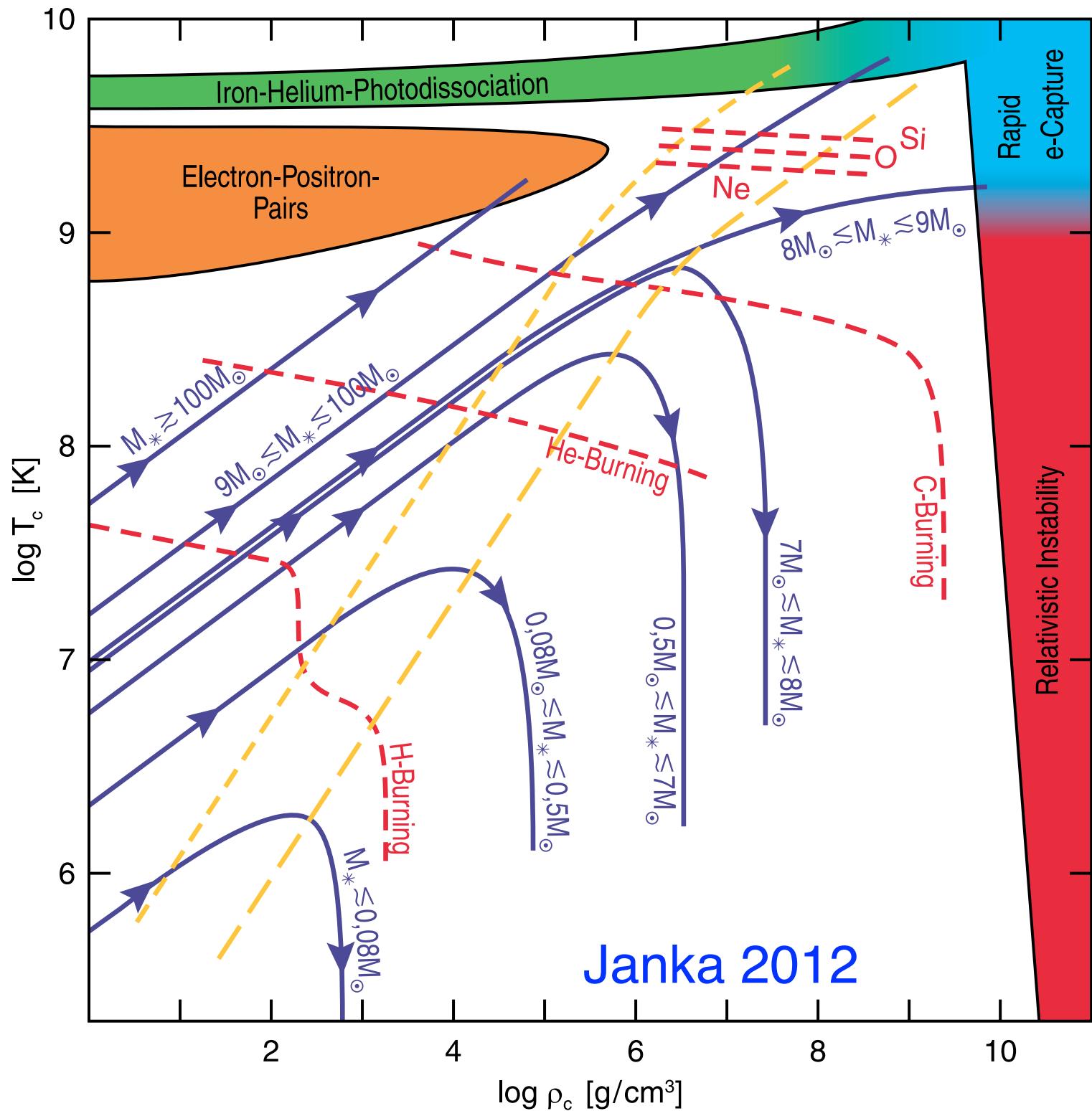
radiation pressure becomes important for $M \gtrsim 10 M_\odot$

$$E_{\text{tot}} = \frac{3\gamma - 4}{3(\gamma - 1)} E_G \sim -\frac{GM^2}{R} \text{ for } \gamma > \frac{4}{3}$$

Stability of nuclear burning

fluctuation $\delta T_c > 0 \Rightarrow E_{\text{tot}} \sim -\frac{GM^2}{R} \uparrow \Rightarrow R \uparrow$



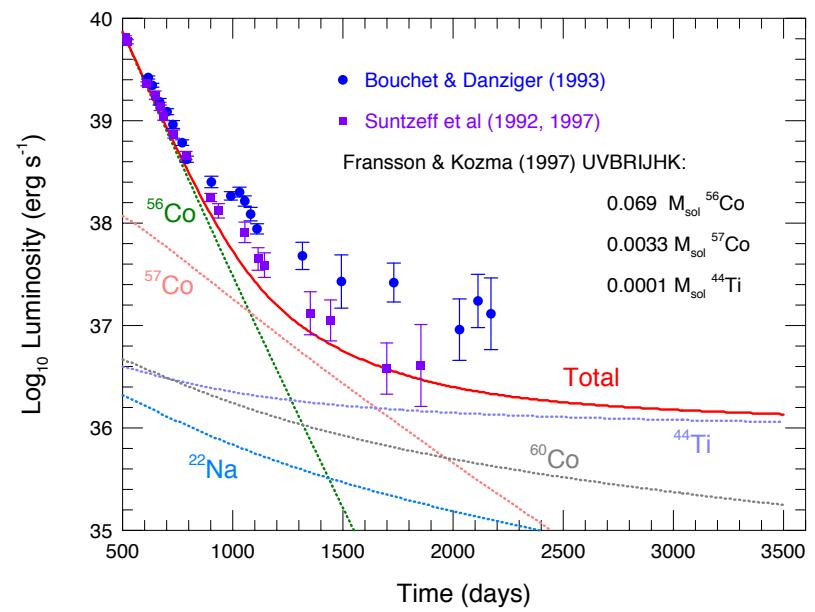
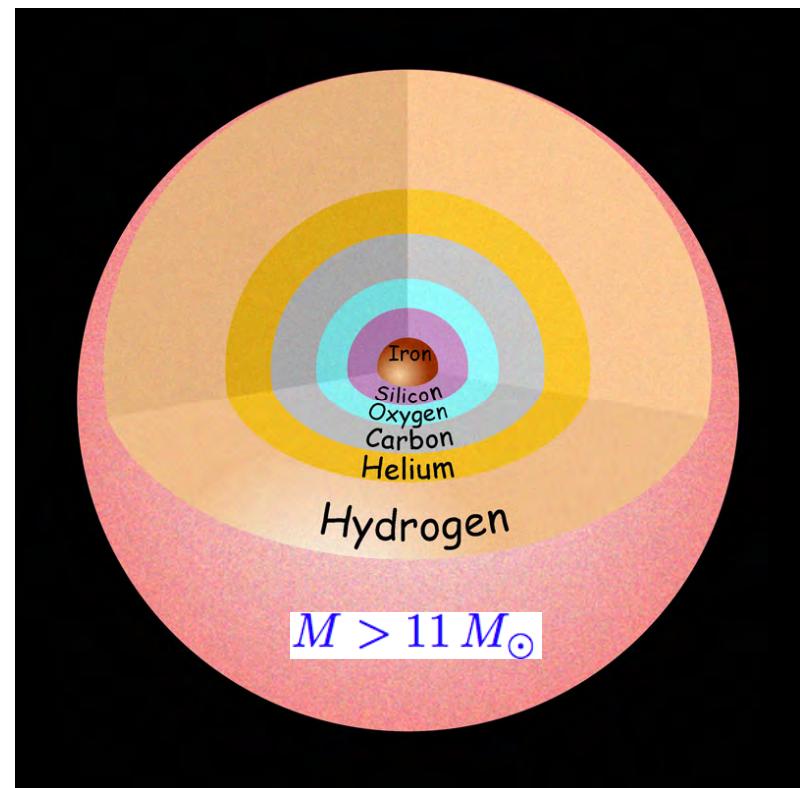
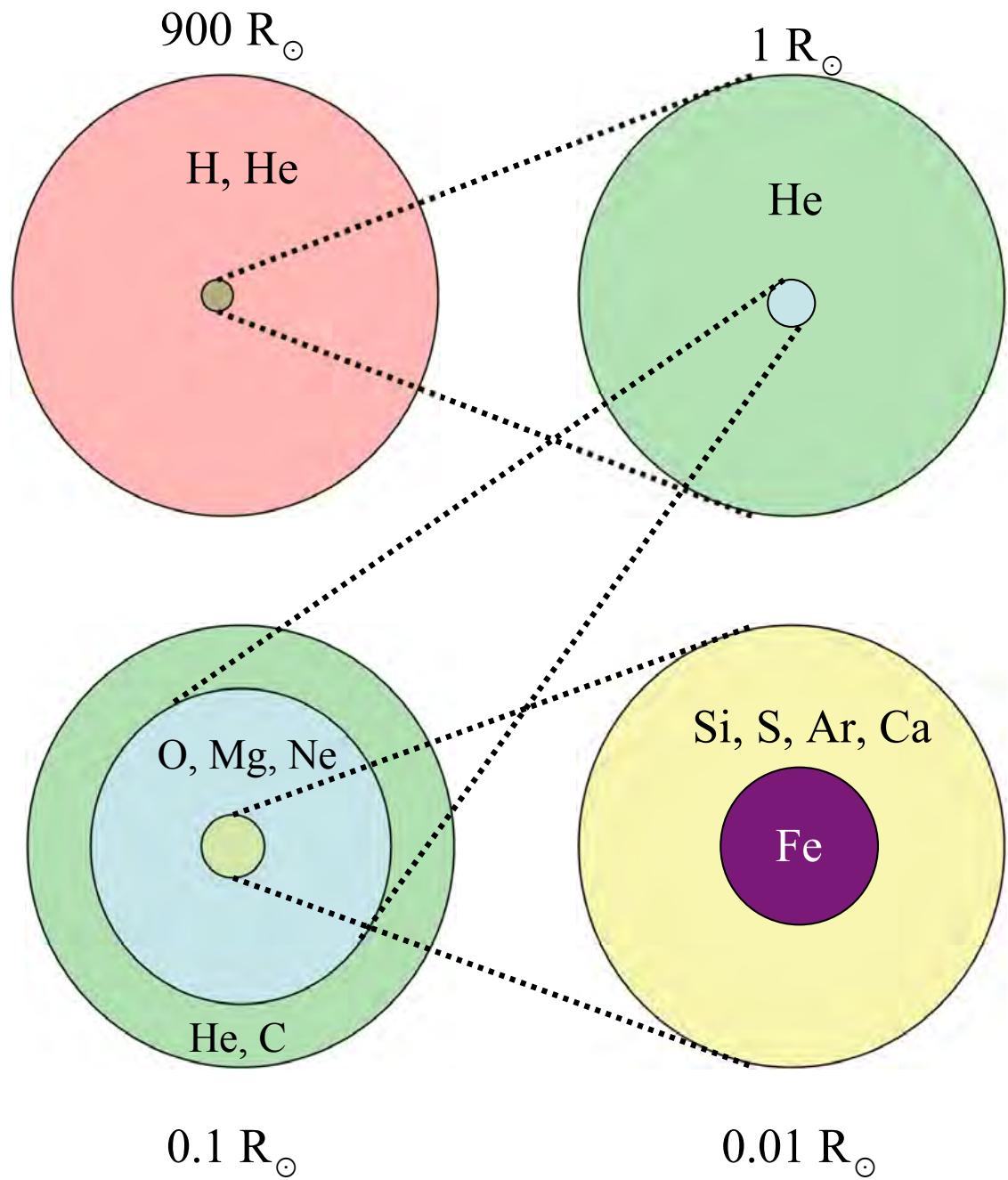


Nuclear burning stages

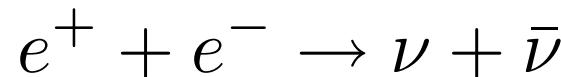
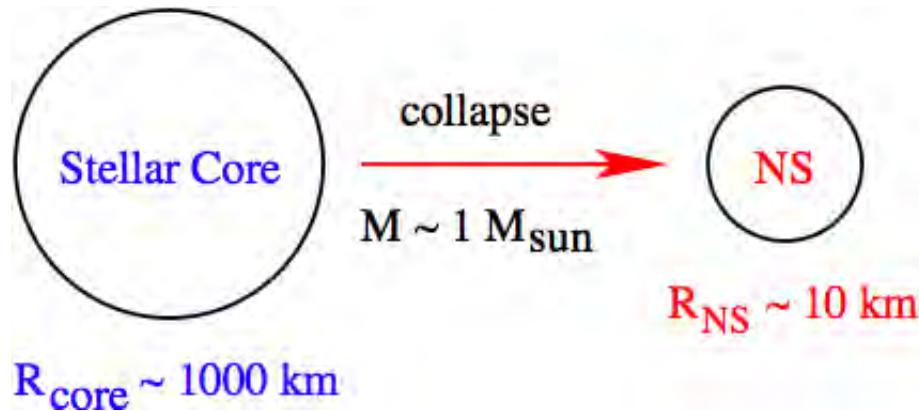
($20 M_{\odot}$ stars)

Fuel	Main Product	Secondary Product	T (10^9 K)	Time (yr)	Main Reaction
H	He	^{14}N	0.02	10^7	$4 \text{ H} \xrightarrow{\text{CNO}} {}^4\text{He}$
He	O, C	^{18}O , ^{22}Ne s-process	0.2	10^6	$3 \text{ He}^4 \rightarrow {}^{12}\text{C}$ ${}^{12}\text{C}(\alpha, \gamma){}^{16}\text{O}$
C	Ne, Mg	Na	0.8	10^3	${}^{12}\text{C} + {}^{12}\text{C}$
Ne	O, Mg	Al, P	1.5	3	${}^{20}\text{Ne}(\gamma, \alpha){}^{16}\text{O}$ ${}^{20}\text{Ne}(\alpha, \gamma){}^{24}\text{Mg}$
O	Si, S	Cl, Ar, K, Ca	2.0	0.8	${}^{16}\text{O} + {}^{16}\text{O}$
Si, S	Fe	Ti, V, Cr, Mn, Co, Ni	3.5	0.02	${}^{28}\text{Si}(\gamma, \alpha)\dots$

25 M_{\odot} Presupernova Star

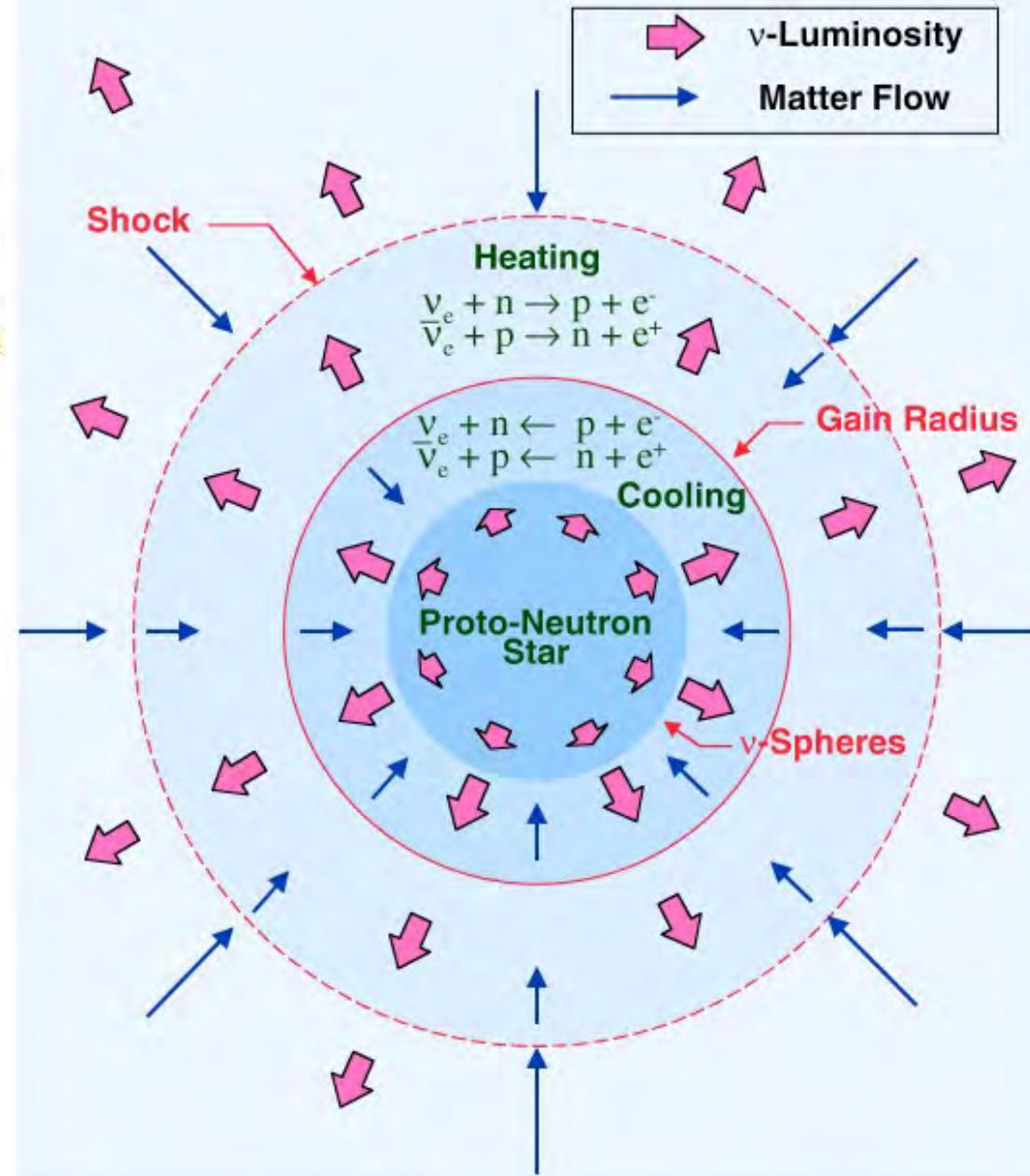


Supernovae as a neutrino phenomenon

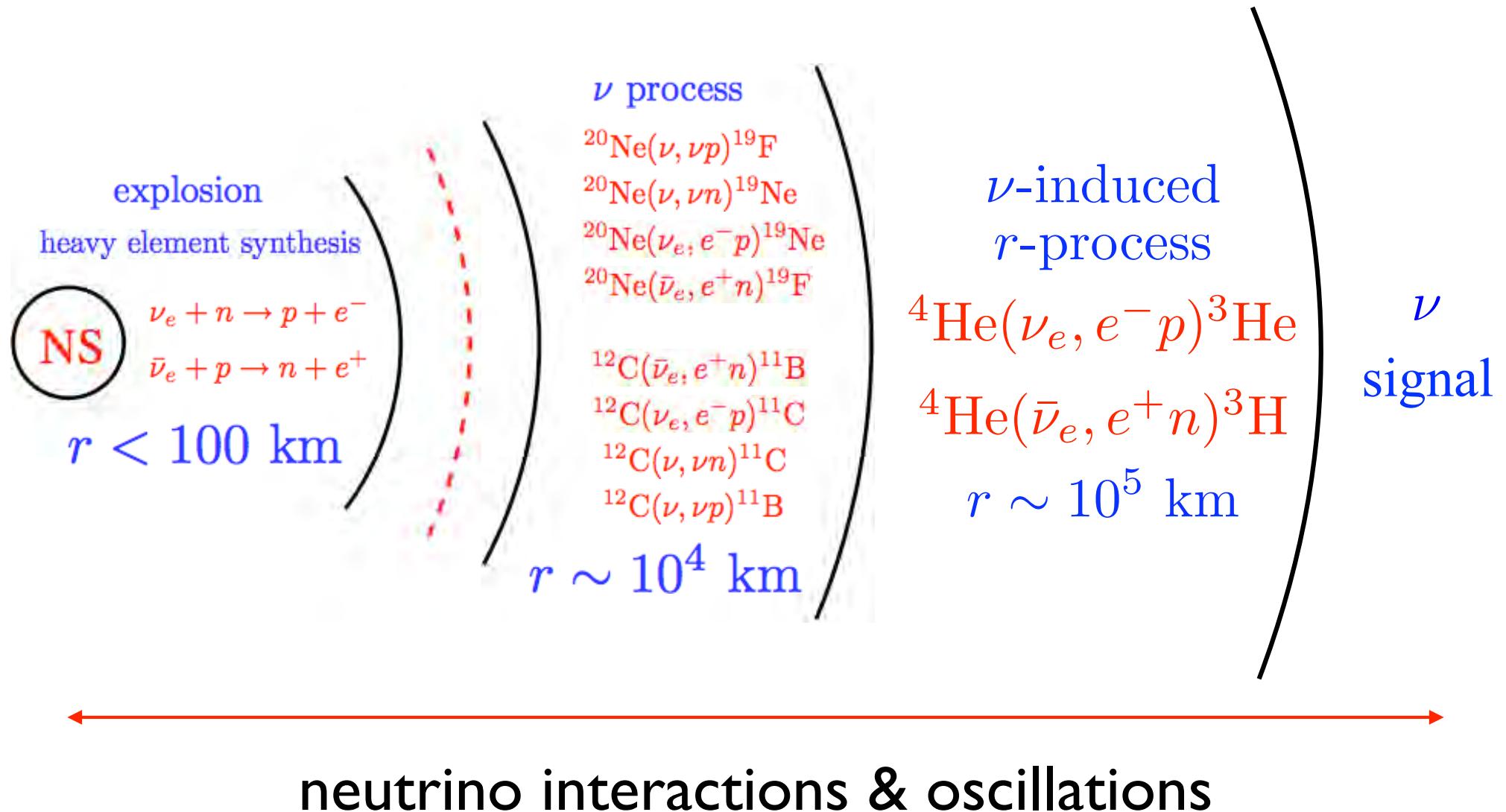


$$\frac{GM^2}{R_{\text{NS}}} \sim 3 \times 10^{53} \text{ erg}$$

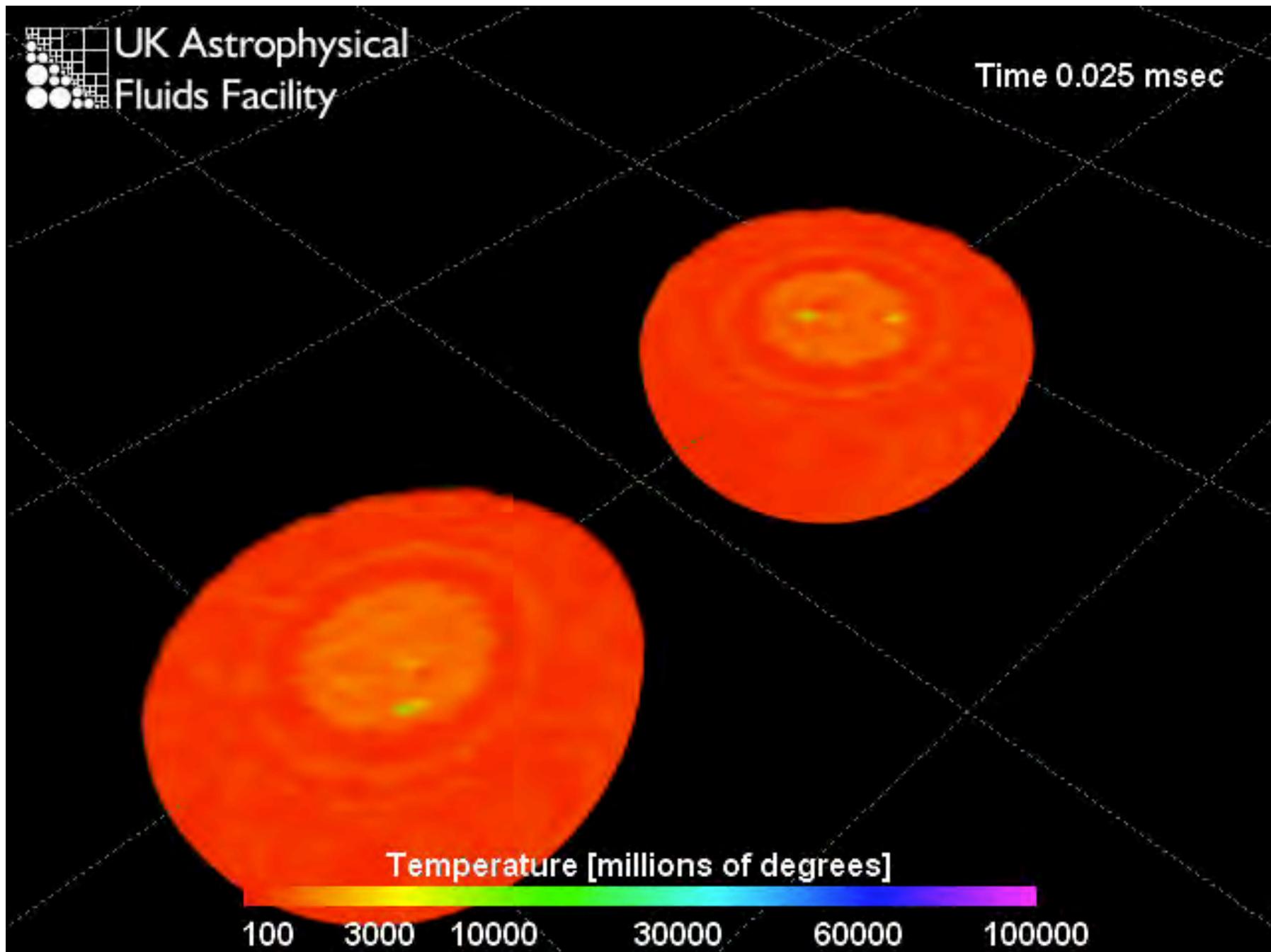
$\Rightarrow \nu_e, \bar{\nu}_e, \nu_\mu, \bar{\nu}_\mu, \nu_\tau, \bar{\nu}_\tau$



neutrino processes in typical core-collapse supernovae



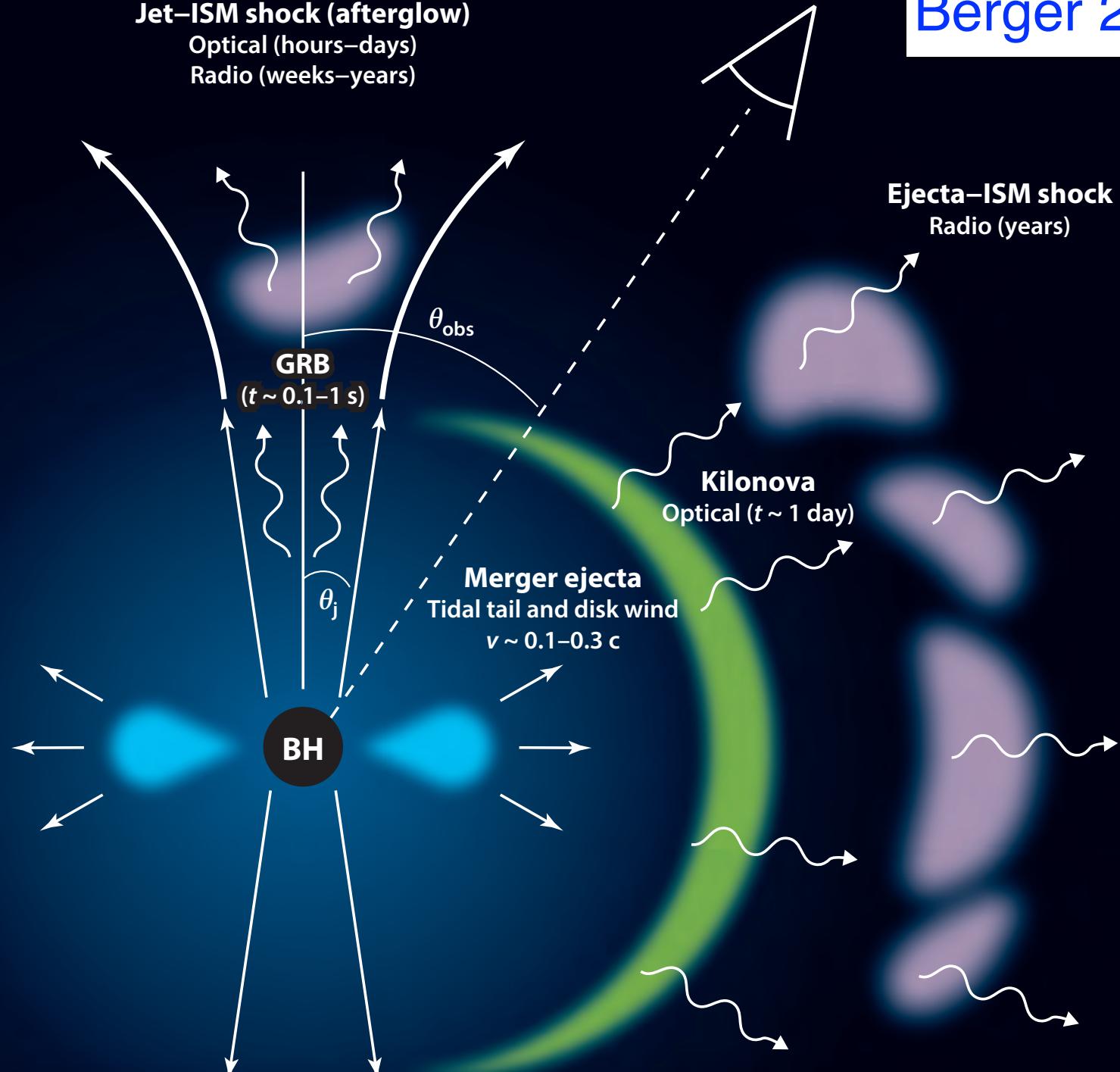
r-process in ejecta from neutron star merger



Jet-ISM shock (afterglow)

Optical (hours–days)

Radio (weeks–years)



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1 Atomic # Sybmld Name Atomic Mass	2 C Solid	3	4 Hg Liquid	5	6 H Gas	7	8	9	10	11	12	13	14	15	16	17	18																																																					
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THE BIG BANG

INFLATION

COSMIC MICROWAVE
BACKGROUND
400,000 YEARS AFTER
BIG BANG

THE DARK AGES

FIRST STARS
400,000,000 YEARS
AFTER BIG BANG

GALAXY EVOLUTION
CONTINUES...

