

CORRELATIONS AND DISTINGUISHABILITY CHALLENGES IN SUPERNOVA MODELS: INSIGHTS FROM FUTURE NEUTRINO DETECTOR

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Core-Collapse Supernovae

1. Core increases T without degenerating: Hydrodynamic burning of H, He, O, C, Ne, Si. Layers with material from a different burning stage.

2. Silicon burning shell: core reaches Chandrasekhar mass, and the material degenerates.

3. "Pre-supernova" stage: $T \sim 10^{10}\text{K}$ and $\rho \sim 10^{10}\text{g/cm}^3$

4. Pressure of the degenerate electrons can no longer stabilize the core → **COLLAPSE**

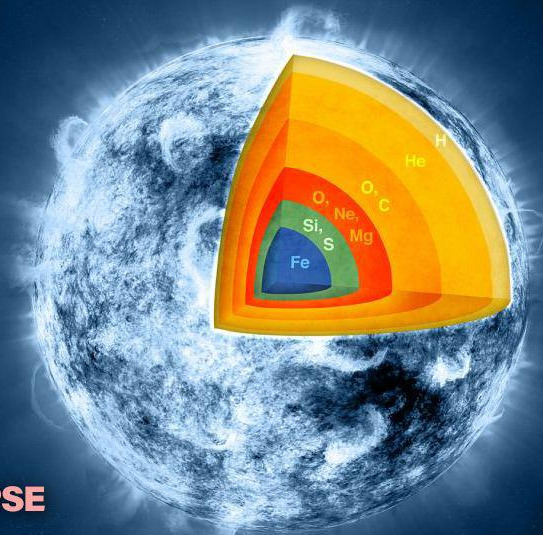
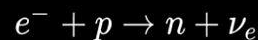
5. Mechanisms that accelerate the collapse:



6. Neutrino trapping and the nuclear saturation density is reached: STOPS THE IMPLOSION, Interactions begin to be repulsive → **EXPLOSION**

7. Collision between the outer layer that continue falling to the center with the inner core → **SHOCK WAVE**

8. Neutrinos of all flavors are produced



INITIAL FLUXES



QKE
NEUTRINO OSCILLATION
MODELLING



**FLUXES THAT
REACH THE EARTH**



INITIAL FLUXES



The initial spectral distribution is often parametrized by a parameter-fit that allows for deviations from a strictly thermal spectrum.

**SUPERNOVA SIMULATIONS
PROVIDE THE INDICATIVE VALUES
OF THE PARAMETERS**

(EoS?, unsuccessful explosions,
different approx, etc.)

QKE

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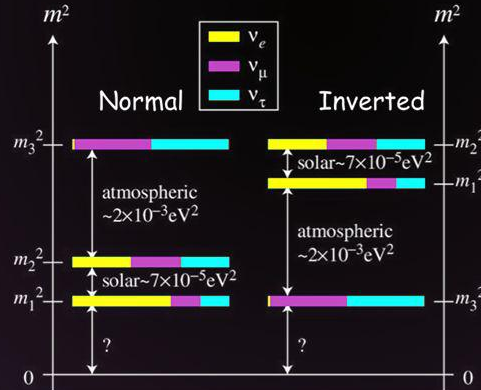
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QKE

NEUTRINO OSCILLATION MODELLING



NEUTRINO MASS ORDERING PROBLEM



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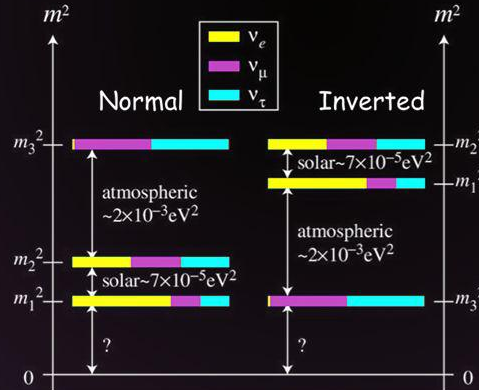
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QKE

NEUTRINO OSCILLATION MODELLING



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FLUXES THAT REACH THE EARTH



DETECTOR CHARACTERISTICS

DETECTION CHANNEL

RELEVANT CROSS-SECTIONS

UPCOMING GALACTIC SUPERNOVA: OPPORTUNITY TO COMPARE MODEL PREDICTIONS WITH OBSERVATIONS.

- ▶ **Current research heavily relies on computer simulations.**
- ▶ **Research groups worldwide conduct SN simulations**
→ different dimensionalities, progenitor masses ,
compositions, rotational velocities, EoS, approximations.

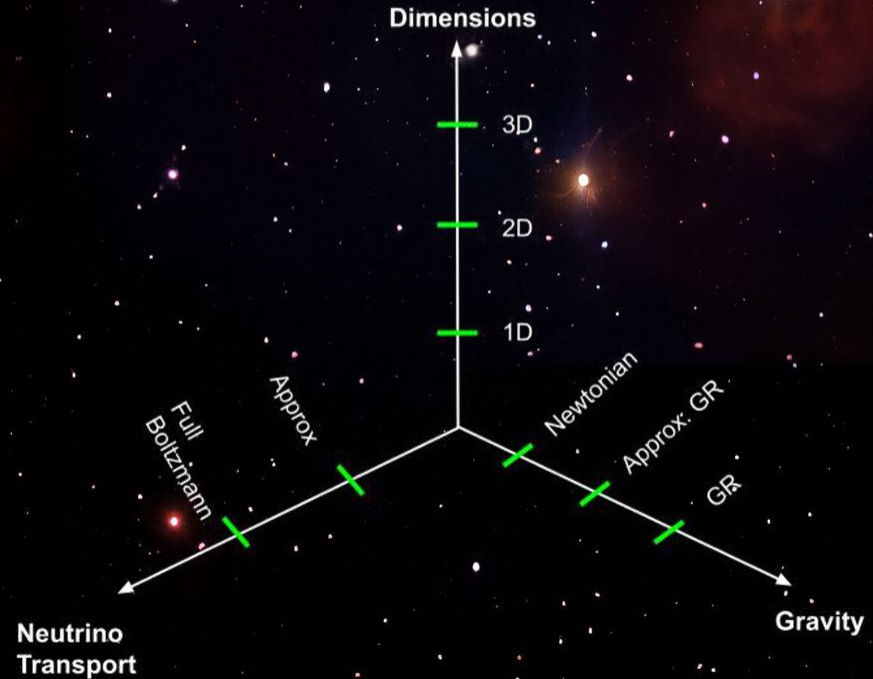
- ▶ **Vary in complexity:**

- Highly detailed models:**

- Substantial computational resources.
 - Simulate individual or few progenitors accurately

- Simplified studies:**

- Spherical symmetry
 - Approximations in neutrino transport
 - Reduced computational demands
 - Less realistic but allows large number of simulations



Goal

- ▶ **Conduct a comprehensive analysis of the expected neutrino signal in three future neutrino detectors for a large number of initial SN models incorporating simulations from various groups.**
- ▶ **Expected neutrino signal derived from Boltzmann-radiation-hydrodynamics models studied for the very first time.**
- ▶ **Examine correlation and patterns in signals.**
- ▶ **Assess the potential of planned neutrino detectors in discriminating between different SN models.**

SN Fluxes at the Earth

$$F_{\nu_\alpha}^0 = \frac{L_{\nu_\alpha}(t)}{4\pi d^2 \langle E_{\nu_\alpha}(t) \rangle} f_{\nu_\alpha}(E_{\nu_\alpha}, t) \quad \text{Un-oscillated fluxes}$$

$$f_{\nu_\alpha}(E_\nu, t) = \mathcal{N} \left(\frac{E_\nu}{\langle E_{\nu_\alpha}(t) \rangle} \right)^{\beta_{\nu_\alpha}(t)} \exp \left[-(\beta_{\nu_\alpha}(t) + 1) \frac{E_\nu}{\langle E_{\nu_\alpha}(t) \rangle} \right] \quad \frac{\langle E_{\nu_\alpha}^2 \rangle}{\langle E_{\nu_\alpha} \rangle^2} = \frac{2 + \beta_{\nu_\alpha}}{1 + \beta_{\nu_\alpha}} \quad \text{Normalized distribution}$$

$$F_{\nu_e} = P_e F_{\nu_e}^0 + (1 - P_e) F_{\nu_x}^0 ,$$

$$F_{\bar{\nu}_e} = \bar{P}_e F_{\bar{\nu}_e}^0 + (1 - \bar{P}_e) F_{\bar{\nu}_x}^0 ,$$

$$F_{\nu_x} = \frac{1}{2}(1 - P_e) F_{\nu_e}^0 + \frac{1}{2}(1 + P_e) F_{\nu_x}^0 ,$$

$$F_{\bar{\nu}_x} = \frac{1}{2}(1 - \bar{P}_e) F_{\bar{\nu}_e}^0 + \frac{1}{2}(1 + \bar{P}_e) F_{\bar{\nu}_x}^0 .$$

Neutrino fluxes at Earth (MSW)

$$\text{NMO} \rightarrow \begin{aligned} P_e &= |U_{e1}|^2 P_H P_L + |U_{e2}|^2 P_H (1 - P_L) + |U_{e3}|^2 (1 - P_H) , \\ \bar{P}_e &= |U_{e1}|^2 , \end{aligned}$$

$$\text{IMO} \rightarrow \begin{aligned} P_e &= |U_{e1}|^2 P_L + |U_{e2}|^2 (1 - P_L) , \\ \bar{P}_e &= |U_{e1}|^2 \bar{P}_H + |U_{e3}|^2 (1 - \bar{P}_H) , \end{aligned}$$

Expected signal in underground detectors

$$\frac{dN(t)}{dt} = N_{tar} \int_{E_{th_d}}^{E_{max}} dE \int_{E_{th_r}}^{\infty} dE_{\nu} \int_0^{\infty} dE' \epsilon(E') \frac{dF_{\nu}}{dE_{\nu}}(E_{\nu}, t) \frac{d\sigma}{dE'}(E_{\nu}, E') G(E, E', \delta),$$

Expected signal in underground detectors

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Detector info

Thresholds, Fiducial Vols., efficiency, smearing, resolution, quenchings, etc.

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Detector info

Thresholds, Fiducial Vols., efficiency, smearing, resolution, quenchings, etc.

