Neutron star crusts and multi-messenger nuclear-astrophysics William G. Newton

The work presented in this talk would not be possible without an amazing team of undergraduates and Master's students, including

Rebecca Preston, Amber Stinson, Lauren Balliet, Michael Ross, Gabriel Crocombe, Blake Head, Josh Sanford, Zachary Langford

Texas A&M University-Commerce

Duncan Neill, David Tsang – University of Bath



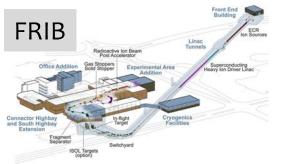




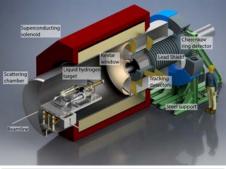




Strong, Weak, EM signals



Elliptic flow p/n ratios Pion production Resonance widths, Centroid energies **Optical potentials** Scattering X-sections



Computation

PREX/CREX/MREX

Multi-messenger Nuclear & **Astro Physics**

Weak, EM, Grav signals

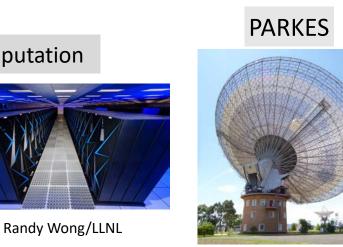






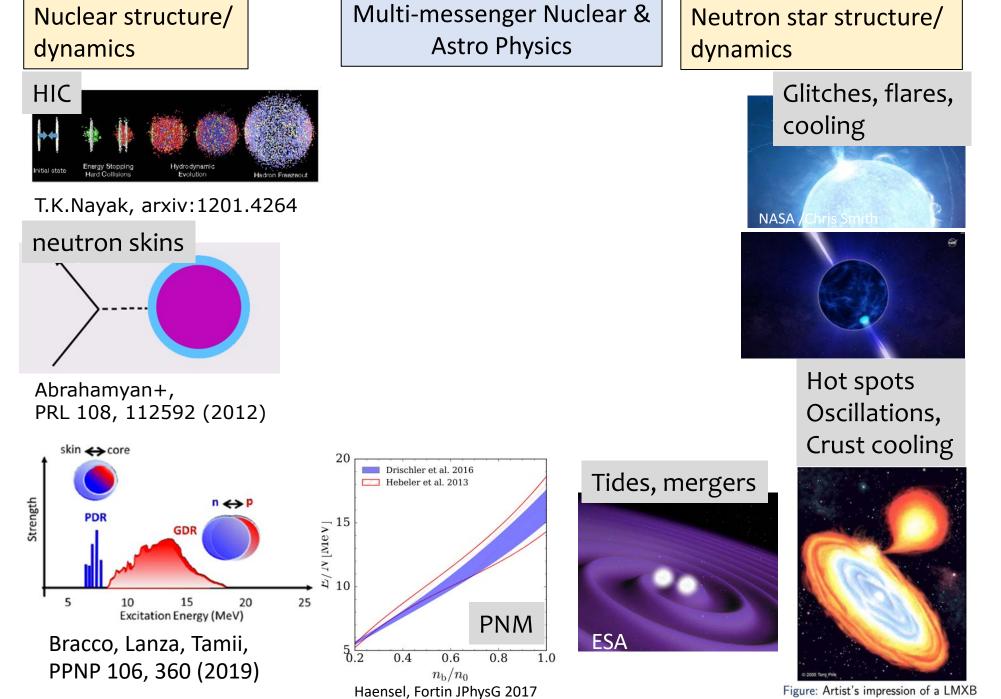
NICER

X-ray flux and light curves Gravitational waveforms Pulsar timing

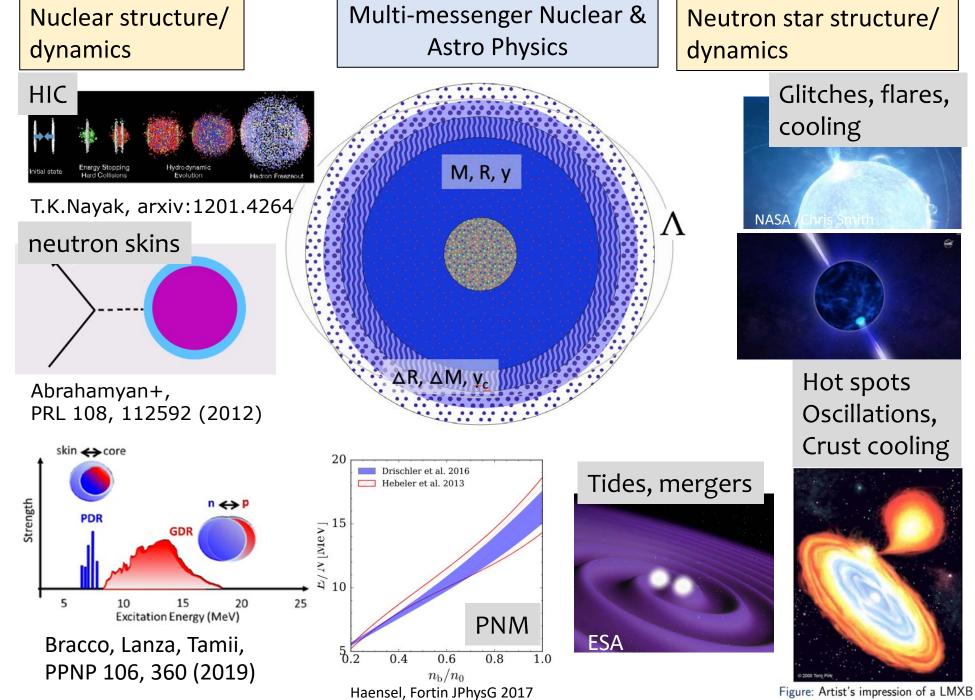




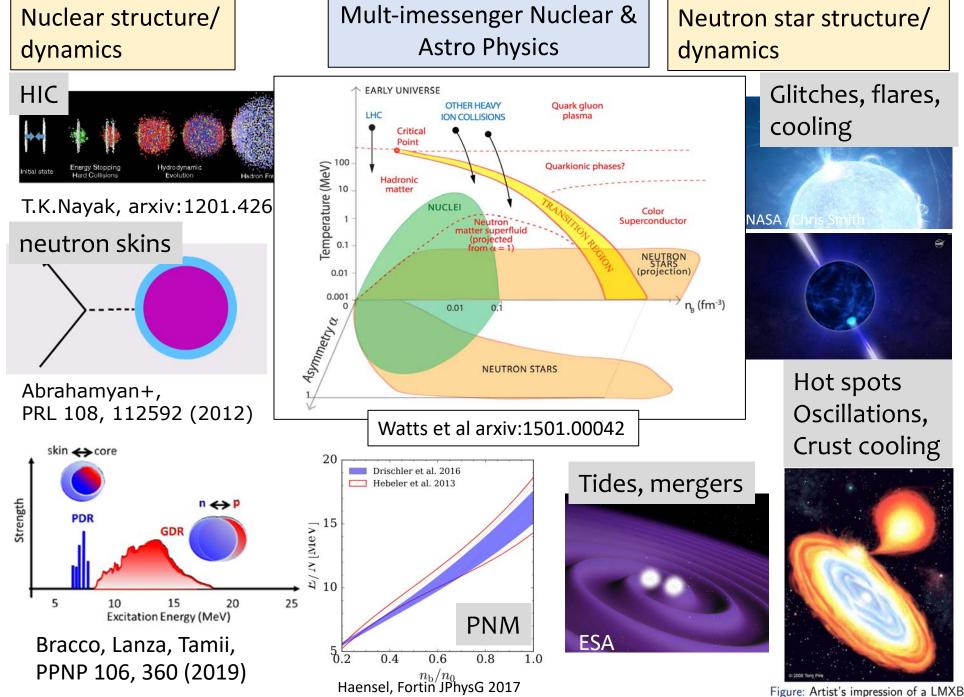




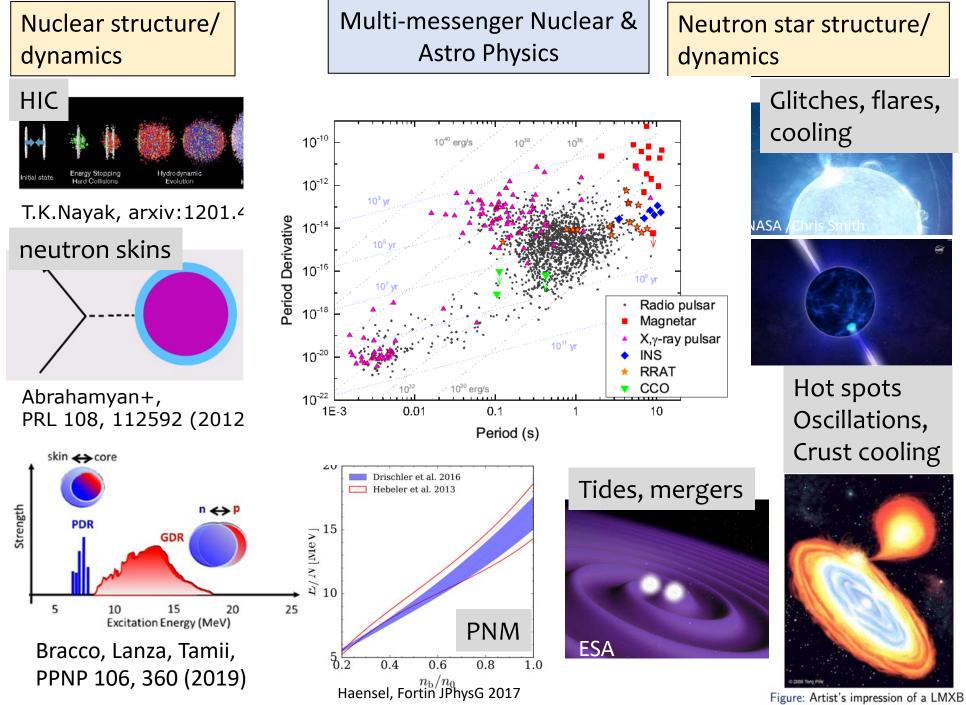
- credit Tony Piro, 2005.



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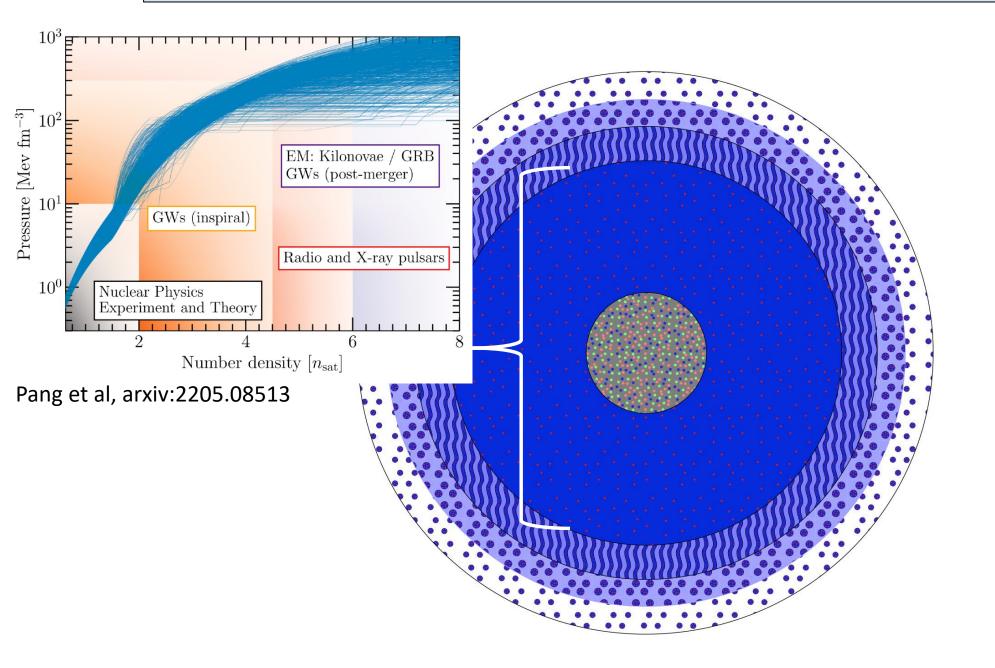
- credit Tony Piro, 2005.