FIRST GALAXIES
1000,000,000 YEARS
AFTER BIG BANG

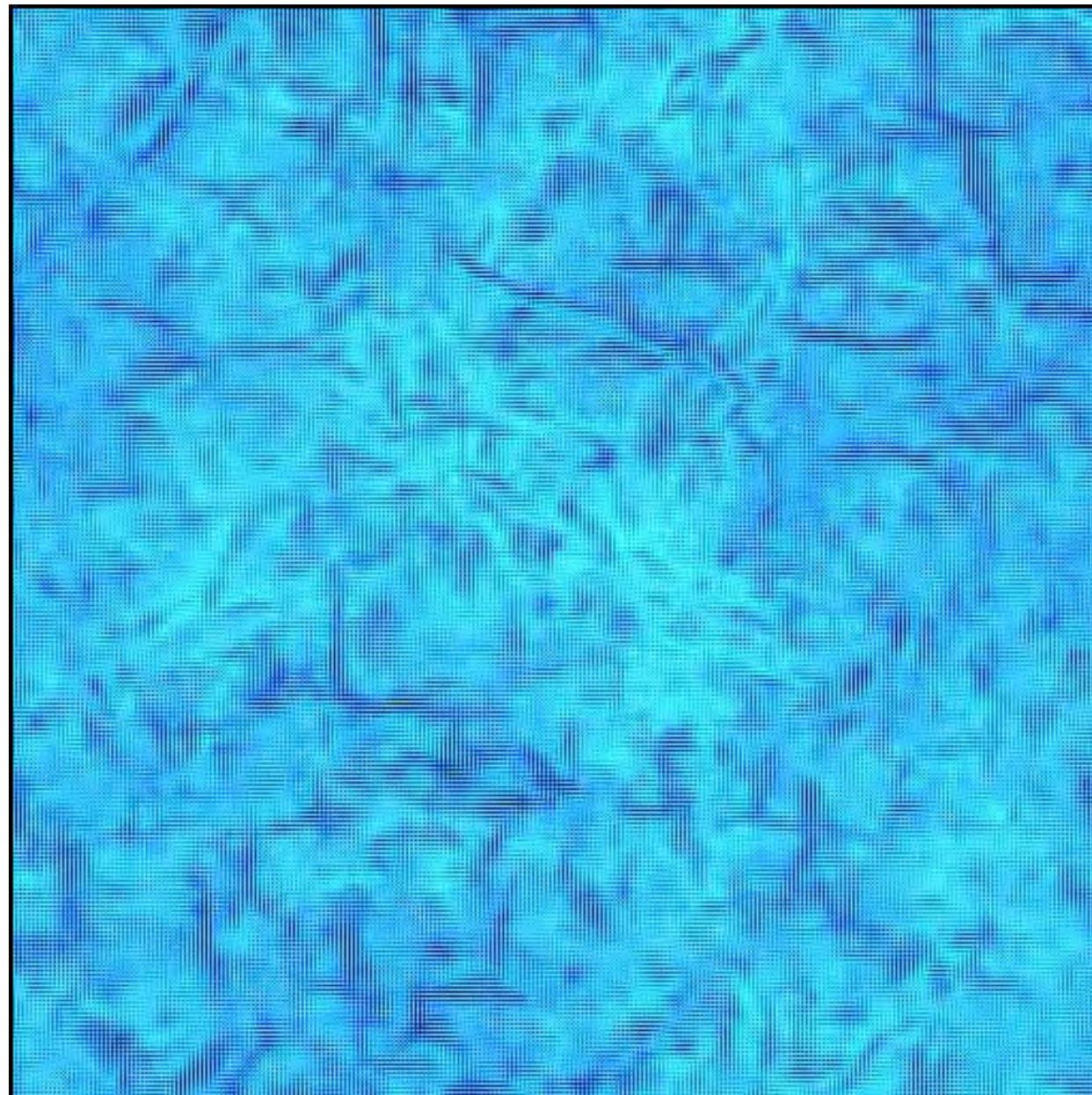
FORMATION OF
THE SOLAR SYSTEM
8,700,000,000 YEARS
AFTER BIG BANG

NOW
13,700,000,000 YEARS
AFTER BIG BANG

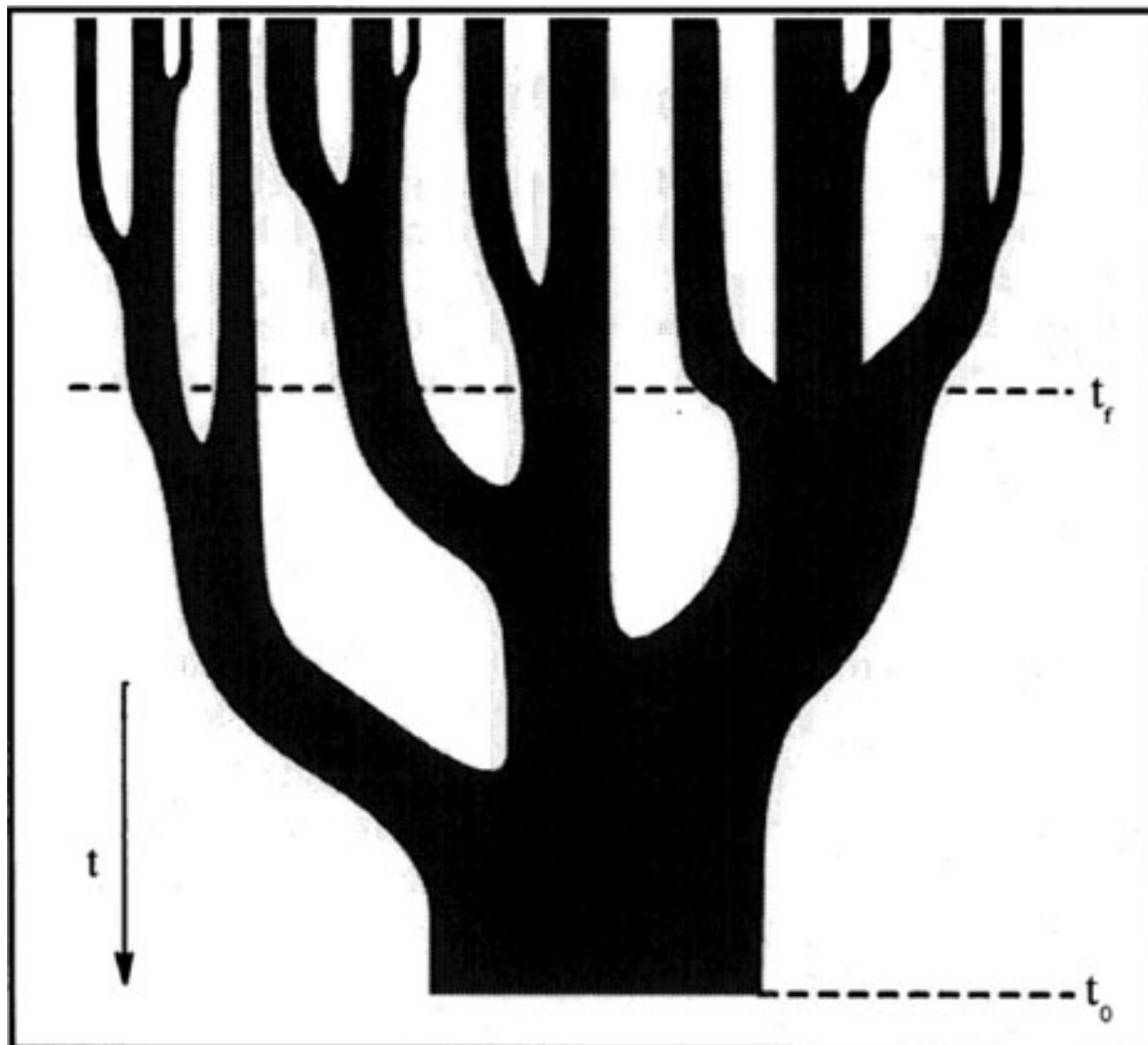
DARK ENERGY?

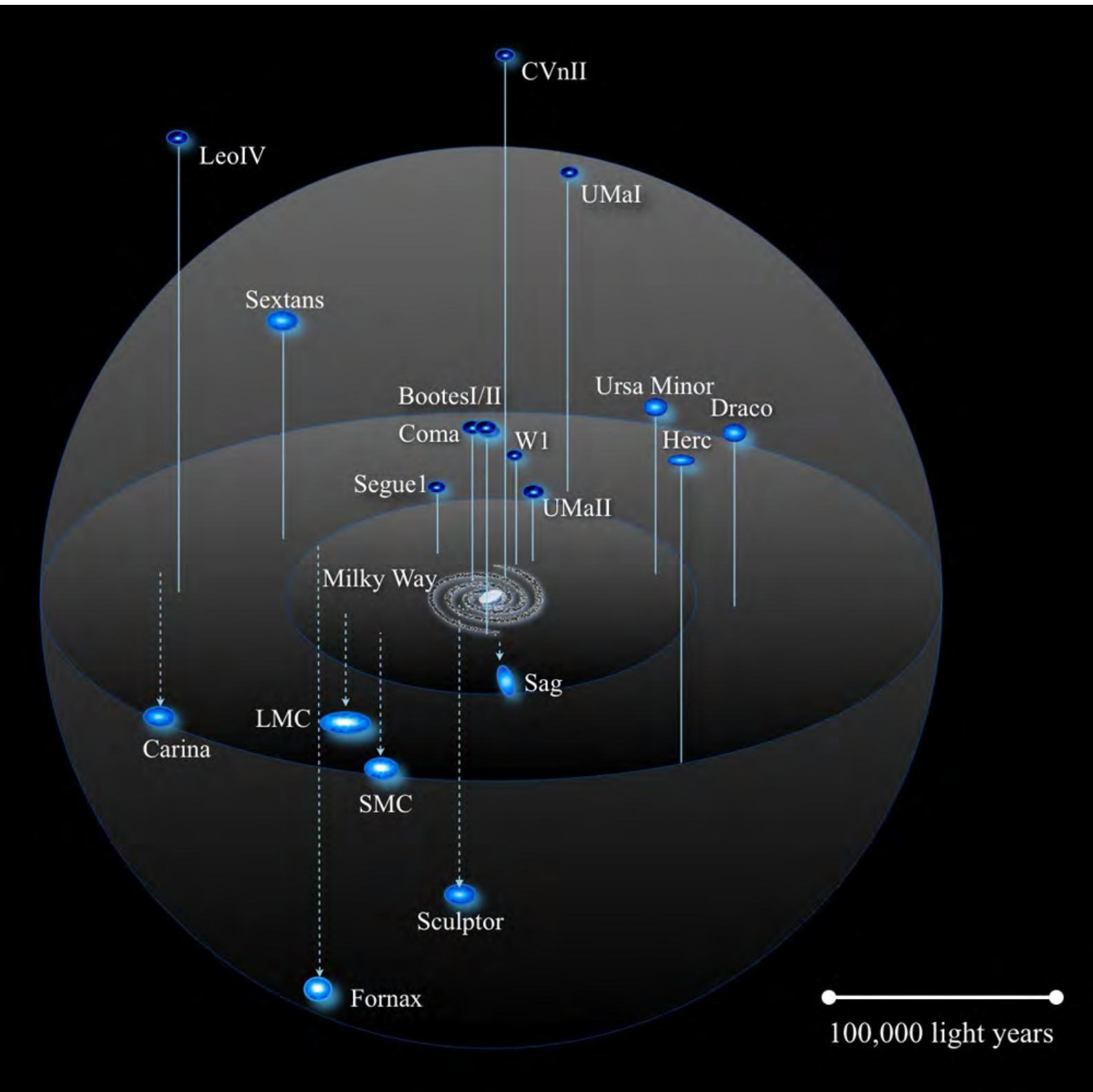


Hierarchical Structure Formation

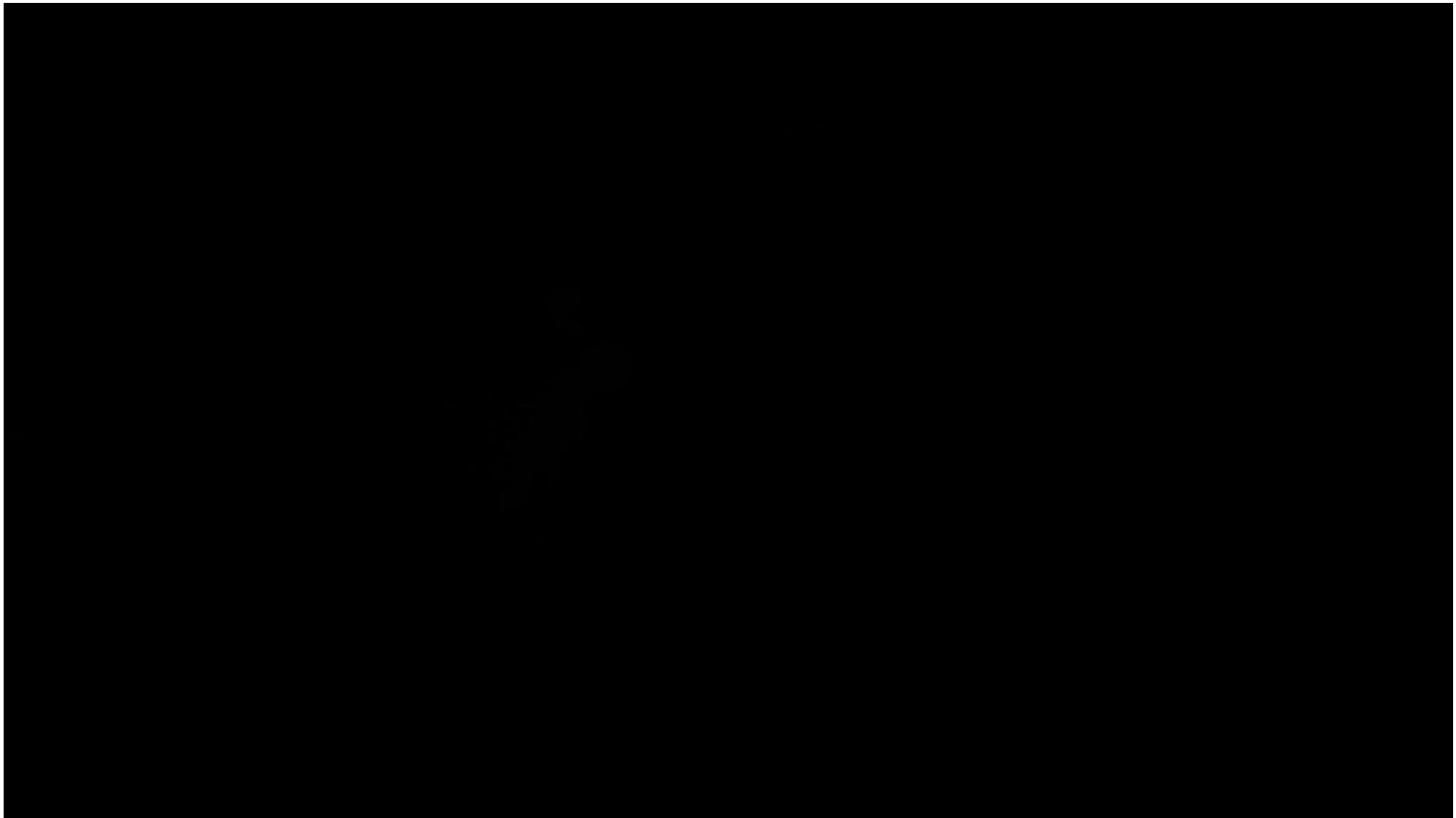


Merger Tree

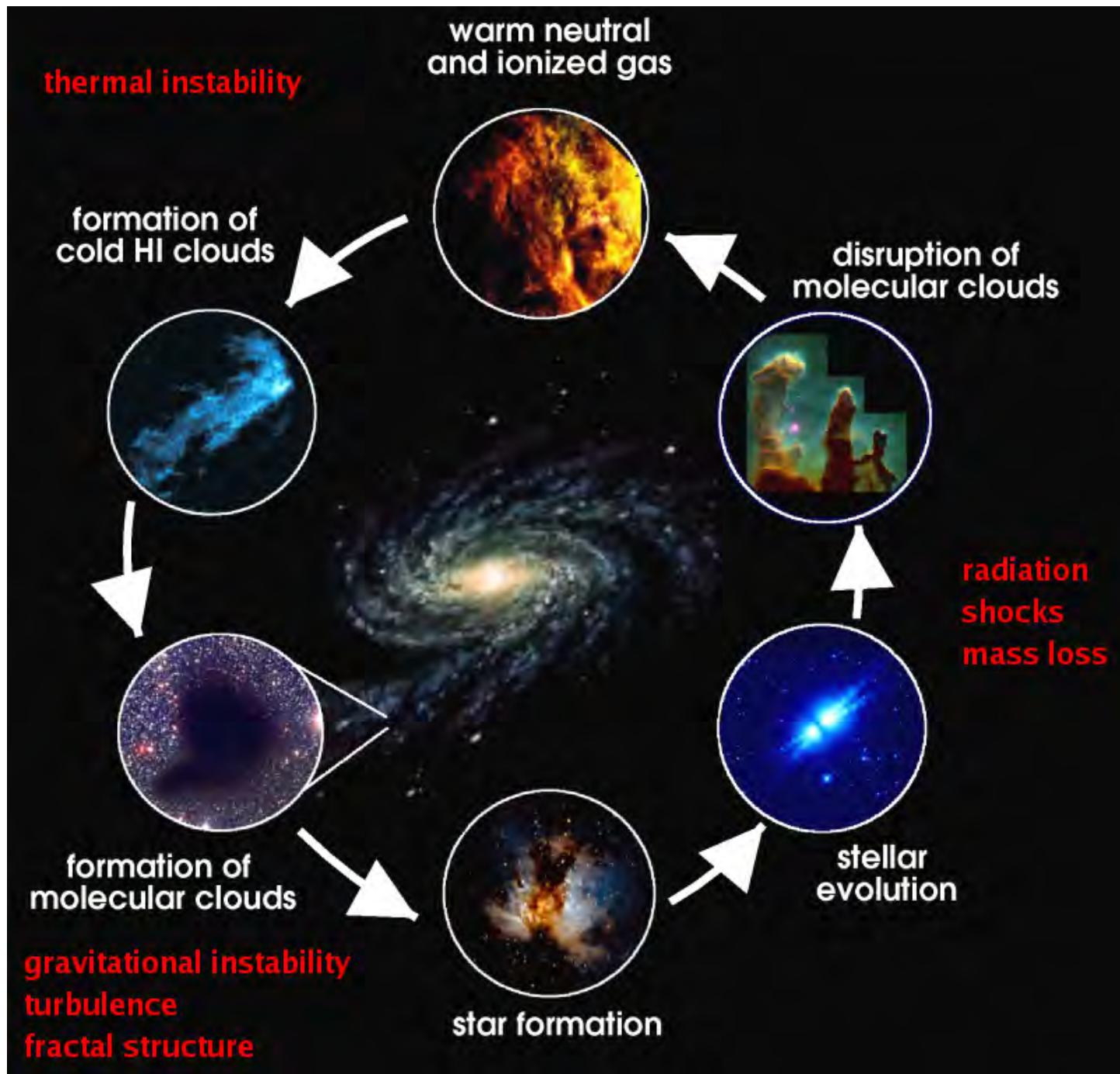




Galaxy formation with gas dynamics

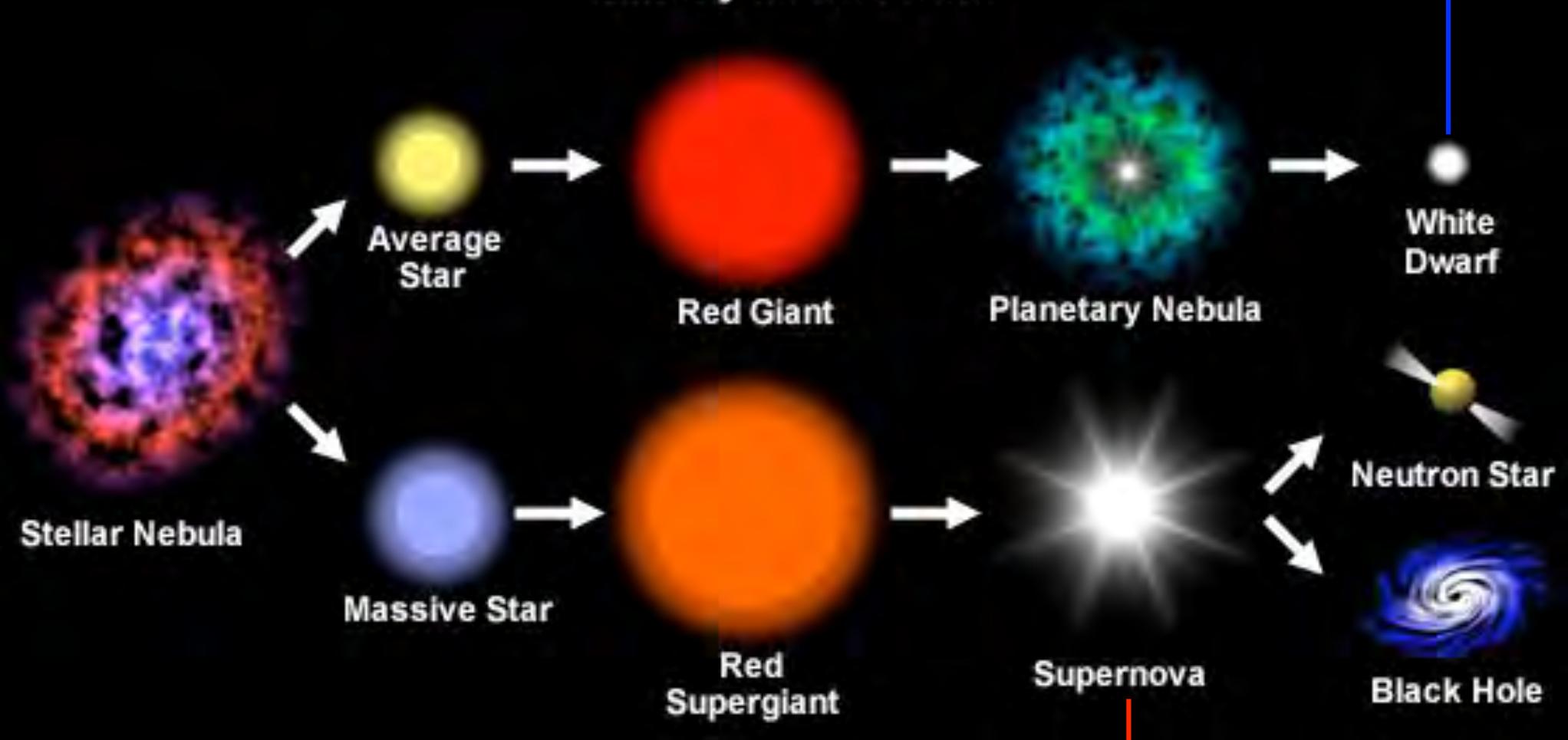


Life Cycle of Interstellar Medium



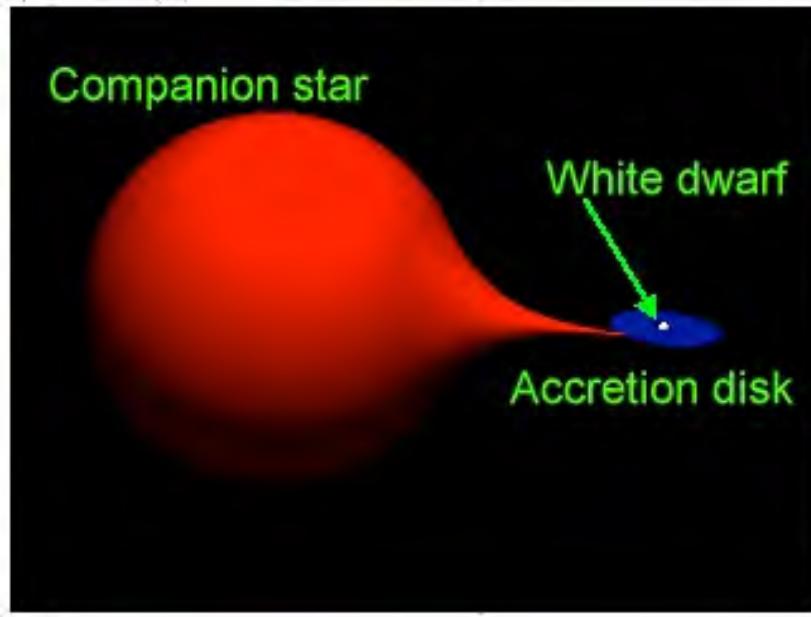
Type Ia SNe

Life Cycle of a Star

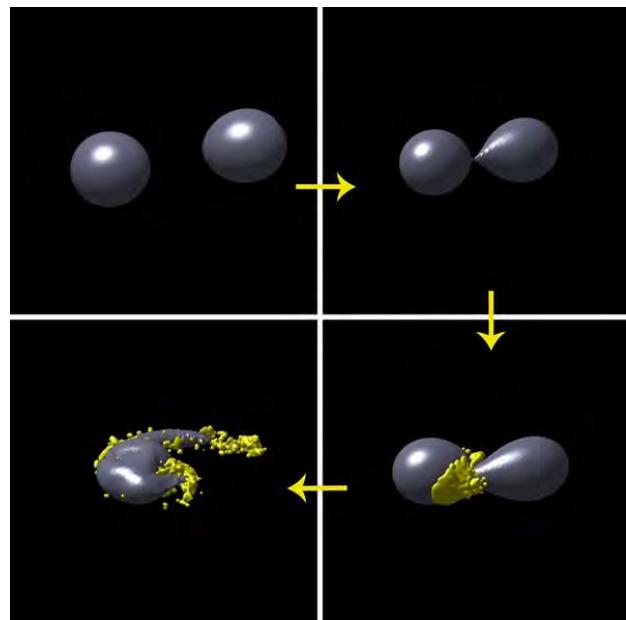


core-collapse SNe (mostly Type II)
~ 10 Myr

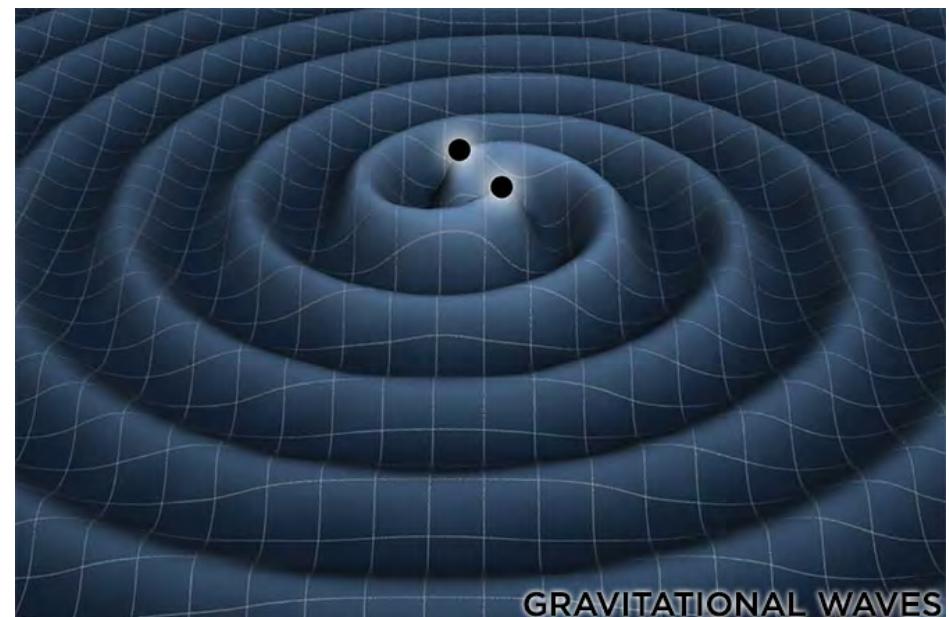
SNe Ia: Accreting and Merging White Dwarfs