Neutrino Fluxes

Oscillation schemes, different flavors for different reactions

Expected signal in underground detectors

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Detector info

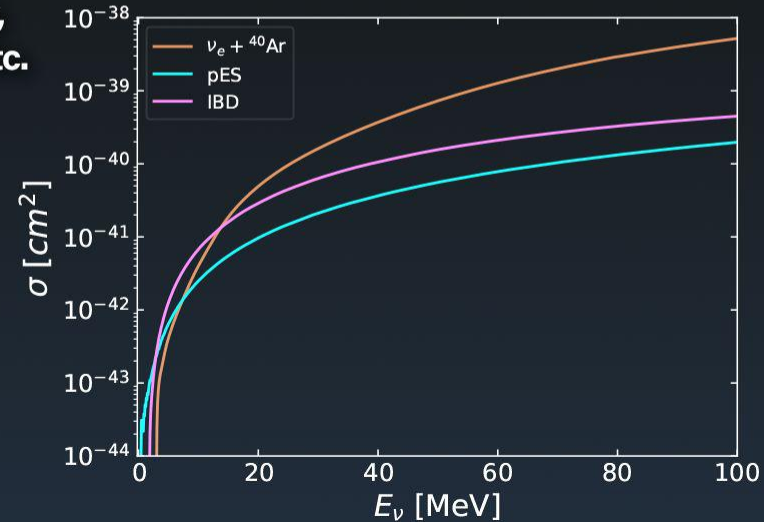
Thresholds, Fiducial Vols., efficiency, smearing, resolution, quenchings, etc.

Neutrino Fluxes

Oscillation schemes, different flavors for different reactions

Nuclear physics → XS

CC and NC cross-sections



HYPER KAMIOKANDE

Next-generation water Cherenkov detector.
188kt of ultra-pure water.

Dominant channel: $\bar{\nu}_e + p \rightarrow n + e^+$

$$E_{th_r} = 1.806 \text{ MeV} \quad E_{th_d} = 7 \text{ MeV} \quad E_{max} = 100 \text{ MeV}$$

$$\delta/\text{MeV} = -0.0839 + 0.349\sqrt{E'/\text{MeV}} + 0.0397E'/\text{MeV}$$

DUNE

Four 10-kt liquid argon time projection chambers
(LArTPCs) Dominant channel: $\nu_e + {}^{40}\text{Ar} \rightarrow e^- + {}^{40}\text{K}^*$

$$E_{th_d} = 5 \text{ MeV} \quad E_{max} = 100 \text{ MeV}$$

JUNO

Next-generation liquid scintillator
C₆H₅CnH_{2n+1} where n is 95% 9-12, 5% 13-14



$$\frac{d\sigma}{dE'}(E_\nu, E') = \frac{G_F^2 m_p}{2\pi E_\nu^2} [(C_V \pm C_A)^2 E_\nu^2 + (C_V \mp C_A)(E_\nu - E')^2 - (C_V^2 - C_A^2)m_p E']$$



$$E_\nu^{min} = \frac{E' + \sqrt{E'(E' + 2m_p)}}{2} \quad \delta = 3\% \sqrt{E'} \quad E_{th_d} = 0.2 \text{ MeV} \quad E_{max} = 60 \text{ MeV}$$

NUMERICAL TOOLS DEVELOPED

BASED ON SNTTOOLS AND SNEWPY + OUR CONTRIBUTIONS

Neutrino signal calculations:

Neutrino signal as function of E and t

- + new detectors
- + new cross sections
- + smearing
- + detector characteristics
(efficiency, thresholds, quenchings,
fiducial volumes, etc)
- + different oscillation schemes
- + different initial SN models



Statistical analysis:

Montecarlo simulation

- + likelihoods
- + bayes factors
- + model distinguishability analysis

SN Models

2D Boltzmann-radiation-hydrodynamics models (9)

[suyimoshi, yamada, Nagakura]

EoS:

Furusawa and Togashi (FT),

Furusawa and Shen (FS),

Lattimer and Swesty (LS)

11.2, 15, 27 M_{\odot}

With and without rotation

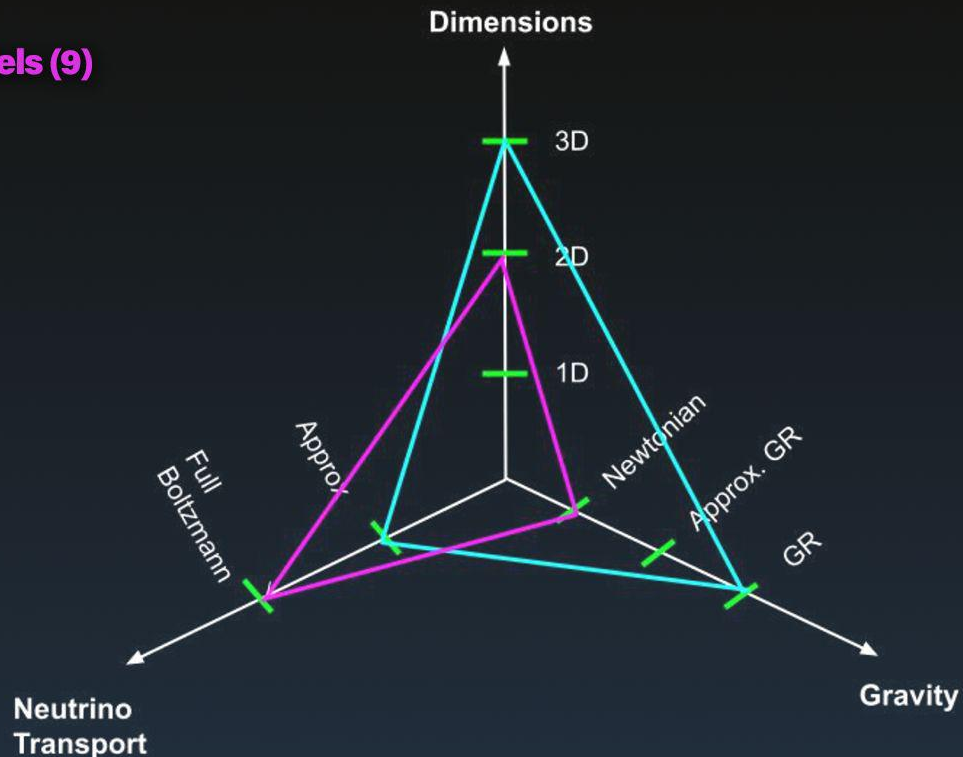
3D Princeton models (9)

[Burrows, Vartanyan]

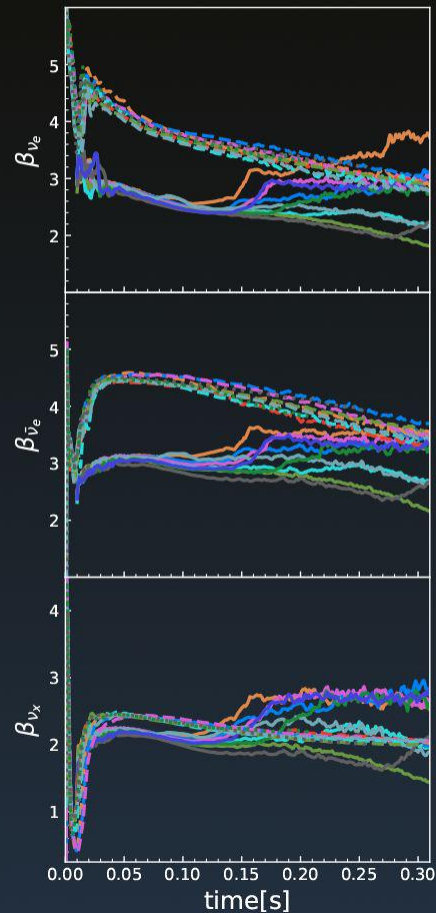
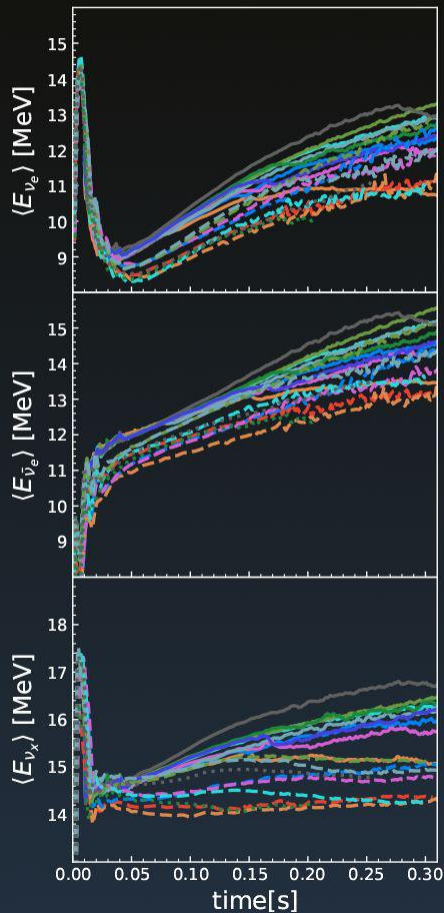
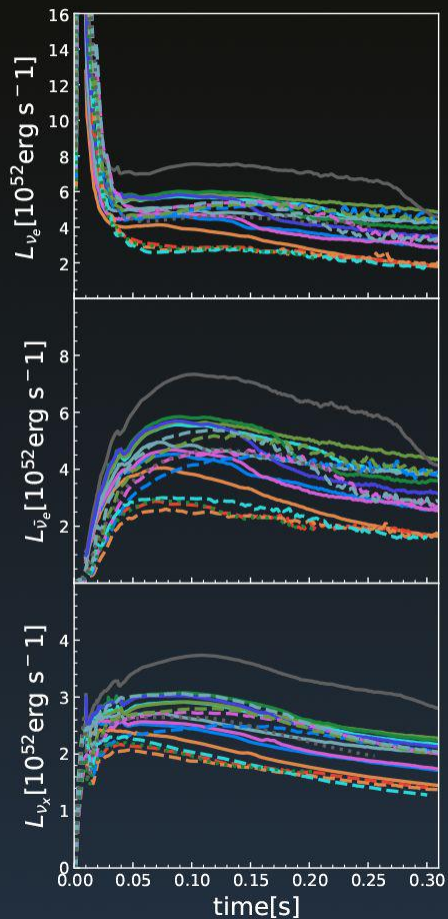
EoS: SFHo

9,10,12,13,14,15,19,25,60 M_{\odot}

Non rotationals



SN Models spectral parameters



Neutrino signal results

18 SN models:

3D Princeton (Fornax)

2D Boltzmann

Detectors:

HK - IBD channel

DUNE - nue + Ar channel

JUNO - IBD + pES channel

Neutrino mixing:

No oscillation(NO)