Putting the Multi in Multi-messenger

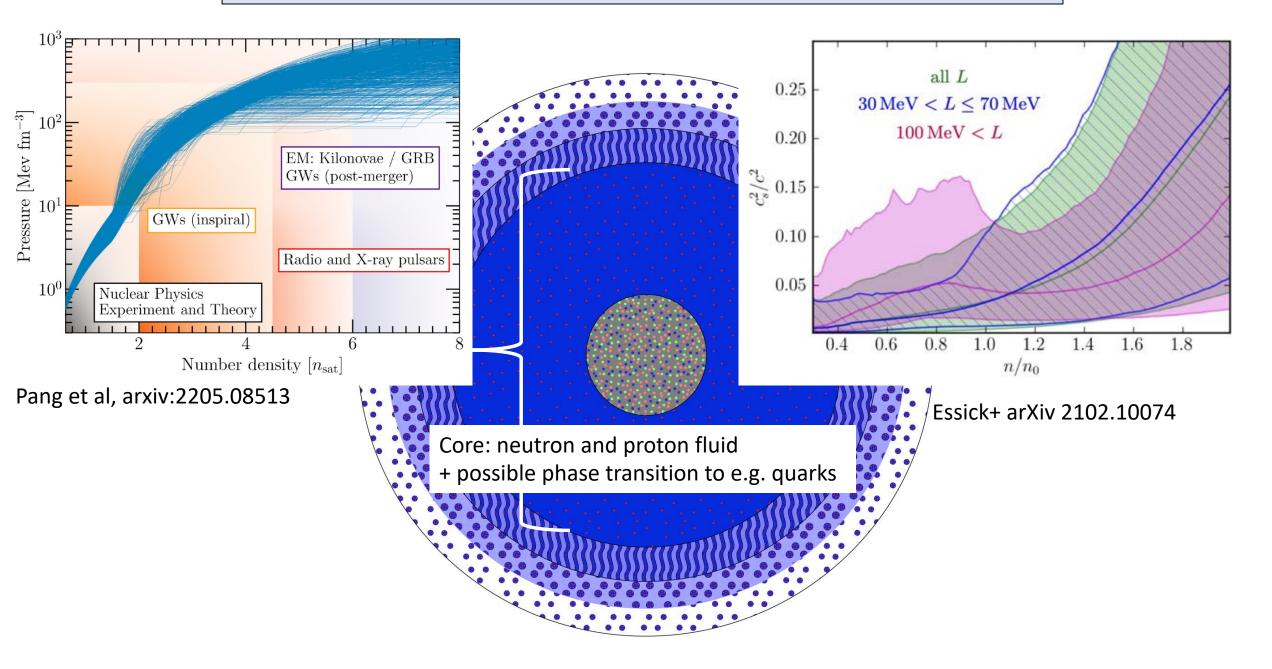
Nuclear	Neutron star
Isospin diffusion in HICs	Masses and radii
Dipole polarizability	Tidal deformability
Spectral ratios of light clusters	Moment of inertia
Nuclear masses and radii	Gravitational binding energy
Isobaric analog states	Cooling of young neutron stars
n/p ratios in HICs	Bulk oscillation modes
Neutron skins	Crust cooling
Mirror nuclei	Pulsar glitches
Giant resonances	Lower and upper limits on neutron star spin periods
Flow of particles in HICs	Torsional crust oscillations
Charged pion ratios in HICs	Crust-core interface modes

What do we want to do with this (potential data)?

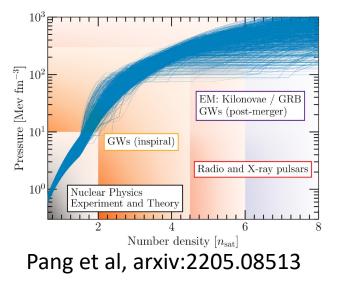
Modern approach: create ensembles of EOSs/neutron star models for statistical inference