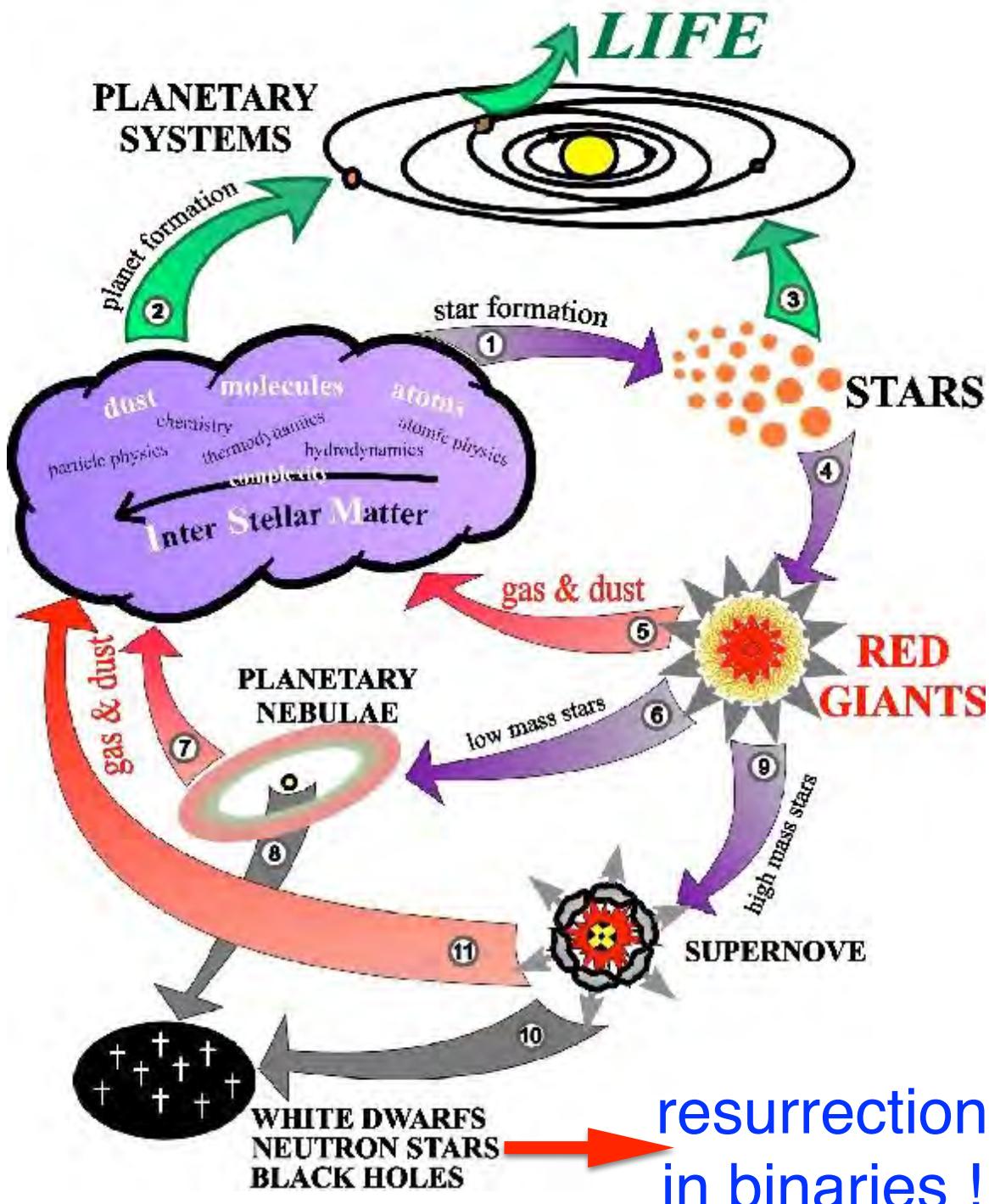


~1 Gyr

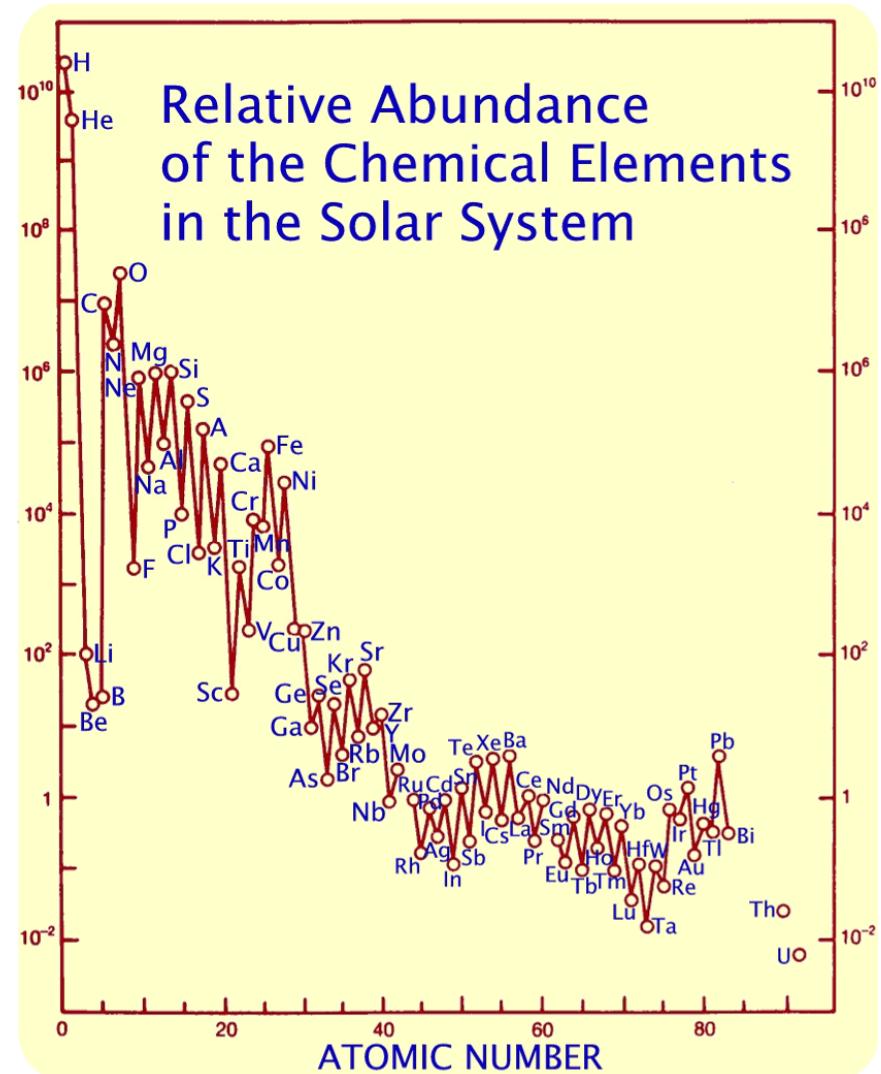


Neutron
Star
Mergers
~100 Myr





Arise from the Ashes



resurrection
in binaries !

Outline (7/16 & 7/17)

1. Overview of nucleosynthesis
2. Big Bang Nucleosynthesis
3. r-Process Nucleosynthesis
4. Neutrino-induced Nucleosynthesis
5. Cosmological Structure Formation
6. Galactic Chemical Evolution

Textbooks

Clayton: Principles of Stellar Evolution & Nucleosynthesis

Arnett: Supernovae & Nucleosynthesis

Pagel: Nucleosynthesis & Chemical Evolution of Galaxies

Basics of Big Bang Nucleosynthesis

initial state ($T > 1$ MeV): n, p

$$X_n + X_p = 1 \Rightarrow \text{need } n/p$$

rate of change in abundance:

$$\frac{dY_i}{dt} = P(t) - D(t)Y_i, \quad Y_i = \frac{X_i}{A_i}, \quad n_i = \rho_b N_A Y_i$$

$P(t)$: production rate
 $D(t)$: destruction rate } both depend on $T(t)$ and $\rho_b(t)$

$T(t)$ specified by dynamics of expansion

$\rho_b(t)$ specified by conservation of entropy per baryon

$$s \propto g_{\text{eff}}^*(t) \frac{T^3}{\rho_b} \propto g_{\text{eff}}^*(t) \frac{n_\gamma}{n_b} = \text{const.}$$

baryon-to-photon ratio: $\eta = \frac{n_{b,0}}{n_{\gamma,0}} \Rightarrow s \approx \frac{3.6}{\eta}$

entropy conservation \Rightarrow evolution of ρ_{rel} at $100 > T > 1 \text{ MeV}$

$$TS = E + PV - \mu N \Rightarrow S = \frac{E + PV - \mu N}{T}$$

fully relativistic: $S_{\text{rel}} = \frac{\rho_{\text{rel}} V + (\rho_{\text{rel}}/3)V}{T} \propto g_{\text{eff}} T(t)^3 R(t)^3$

$$g_{\text{eff}} = \text{const.} \Rightarrow T(t) \propto R(t)^{-1}, \dot{T}/T = -\dot{R}/R$$

$$\left(\frac{\dot{R}}{R}\right)^2 = \left(\frac{\dot{T}}{T}\right)^2 = \frac{8\pi}{3} G \rho_{\text{rel}} = \left(\frac{8\pi}{3} G\right) g_{\text{eff}} \frac{\pi^2}{15} T^4$$

$$T \rightarrow \infty \text{ as } t \rightarrow 0 \Rightarrow \frac{\dot{T}}{T} = -\sqrt{\frac{8\pi^3}{45} g_{\text{eff}} G T^4}$$

$$t \approx \frac{1}{2} \sqrt{\frac{45}{8\pi^3}} \frac{1}{\sqrt{g_{\text{eff}} G}} \frac{1}{T^2} = \frac{1.71}{\sqrt{g_{\text{eff}}}} \left(\frac{\text{MeV}}{T}\right)^2 \text{ s}$$

$$N_\nu = 3 \Rightarrow g_{\text{eff}} = \frac{43}{8}, \quad t \approx 0.74 \left(\frac{\text{MeV}}{T}\right)^2 \text{ s}$$

BBN and Neutrinos

freeze-out of n/p : $\nu_e + n \rightleftharpoons p + e^-$, $\bar{\nu}_e + p \rightleftharpoons n + e^+$

$$\sigma_{\nu_e n} \approx \frac{G_F^2}{\pi} \cos^2 \theta_C (f^2 + 3g^2) (E_{\nu_e} + \Delta)^2$$

$$\sigma_{\bar{\nu}_e p} \approx \frac{G_F^2}{\pi} \cos^2 \theta_C (f^2 + 3g^2) (E_{\bar{\nu}_e} - \Delta)^2$$

$$\cos^2 \theta_C = 0.95, \quad f = 1, \quad g = 1.26, \quad \Delta = M_n - M_p = 1.293 \text{ MeV}$$

$$\begin{aligned} \text{rate per nucleon: } \lambda_{\nu N} &\approx \frac{4\pi}{(2\pi)^3} \int_0^\infty \frac{\sigma_{\nu N} E_\nu^2}{\exp(E_\nu/T) + 1} dE_\nu \\ &\approx 0.4 \left(\frac{T}{\text{MeV}} \right)^5 \text{ s}^{-1} \end{aligned}$$

$$\begin{aligned} \int_{t_{\text{FO}}}^\infty \lambda_{\nu N} dt &\sim \int_0^{T_{\text{FO}}} 0.4 \left(\frac{T}{\text{MeV}} \right)^5 \times 2 \times 0.74 \left(\frac{\text{MeV}}{T} \right)^3 dT \\ &\sim 0.2 \left(\frac{T_{\text{FO}}}{\text{MeV}} \right)^3 \sim 1 \Rightarrow T_{\text{FO}} \sim 1.7 \text{ MeV} \end{aligned}$$

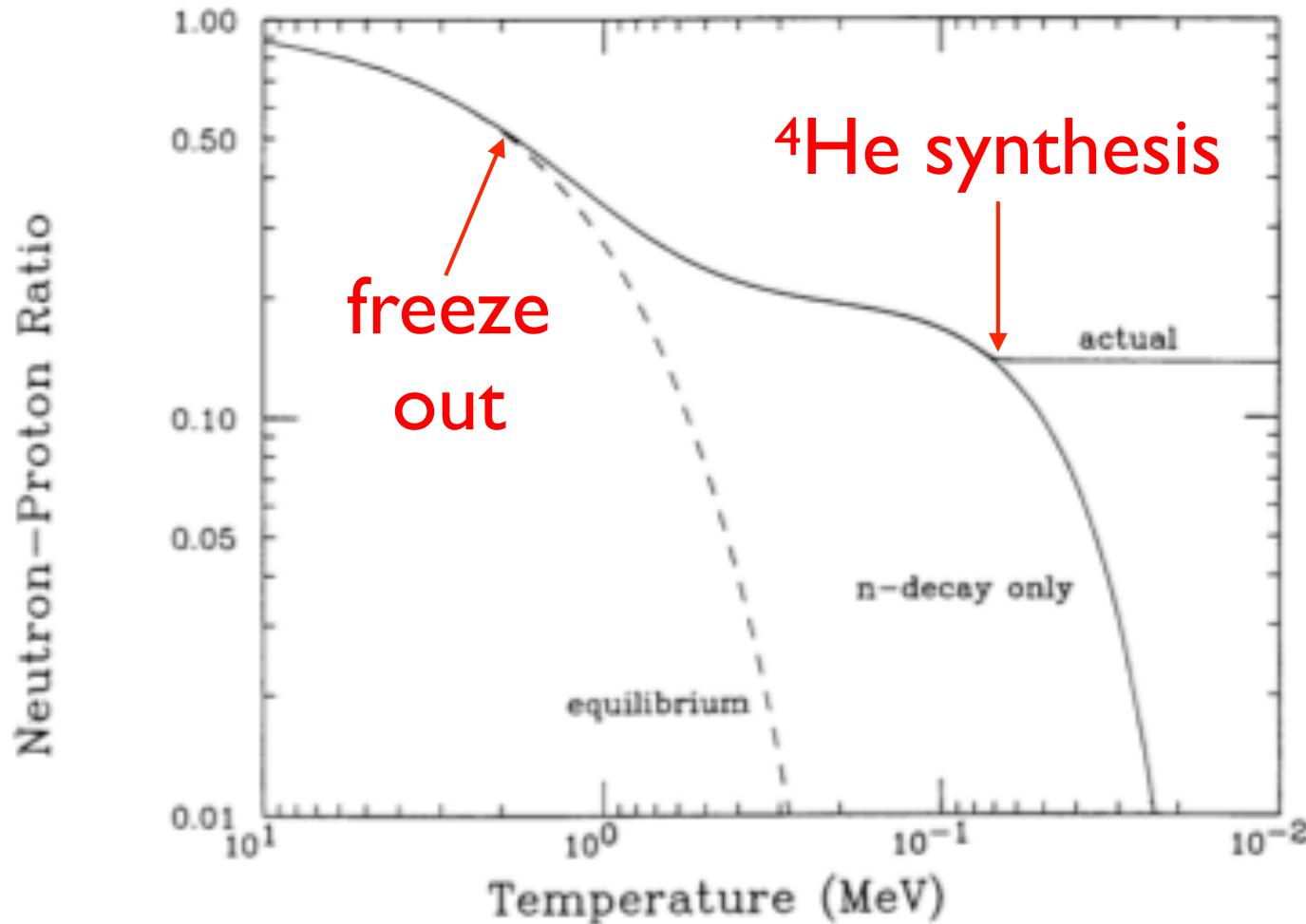


FIG. 1.—Evolution of the neutron-proton ratio with temperature. The NSE ratio is given by the dashed curve. If neutron decay is the only reaction (all other reactions are shut off), the n/p ratio follows the solid curve. The actual final value of the ratio is shown by the straight horizontal line.

$$\frac{n}{p} = \left(\frac{n}{p} \right)_{FO} \exp \left(-\frac{t - t_{FO}}{\tau_n} \right) \sim \exp \left(-\frac{\Delta}{T_{FO}} - \frac{t - t_{FO}}{\tau_n} \right)$$

$$\frac{n}{p} = \left(\frac{n}{p}\right)_{\text{FO}} \exp\left(-\frac{t - t_{\text{FO}}}{\tau_n}\right) \sim \exp\left(-\frac{\Delta}{T_{\text{FO}}} - \frac{t - t_{\text{FO}}}{\tau_n}\right)$$

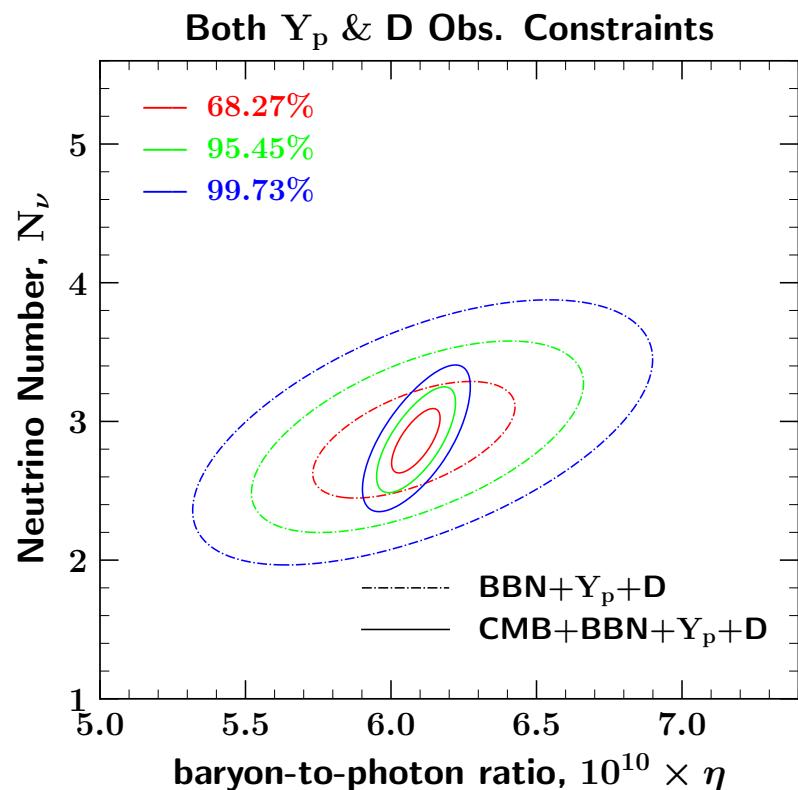
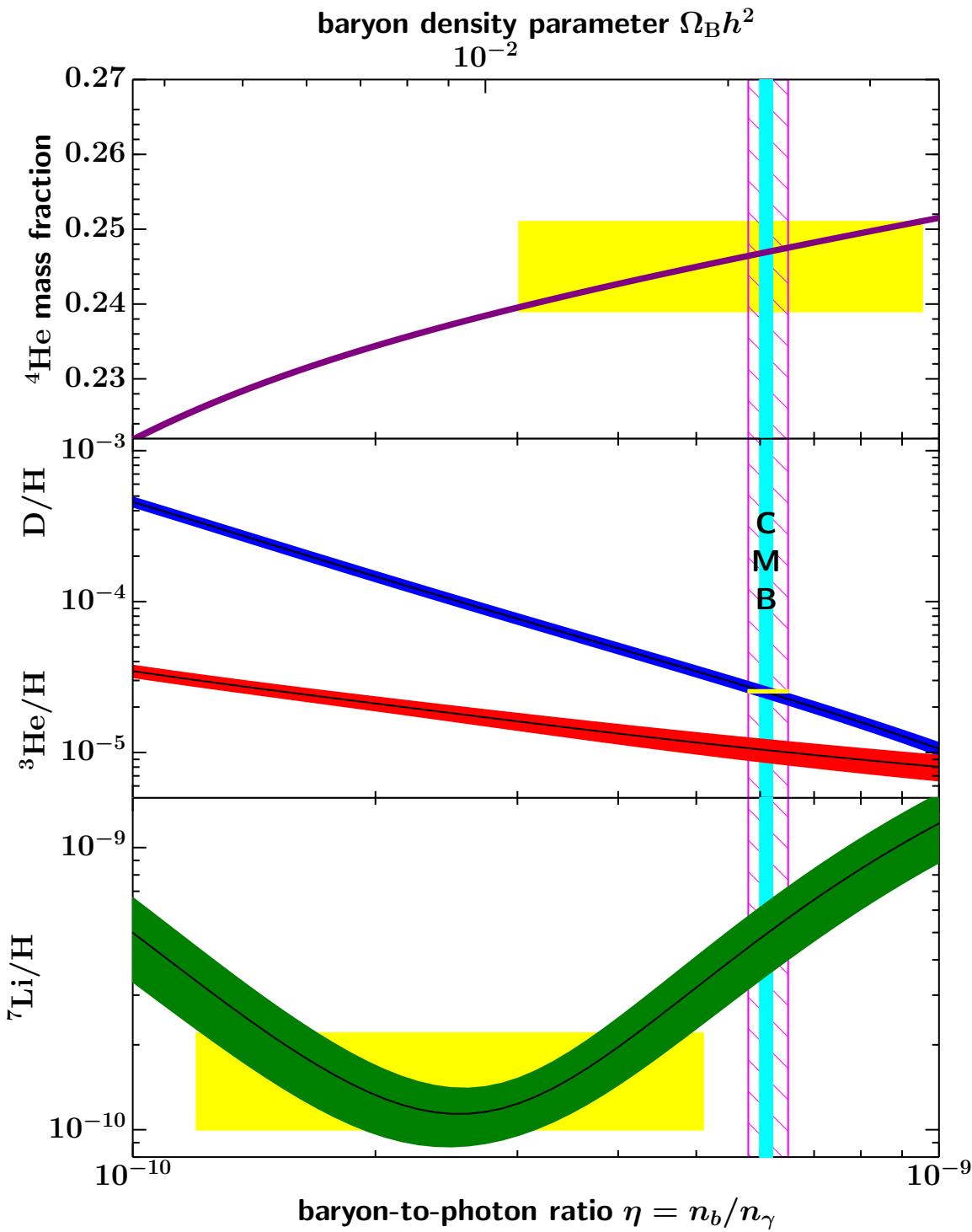
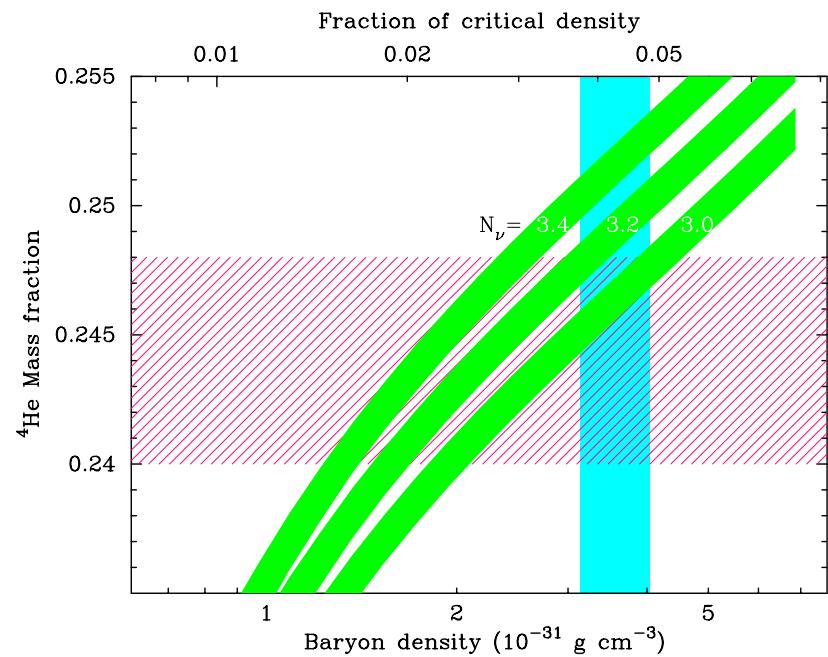
$${}^4\text{He} \text{ production: } X({}^4\text{He}) \sim \frac{2n}{n+p} = \frac{2(n/p)}{(n/p)+1}$$

$$t \sim \frac{1}{\sqrt{a+bN_\nu}} \frac{1}{T^2}, \quad \lambda_{\nu N} \propto T^5, \quad \int_{t_{\text{FO}}}^{\infty} \lambda_{\nu N} dt \propto \frac{T_{\text{FO}}^3}{\sqrt{a+bN_\nu}} \sim \text{const.}$$

$$N_\nu \uparrow \Rightarrow T_{\text{FO}} \uparrow \Rightarrow \left(\frac{n}{p}\right)_{\text{FO}} \sim \exp\left(-\frac{\Delta}{T_{\text{FO}}}\right) \uparrow$$

$$N_\nu \uparrow \Rightarrow t \downarrow \Rightarrow \frac{n}{p} = \left(\frac{n}{p}\right)_{\text{FO}} \exp\left(-\frac{t - t_{\text{FO}}}{\tau_n}\right) \uparrow$$

$$N_\nu \uparrow \Rightarrow \frac{n}{p} \uparrow \Rightarrow X({}^4\text{He}) \uparrow$$



Nuclear Statistical Equilibrium (NSE)

$$Zp + (A - Z)n \rightleftharpoons (Z, A) + \gamma \Rightarrow Z\mu_p + (A - Z)\mu_n = \mu(Z, A)$$

considering excited states of nuclei:

$$\begin{aligned} n(Z, A) &= \sum_i \frac{2J_i + 1}{(2\pi)^3} \int_0^\infty \frac{4\pi p^2 dp}{\exp\{[(p^2/2M) + M + E_i - \mu]/T\}} \\ &= G(Z, A) \left(\frac{MT}{2\pi}\right)^{3/2} \exp\left[\frac{\mu(Z, A) - M(Z, A)}{T}\right] \end{aligned}$$

$$\text{nuclear partition function: } G(Z, A) = \sum_i (2J_i + 1) \exp\left(-\frac{E_i}{T}\right)$$

$$\begin{aligned} \Rightarrow X(Z, A) &= X_p^Z X_n^{A-Z} \frac{G(Z, A)}{2^A} A^{5/2} \\ &\times \left(\frac{\rho_b}{M_N}\right)^{A-1} \left(\frac{2\pi}{M_N T}\right)^{3(A-1)/2} \exp\left[\frac{B(Z, A)}{T}\right] \end{aligned}$$