NMO

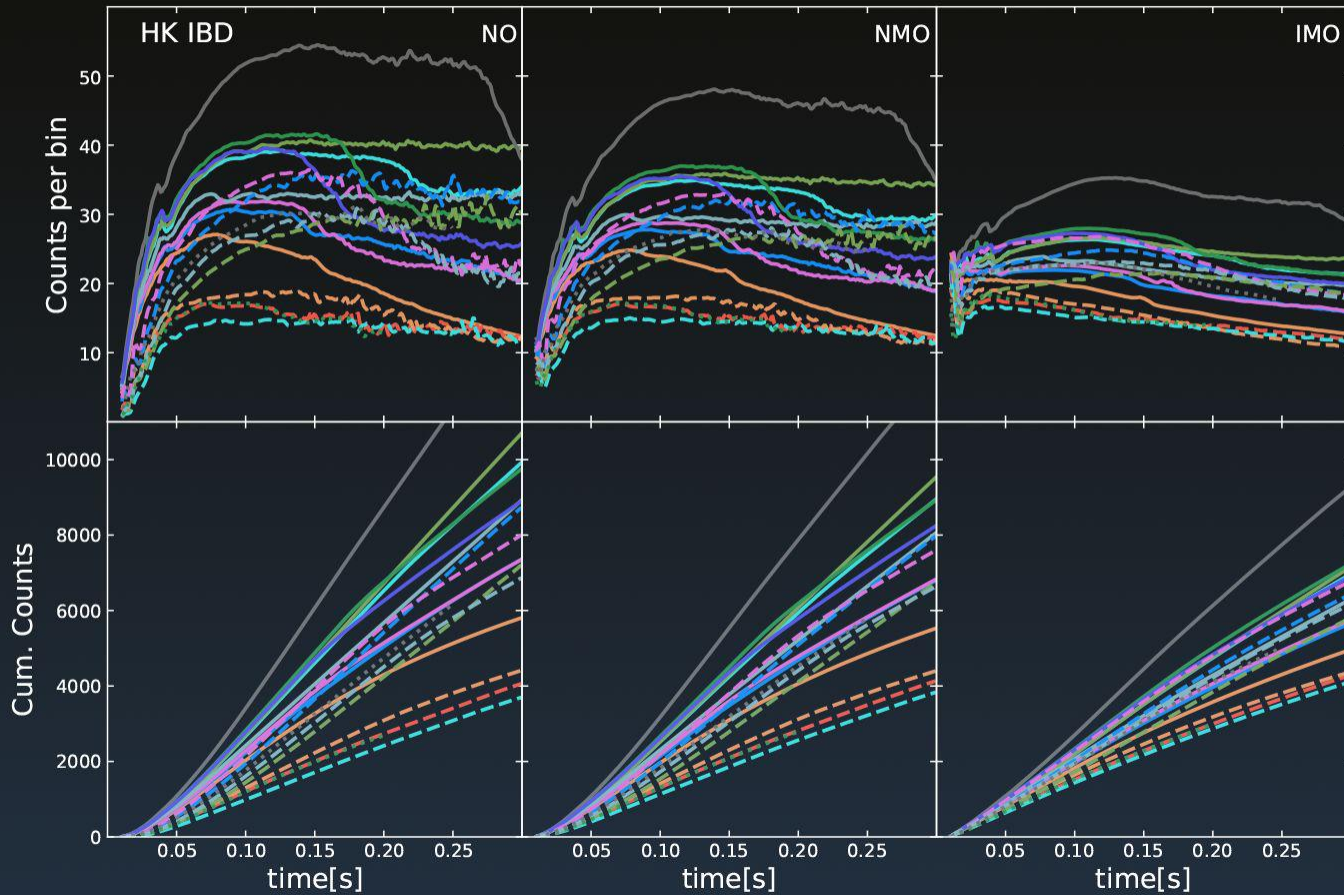
IMO

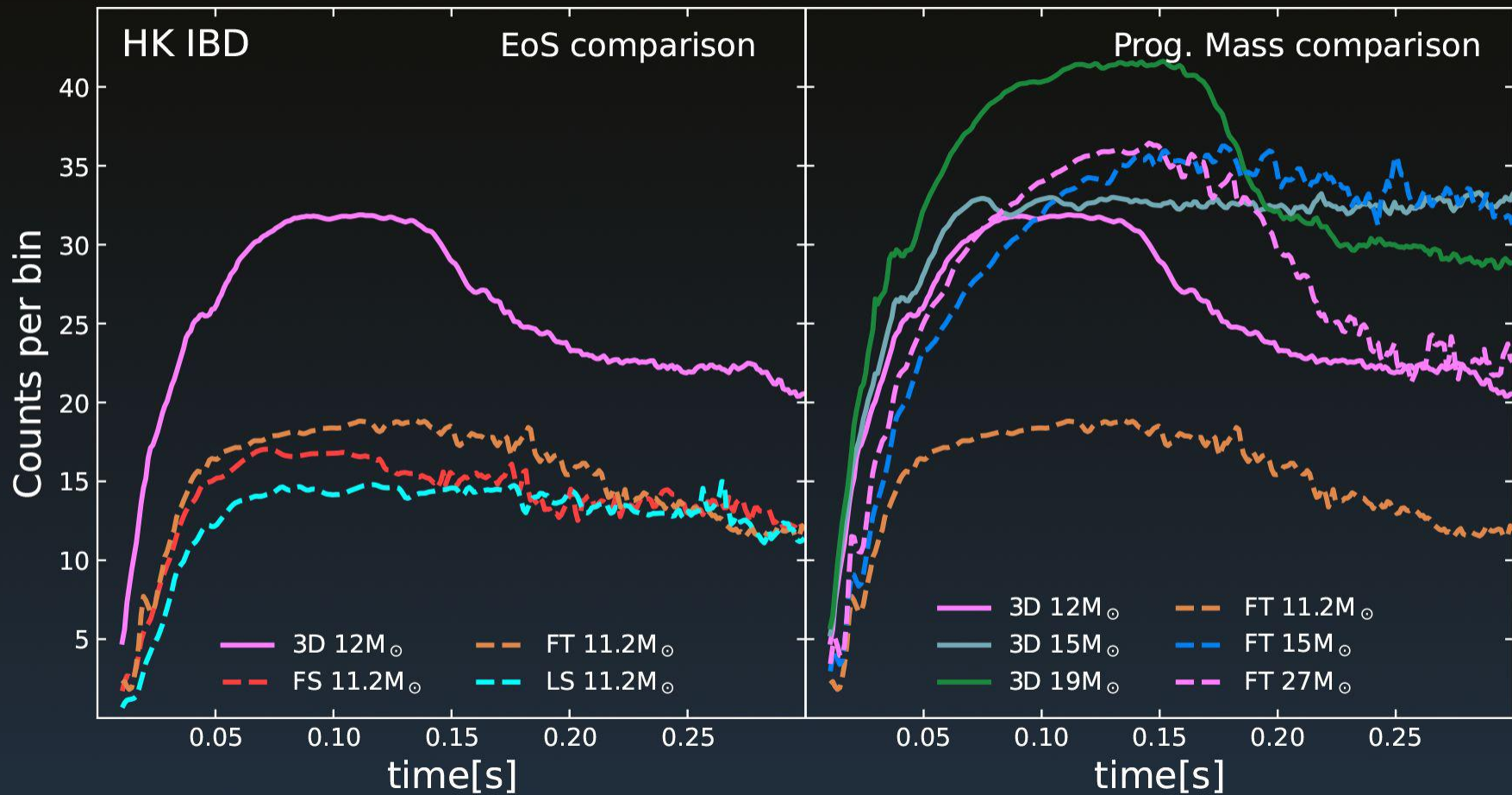
Event distribution with time or E

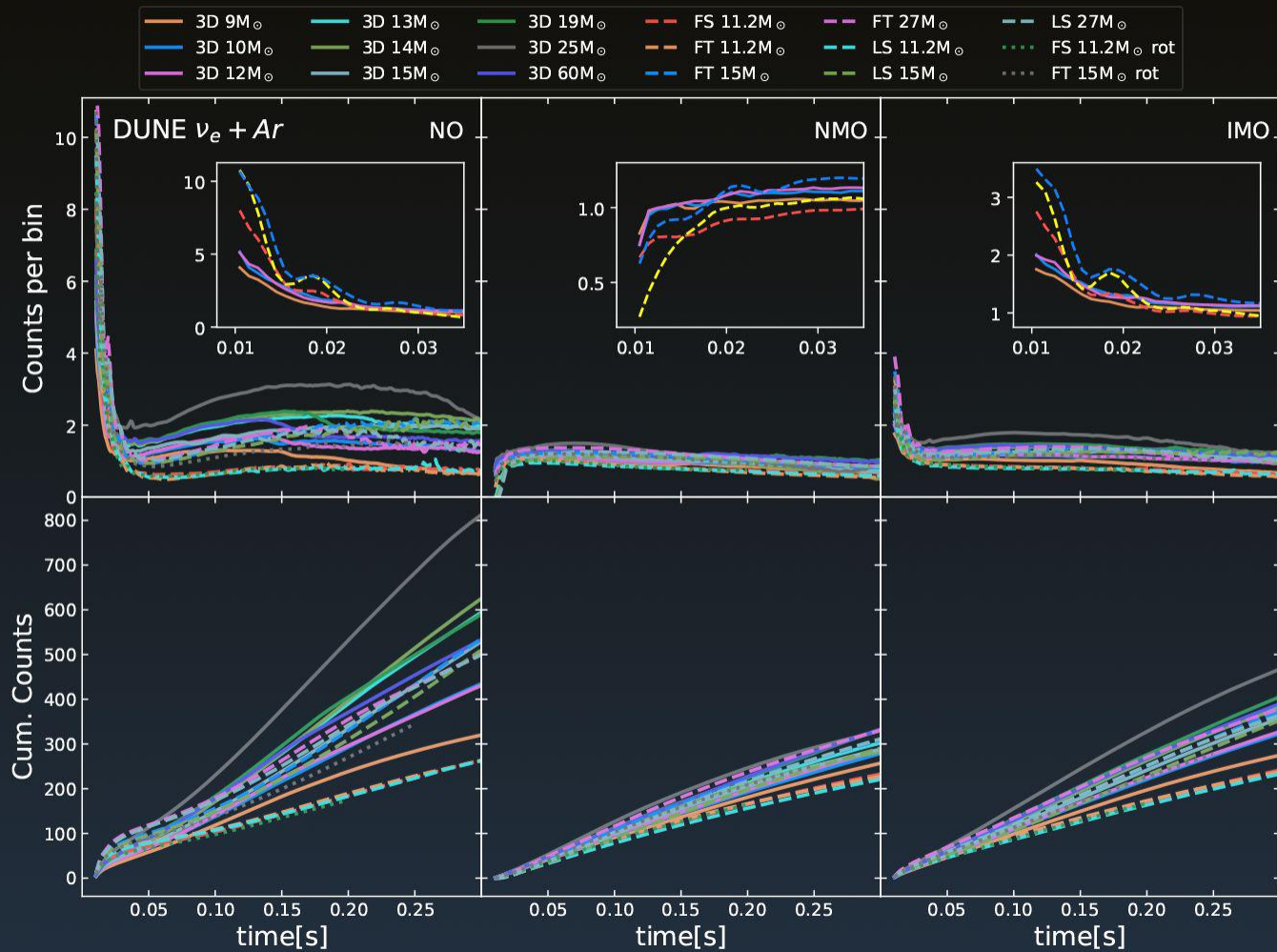
Cumulative number of events

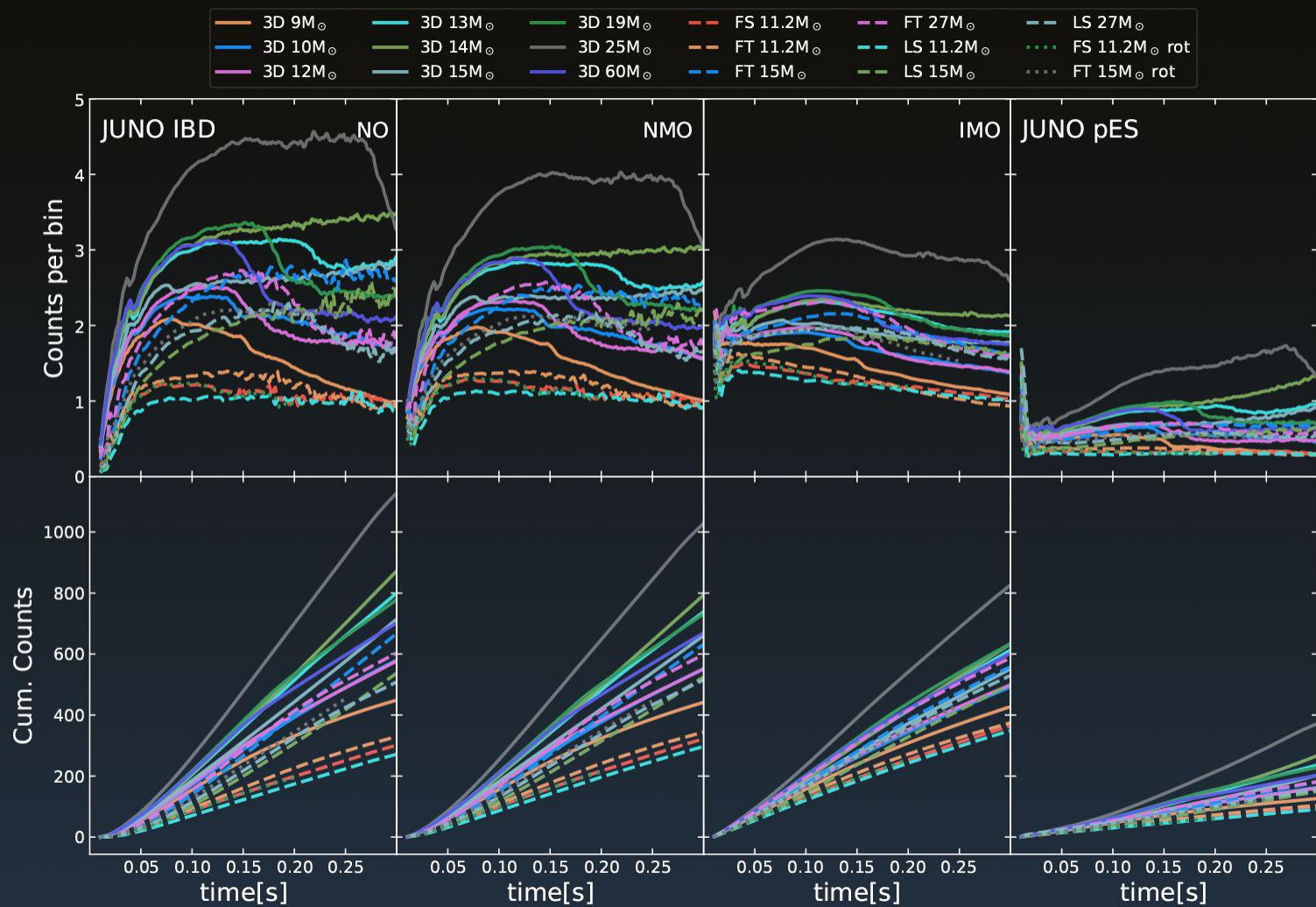
Total number of events

($T_{fin}=300\text{ms}$ and $d=10\text{kpc}$)









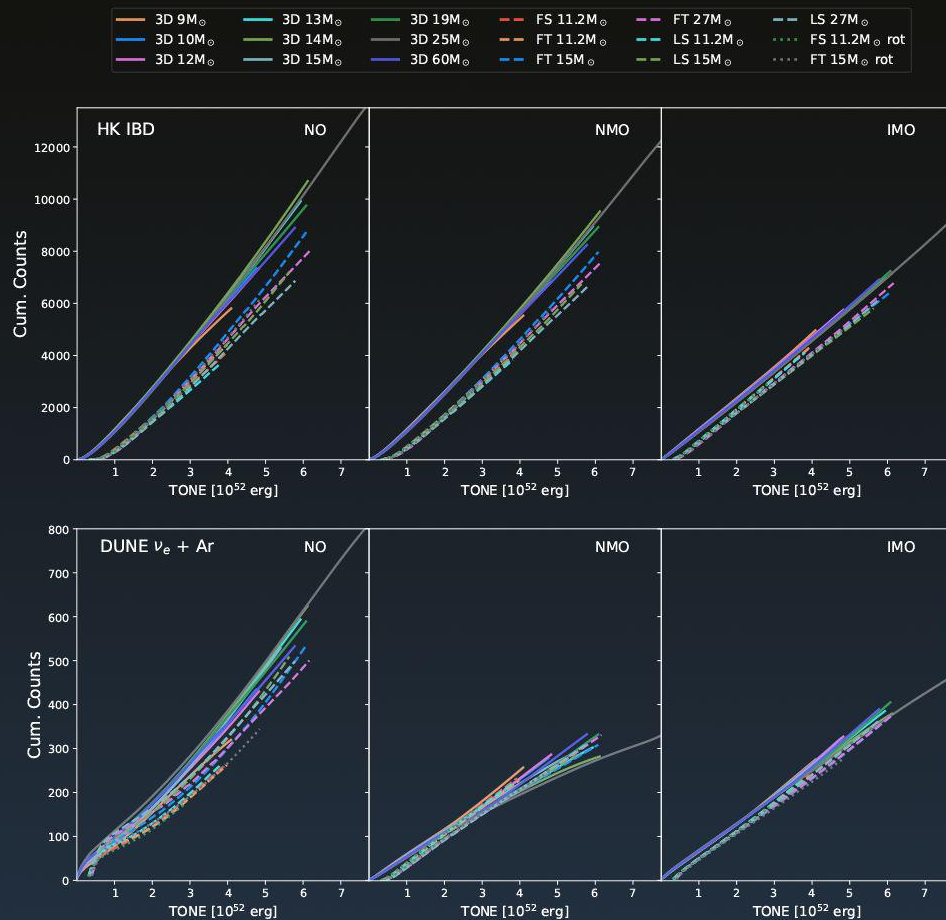