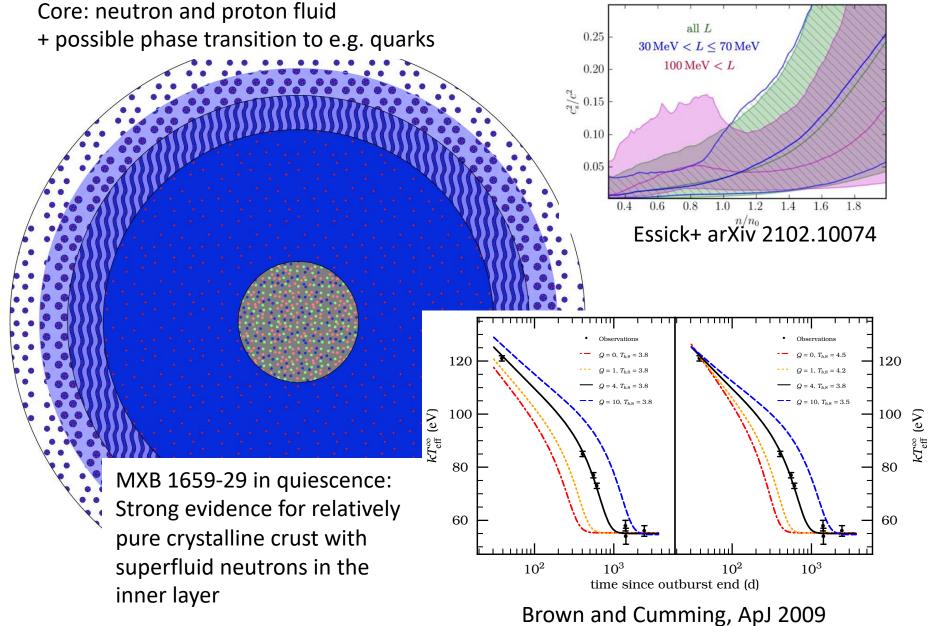


And go ahead and infer! To date, emphasis has been on the EOS of the core

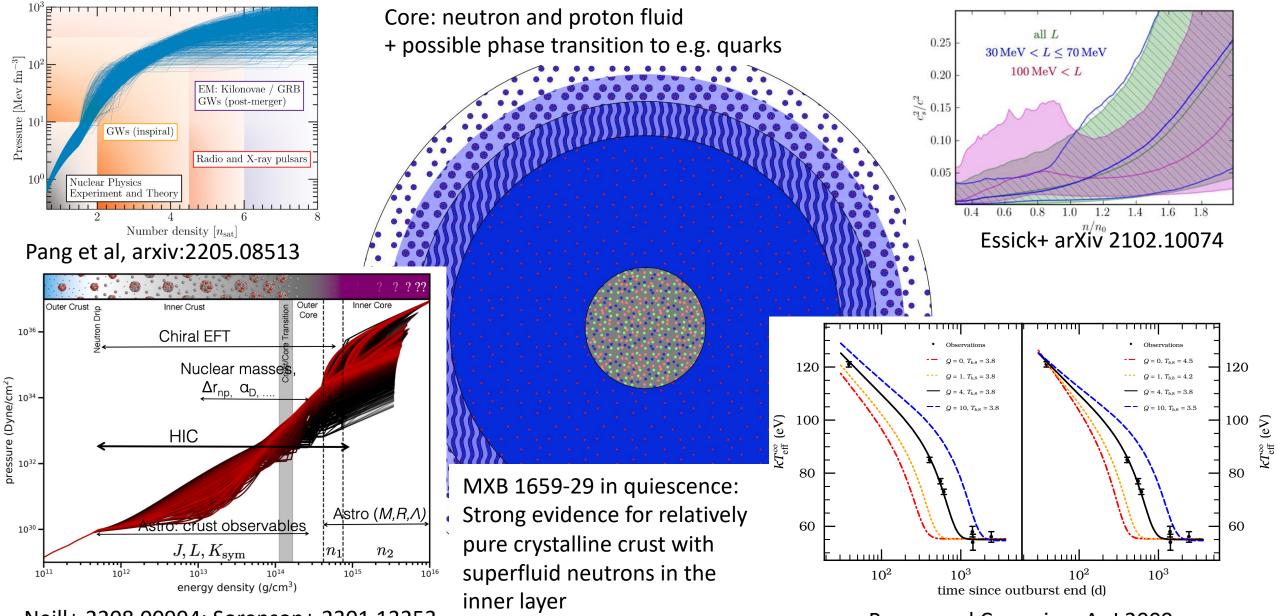


But the crust is there too, and several observables are sensitive to it





So let's include the crust when we build our ensembles



Neill+ 2208.00994; Sorenson+ 2301.13253 '

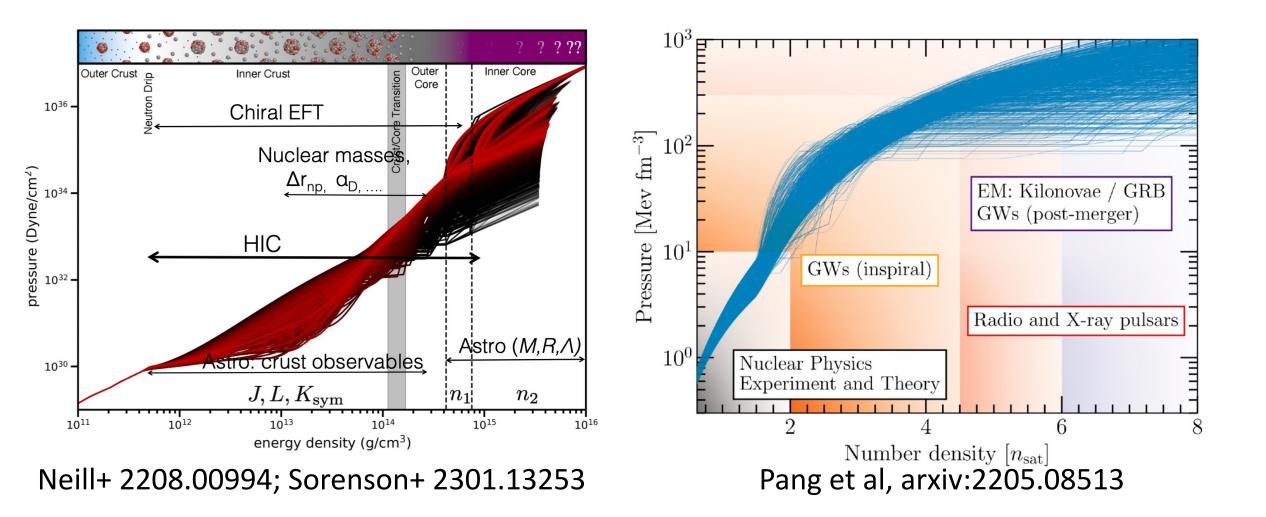
Brown and Cumming, ApJ 2009

Putting the Multi in Multi-messenger

Nuclear	Neutron star
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arxiv:2301.13253 What do we want to do with this (potential data)?