In NSE, no rates are needed to calculate abundances:

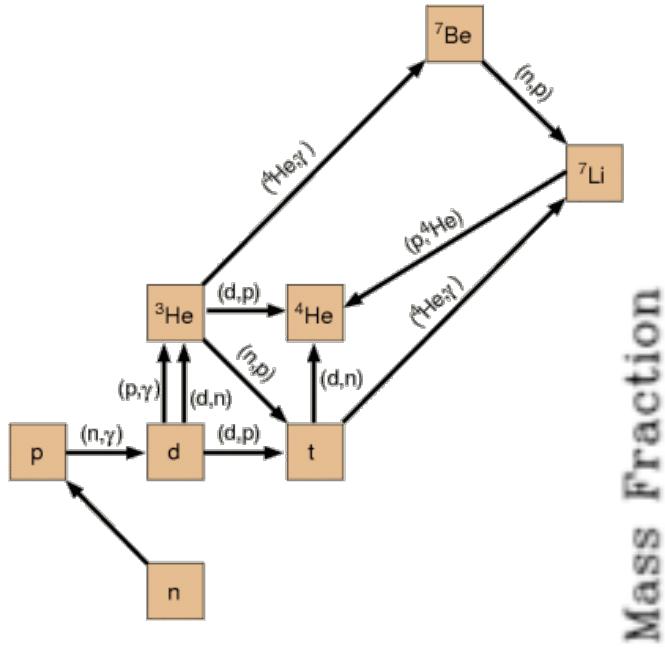
$$1 = X_n + X_p + \sum_{(Z,A)} X(Z, A)$$

$$Y_e = X_p + \sum_{(Z,A)} \frac{Z}{A} X(Z, A)$$

$$\begin{aligned} X(Z, A) &= X_p^Z X_n^{A-Z} \frac{G(Z, A)}{2^A} A^{5/2} \\ &\times \left(\frac{\rho_b}{M_N} \right)^{A-1} \left(\frac{2\pi}{M_N T} \right)^{3(A-1)/2} \exp \left[\frac{B(Z, A)}{T} \right] \end{aligned}$$

$$\left[\eta \left(\frac{T}{M_N} \right)^{3/2} \right]^{A-1} \exp \left[\frac{B(Z, A)}{T} \right] \sim 1$$

$\Rightarrow (Z, A)$ drops out of NSE at lower T



$p(n, \gamma)d$

$d(d, n)^3\text{He}$

$d(d, p)t$

$^3\text{He}(d, p)^4\text{He}$

$t(d, n)^4\text{He}$

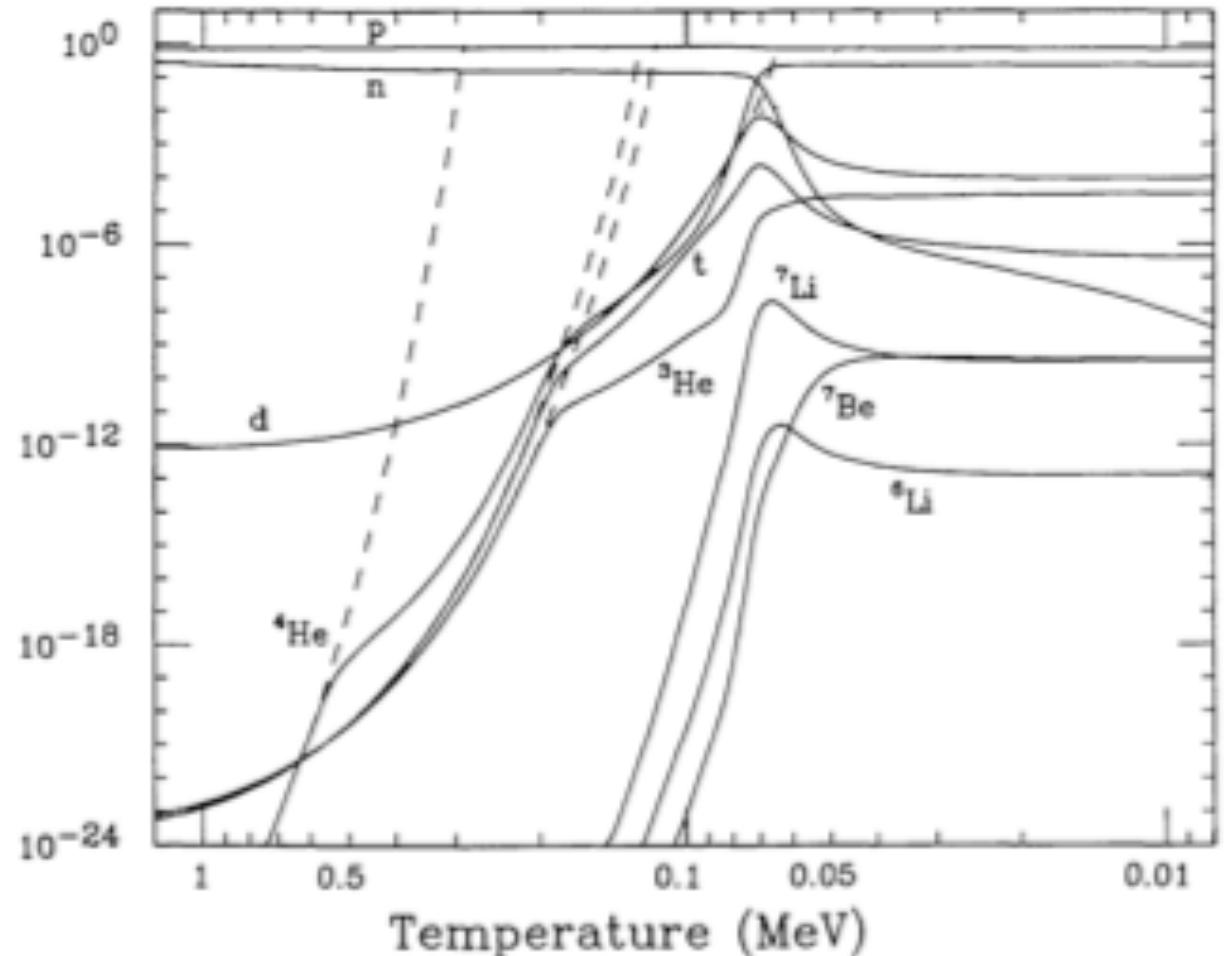
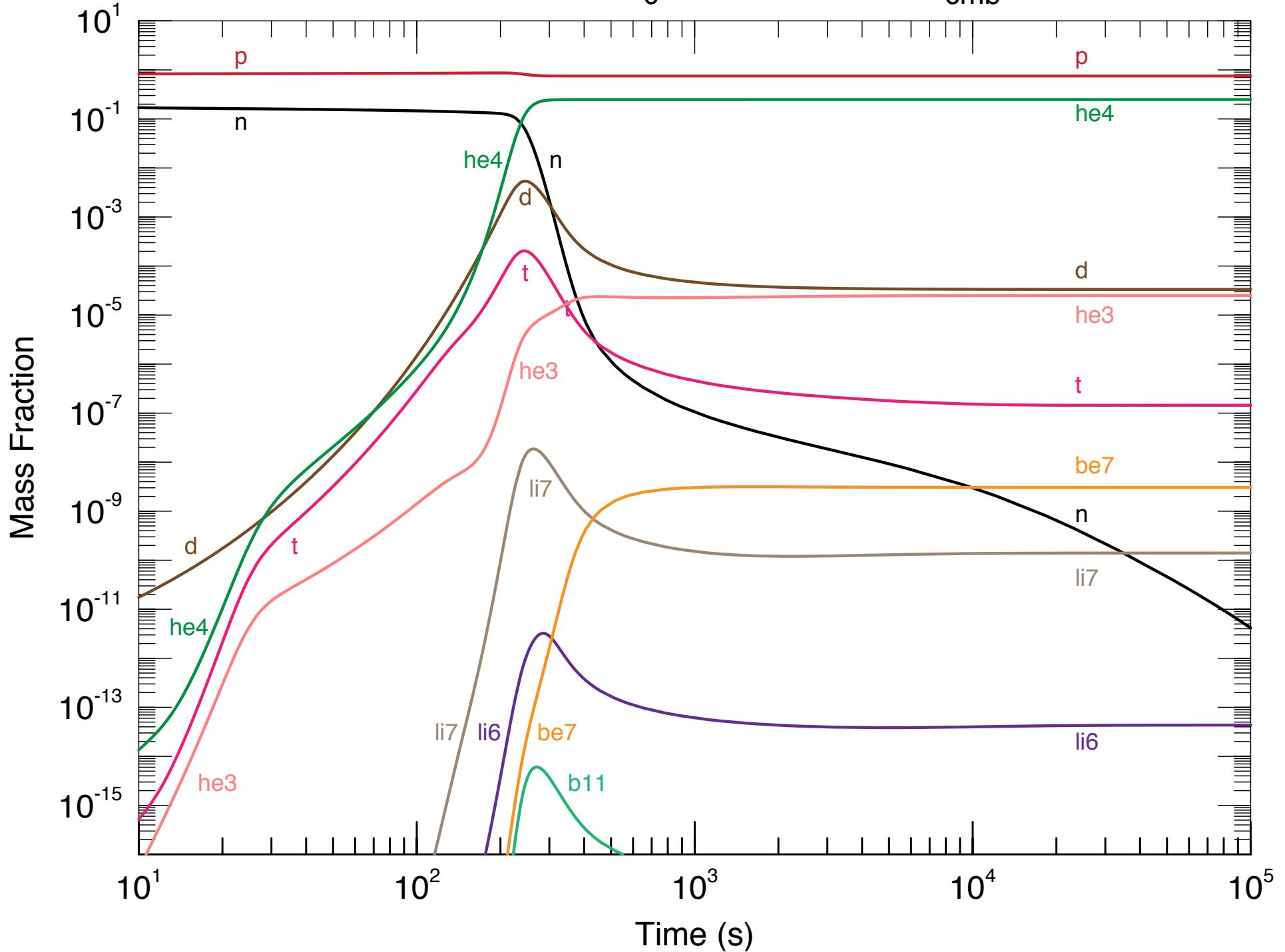


FIG. 2.—Evolution of light-element abundances with temperature, for a baryon-to-photon ratio $\eta_{10} = 3.16$. The dashed curves give the NSE curves of ^4He , t , ^3He , and d , respectively. The dotted curve is explained in the text.

$$T_{\text{NSE}} \sim \frac{B(Z, A)/(A - 1)}{\ln \eta^{-1} + (3/2) \ln(M_N/T)}$$

$$\eta = 6.23 \times 10^{-10} \quad N_V = 3.0 \quad H_0 = 70.50 \text{ km/s/Mpc} \quad T_{\text{cmb}} = 2.725 \text{ K}$$



Expansion from high temperature & density

- nuclear statistical equilibrium (NSE)

all strong & electromagnetic reactions in equilibrium



- quasi-statistical equilibrium (QSE)

clusters of nuclei form & reactions involving n, p,
& light nuclei in equilibrium within each cluster



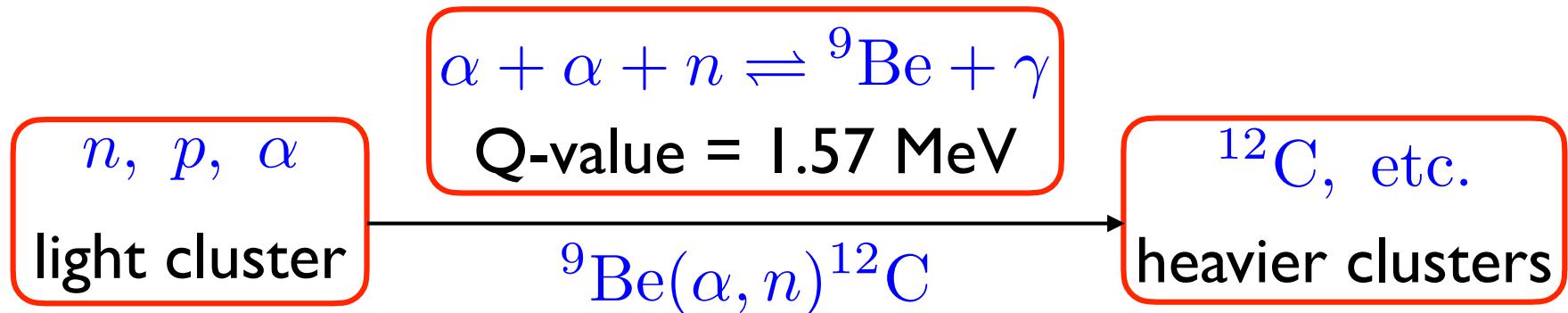
- hot r-process

QSE within each isotopic chain only



quasi-statistical equilibrium (QSE)

typically QSE is achieved for $0.5 \gtrsim T \gtrsim 0.25$ MeV



determination of neutron-to-seed ratio

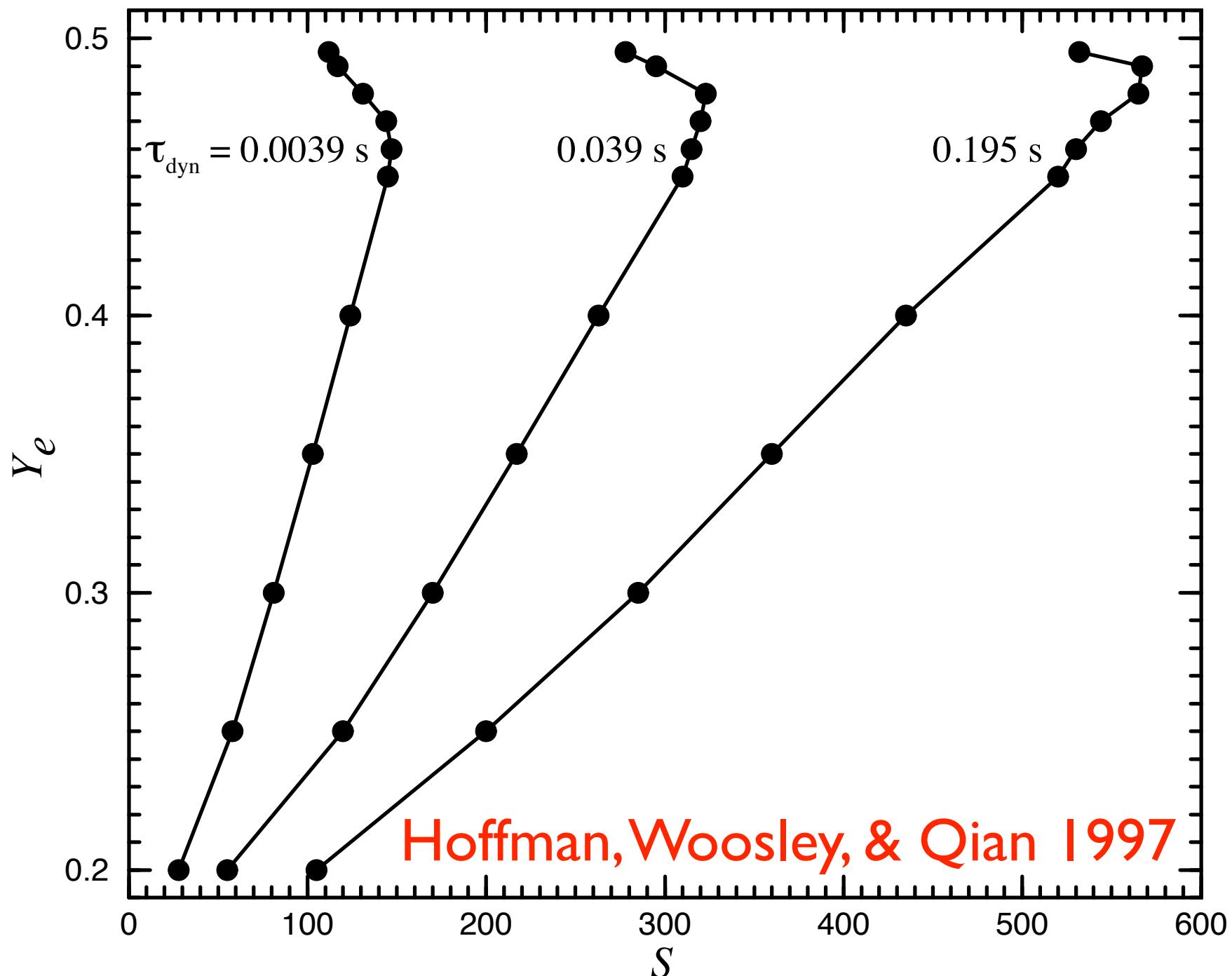
$Y_e \downarrow \Rightarrow$ more neutrons

entropy per baryon $S \propto (T^3/\rho) \uparrow \Rightarrow$

more photons capable of disintegrating ${}^9\text{Be} \Rightarrow$ fewer seeds

$\tau_{\text{dyn}} \downarrow \Rightarrow$ shorter time for producing seeds

conditions for producing r-nuclei with A~200



$(n, \gamma) \rightleftharpoons (\gamma, n)$ equilibrium: $n + (Z, A) \rightleftharpoons (Z, A + 1) + \gamma$

typically achieved for $0.25 \gtrsim T \gtrsim 0.1$ MeV

$$\frac{Y(Z, A + 1)}{Y(Z, A)} = \frac{G(Z, A + 1)}{G(Z, A)} \left(\frac{A + 1}{A} \right)^{3/2}$$

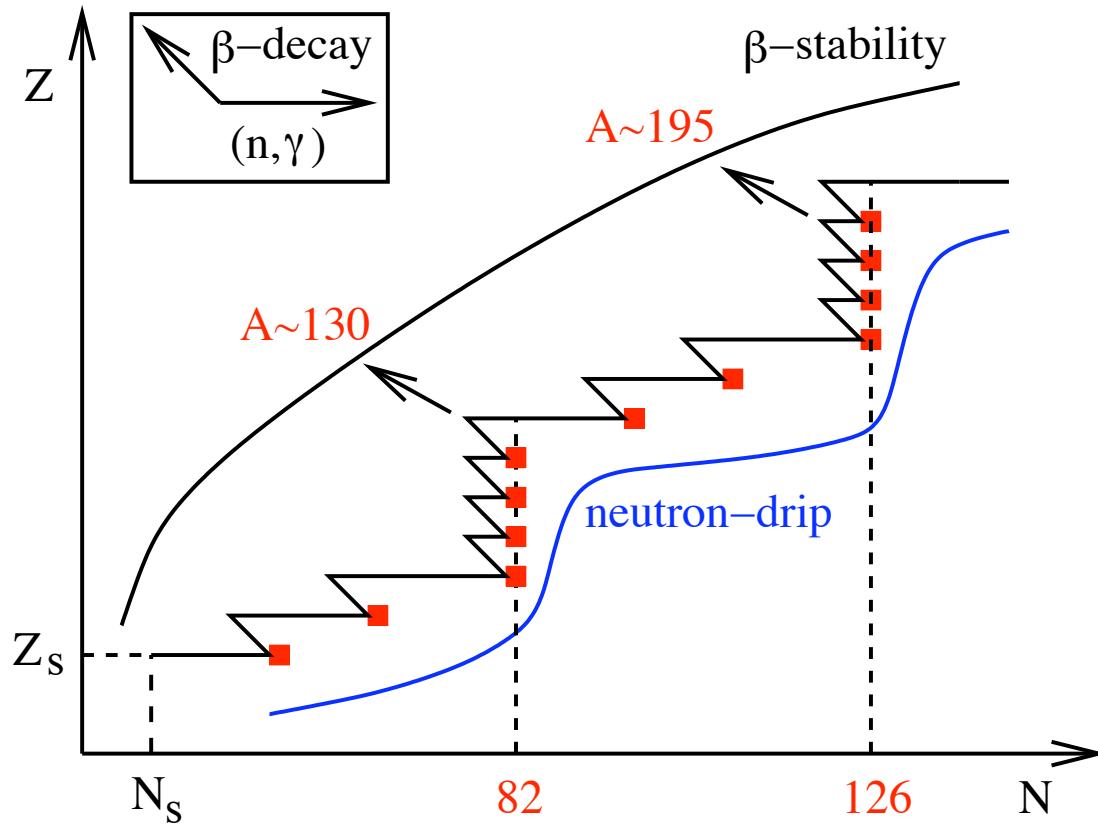
$$\times \boxed{\frac{n_n}{2} \left(\frac{2\pi}{M_N T} \right)^{3/2} \exp \left[\frac{S_n(Z, A + 1)}{T} \right]}$$

$\sim 1 \Rightarrow$ waiting-point nuclei

$$\begin{aligned} S_n^{\text{WP}} &\sim T \ln \left[\frac{2}{n_n} \left(\frac{M_N T}{2\pi} \right)^{3/2} \right] \\ &= T_9 [2.79 + 0.198(1.5 \log T_9 - \log n_{n,20})] \text{ MeV} \end{aligned}$$

steady β -flow in $(n, \gamma) \rightleftharpoons (\gamma, n)$ equilibrium

Rapid neutron capture: the r-process



$$\frac{dY_{\text{WP}}(Z+1)}{dt} = \lambda_{\beta}^{\text{WP}}(Z)Y_{\text{WP}}(Z) - \lambda_{\beta}^{\text{WP}}(Z+1)Y_{\text{WP}}(Z+1) = 0$$

$$\Rightarrow Y_{\text{WP}}(Z) \propto 1/\lambda_{\beta}^{\text{WP}}(Z)$$

$(n, \gamma) \rightleftharpoons (\gamma, n)$ equilibrium: $n + (Z, A) \rightleftharpoons (Z, A + 1) + \gamma$

$$\frac{Y(Z, A + 1)}{Y(Z, A)} = \frac{G(Z, A + 1)}{2G(Z, A)} \left(\frac{A + 1}{A} \right)^{3/2}$$

$$\times n_n \left(\frac{2\pi}{M_N T} \right)^{3/2} \exp \left[\frac{S_n(Z, A + 1)}{T} \right]$$

$$n_n \langle \sigma_{n,\gamma}(Z, A) v \rangle Y(Z, A) = \lambda_{\gamma,n}(Z, A + 1) Y(Z, A + 1)$$

$$\Rightarrow \lambda_{\gamma,n}(Z, A + 1) = n_n \frac{Y(Z, A)}{Y(Z, A + 1)} \langle \sigma_{n,\gamma}(Z, A) v \rangle$$

$$= \frac{2G(Z, A)}{G(Z, A + 1)} \left(\frac{A}{A + 1} \right)^{3/2} \langle \sigma_{n,\gamma}(Z, A) v \rangle$$

$$\times \left(\frac{M_N T}{2\pi} \right)^{3/2} \exp \left[- \frac{S_n(Z, A + 1)}{T} \right]$$

typical nuclear properties of waiting-point nuclei with $N = 82$ & 126

$$S_n \sim 2 \text{ MeV}, \langle \sigma_{n,\gamma} v \rangle \sim 10^{-20} \text{ cm}^3 \text{ s}^{-1}, \lambda_\beta \sim 10 \text{ s}^{-1}$$

$$n_n \langle \sigma_{n,\gamma} v \rangle \gg \lambda_\beta \Rightarrow n_n \gg 10^{21} \text{ cm}^{-3}$$

$$T_9 \sim 1 \Rightarrow \lambda_{\gamma,n} \sim 10^4 \text{ s}^{-1} \gg \lambda_\beta$$

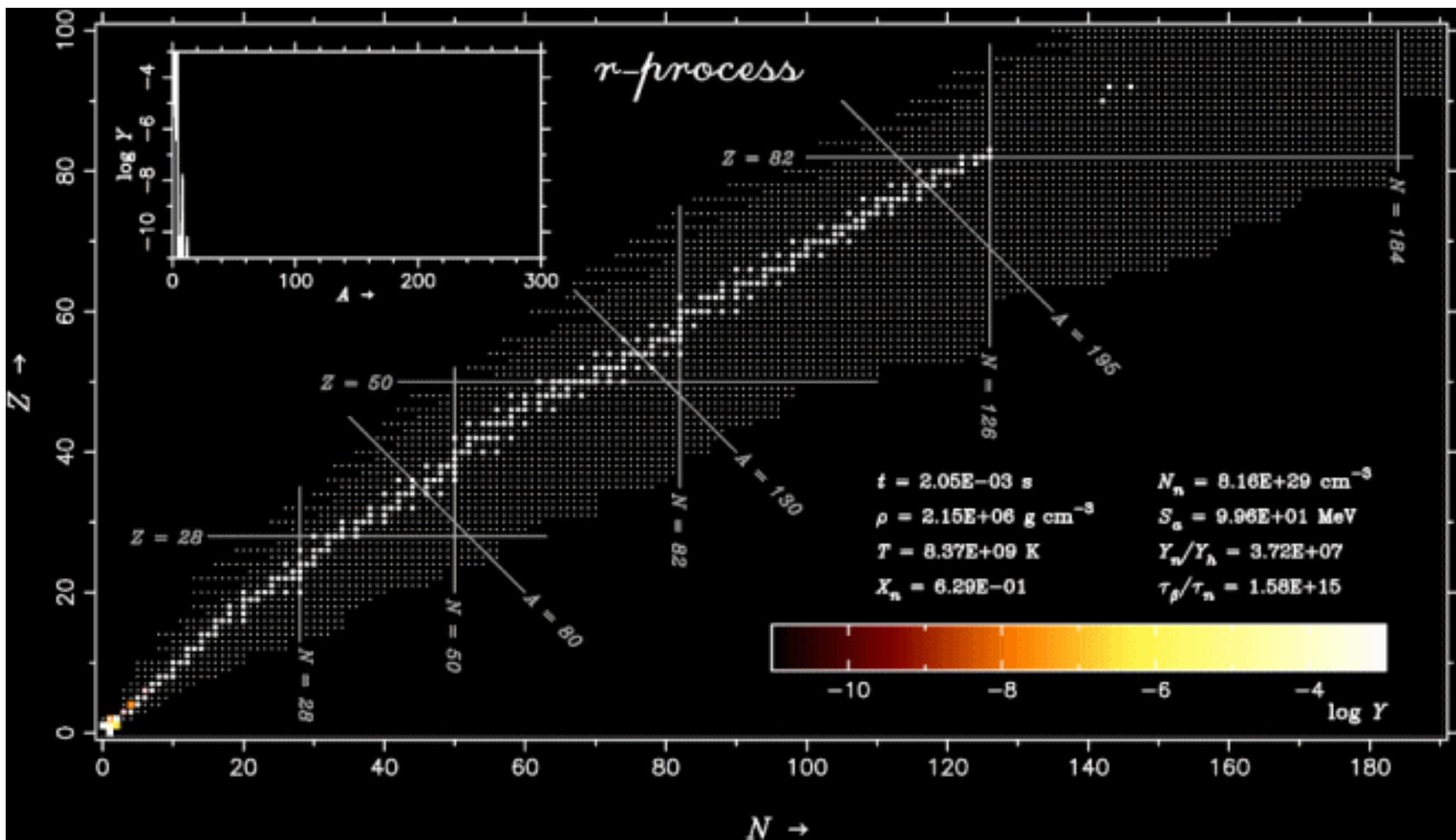
time for β -flow to go from seeds through $N = 82$ and 126

is $\lesssim 1$ s

in contrast to a hot r-process, a cold r-process occurs
when photo-disintegration can be ignored