Neutrino signal results useful correlation

TONE: energy- and flavor-integrated time-cumulative neutrino radiation up to a given post-bounce time.

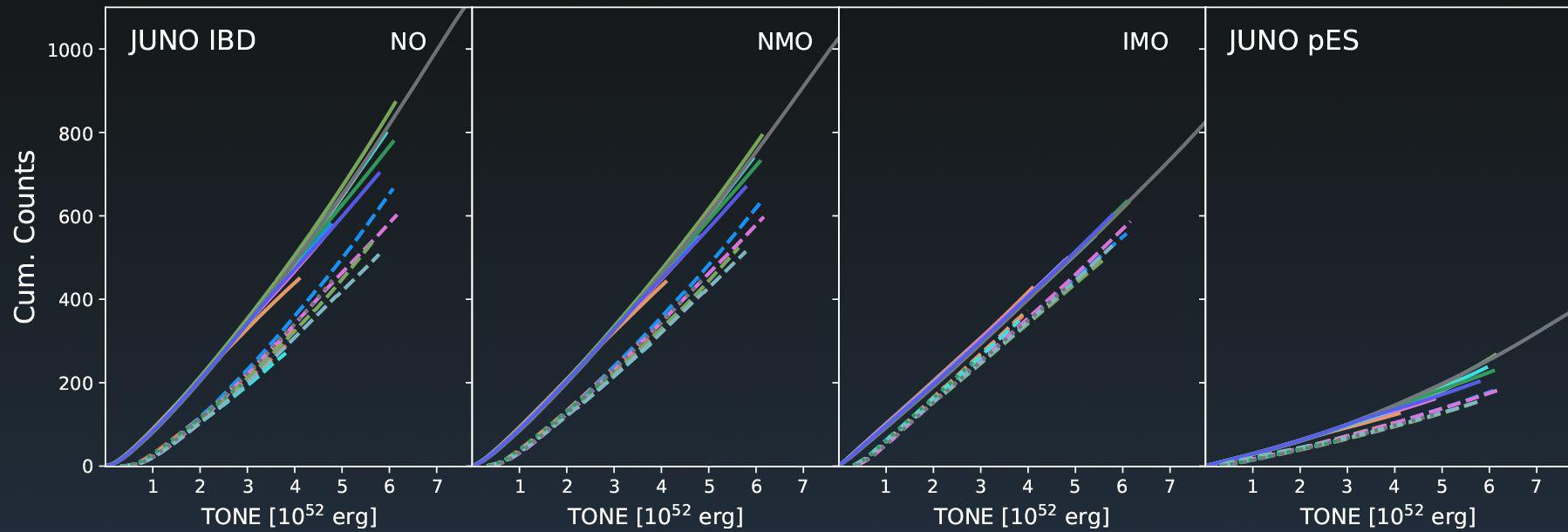
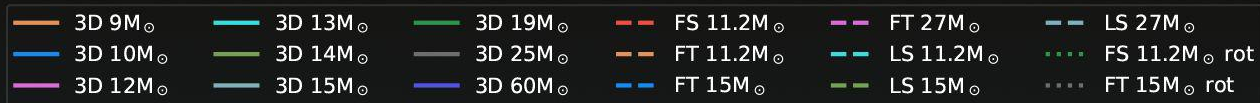
$$F_{\bar{\nu}_e} = \bar{P}_e F_{\bar{\nu}_e}^0 + (1 - \bar{P}_e) F_{\bar{\nu}_x}^0$$

$P_e \sim 0.7$ NMO; 0.3 IMO

Fits with 1st or 2nd order polynomials.