So let's go about the task of creating consistent crust and core EOSs

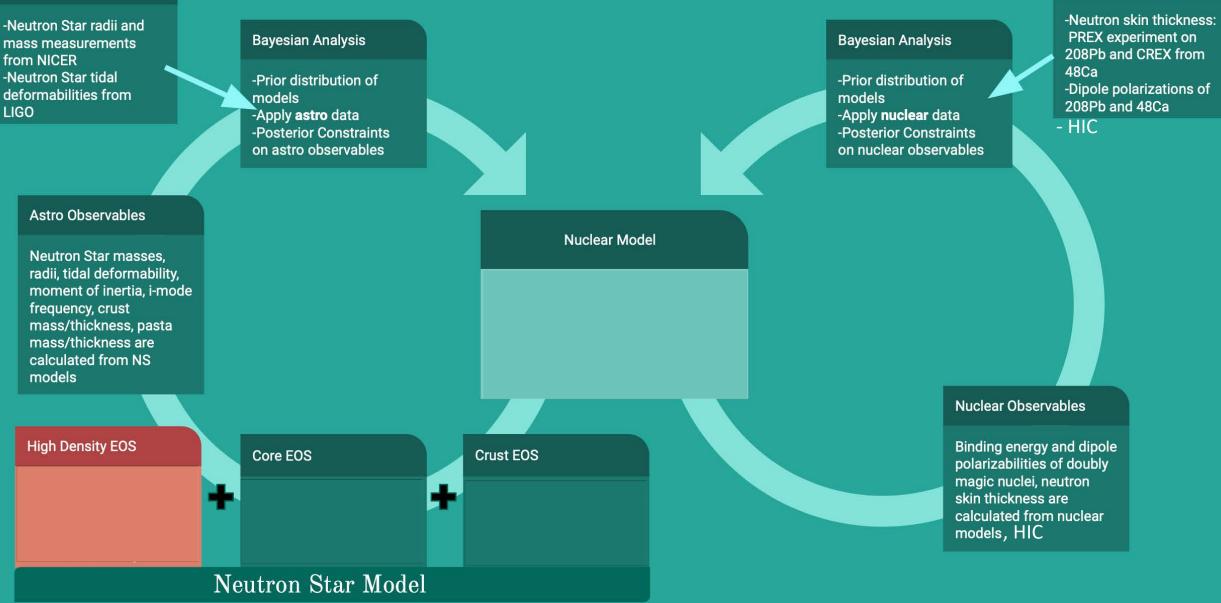


Astro Data

mass measurements from NICER -Neutron Star tidal deformabilities from LIGO

Combining nuclear and astrophysical data: a perspective

Nuclear Data



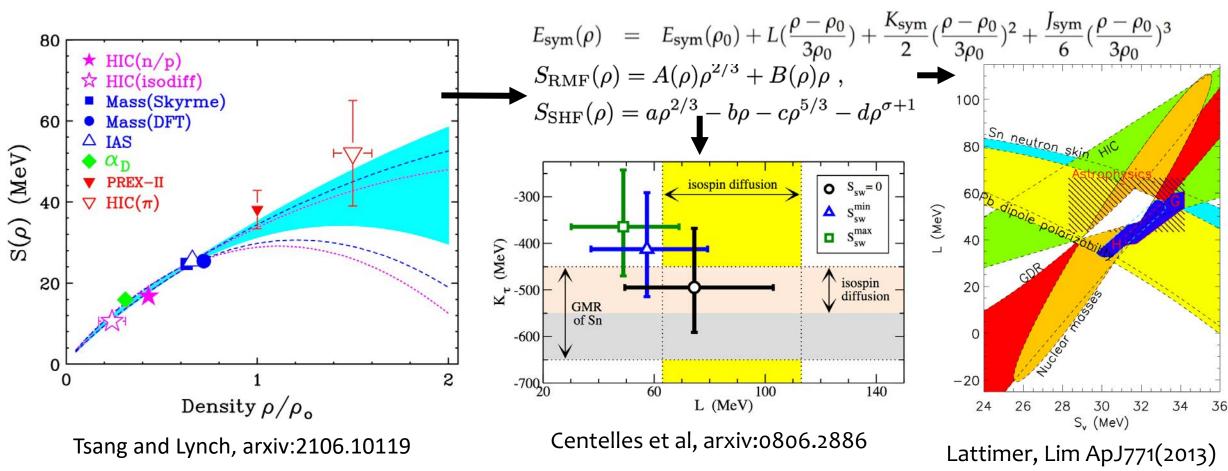
The nuclear symmetry energy: parameterizing our ignorance in a physically meaningful way

$$E_{\text{sym}}(\rho) = E_{\text{sym}}(\rho_0) + L(\frac{\rho - \rho_0}{3\rho_0}) + \frac{K_{\text{sym}}}{2}(\frac{\rho - \rho_0}{3\rho_0})^2 + \frac{Q_{\text{sym}}}{6}(\frac{\rho - \rho_0}{3\rho_0})^3$$

$$\int_{0}^{20} \int_{0}^{10} \frac{1}{\rho_0} \frac{1}{\rho_0$$

Different observables constrain at different densities...

... so resulting constraints on nuclear matter parameters at saturation density involve model-dependent extrapolation



Lattimer, Steiner EPJA50 (2013)

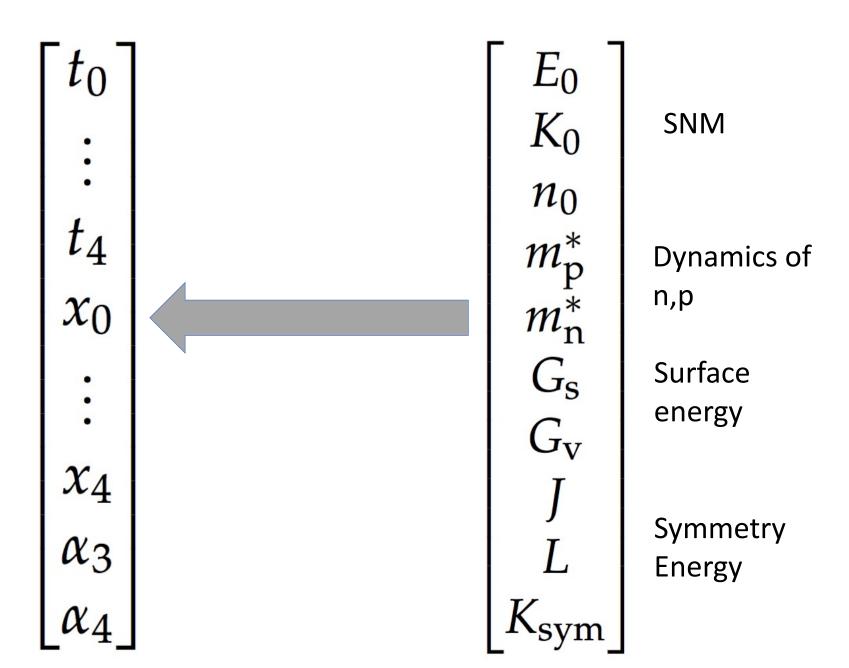
Our choice of model: Skyrme-Hartree-Fock