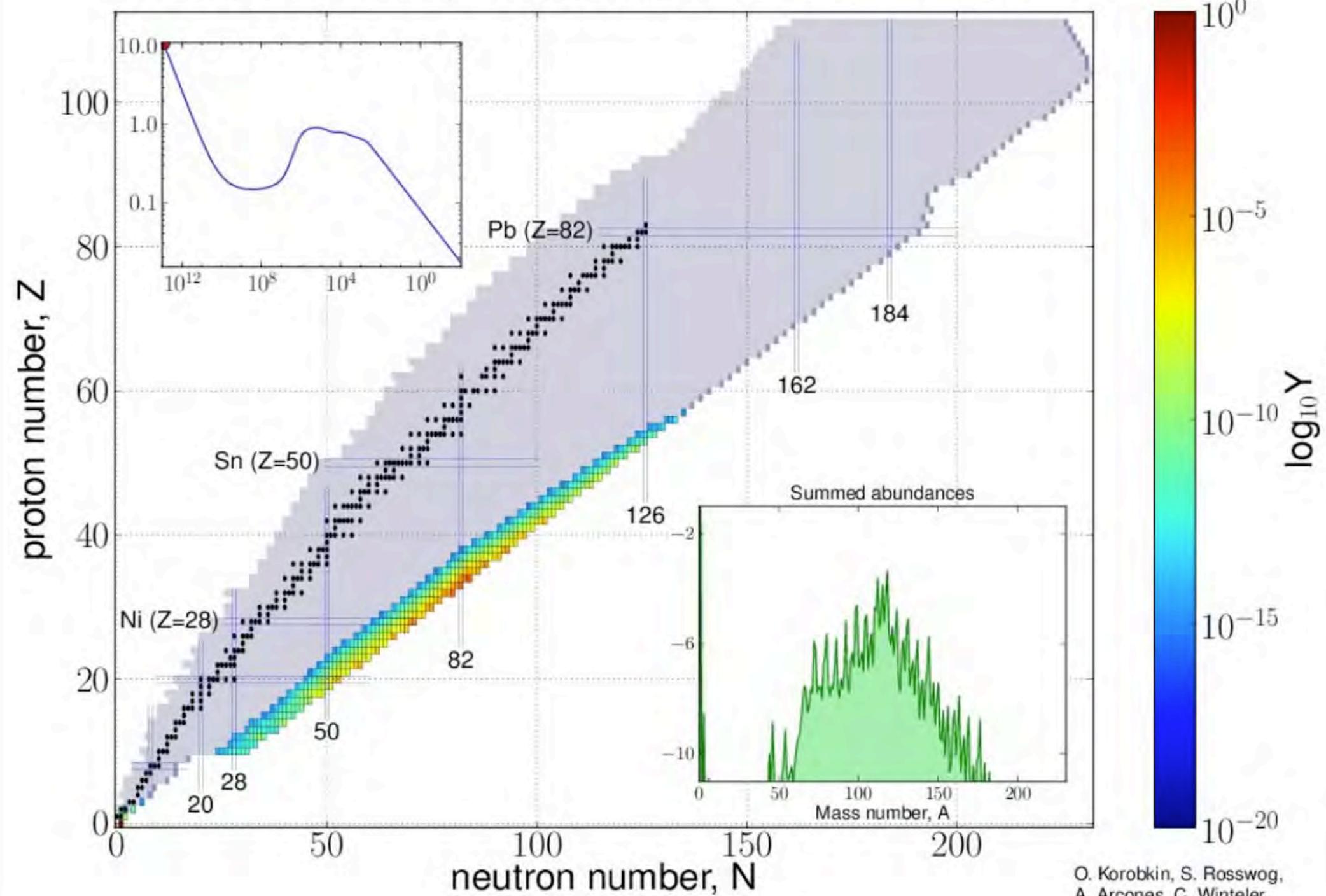
$T_9 \ll 1 \Rightarrow$ cold r-process

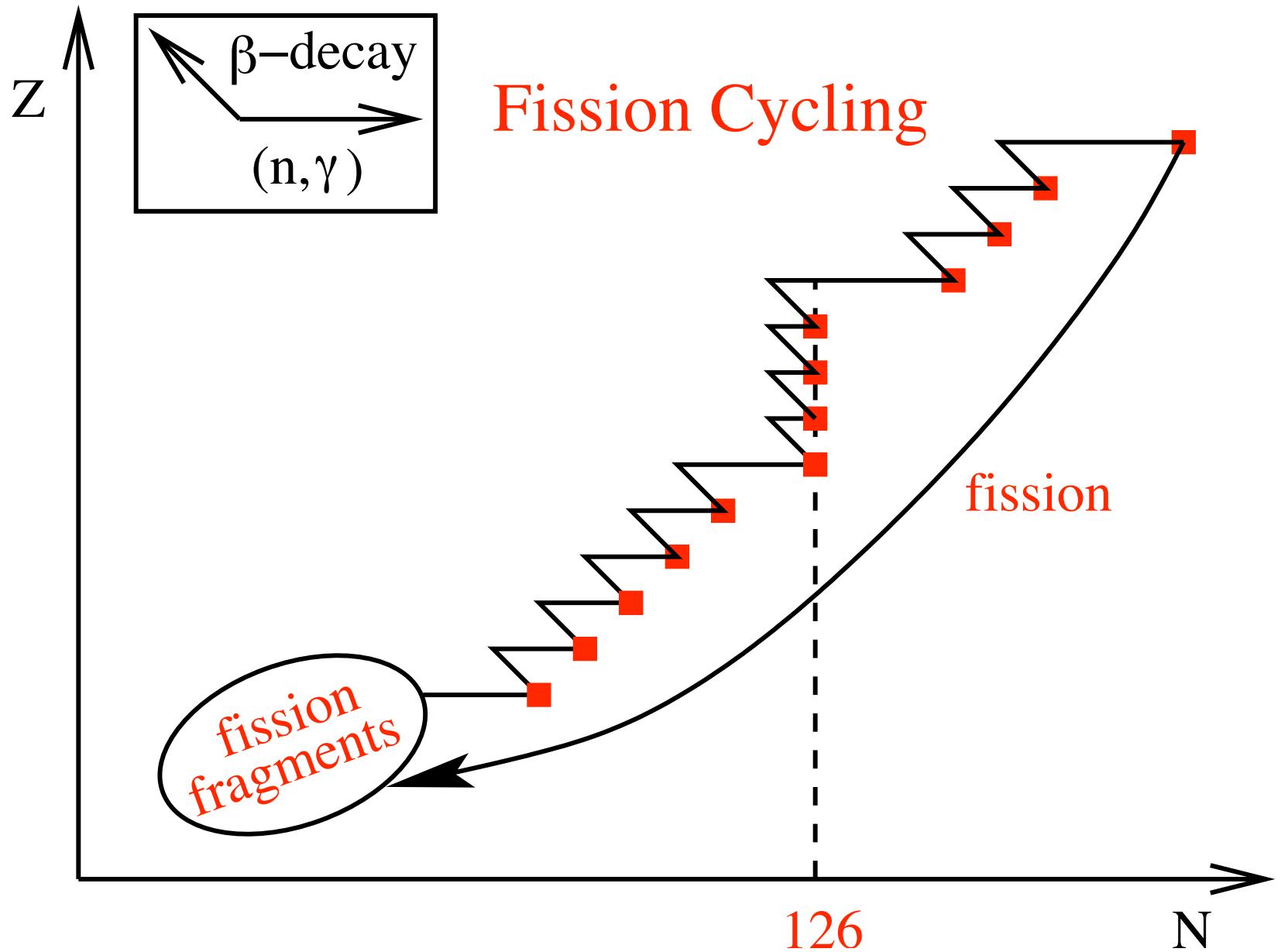
Hot r-process starting with free nucleons



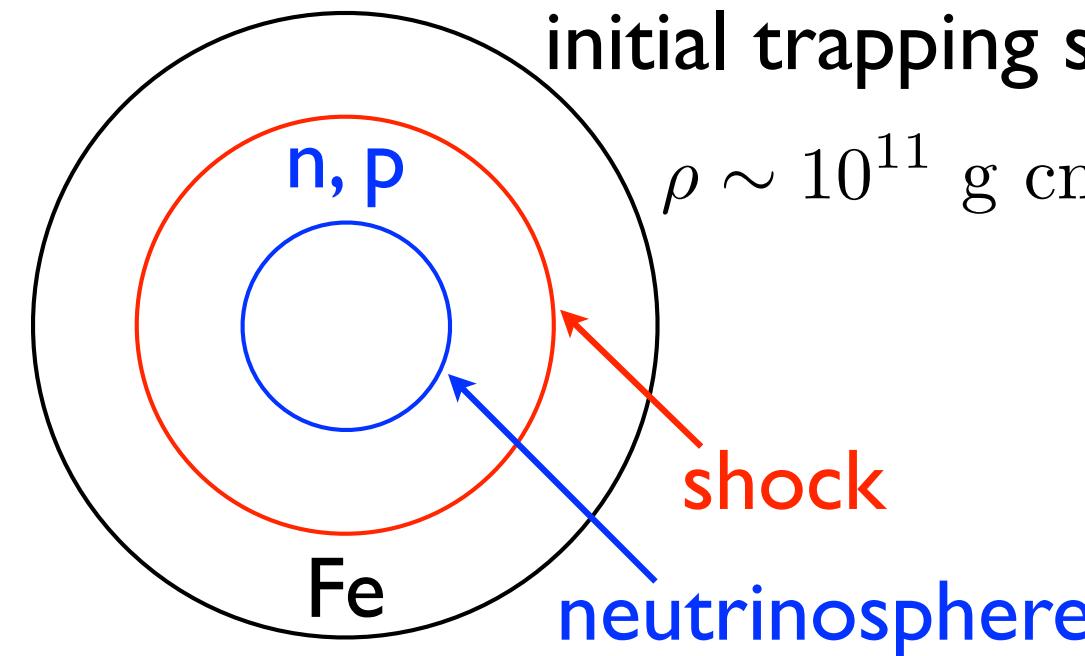
Wanajo et al. 2004

$t : 0.00e+00 \text{ s} / T : 10.96 \text{ GK} / \rho_b : 8.71e+12 \text{ g/cm}^3$





“neutronization” pulse at shock breakout

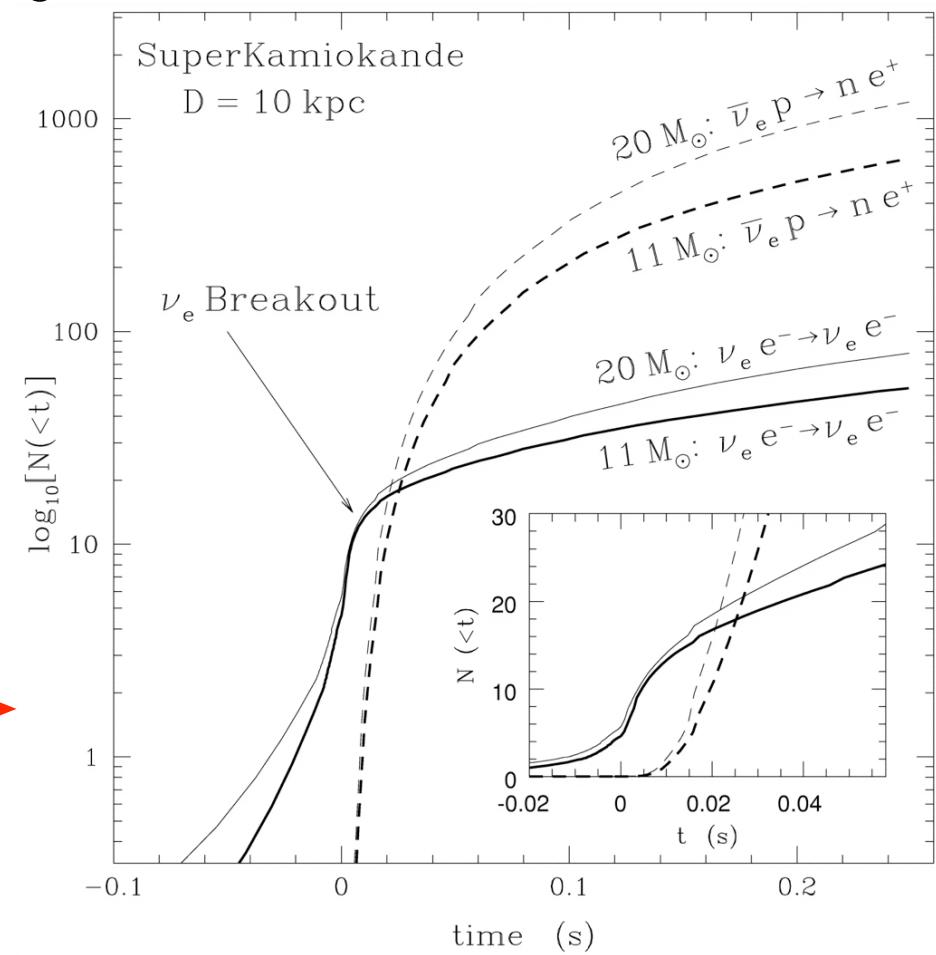


initial trapping surface

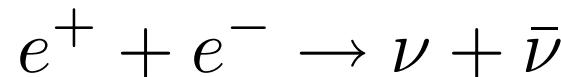
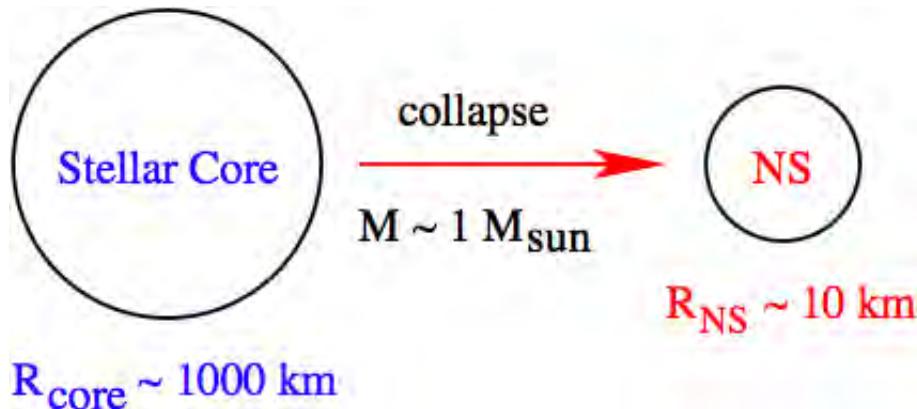
$$\rho \sim 10^{11} \text{ g cm}^{-3}$$

without oscillations

(Thompson et al. 2003)

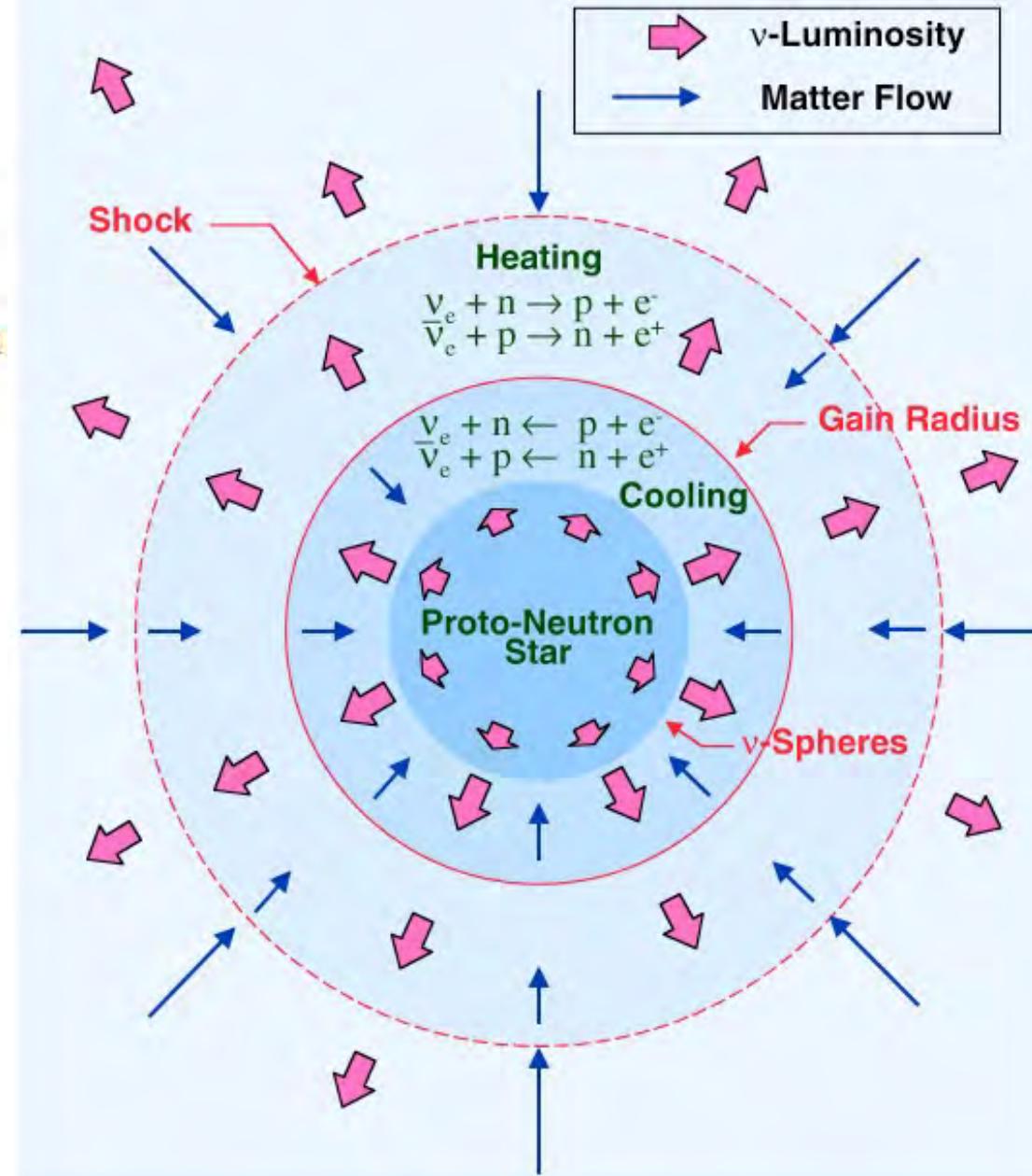


Supernovae as a neutrino phenomenon

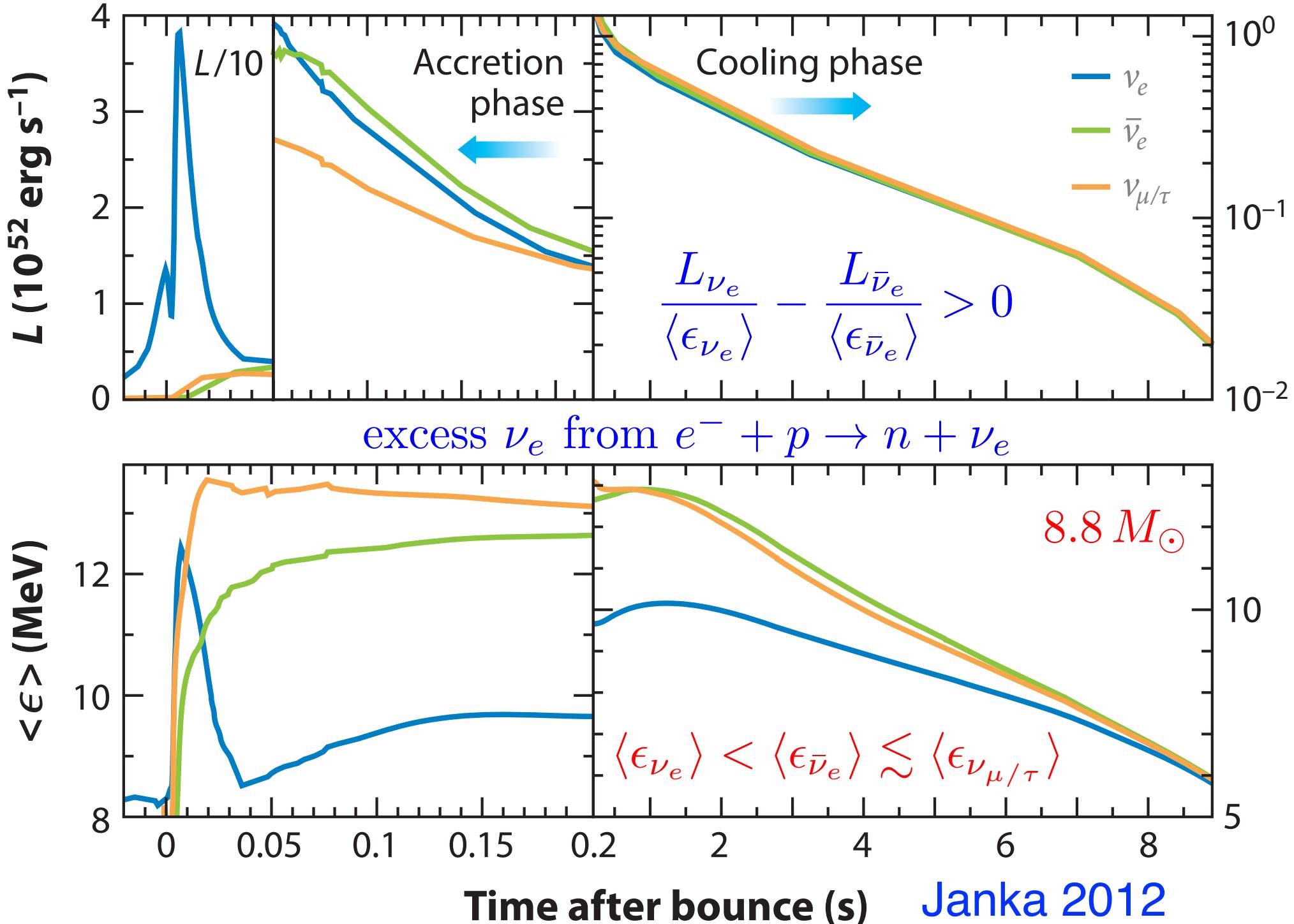


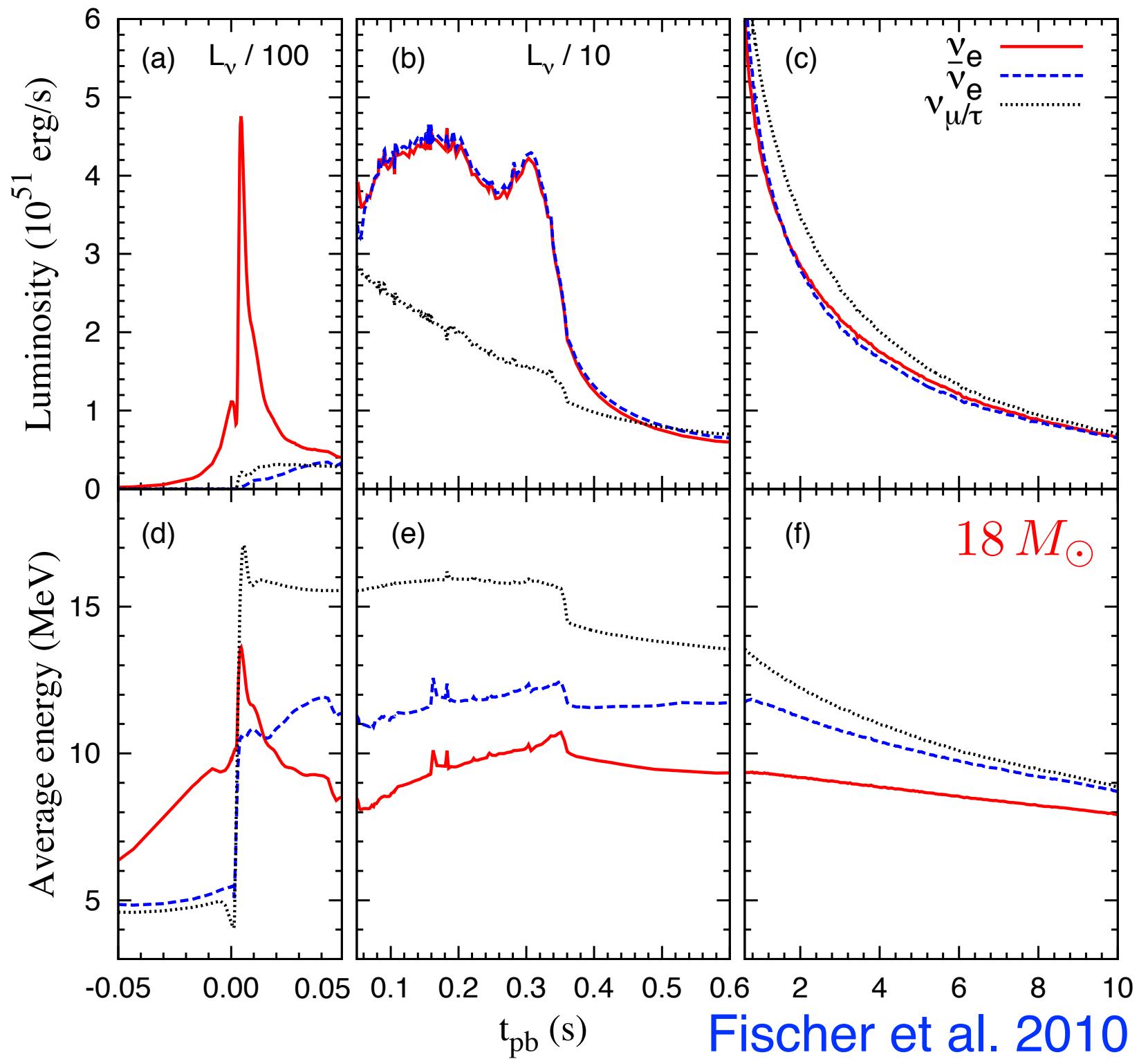
$$\frac{GM^2}{R_{\text{NS}}} \sim 3 \times 10^{53} \text{ erg}$$

$\Rightarrow \nu_e, \bar{\nu}_e, \nu_\mu, \bar{\nu}_\mu, \nu_\tau, \bar{\nu}_\tau$



neutrino emission from a low-mass SN





Fischer et al. 2010

setting n/p in the neutrino-driven wind

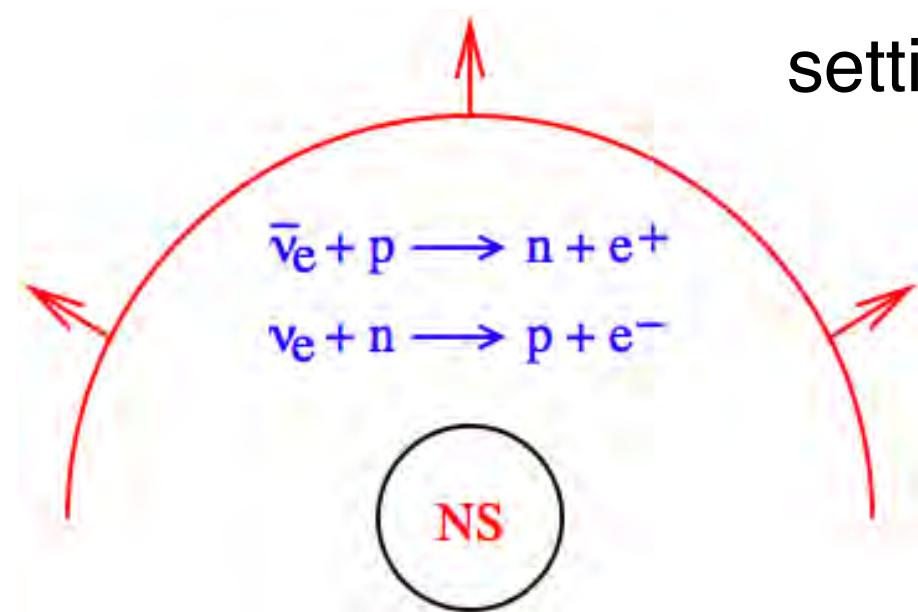
$$n/p > 1 \Rightarrow Y_e < 0.5$$

Qian et al. 1993

Qian & Woosley 1996

McLaughlin et al. 1996

Horowitz & Li 1999



$$\sigma_{\nu N} \propto (E_\nu \mp \Delta_{np})^2$$

$$\lambda_{\bar{\nu}_e p} = \frac{L_{\bar{\nu}_e}}{4\pi r^2} \frac{\langle \sigma_{\bar{\nu}_e p} \rangle}{\langle E_{\bar{\nu}_e} \rangle} \propto L_{\bar{\nu}_e} \left(\frac{\langle E_{\bar{\nu}_e}^2 \rangle}{\langle E_{\bar{\nu}_e} \rangle} - 2\Delta_{np} \right)$$