Neutrino signal results useful correlation

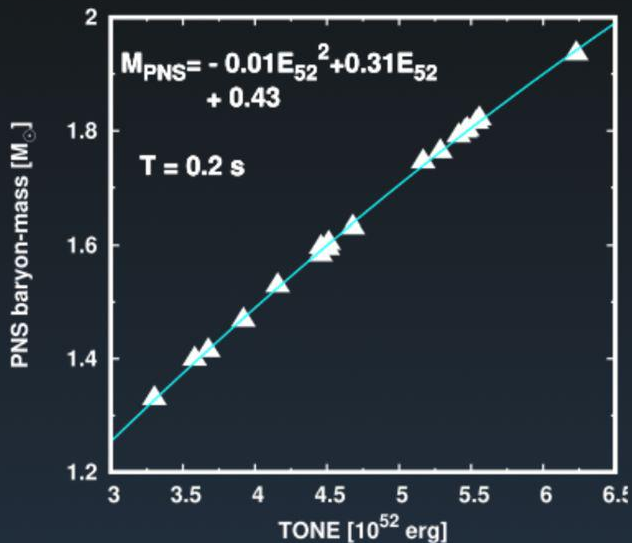


Neutrino signal results useful correlation

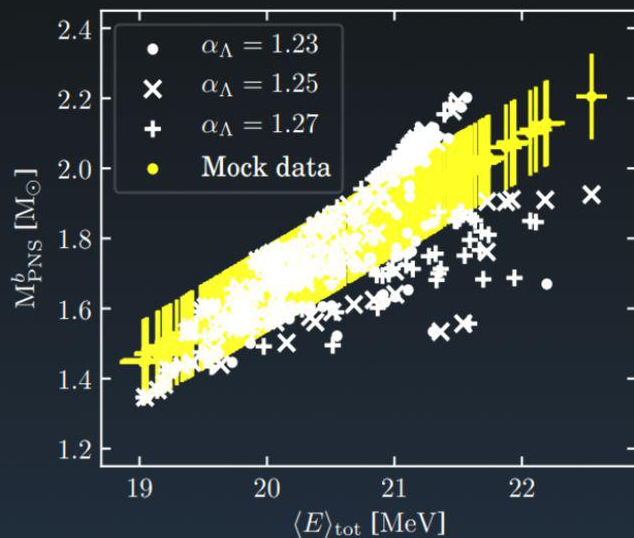
Cumulative events \rightarrow TONE \rightarrow PNS properties

Correlation between TONE and PNS mass/radius

“Efficient method for estimating the time evolution of the proto-neutron star mass and radius from a supernova neutrino signal” [Nagakura & Vartanyan 2022]



Nagakura & Vartanyan 2022



Warren et al. 2020

SN model discrimination

- ▶ Following [Olsen & Qian 2021]: Bayes factors as strength of evidence for model distinguishability
- ▶ Comparison between models of our dataset
- ▶ Detectors HK, DUNE and JUNO (pES channel did not present sensitivity)
- ▶ Different mass orderings

We explored discrimination power regarding 3 aspects:

1. EoS

2. Progenitor Mass

3. Neutrino mass ordering

SN model discrimination

Probability distribution for an event to be observed at time t with energy E , for an specific emission model M_j :

$$p(E, t|M_j) = \frac{1}{\langle N \rangle} \frac{d^2 N}{dt dE}$$

Extended maximum likelihood [Barlow 1990]:

$$\mathcal{L}(D|M_j) = \frac{e^{-\langle N \rangle} \langle N \rangle^N}{N!} \prod_{i=1}^N p(E_i, t_i|M_j)$$

$$\mathcal{B}_{\alpha\beta} = \frac{\mathcal{L}(D|M_\alpha)}{\mathcal{L}(D|M_\beta)}$$

$\ln \mathcal{B}_{\alpha\beta}$	Interpretation
0 to 1	Not worth more than a bare mention
1 to 3	Positive evidence favoring M_α
3 to 5	Strong evidence favoring M_α
> 5	Very Strong evidence favoring M_α

By executing this procedure a total of 10^4 times and leveraging our Monte Carlo simulated signals, we are able to compute both the mean $\langle \ln \mathcal{B}_{\alpha\beta} \rangle$ and the standard deviation $\sigma(\ln \mathcal{B}_{\alpha\beta})$ across all pairs of our model instances. To discern a substantial preference for M_α over M_β , we establish a threshold of $\ln \mathcal{B}_{\alpha\beta} > 5$

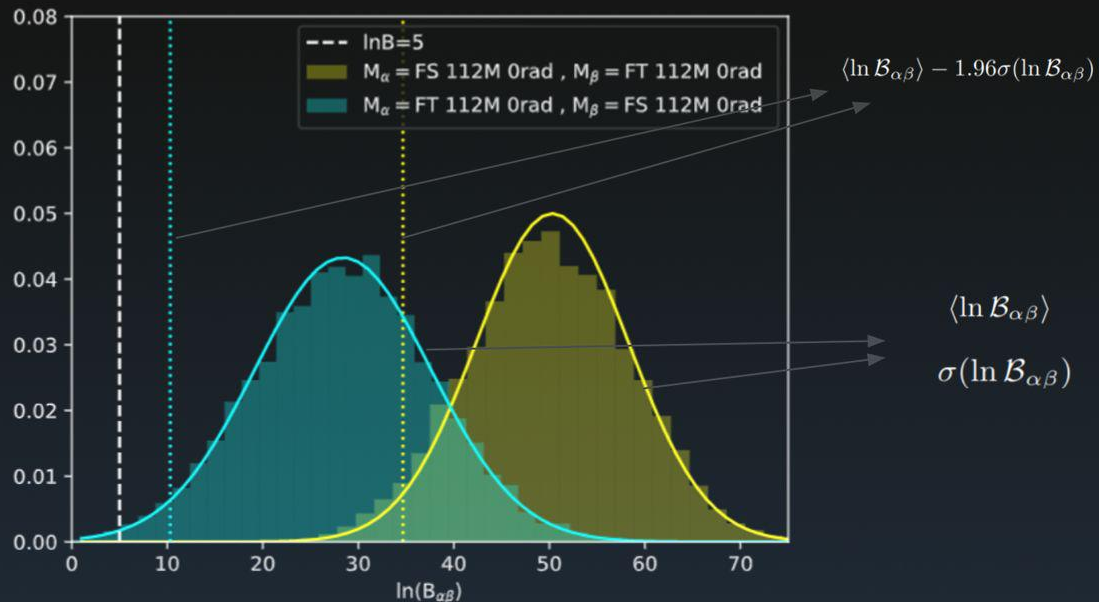
Example case: EoS discrimination

IBD channel at HK

Models: 2D, $M=11.2 M_{\odot}$

1. True model: FS, other: FT

2. True model: FT, other: FS



Binned outcomes of the Monte Carlo sample for both model pairs

Models distinguishable at a 95% CL if:

$$\langle \ln \mathcal{B}_{\alpha\beta} \rangle - 1.96\sigma(\ln \mathcal{B}_{\alpha\beta}) > 5$$

EoS discrimination

HK IBD			DUNE $\nu_e + \text{Ar}$		JUNO IBD	
M_α/M_β	FT 11.2M $_\odot$	LS 11.2M $_\odot$	FT 11.2M $_\odot$	LS 11.2M $_\odot$	FT 11.2M $_\odot$	LS 11.2M $_\odot$
FS 11.2M $_\odot$						
NO	39.34 \pm 8.62	56.65 \pm 10.29	3.49 \pm 2.69	1.81 \pm 1.89	3.71 \pm 5.72	4.52 \pm 6.57
NMO	27.38 \pm 6.97	27.42 \pm 8.14	0.55 \pm 1.18	0.82 \pm 0.86	2.05 \pm 5.46	2.28 \pm 6.06
IMO	22.17 \pm 6.81	18.45 \pm 7.44	0.65 \pm 1.49	0.94 \pm 0.93	2.02 \pm 4.48	1.03 \pm 3.38
FT 11.2M $_\odot$						
NO		141.75 \pm 16.08		2.37 \pm 2.20		11.67 \pm 7.85
NMO		76.36 \pm 12.53		0.64 \pm 1.57		6.40 \pm 6.67
IMO		29.58 \pm 7.78		0.84 \pm 1.59		2.72 \pm 5.30
M_α/M_β	LS 15M $_\odot$		LS 15M $_\odot$		LS 15M $_\odot$	
FT 15M $_\odot$						
NO	278.94 \pm 22.49		2.95 \pm 2.63		23.97 \pm 10.48	
NMO	171.71 \pm 17.93		1.01 \pm 2.15		15.26 \pm 9.64	
IMO	73.47 \pm 12.40		1.71 \pm 2.16		6.42 \pm 6.82	
M_α/M_β	LS 27M $_\odot$		LS 27M $_\odot$		LS 27M $_\odot$	
FT 27M $_\odot$						
NO	216.62 \pm 20.14		2.44 \pm 2.25		18.39 \pm 9.95	
NMO	139.11 \pm 17.05		1.24 \pm 2.07		11.81 \pm 9.28	
IMO	71.25 \pm 12.41		1.28 \pm 2.04		6.26 \pm 6.78	

Prog. Mass discrimination

HK IBD			DUNE $\nu_e + \text{Ar}$		JUNO IBD	
M_α/M_β	FT 15M $_\odot$	FT 27M $_\odot$	FT 15M $_\odot$	FT 27M $_\odot$	FT 15M $_\odot$	FT 27M $_\odot$
FT 11.2M $_\odot$						
NO	1572.51 \pm 59.41	1071.60 \pm 47.85	114.30 \pm 16.75	83.90 \pm 13.52	126.98 \pm 17.94	83.41 \pm 15.15
NMO	1159.26 \pm 48.98	882.97 \pm 42.88	13.84 \pm 6.31	20.24 \pm 7.16	94.51 \pm 15.84	70.99 \pm 13.63
IMO	467.46 \pm 31.33	559.28 \pm 34.08	34.74 \pm 8.89	34.30 \pm 8.75	43.18 \pm 11.11	50.89 \pm 11.82
FT 15M $_\odot$						
NO		178.36 \pm 19.33		13.79 \pm 5.46		15.28 \pm 7.85
NMO		97.01 \pm 14.01		2.89 \pm 1.40		8.07 \pm 7.50
IMO		22.03 \pm 6.53		2.06 \pm 2.00		4.13 \pm 3.08
M_α/M_β	LS 15M $_\odot$	LS 27M $_\odot$	LS 15M $_\odot$	LS 27M $_\odot$	LS 15M $_\odot$	LS 27M $_\odot$
LS 11.2M $_\odot$						
NO	1227.67 \pm 52.72	994.21 \pm 46.24	95.42 \pm 14.77	87.39 \pm 14.56	94.90 \pm 15.96	75.47 \pm 14.02
NMO	885.34 \pm 44.49	785.81 \pm 40.99	11.04 \pm 5.59	16.73 \pm 6.47	69.47 \pm 14.02	61.21 \pm 13.04
IMO	342.00 \pm 26.36	443.39 \pm 29.72	28.49 \pm 7.95	30.66 \pm 8.33	31.29 \pm 10.23	39.74 \pm 11.10
LS 15M $_\odot$						
NO		107.34 \pm 14.90		11.42 \pm 4.86		8.63 \pm 6.66
NMO		64.12 \pm 11.33		1.49 \pm 1.04		6.18 \pm 5.07
IMO		25.66 \pm 7.07		2.07 \pm 1.99		4.04 \pm 2.21