Density Functional Theory (e.g. Skyrme) $\mathcal{H}_{\delta} = \frac{1}{4} t_0 \rho^2 [(2+x_0) - (2x_0+1)(y_p^2+y_n^2)]$ Local interaction $\mathcal{H}_{\rho} = \frac{1}{4} t_3 \rho^{2+\alpha_3} [(2+x_3) - (2x_3+1)(y_p^2+y_n^2)]$ Density dependent $+ \frac{1}{4} t_4 \rho^{2+\alpha_4} [(2+x_4) - (2x_4+1)(y_p^2+y_n^2)]$

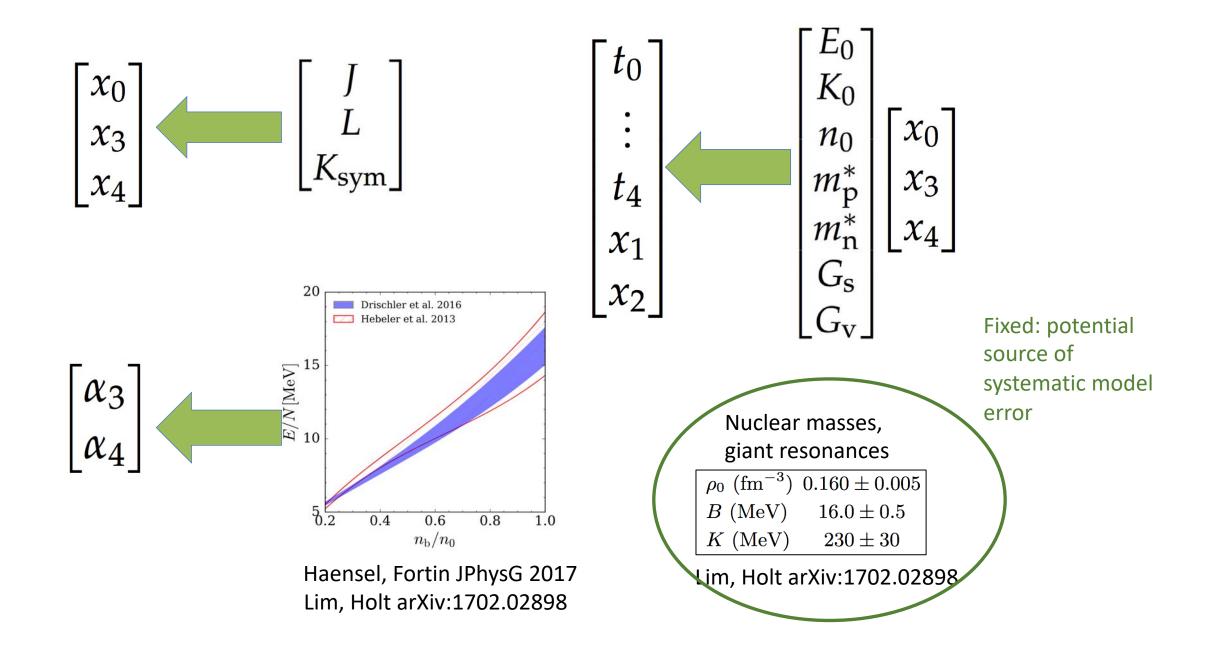
$$egin{aligned} \mathcal{H}_{ ext{eff}} &= rac{1}{8}
ho[t_1(2+x_1)+t_2(2+x_2)] au \ &+ rac{1}{8}
ho[t_1(2x_1+1)+t_2(2x_2+1)](au_p y_p+ au_n y_n) \end{aligned}$$
 3 body

$$\mathcal{H}_{\text{grad}} = \frac{1}{32} (\nabla \rho)^2 [3t_1(2+x_1) - t_2(2+x_2)]$$
Gradient...
$$-\frac{1}{32} [3t_1(2x_1+1) + t_2(2x_2+1)] [(\nabla \rho_p)^2 + (\nabla \rho_n)^2)$$

Used in a variational principle on total energy leads to coupled Schrödinger-like equations for the wavefunctions. Solutions converge to ground state (Hohenberg-Kohn theorem) Map nuclear matter parameters to model parameters and systematically generate models

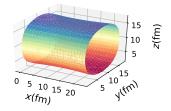


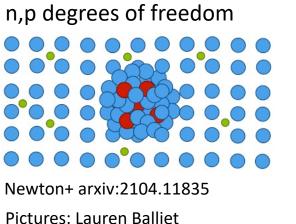
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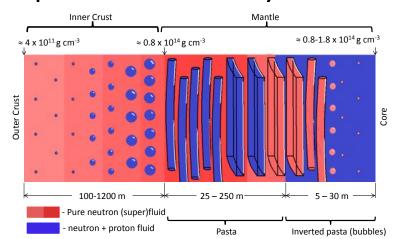
Modeling the crust

3D Skyrme HF: n,p degrees of freedom

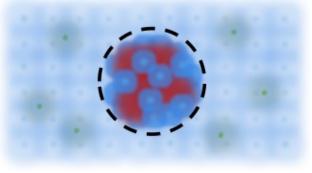




 $\mathcal{H}_{\delta} + \mathcal{H}_{\rho} + \mathcal{H}_{eff} + \mathcal{H}_{grad} + \mathcal{H}_{Coul}$ Nuclear EDF: Bulk+Gradient Specific model: Skyrme



CLDM:Bulk fluid and surface degrees of freedom



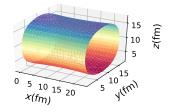
Newton et al arxiv: 1110.4043 Balliet+; arxiv:2009.07696

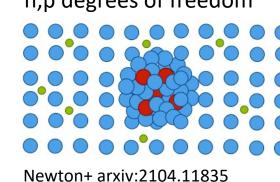
$$\begin{split} \mathcal{H}_{\delta} + \mathcal{H}_{\rho} + \mathcal{H}_{eff} & \sigma(y_p) \\ \text{Nuclear EDF: Bulk +} \\ \text{separate surface energy function} \\ \text{specific model: LLPR 1985} \end{split}$$