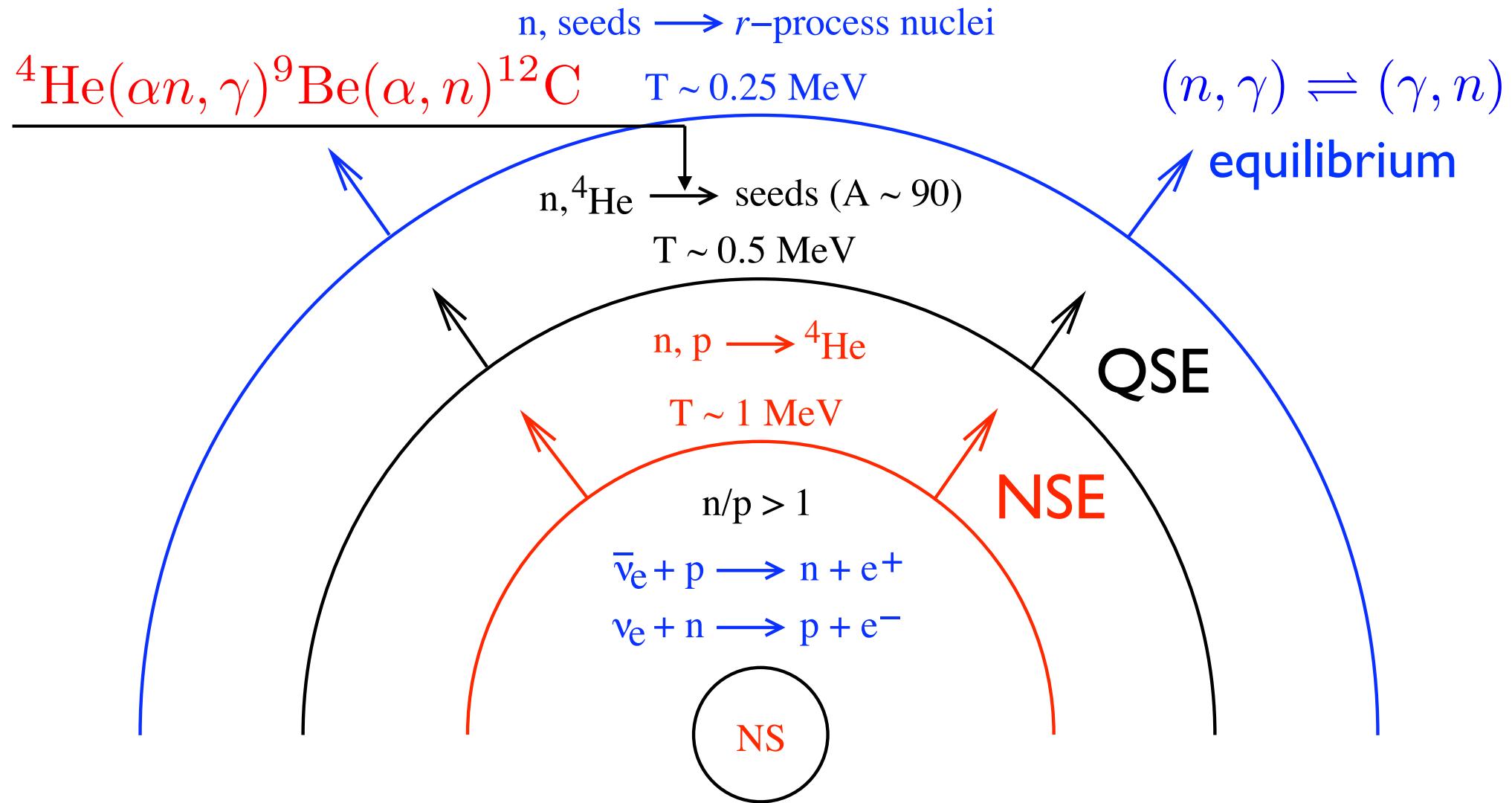
$$\lambda_{\nu_e n} = \frac{L_{\nu_e}}{4\pi r^2} \frac{\langle \sigma_{\nu_e n} \rangle}{\langle E_{\nu_e} \rangle} \propto L_{\nu_e} \left(\frac{\langle E_{\nu_e}^2 \rangle}{\langle E_{\nu_e} \rangle} + 2\Delta_{np} \right)$$

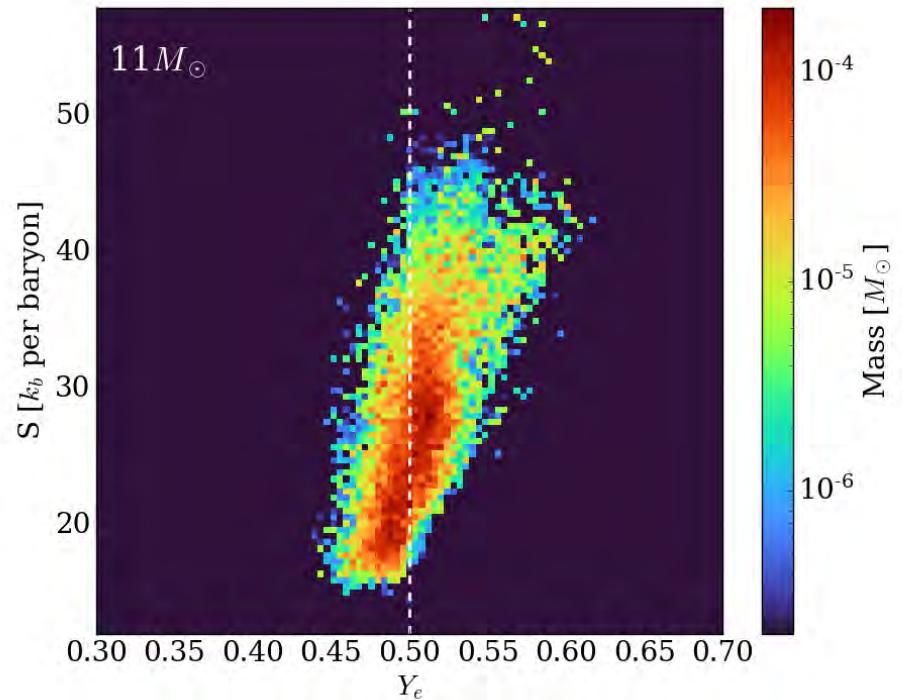
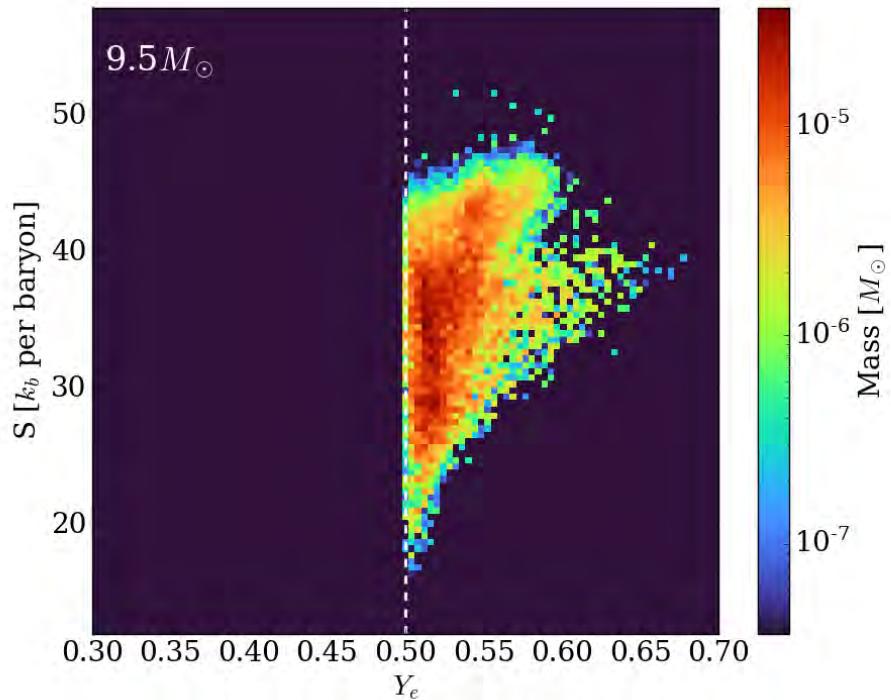
$$\frac{\langle E_{\bar{\nu}_e}^2 \rangle}{\langle E_{\bar{\nu}_e} \rangle} - \frac{\langle E_{\nu_e}^2 \rangle}{\langle E_{\nu_e} \rangle} > 4\Delta_{np} \approx 5.2 \text{ MeV} \Rightarrow \frac{n}{p} > 1$$

neutrino opacities!

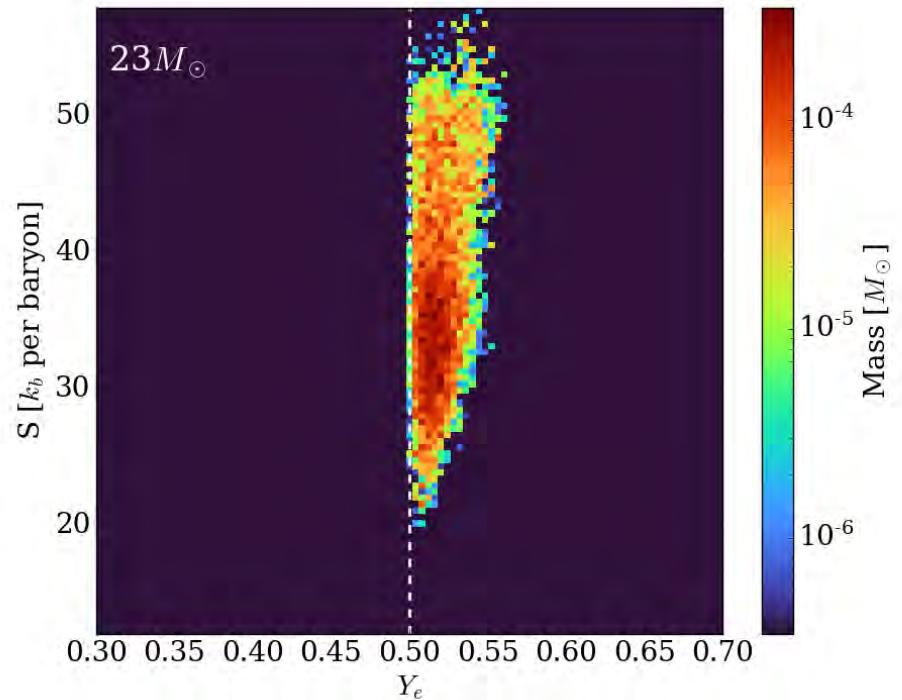
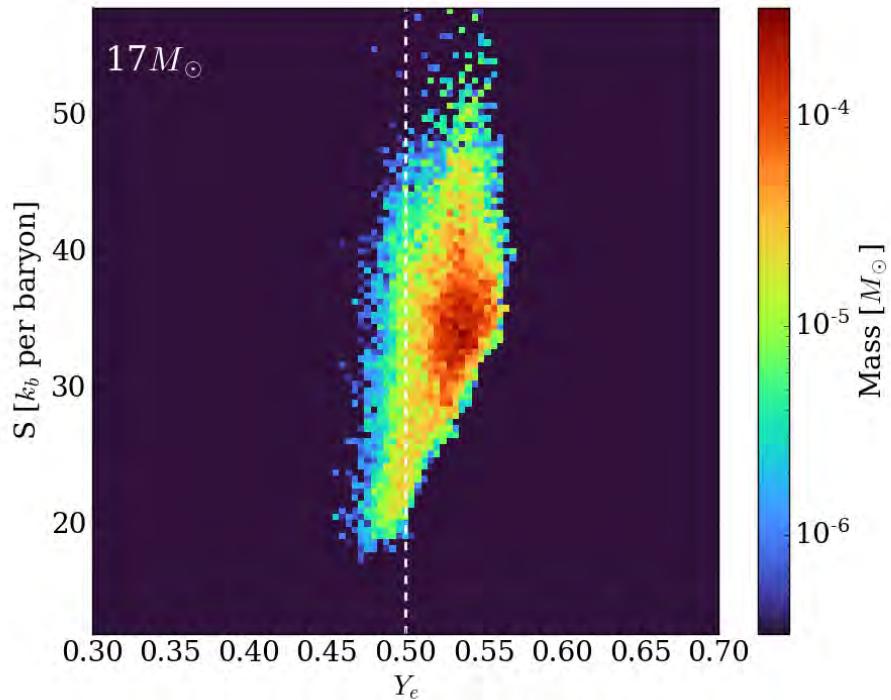
Martinez-Pinedo et al. 2012; Roberts & Reddy 2012

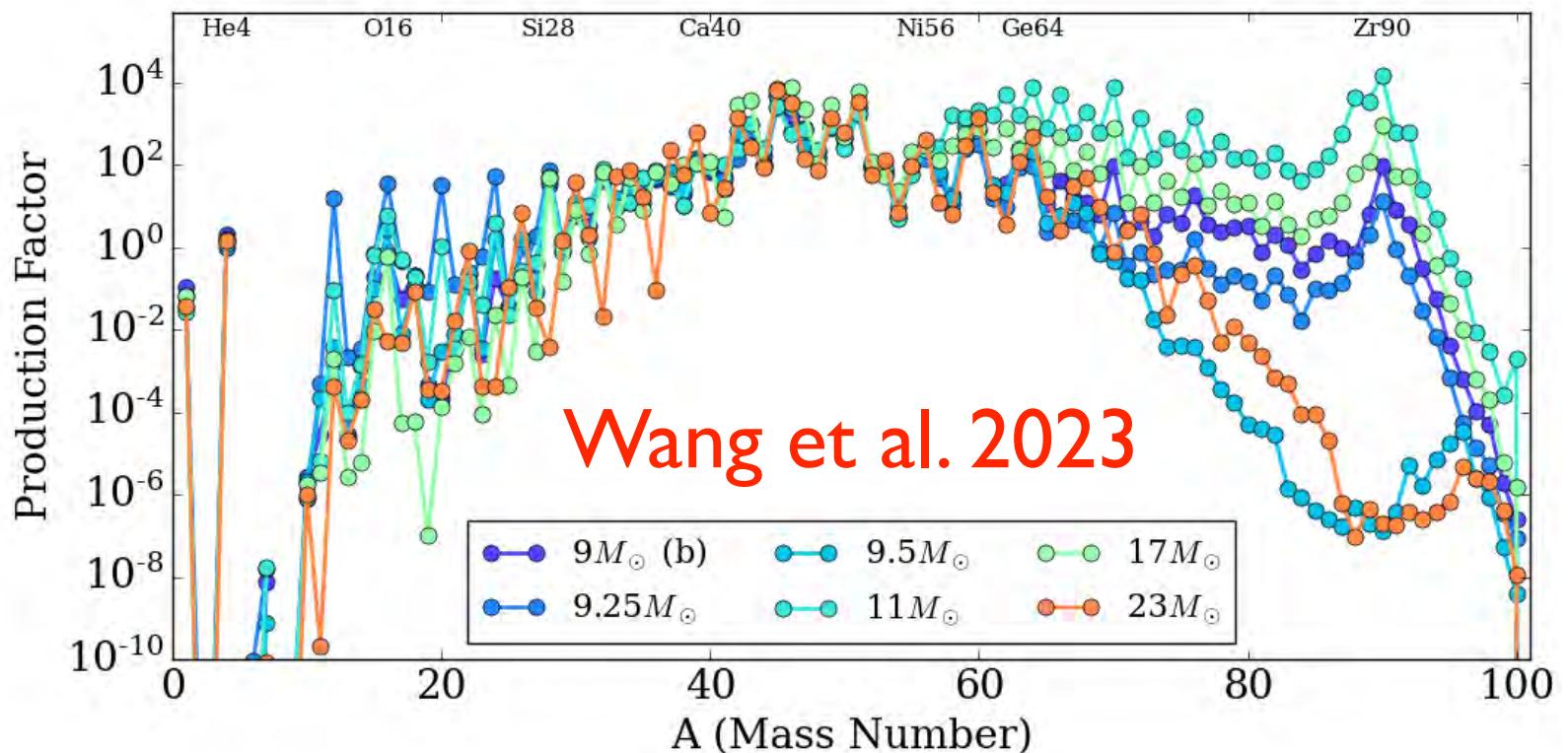
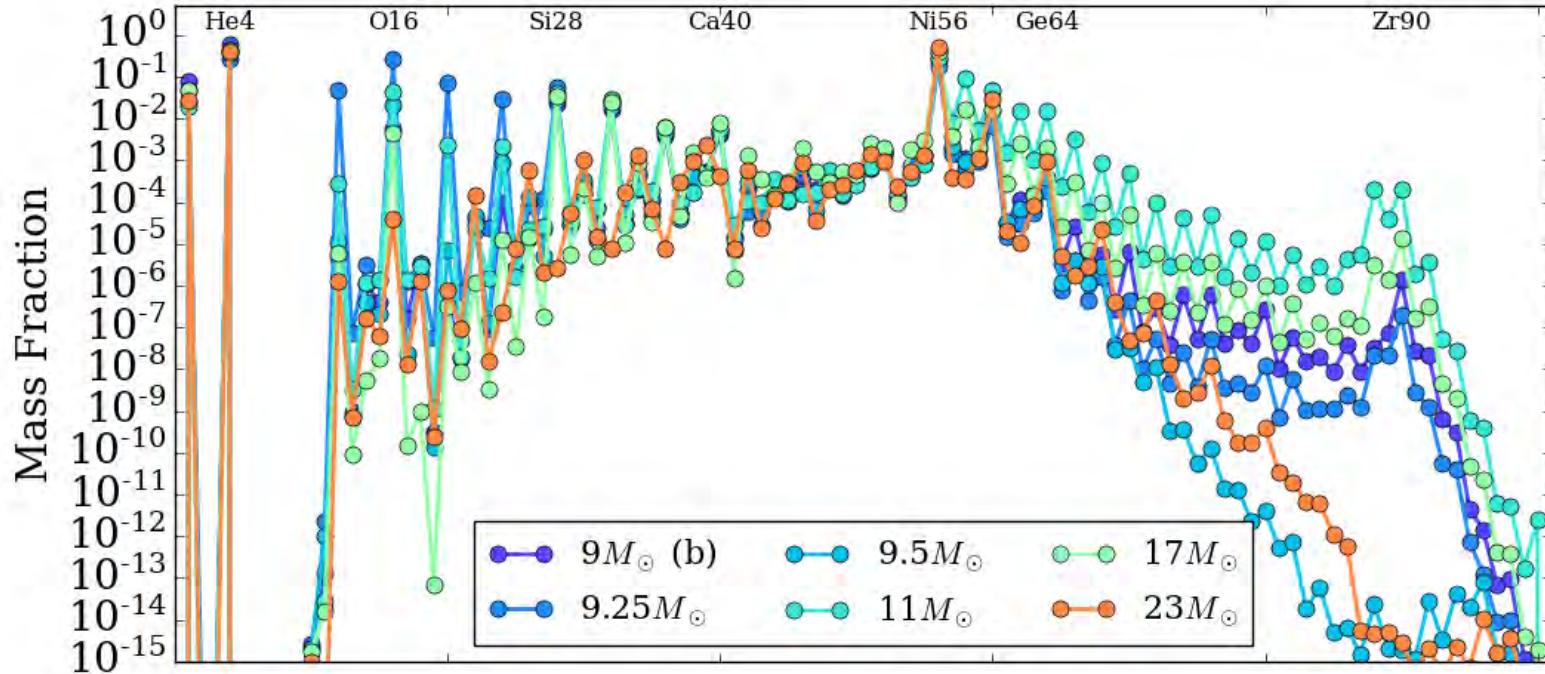
r-Process in Neutrino-driven Wind
 (e.g., Woosley & Baron 1992; Meyer et al. 1992; Woosley et al. 1994)





neutrino-driven winds in 3D models (Wang et al. 2023)

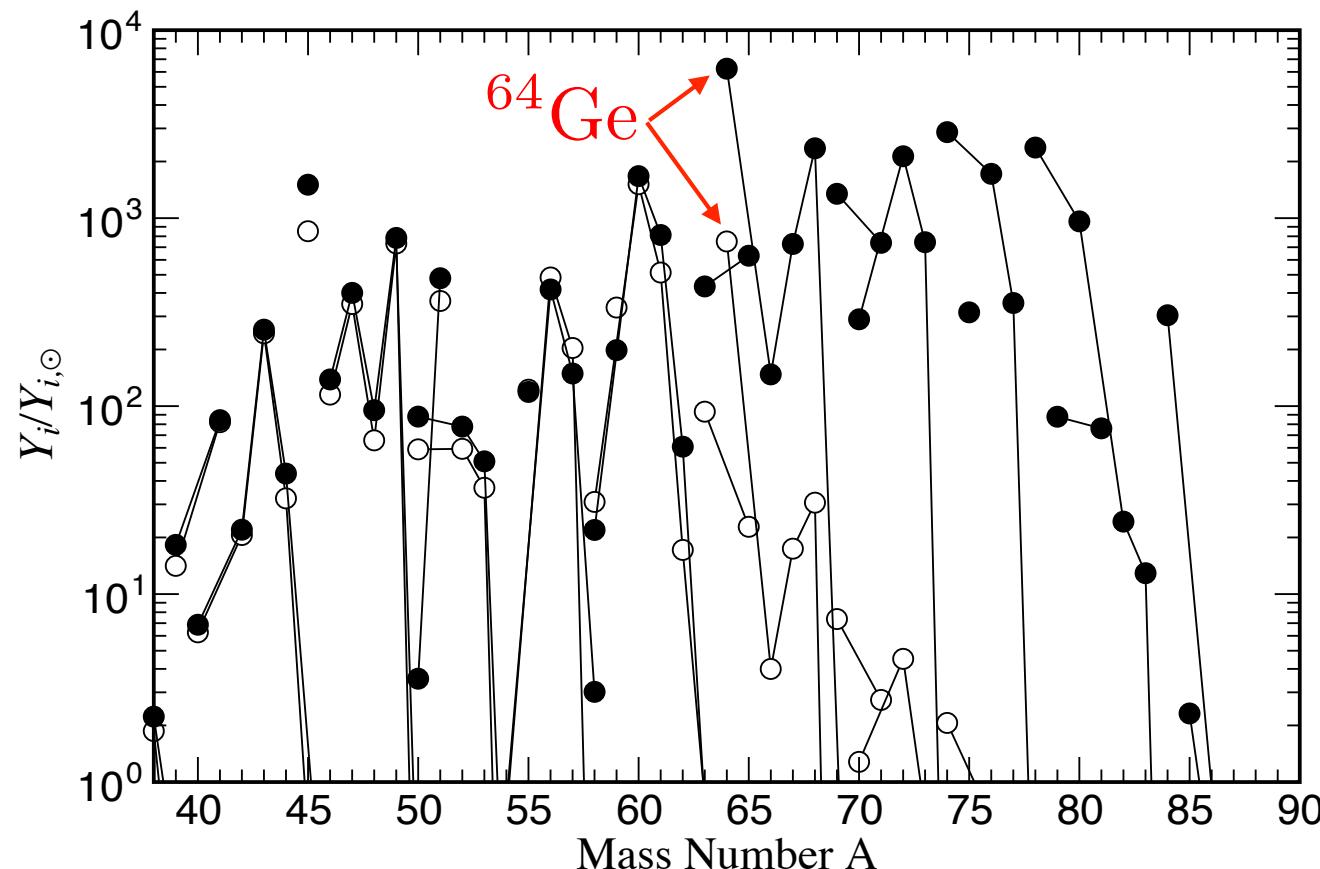
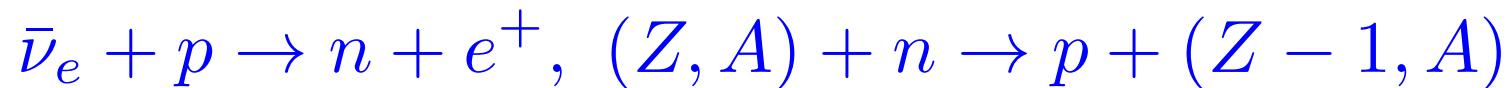




the vp-process in p-rich v-driven winds (Frohlich et al. 2006a,b; Prael et al. 2005,2006)

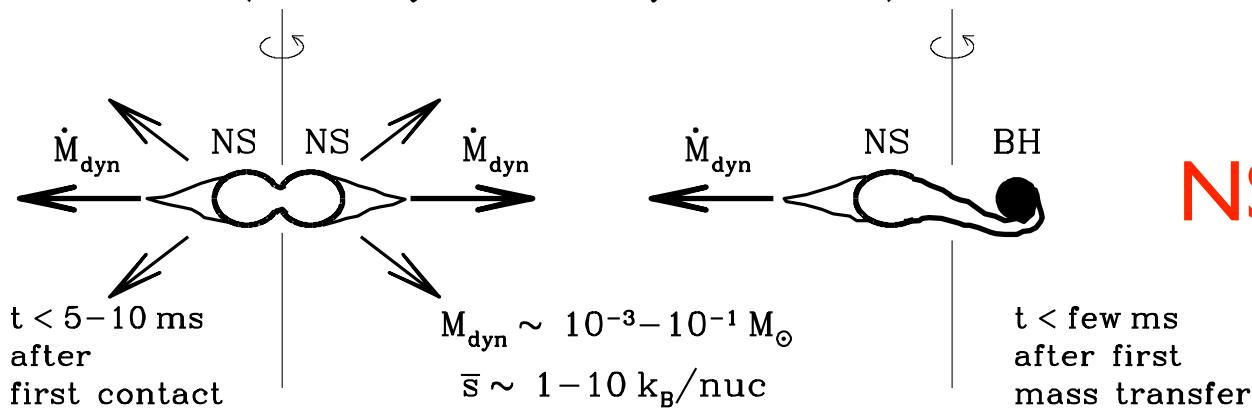
$(p, \gamma) \rightleftharpoons (\gamma, p)$ equilibrium \Rightarrow waiting point

break through waiting-point nuclei with slow beta decay:



Mass Loss Phases During NS-NS and NS-BH Merging

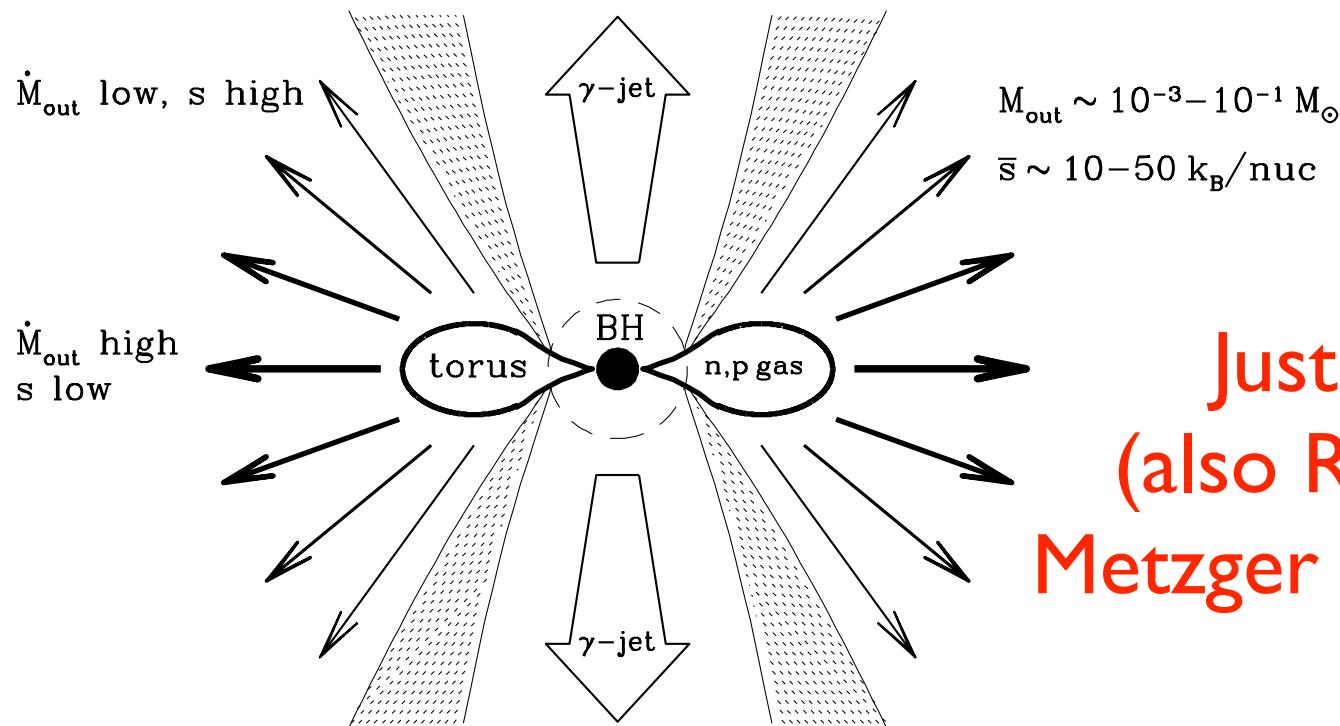
Merger Phase: Prompt/dynamical ejecta
(due to dynamic binary interaction)



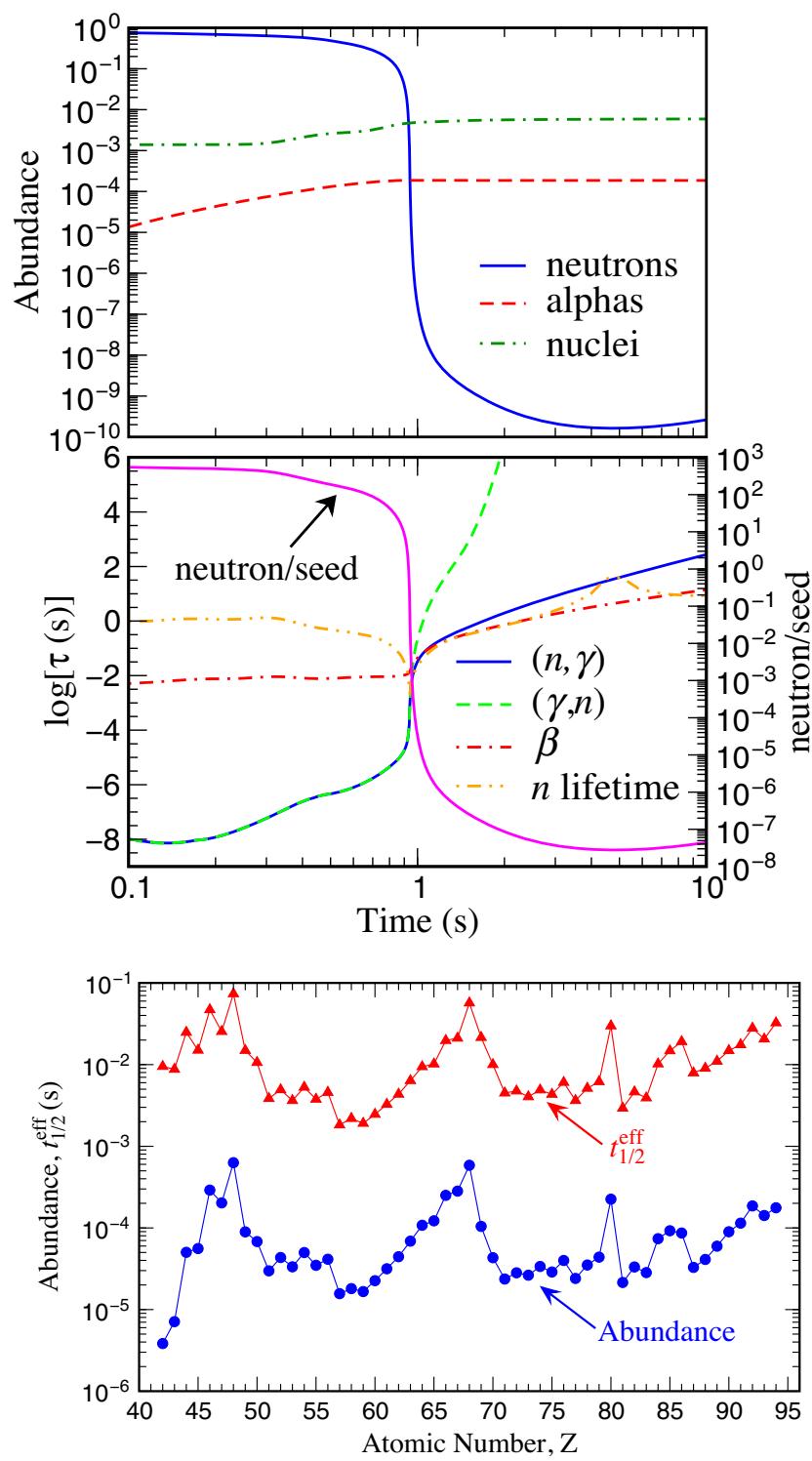
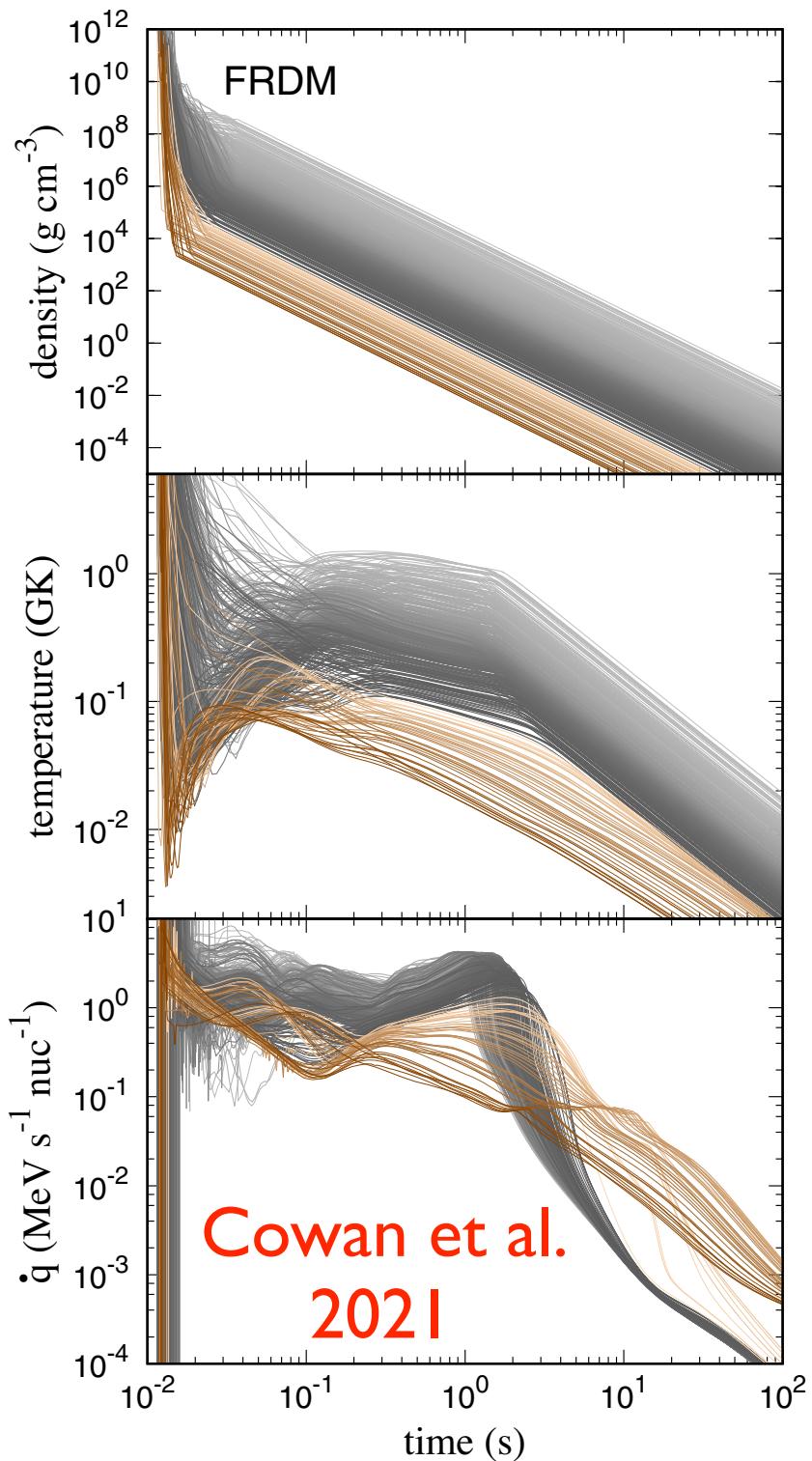
NS matter

BH-Torus Phase: Disk ejecta

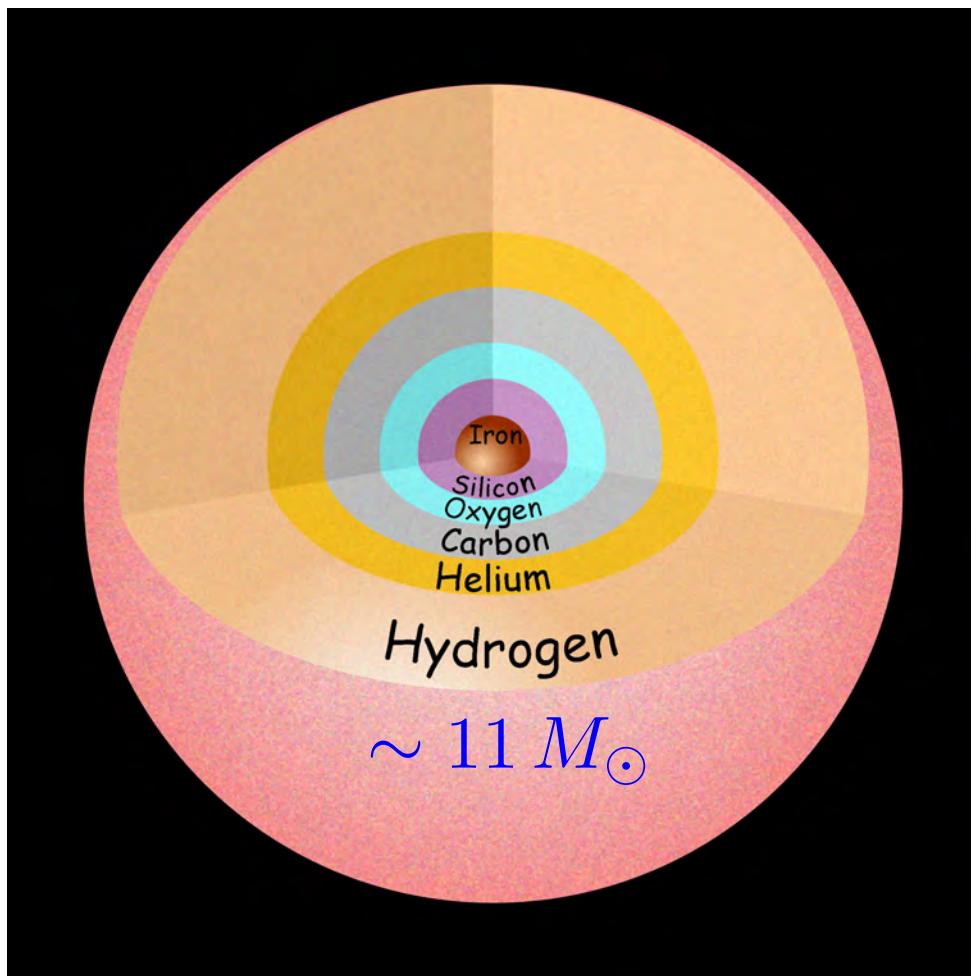
(due to ν heating, viscosity/magn. fields, recombination)



Just + 2014
(also Rosswog +;
Metzger +; Wanajo +)



Neutrino-Induced n Capture in He Shell of early SNe



neutron production by

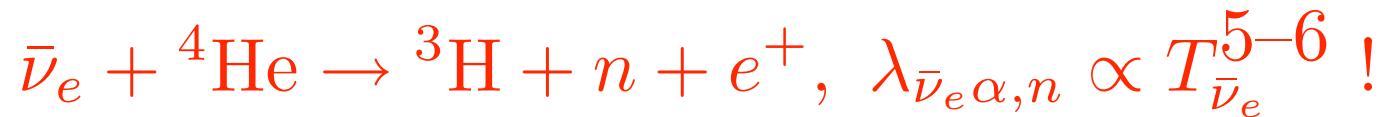


Epstein, Colgate, & Haxton 1988

neutron capture by ^{56}Fe

high n_n requires few ^{56}Fe

→ early SNe



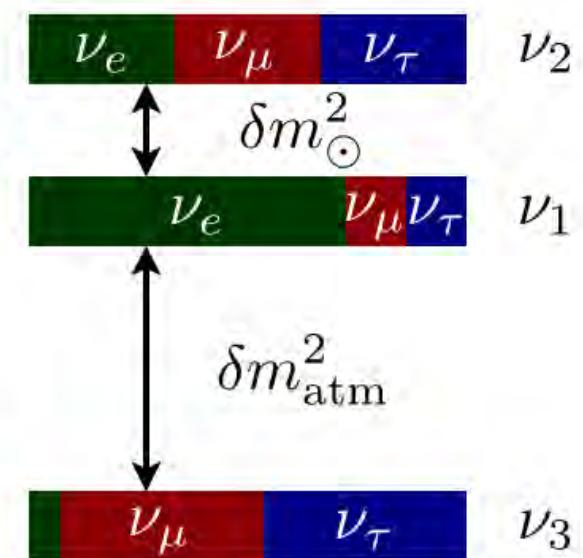
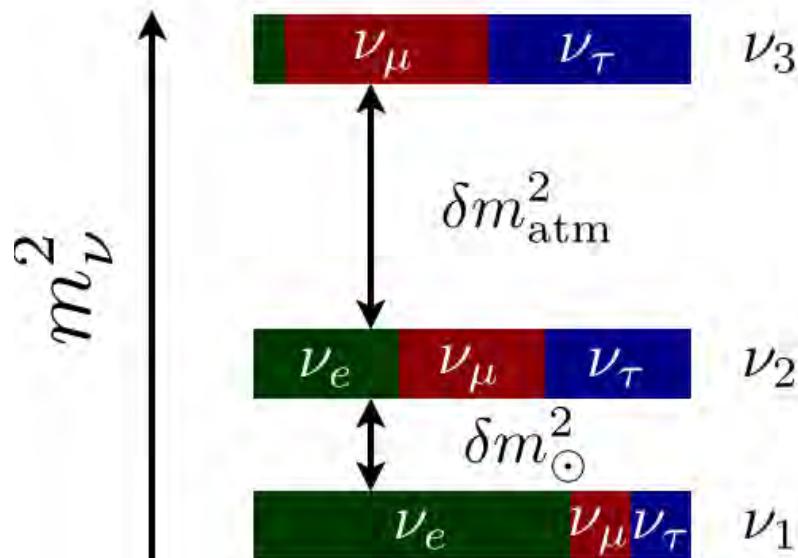
Banerjee, Haxton, & Qian 2011

Neutrino Spectra & Flavor Oscillations

$$T_{\nu_e} \sim 3\text{--}4 \text{ MeV}, T_{\bar{\nu}_e} \sim 4\text{--}5 \text{ MeV}, T_{\nu_{\mu,\tau}} = T_{\bar{\nu}_{\mu,\tau}} \sim 6\text{--}8 \text{ MeV}$$

normal mass hierarchy

inverted mass hierarchy



in supernovae

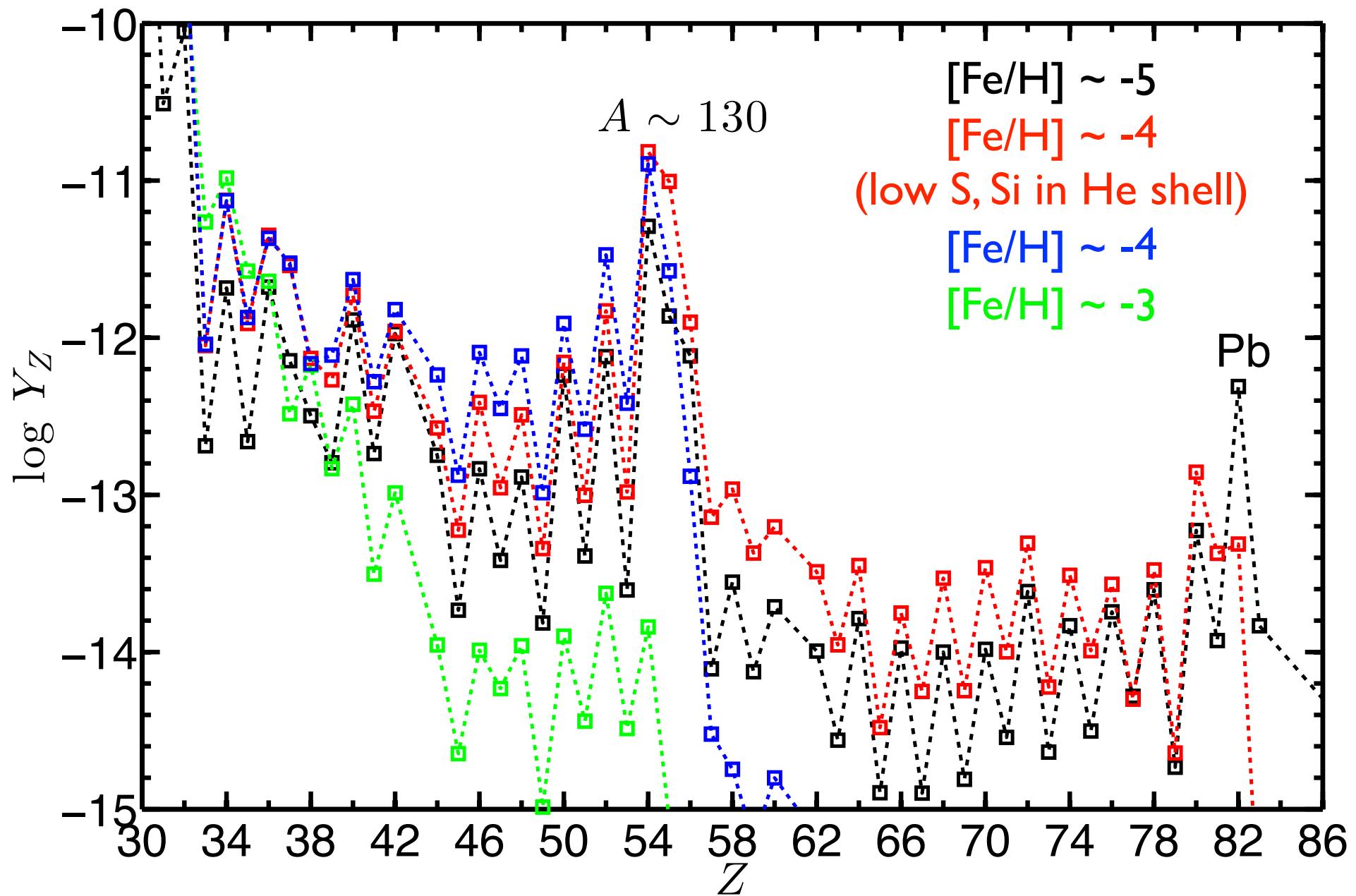


$$\nu_e \rightleftharpoons \nu_{\mu,\tau}$$

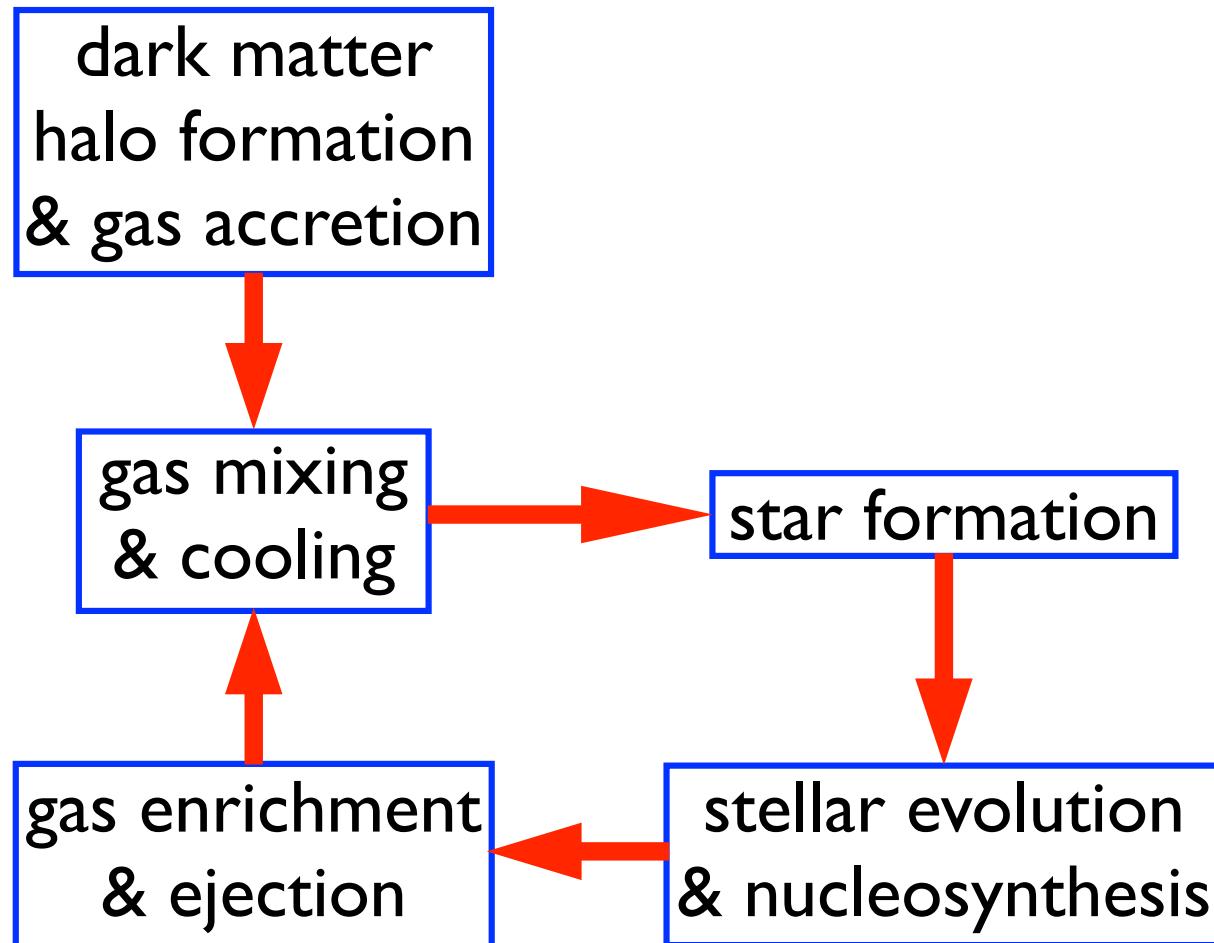
$$\bar{\nu}_e \rightleftharpoons \bar{\nu}_{\mu,\tau}$$