Mass ordering discrimination

Neutronization burst only

	HK IBD	DUNE $\nu_e + \text{Ar}$	JUNO IBD
M_α/M_β	IMO	IMO	IMO
FS 11.2M \odot			
NMO	421.25 \pm 24.84	13.17 \pm 4.78	36.64 \pm 9.85
FT 11.2M \odot			
NMO	484.19 \pm 29.37	17.17 \pm 5.14	41.80 \pm 10.30
FT 15M \odot			
NMO	1156.26 \pm 48.89	30.76 \pm 6.50	74.37 \pm 15.90
FT 27M \odot			
NMO	953.31 \pm 42.51	30.05 \pm 6.44	68.66 \pm 14.65
LS 11.2M \odot			
NMO	473.83 \pm 27.32	18.18 \pm 5.24	41.61 \pm 10.49
LS 15M \odot			
NMO	1021.29 \pm 44.53	34.29 \pm 6.67	72.24 \pm 14.32
LS 27M \odot			
NMO	914.57 \pm 40.71	35.79 \pm 6.82	69.24 \pm 14.33
3D 9M \odot			
NMO	477.65 \pm 30.86	9.57 \pm 3.85	30.11 \pm 10.42
3D 10M \odot			
NMO	626.42 \pm 37.26	14.59 \pm 4.31	39.04 \pm 11.89
3D 12M \odot			
NMO	626.27 \pm 35.97	14.39 \pm 4.35	37.32 \pm 12.22
3D 13M \odot			
NMO	946.44 \pm 45.47	25.37 \pm 5.15	56.49 \pm 14.59
3D 14M \odot			
NMO	1099.72 \pm 50.60	30.97 \pm 5.39	64.07 \pm 15.45
3D 15M \odot			
NMO	849.89 \pm 43.69	20.82 \pm 4.86	50.80 \pm 13.97
3D 19M \odot			
NMO	868.20 \pm 42.97	20.27 \pm 3.58	52.55 \pm 13.73
3D 25M \odot			
NMO	1384.53 \pm 55.17	47.57 \pm 6.33	80.95 \pm 17.60
3D 60M \odot			
NMO	766.91 \pm 39.89	19.09 \pm 4.79	45.63 \pm 13.26

DUNE $\nu_e + \text{Ar}$	
M_α/M_β	IMO
FS 11.2M \odot	
NMO	8.49 \pm 2.97
FT 11.2M \odot	
NMO	12.55 \pm 3.08
FT 15M \odot	
NMO	13.80 \pm 3.45
FT 27M \odot	
NMO	17.80 \pm 3.60
LS 11.2M \odot	
NMO	13.06 \pm 2.77
LS 15M \odot	
NMO	17.08 \pm 2.86
LS 27M \odot	
NMO	21.70 \pm 2.98
3D 9M \odot	
NMO	6.64 \pm 2.32
3D 10M \odot	
NMO	8.14 \pm 2.42
3D 12M \odot	
NMO	8.20 \pm 2.45
3D 13M \odot	
NMO	10.58 \pm 2.64
3D 14M \odot	
NMO	9.88 \pm 2.67
3D 15M \odot	
NMO	9.13 \pm 2.56
3D 19M \odot	
NMO	11.68 \pm 2.76
3D 25M \odot	
NMO	15.24 \pm 3.07
3D 60M \odot	
NMO	11.40 \pm 2.74

Conclusions:

- ▶ Nu-signal from Boltzmann-radiation-hydrodynamics models studied for the very first time
- ▶ Systematic exploration across 3 detectors and comparison with Fornax 3D simulations
- ▶ Correlation between the TONE and the cumulative count statistics:
primarily influenced by the massive lepton neutrinos that originate from the source.
- ▶ **The strength of this correlation is shaped by the surviving probabilities,**
which highlights the crucial aspect of accounting for neutrino oscillations in the analysis
- ▶ The correlation between all models is notably **narrower at DUNE**
- ▶ The forthcoming Galactic SN's neutrino signal will provide us with the means
to differentiate among a diverse array of models
- ▶ **HK: optimal detector for distinguishing between models** (mass ordering, EoS and Mprog)

Future work:

- ▶ **Include more SN initial models in the study**
- ▶ **Further explore the correlation. Incorporate errors, Y_{μ} ?, supplement the study with hydrodynamic information, etc**
- ▶ **Improve the neutrino oscillation treatment (collective effects?, sterile neutrinos?, non-adiabatic transformations? etc)**
- ▶ **Explore other neutral channels**
- ▶ **DUNE cross sections \rightarrow most of the predictions (99% in literature) are made using SNOwglOBes Xs [E. Kolbe et al. 2003] \rightarrow several new argon XS calculations. Effects on expected DUNE signal?**
- ▶ **Hybrid EoS, hadron-quark phase transition: signatures in the neutrino and GW signals?**
- ▶ **New observable definitions:
Ratios of events in different detection channels (CC/NC). [Saez et al. 2310.19939]**

Backup slides

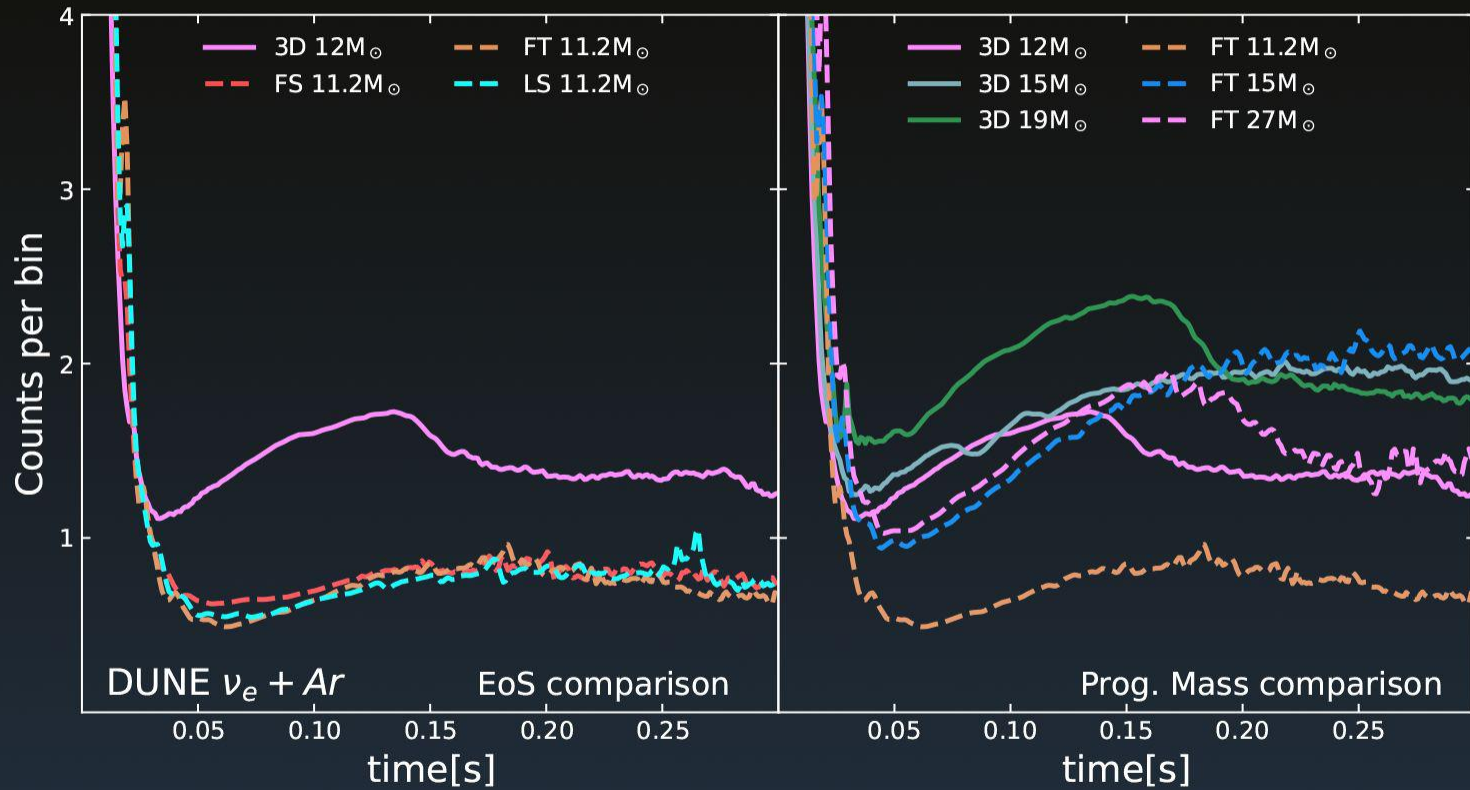
-)**Furusawa and Shen (FS)**: based on the Shen EoS [Shen:1998]. Nuclear statistical equilibrium (NSE) is considered for the ensemble of nuclei in order to calculate the thermodynamical and statistical properties of nonuniform matter.

The Shen EoS, models the strong interaction by the RMF theory with the TM1 parameter set.

-)**Lattimer and Swesty (LS)**: based on the liquid drop model of the nuclei and the Skyrme-type interaction. As for the composition, Lattimer and Swesty [Lattimer:1991] assume that the heavy nuclei are represented by a single nuclear species (single nuclear approximation, SNA), and only the alpha particle is considered as the light nuclei.

-)**Furusawa and Togashi (FT)**: based on Togashi variational method but NSE is assumed [Togashi:2017].

-)**SFHo**: based on the relativistic mean field (RMF) theory whose parameters are tuned to fit the observations[Steiner:2013]. (do not employ the SNA in the vicinity of the saturation density. Rather, several thousand nuclei are taken into account, whose masses and binding energies, when available, are taken from experiment.)