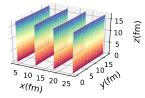
$$\sigma_{s}(y_{p}) = \sigma_{0} \frac{2^{p+1} + b}{\frac{1}{y_{p}^{p}} + b + \frac{1}{(1-y_{p})^{p}}}$$

Modeling the crust

3D Skyrme HF: n,p degrees of freedom



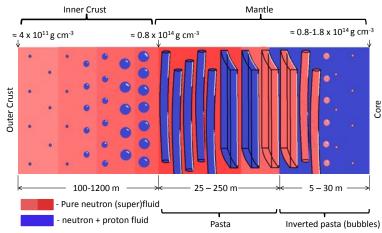




 $\mathcal{H}_{\delta} + \mathcal{H}_{\rho} + \mathcal{H}_{eff} + \mathcal{H}_{grad} + \mathcal{H}_{Coul}$ Nuclear EDF: Bulk+Gradient

Specific model: Skyrme

Pictures: Lauren Balliet

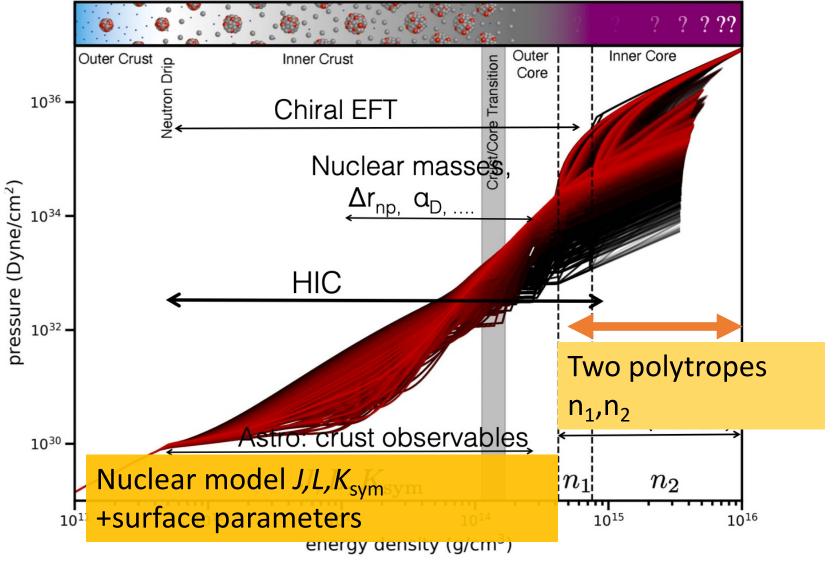


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 $\mathcal{H}_{\delta} + \mathcal{H}_{\rho} + \mathcal{H}_{eff} \quad \sigma(y_p)$ Nuclear EDF: Bulk + separate surface energy function specific model: LLPR 1985 $\sigma_{s}(y_p) = \sigma_0 \frac{2^{p+1} + b}{\frac{1}{y_p^p} + b + \frac{1}{(1-y_p)^p}}$

Model parameters



Neill+ 2208.00994; Sorenson+ 2301.13253

Astro Data

-Neutron Star radii and mass measurements from NICER -Neutron Star tidal deformabilities from LIGO

Bayesian Analysis

-Prior distribution of models -Apply **astro** data -Posterior Constraints on astro observables

Astro Observables

Neutron Star masses, radii, tidal deformability, moment of inertia, i-mode frequency, crust mass/thickness, pasta mass/thickness are calculated from NS models

High Density EOS

Polytropic model is used for high density inner core of neutron star at 1.5 and 2.7 times saturation density

Core EOS

Skyrme is used as input to core EOS up to 1.5 times saturation density

This work

Nuclear Model

Skyrme Hartree Fock energy

density functionals

parameterized by symmetry

energy values:

Bayesian Analysis

-Prior distribution of models -Apply **nuclear** data -Posterior Constraints on nuclear observables -Neutron skin thickness: PREX experiment on 208Pb and CREX from 48Ca -Dipole polarizations of 208Pb and 48Ca

Nuclear Data

Nuclear Observables

Binding energy and dipole polarizabilities of doubly magic nuclei, neutron skin thickness are calculated from nuclear models

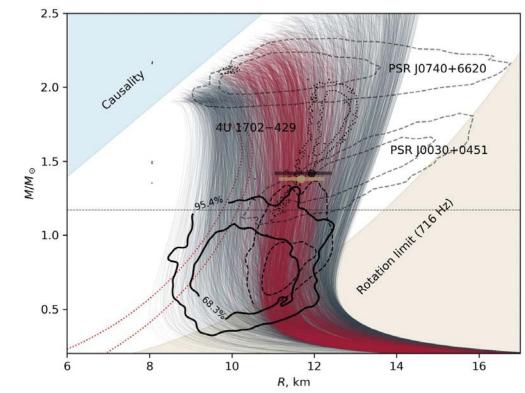
Neutron Star Model

Crust EOS

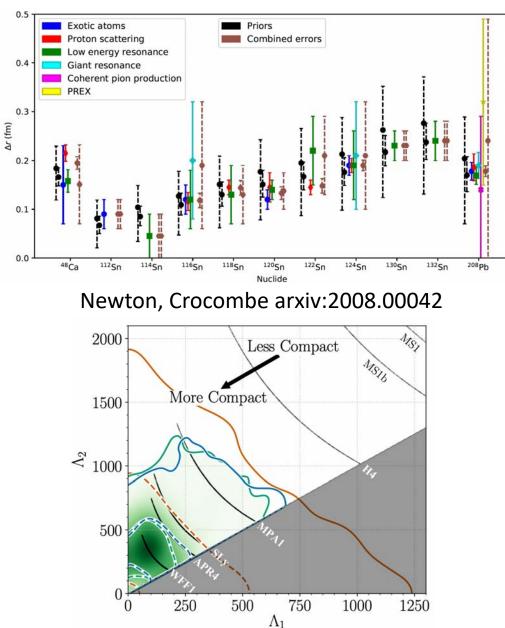
Skyrme + Compressible Liquid Drop Model is input to crust EOS