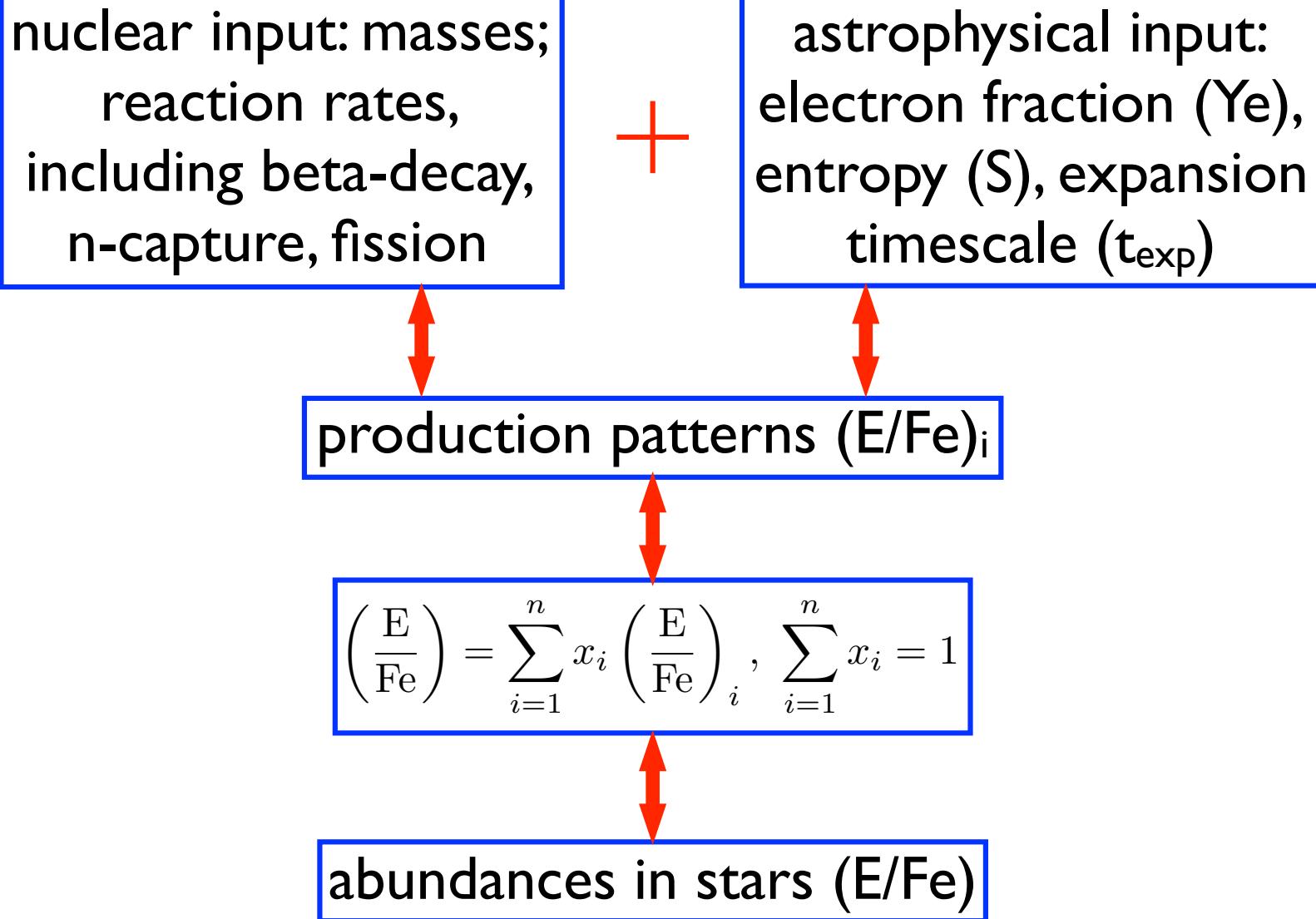
Banerjee, Qian, Heger, & Haxton 2016

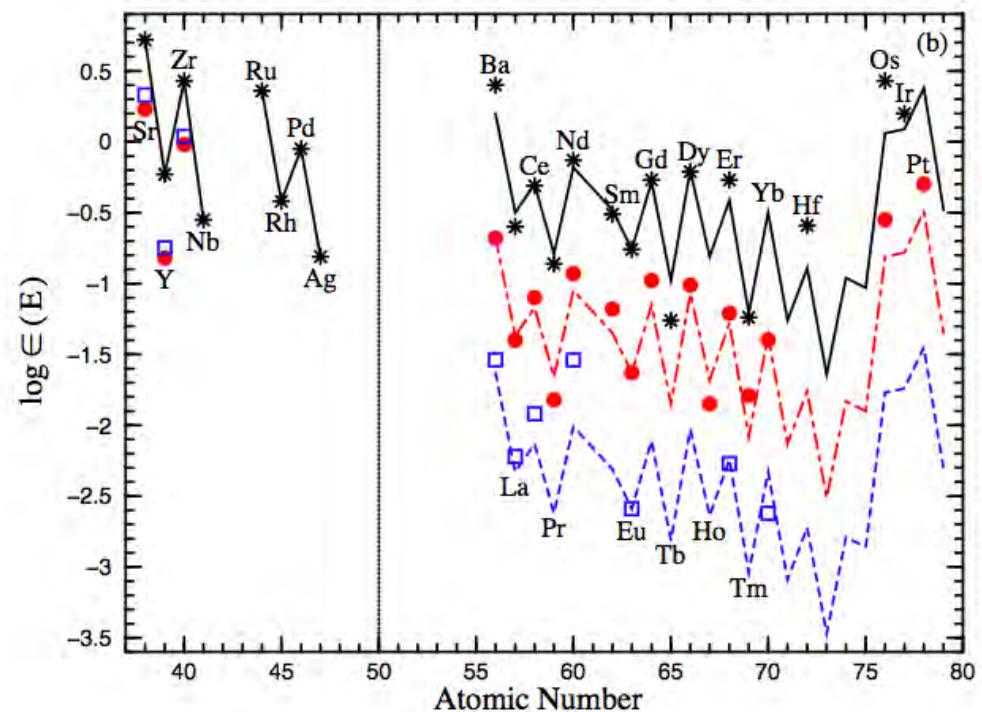
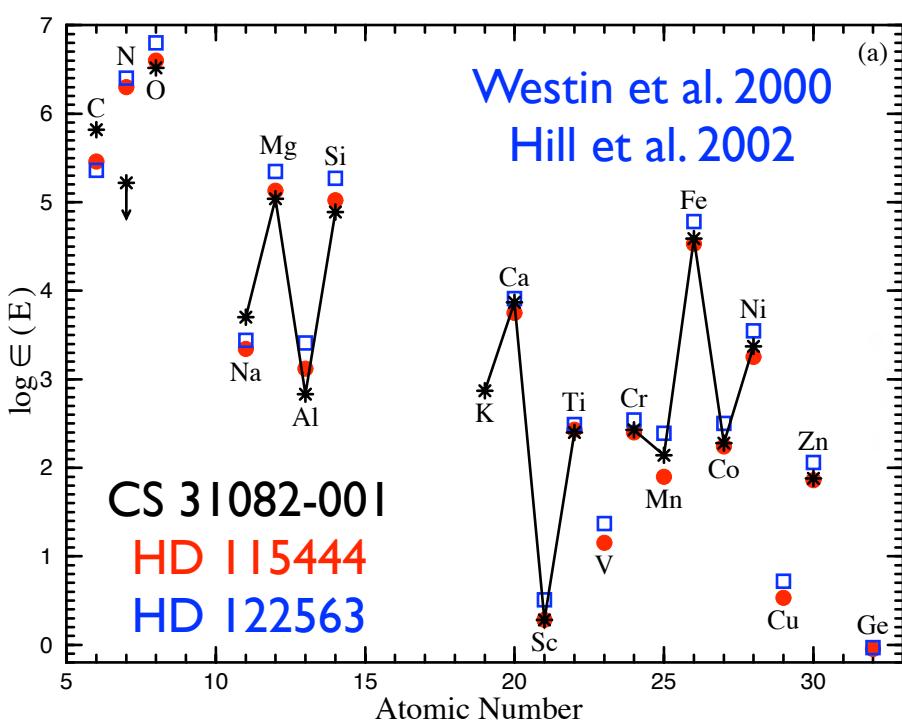
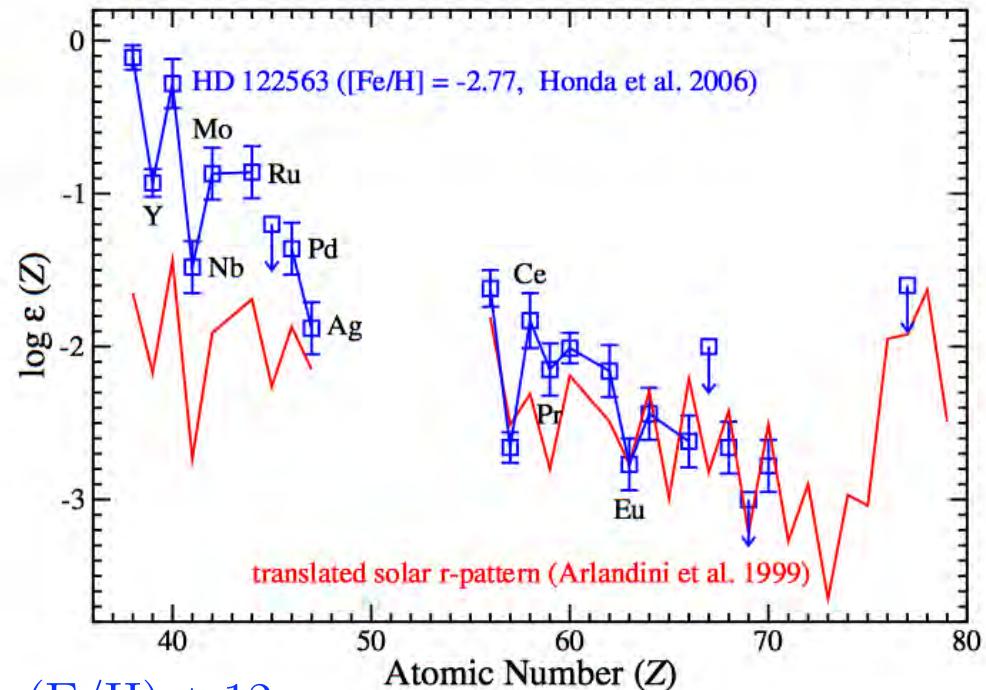
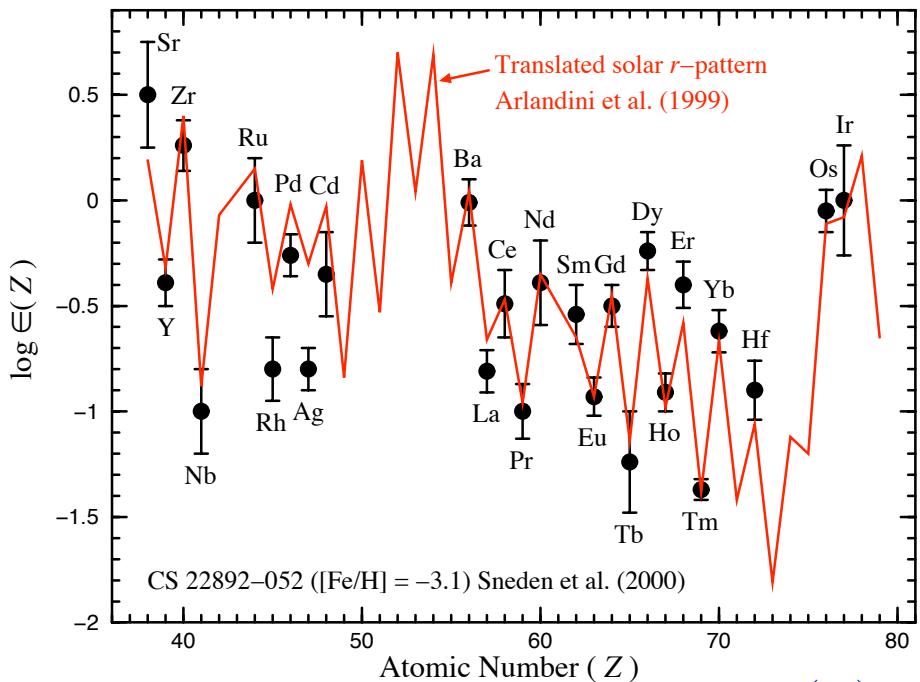


Ab Initio Model of Chemical Evolution



Data-Driven Nuclear Forensics





Early Chemical Evolution (\sim First Gyr)

dominant sources: Core-Collapse SNe & NS Mergers

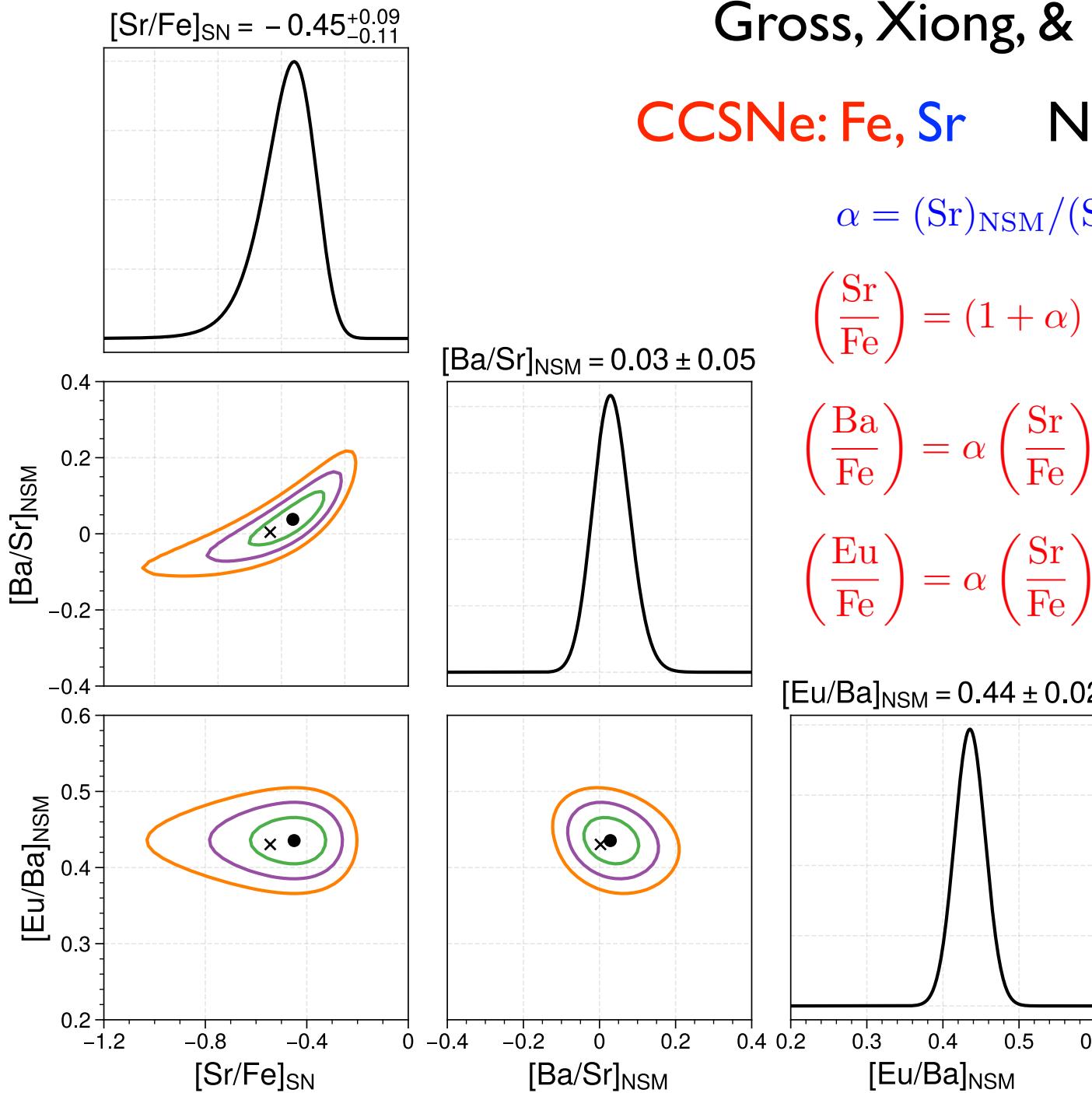
$$\left(\frac{E}{Fe} \right) = x \left(\frac{E}{Fe} \right)_1 + (1 - x) \left(\frac{E}{Fe} \right)_2$$

R-Process Alliance search for r-process-enhanced stars in the Galactic Halo: data on (Sr/Fe), (Ba/Fe), (Eu/Fe) for 195 stars
(Holmbeck et al. 2020)

$$[E/Fe] = \log(E/Fe) - \log(E/Fe)_{\text{sun}}$$

$$[Sr/Fe]_1 = -0.49, [Ba/Fe]_1 = -3.00, [Eu/Fe]_1 = -0.77 \\ [Sr/Fe]_2 = 0.90, [Ba/Fe]_2 = 0.91, [Eu/Fe]_2 = 1.30$$

CCSNe: Fe, Sr NSMs: Sr, Ba, Eu



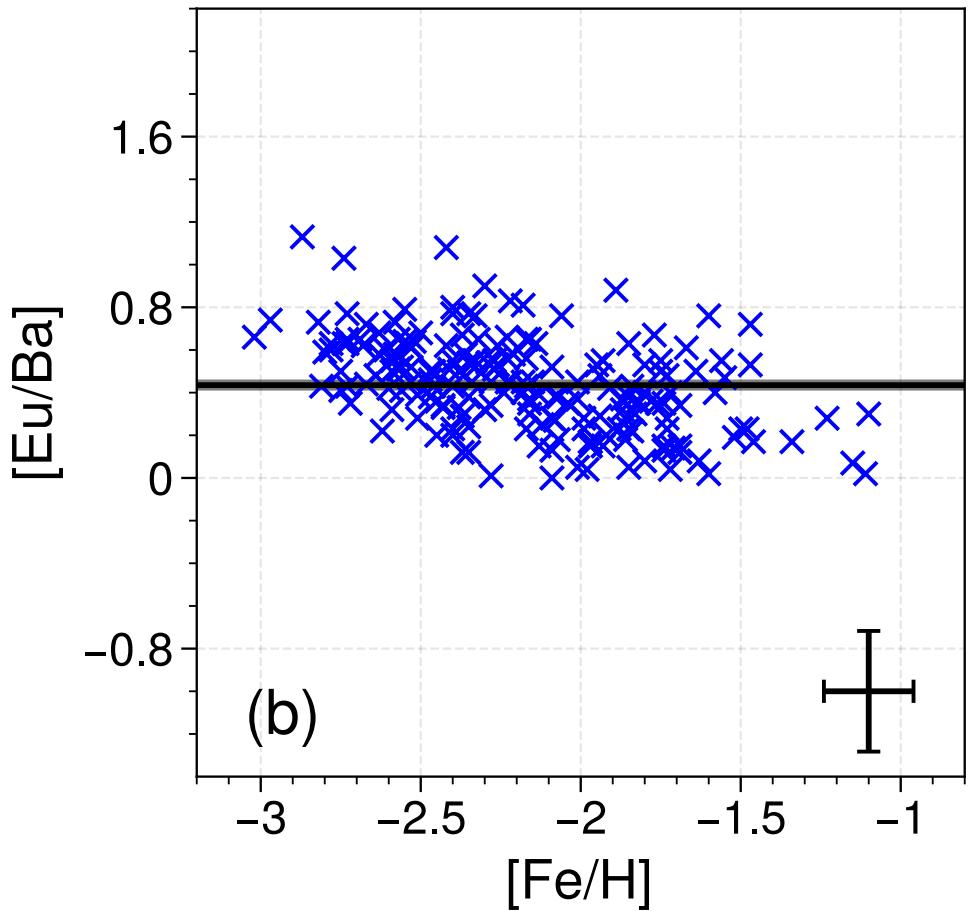
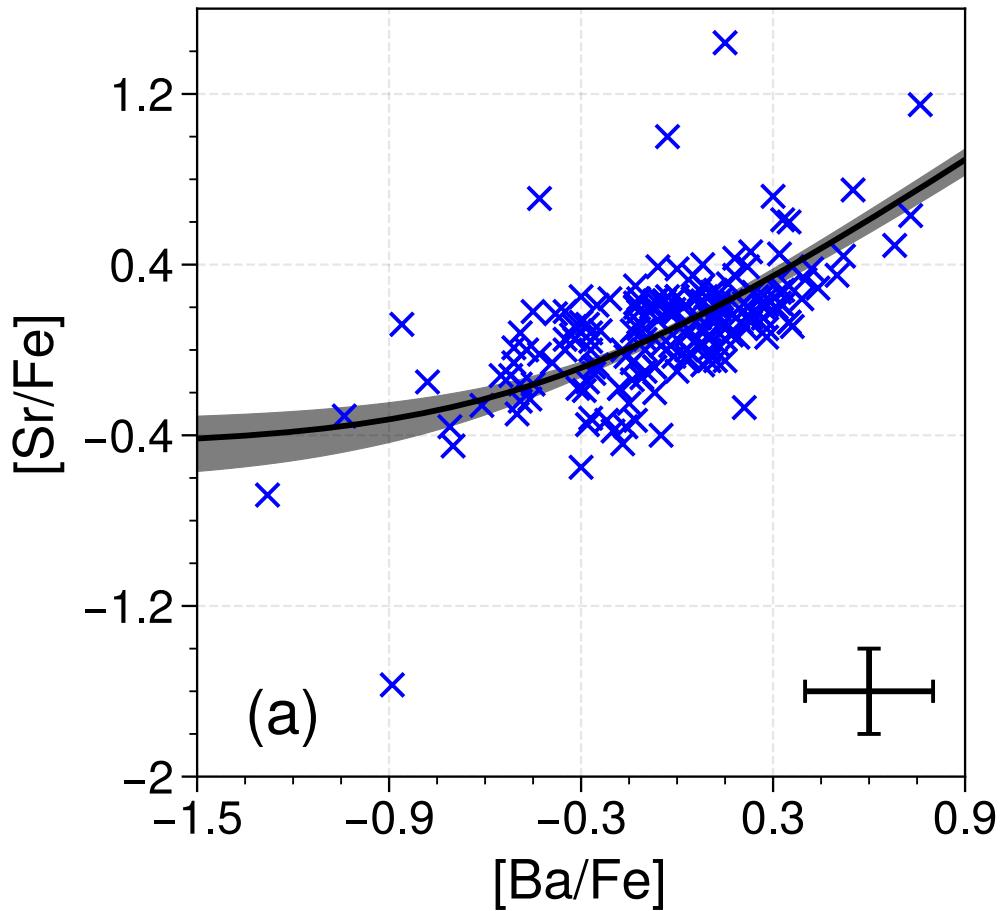
$$\alpha = (\text{Sr})_{\text{NSM}} / (\text{Sr})_{\text{SN}}$$

$$\left(\frac{\text{Sr}}{\text{Fe}} \right) = (1 + \alpha) \left(\frac{\text{Sr}}{\text{Fe}} \right)_{\text{SN}}$$

$$\left(\frac{\text{Ba}}{\text{Fe}} \right) = \alpha \left(\frac{\text{Sr}}{\text{Fe}} \right)_{\text{SN}} \left(\frac{\text{Ba}}{\text{Sr}} \right)_{\text{NSM}}$$

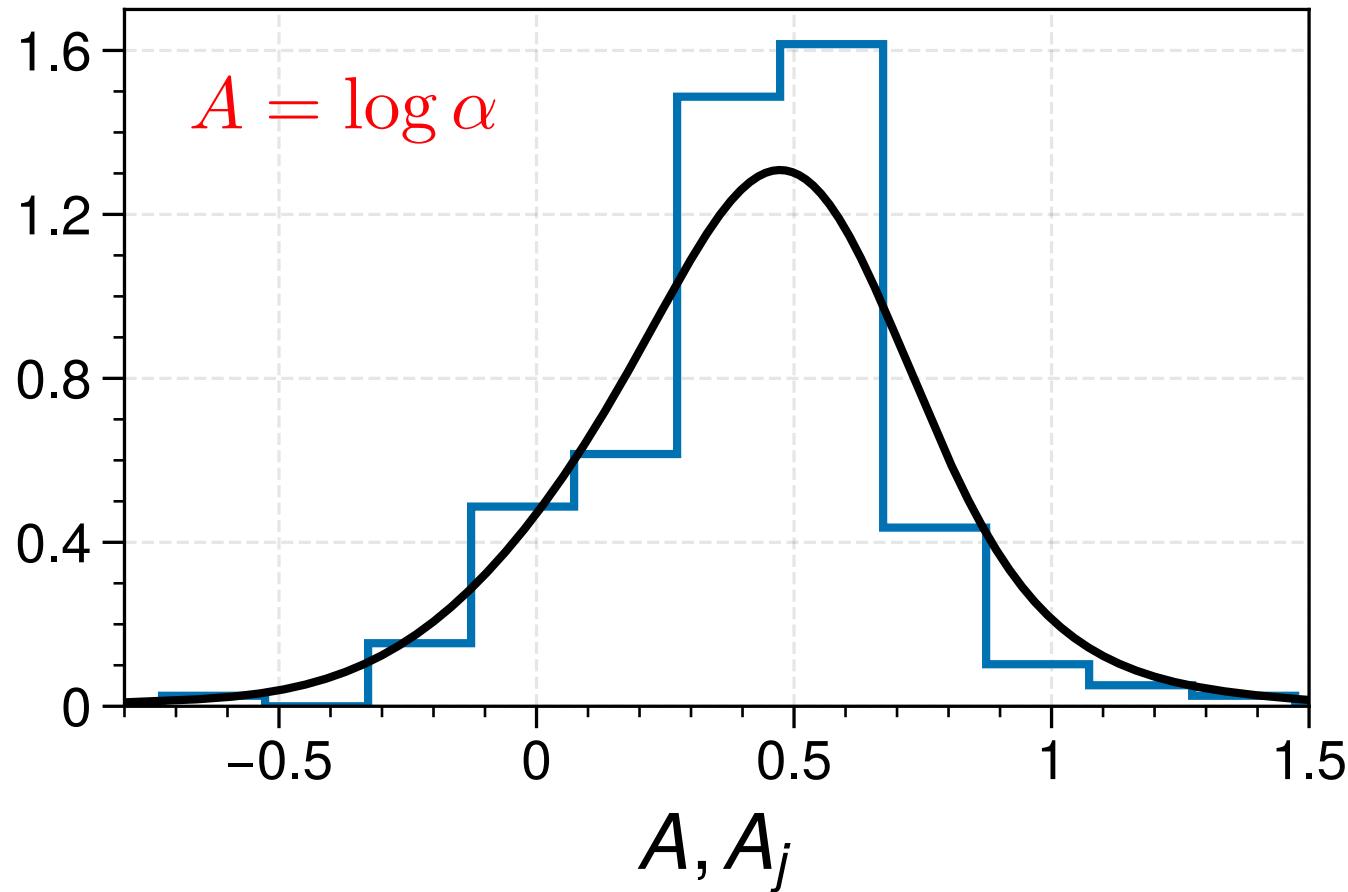
$$\left(\frac{\text{Eu}}{\text{Fe}} \right) = \alpha \left(\frac{\text{Sr}}{\text{Fe}} \right)_{\text{SN}} \left(\frac{\text{Ba}}{\text{Sr}} \right)_{\text{NSM}} \left(\frac{\text{Eu}}{\text{Ba}} \right)_{\text{NSM}}$$

Gross, Xiong, & Qian 2023



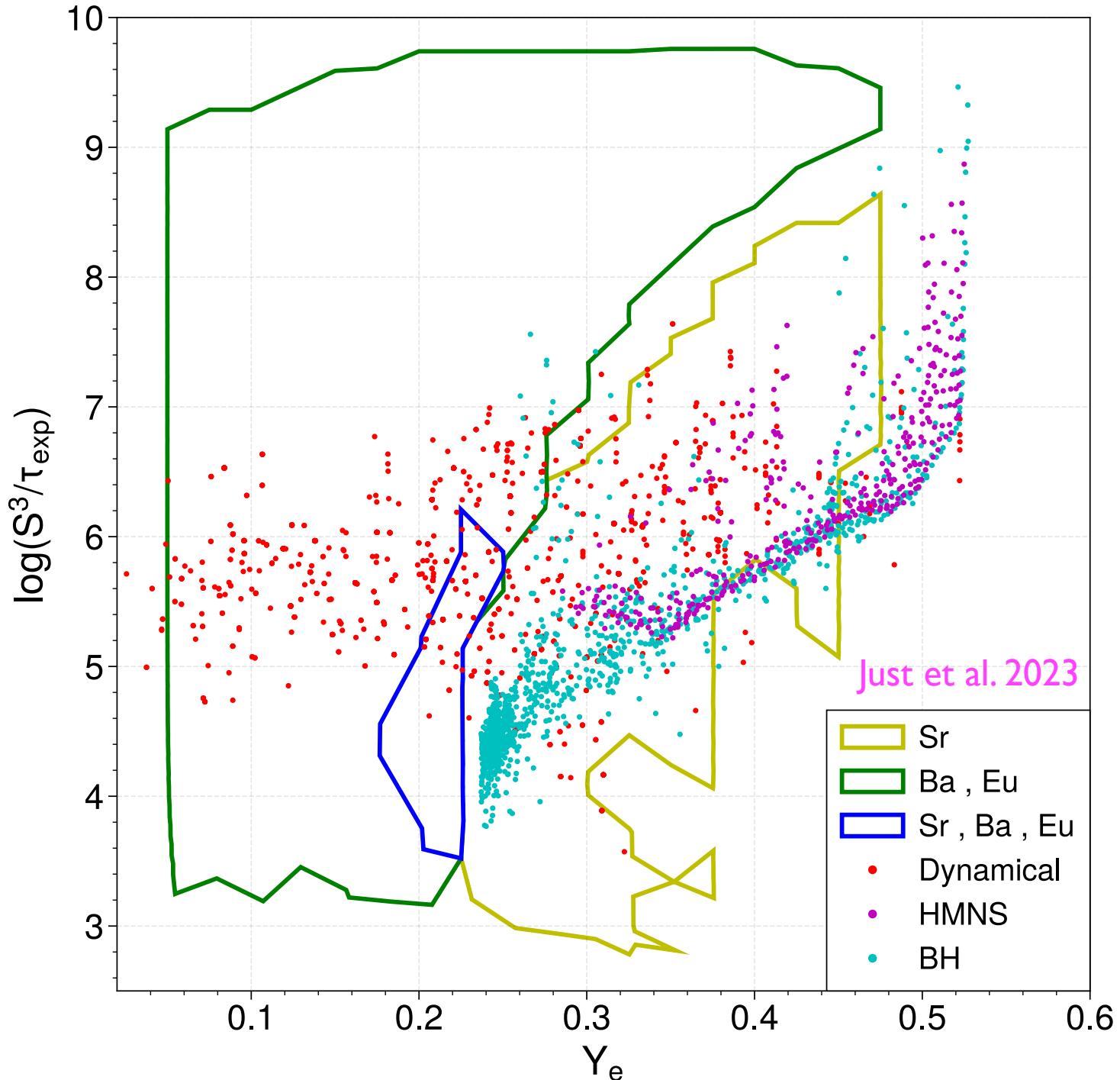
$$\left(\frac{{\rm Sr}}{{\rm Fe}} \right) = \left(\frac{{\rm Sr}}{{\rm Fe}} \right)_{{\rm SN}} + \left(\frac{{\rm Sr}}{{\rm Ba}} \right)_{{\rm NSM}} \left(\frac{{\rm Ba}}{{\rm Fe}} \right)$$

Gross, Xiong, & Qian 2023



$$\alpha_{\odot} \approx \frac{(\text{Sr/Eu})_{\text{NSM}} (\text{Eu})_{\odot}}{(\text{Sr/Fe})_{\text{SN}} (\text{Fe})_{\odot, \text{SN}}} \approx 3 \Rightarrow (\text{Fe})_{\odot, \text{SN}} \approx (\text{Fe})_{\odot}/3$$

Gross, Xiong, & Qian 2023



Gross, Xiong, & Qian 2023