Cherenkov detectors: Charged particles are detected via their Cherenkov light emission. IBD is overwhelmingly dominant in the supernova neutrino energy regime: water Cherenkov detectors are primarily sensitive to the $\bar{\nu}_e$ component of the flux. The primary observable is the Cherenkov radiation of the IBD positron.

Scintillator detectors are composed of hydrocarbons, which have the approximate chemical formula C_nH_{2n} . The energy loss of charged particles is observed via light emitted from deexcitation of molecular energy levels, and a very large number of photons may be released.

LARTPC: Liquid argon has a particular sensitivity to the ν_e component of a supernova neutrino burst, via the dominant interaction, CC absorption of ν_e on ^{40}Ar , for which the observable is the e^- plus deexcitation products from the excited $^{40}\text{K}^*$ final state.

$$\ln B_{\alpha\beta} = N \ln \frac{\langle N \rangle_\alpha}{\langle N \rangle_\beta} - \Delta_{\alpha\beta} + \sum_{i=1}^N \ln \frac{p(E_i, t_i | M_\alpha)}{p(E_i, t_i | M_\beta)},$$

$$\frac{\langle \ln B_{\alpha\beta} \rangle}{\langle N \rangle_\alpha} = \ln \frac{\langle N \rangle_\alpha}{\langle N \rangle_\beta} - \frac{\Delta_{\alpha\beta}}{\langle N \rangle_\alpha} + \left\langle \ln \frac{p(E, t | M_\alpha)}{p(E, t | M_\beta)} \right\rangle_\alpha$$

$$\frac{\sigma[\ln B_{\alpha\beta}]}{\sqrt{\langle N \rangle_\alpha}} = \left[\left(\ln \frac{\langle N \rangle_\alpha}{\langle N \rangle_\beta} \right)^2 + 2 \ln \frac{\langle N \rangle_\alpha}{\langle N \rangle_\beta} \left\langle \ln \frac{p(E, t | M_\alpha)}{p(E, t | M_\beta)} \right\rangle_\alpha + \left\langle \left(\ln \frac{p(E, t | M_\alpha)}{p(E, t | M_\beta)} \right)^2 \right\rangle_\alpha \right]^{1/2},$$

$$\left\langle \ln \frac{p(E, t | M_\alpha)}{p(E, t | M_\beta)} \right\rangle_\alpha = \int_0^{9^s} dt \int_{E_{\min}}^{\infty} dE p(E, t | M_\alpha) \times \ln \frac{p(E, t | M_\alpha)}{p(E, t | M_\beta)}.$$

DUNE $\nu_e + \text{Ar}$								
M_α/M_β	3D 10M $_\odot$	3D 12M $_\odot$	3D 13M $_\odot$	3D 14M $_\odot$	3D 15M $_\odot$	3D 19M $_\odot$	3D 25M $_\odot$	3D 60M $_\odot$
3D 9M $_\odot$								
NO	28.67 \pm 8.27	23.96 \pm 7.36	104.20 \pm 16.98	126.69 \pm 19.80	75.79 \pm 14.60	94.33 \pm 15.26	243.06 \pm 27.15	62.02 \pm 11.84
NMO	2.27 \pm 1.01	2.48 \pm 1.85	5.20 \pm 4.02	5.22 \pm 0.12	3.71 \pm 2.05	4.46 \pm 2.20	9.63 \pm 4.17	10.41 \pm 5.53
IMO	5.15 \pm 3.57	5.52 \pm 3.54	20.68 \pm 7.91	18.45 \pm 8.49	12.75 \pm 6.26	27.01 \pm 8.41	50.71 \pm 13.81	20.98 \pm 7.09
3D 10M $_\odot$								
NO		1.54 \pm 1.99	27.75 \pm 8.21	40.97 \pm 10.53	13.47 \pm 5.91	25.30 \pm 7.45	122.87 \pm 17.79	12.18 \pm 4.94
NMO		0.60 \pm 0.25	3.13 \pm 1.42	3.23 \pm 0.55	1.65 \pm 0.46	5.02 \pm 4.07	7.20 \pm 5.45	4.92 \pm 3.43
IMO		0.78 \pm 0.29	5.63 \pm 4.40	4.19 \pm 5.01	2.75 \pm 1.13	9.53 \pm 4.93	24.92 \pm 10.21	6.37 \pm 3.72
3D 12M $_\odot$								
NO			33.35 \pm 9.17	47.51 \pm 11.54	18.27 \pm 6.95	28.07 \pm 7.95	131.66 \pm 18.74	11.82 \pm 4.97
NMO			2.92 \pm 0.74	3.03 \pm 0.96	1.51 \pm 0.25	3.78 \pm 3.82	6.96 \pm 4.01	3.64 \pm 3.14
IMO			4.99 \pm 4.41	5.04 \pm 3.56	2.82 \pm 1.38	8.60 \pm 4.84	23.30 \pm 10.16	5.54 \pm 3.55
3D 13M $_\odot$								
NO				2.87 \pm 2.55	5.45 \pm 3.41	3.18 \pm 2.92	36.61 \pm 9.52	10.81 \pm 5.58
NMO				0.73 \pm 0.67	1.62 \pm 0.49	1.84 \pm 1.21	4.17 \pm 1.51	1.29 \pm 1.12
IMO				1.09 \pm 0.39	2.03 \pm 1.20	0.66 \pm 0.8	7.15 \pm 5.82	1.81 \pm 0.29
3D 14M $_\odot$								
NO					9.31 \pm 4.89	7.94 \pm 5.22	26.73 \pm 7.84	20.37 \pm 7.98
NMO					1.75 \pm 0.64	3.26 \pm 2.59	4.06 \pm 3.67	7.81 \pm 1.44
IMO					2.57 \pm 0.87	1.38 \pm 1.66	8.66 \pm 5.23	2.52 \pm 0.48
3D 15M $_\odot$								
NO						10.07 \pm 4.58	64.36 \pm 12.45	10.64 \pm 5.12
NMO						3.26 \pm 2.59	5.67 \pm 3.06	3.45 \pm 2.09
IMO						3.44 \pm 2.77	14.61 \pm 7.75	2.15 \pm 2.14
3D 19M $_\odot$								
NO							55.53 \pm 10.79	4.13 \pm 3.51
NMO							3.43 \pm 0.62	0.99 \pm 0.13
IMO							5.53 \pm 4.42	1.57 \pm 0.45
3D 25M $_\odot$								
NO								63.21 \pm 14.13
NMO								4.15 \pm 1.29
IMO								7.06 \pm 3.85

JUNO IBD

M_α / M_β	3D 10M $_\odot$	3D 12M $_\odot$	3D 13M $_\odot$	3D 14M $_\odot$	3D 15M $_\odot$	3D 19M $_\odot$	3D 25M $_\odot$	3D 60M $_\odot$
3D 9M $_\odot$								
NO	25.87 \pm 9.30	21.88 \pm 8.61	128.39 \pm 17.89	183.91 \pm 22.00	91.73 \pm 15.19	103.62 \pm 16.09	345.09 \pm 30.17	62.77 \pm 13.02
NMO	17.68 \pm 7.70	15.80 \pm 7.70	95.31 \pm 14.79	135.12 \pm 17.98	63.88 \pm 12.65	82.75 \pm 14.10	275.37 \pm 26.02	51.72 \pm 11.46
IMO	6.87 \pm 4.21	7.40 \pm 4.98	37.61 \pm 7.03	61.74 \pm 12.71	24.14 \pm 8.44	48.14 \pm 11.24	158.38 \pm 19.57	33.23 \pm 9.99
3D 10M $_\odot$								
NO		5.51 \pm 2.96	42.21 \pm 11.19	79.36 \pm 14.36	22.15 \pm 8.88	31.87 \pm 9.99	202.43 \pm 22.22	13.75 \pm 7.78
NMO		4.93 \pm 2.13	33.43 \pm 9.76	60.19 \pm 12.48	16.11 \pm 7.61	27.73 \pm 9.20	166.88 \pm 19.95	12.84 \pm 6.85
IMO		3.71 \pm 1.28	18.40 \pm 7.81	30.60 \pm 9.16	5.75 \pm 3.89	20.29 \pm 8.24	105.48 \pm 15.91	11.99 \pm 6.86
3D 12M $_\odot$								
NO			49.57 \pm 11.66	89.60 \pm 15.17	28.48 \pm 9.86	33.92 \pm 10.60	214.60 \pm 22.74	12.35 \pm 6.98
NMO			37.27 \pm 10.02	65.76 \pm 13.07	19.54 \pm 8.29	28.58 \pm 9.70	173.62 \pm 20.31	11.70 \pm 6.41
IMO			18.71 \pm 7.81	31.87 \pm 9.46	6.09 \pm 4.07	18.91 \pm 8.10	105.75 \pm 16.01	10.14 \pm 6.43
3D 13M $_\odot$								
NO				7.23 \pm 4.36	6.93 \pm 4.98	7.14 \pm 4.77	64.52 \pm 13.71	19.72 \pm 7.78
NMO				7.18 \pm 4.19	6.60 \pm 3.94	6.09 \pm 3.96	54.51 \pm 13.21	9.81 \pm 5.35
IMO				5.06 \pm 3.48	5.71 \pm 3.57	4.77 \pm 3.15	38.22 \pm 10.38	6.22 \pm 3.64
3D 14M $_\odot$								
NO					16.23 \pm 7.98	25.13 \pm 9.18	36.47 \pm 11.62	50.74 \pm 11.67
NMO					13.45 \pm 6.55	15.52 \pm 6.44	32.45 \pm 10.92	33.01 \pm 9.72
IMO					7.06 \pm 4.36	5.89 \pm 3.89	27.45 \pm 9.53	8.09 \pm 4.07
3D 15M $_\odot$								
NO						12.66 \pm 8.42	105.66 \pm 16.99	16.02 \pm 8.69
NMO						10.02 \pm 7.18	91.65 \pm 15.69	8.00 \pm 5.59
IMO						8.15 \pm 6.62	65.65 \pm 12.99	6.02 \pm 4.64
3D 19M $_\odot$								
NO							89.22 \pm 15.24	7.29 \pm 4.58
NMO							69.29 \pm 13.76	6.87 \pm 4.58
IMO							39.85 \pm 10.17	4.81 \pm 3.36
3D 25M $_\odot$								
NO								139.92 \pm 18.44
NMO								106.78 \pm 16.44
IMO								57.52 \pm 11.80

SN model	$\langle N \rangle_{\text{HK-IBD}}$	$\langle N \rangle_{\text{DUNE-}\nu_e + \text{Ar}}$	$\langle N \rangle_{\text{JUNO-IBD}}$	$\langle N \rangle_{\text{JUNO-pES}}$					
FS 11.2M _⊙					3D 9M _⊙				
NO	4060.73	263.57	300.86	94.52	NO	5803.79	319.65	448.43	126.17
NMO	4135.24	232.35	321.68		NMO	5526.81	256.74	441.38	
IMO	4287.04	241.24	364.11		IMO	5013	263.57	300.86	
FT 11.2M _⊙					3D 10M _⊙				
NO	4412.90	262.40	328.97	103.67	NO	7349.60	434.14	579.05	162.63
NMO	4399.86	227.22	343.80		NMO	6793.82	278.54	550.26	
IMO	4373.35	237.25	374.0		IMO	5661.89	322.84	328.97	
FT 15M _⊙					3D 12M _⊙				
NO	8729.38	534.59	665.66	180.81	NO	7351.13	430.19	576.43	160.64
NMO	7972.33	308.52	630.12		NMO	6824.39	286.21	550.70	
IMO	6430.36	372.89	557.75		IMO	5751.68	534.59	665.66	
FT 27M _⊙					3D 13M _⊙				
NO	8001.56	500.24	603.00	181.38	NO	9916.66	593.40	798.33	236.19
NMO	7597.67	330.55	597.60		NMO	8944.62	302.16	737.07	
IMO	6775.03	378.87	586.65		IMO	6965.02	385.08	612.31	
LS 11.2M _⊙					3D 14M _⊙				
NO	3698.63	262.58	271.05	90.56	NO	10690.80	624.14	870.83	266.51
NMO	3834.93	220.77	296.67		NMO	9525.98	281.87	792.17	
IMO	4112.51	232.63	348.84		IMO	7153.43	379.32	632.21	
LS 15M _⊙					3D 15M _⊙				
NO	7205.00	509.16	535.97	150.0	NO	8917.34	528.16	713.04	202.60
NMO	6743.71	290.45	523.79		NMO	8054.11	288.13	659.56	
IMO	5804.10	352.57	498.98		IMO	6296.01	356.47	550.63	
LS 27M _⊙					3D 19M _⊙				
NO	6865.32	499.00	507.86	155.38	NO	9747.49	589.19	777.31	228.42
NMO	6633.72	311.00	514.81		NMO	8920.65	331.84	730.03	
IMO	6162.03	364.19	528.99		IMO	7236.70	405.11	633.74	
FS 11.2M _⊙ rot					3D 25M _⊙				
NO	2682.33	174.32	194.45	62.65	NO	13692.27	811.58	1125.61	371.30
NMO	2783.44	164.33	213.83		NMO	12235.23	329.90	1026.43	
IMO	2989.51	167.17	253.31		IMO	9267.60	467.04	825.06	
FT 15M _⊙ rot					3D 60M _⊙				
NO	6156.31	344.28	452.64	124.11	NO	8896.73	532.58	701.21	202.49
NMO	5780.72	251.82	445.79		NMO	8242.20	232.35	668.41	
IMO	5015.67	278.14	431.83		IMO	6909.17	389.45	601.60	

Correlation fittings:

$$\text{Cum}_{(\text{DUNE-NMO})} = (0.43E^2 + 51.19E) \left(\frac{d}{10kpc} \right)^{-2} \left(\frac{V}{40kt} \right)$$

$$\text{Cum}_{(\text{DUNE-IMO})} = (59.63E + 9.91) \left(\frac{d}{10kpc} \right)^{-2} \left(\frac{V}{40kt} \right)$$

$$\text{Cum}_{(\text{JUNO-IBD-NMO})_{3D}} = (5.38E^2 + 92.6E) \left(\frac{d}{10kpc} \right)^{-2} \left(\frac{V}{18.25kt} \right)$$

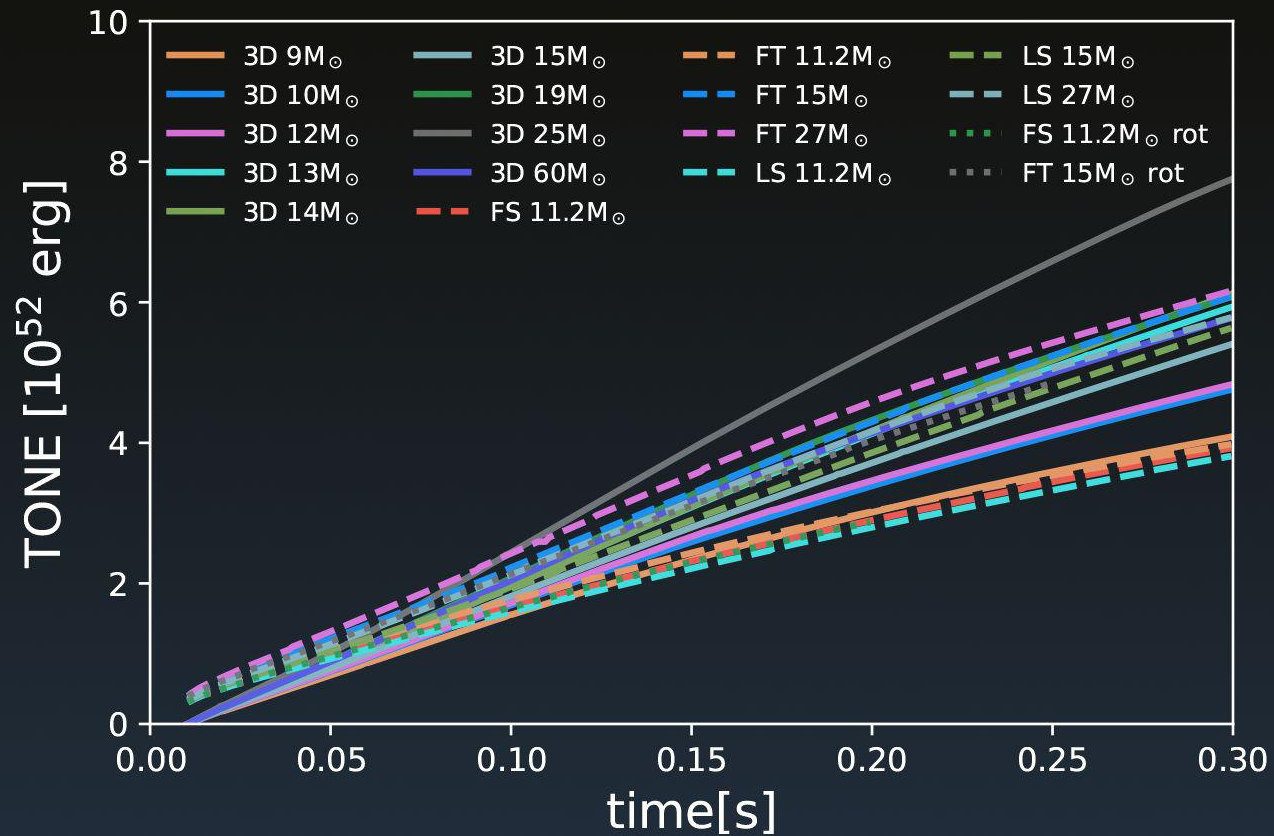
$$\text{Cum}_{(\text{JUNO-IBD-IMO})_{3D}} = (106.8E - 19.37) \left(\frac{d}{10kpc} \right)^{-2} \left(\frac{V}{18.25kt} \right)$$

$$\text{Cum}_{(\text{JUNO-IBD-NMO})_{2D}} = (8.05E^2 + 50.75E) \left(\frac{d}{10kpc} \right)^{-2} \left(\frac{V}{18.25kt} \right)$$

$$\text{Cum}_{(\text{JUNO-IBD-IMO})_{2D}} = (99.26E - 45.5) \left(\frac{d}{10kpc} \right)^{-2} \left(\frac{V}{18.25kt} \right)$$

$$\text{Cum}_{(\text{JUNO-PES})_{3D}} = (2.31E^2 + 26.12E) \left(\frac{d}{10kpc} \right)^{-2} \left(\frac{V}{18.25kt} \right)$$

$$\text{Cum}_{(\text{JUNO-PES})_{2D}} = (1.93E^2 + 17.83E) \left(\frac{d}{10kpc} \right)^{-2} \left(\frac{V}{18.25kt} \right)$$



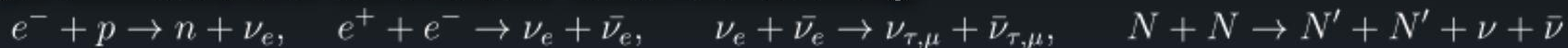
Motivation

- ▶ Various current and next-generation neutrino detection collaborations have assessed the capabilities of their detectors in observing SN neutrinos [Wang et al. 2021, Abi et al. 2021, Abe et al. 2018].
- ▶ These collaborations typically utilize a limited number (one or two) of SN models as benchmarks
- ▶ The neutrino oscillations effects and smearing effects are not always included
- ▶ Not too much analysis has been conducted to date that demonstrates the level of discrimination achievable between different SN models using several realistic detectors and interaction channels. Notable contributions in this direction: [Olsen & Qian 2022] and [Abe et al. 2021].

Conducting a study that takes into account both detector efficiencies, smearing, thresholds, background considerations, and moreover, doesn't solely focus on a limited set of benchmark models, but rather conducts a more generalized study, is a challenging task. But, it is necessary to comprehensively explore the diverse features observed in modern computer simulations.

Core-Collapse Supernovae

- ▶ Final evolutionary stage of stars with masses $M > 8 M_{\odot}$. Represent a long-awaited observation target for neutrino telescopes.
- ▶ About 1% of the gravitational binding energy is released as kinetic energy in the compact object formation, while the remaining 99% is carried out by neutrinos with energies of several MeV.
- ▶ To explain these events, interdisciplinary research that combines nuclear physics, particle physics and astrophysics is needed.
- ▶ The mechanisms leading to neutrino production in the SN core are, mainly, electron capture by nucleons, pair annihilation, flavor-conversion, and nucleon bremsstrahlung:



- ▶ Studying the signals that the neutrino leave in the detectors, with an effective neutrino flavor discrimination, it is possible to infer properties on their physics, since the structure of the neutrino mass spectrum and lepton mixing is imprinted into the detected signal.
- ▶ CCSN neutrinos were already observed for the 1987A SN in the Large Magellanic Cloud (Kamiokande-II, Irvine-Michigan Brookhaven (IMB) and Baksan detectors)
- ▶ At present, several detectors are ready and waiting for the detection of SN neutrinos from the next galactic explosion. SN neutrinos can be detected via weak charged-current (CC) and neutral-current (NC) interactions with electrons and nuclei.
- ▶ Relevant interaction channels for current (and future) detectors: inverse beta decay, neutrino-proton elastic scattering, neutrino-electron elastic scattering, **absorption interaction in liquid Argon**, among others.