J, L, Ksym

Data

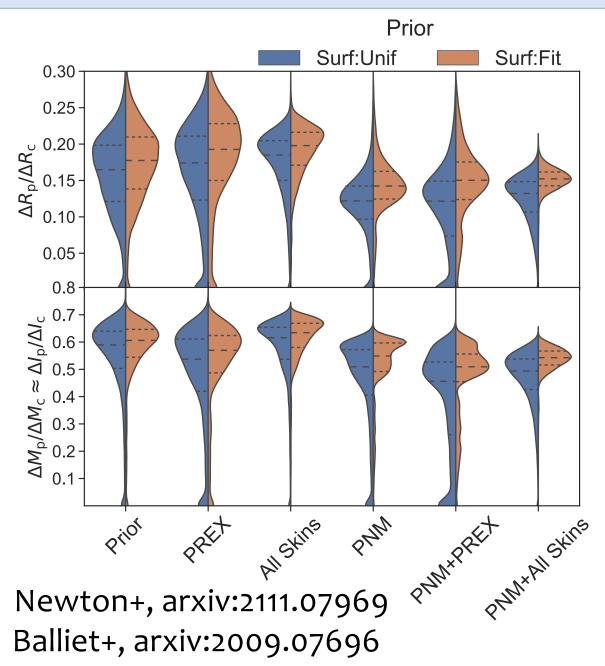


Doroshenko+, Nature Astronomy, **6**, 1444 (2022) Raajimakers arxiv: 1912.05703, 2105.06981

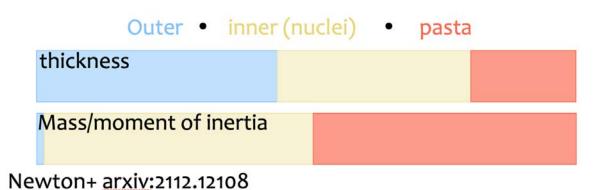


LIGO/Virgo arxiv:1805.11581

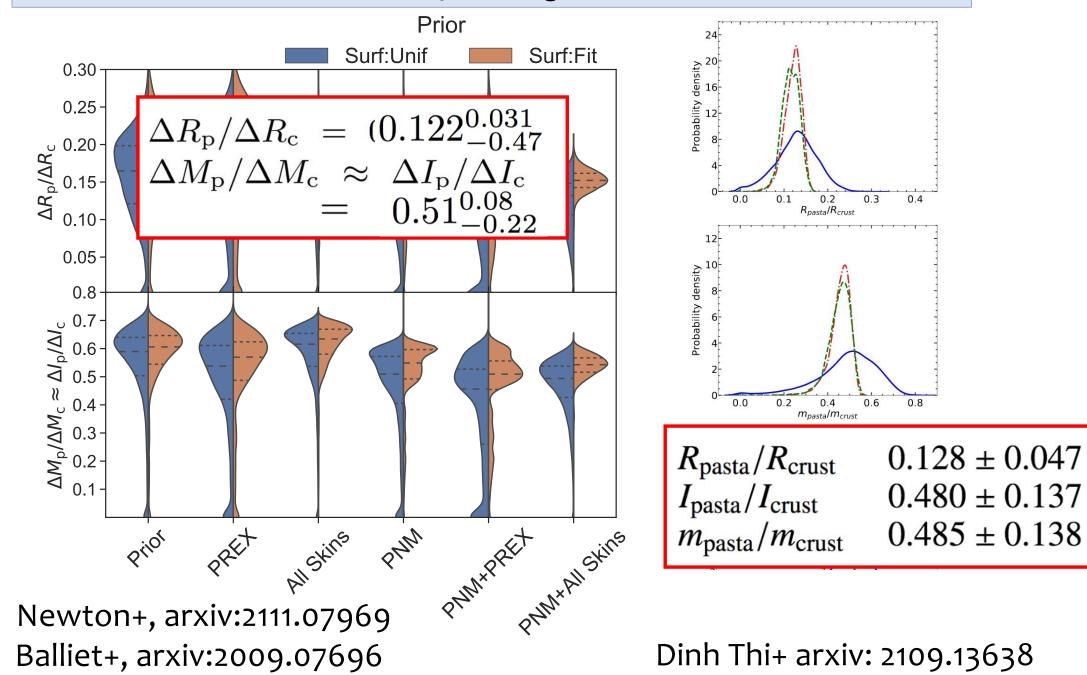
Results: Relative thickness and mass of pasta



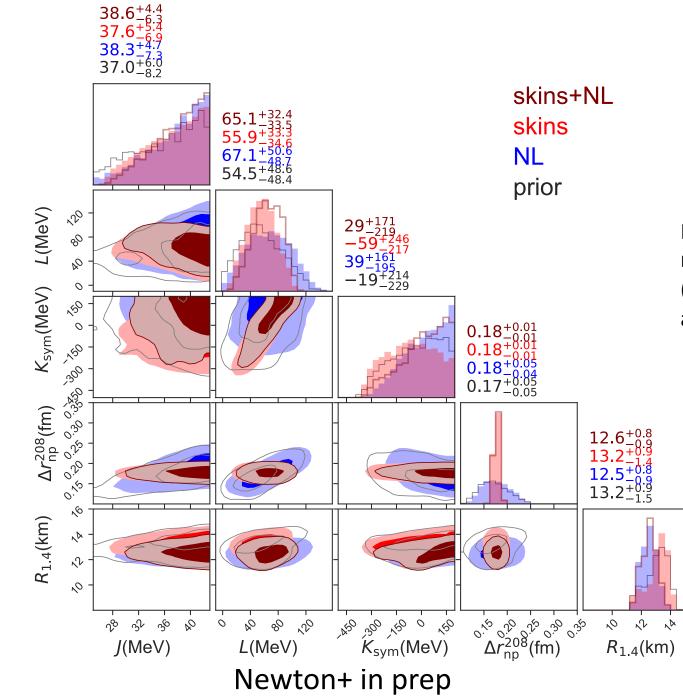
$$\Delta R_{\rm p} / \Delta R_{\rm c} = (0.122^{0.031}_{-0.47} \\ \Delta M_{\rm p} / \Delta M_{\rm c} \approx \Delta I_{\rm p} / \Delta I_{\rm c} \\ = 0.51^{0.08}_{-0.22}$$



Relative thickness and mass of pasta: agreement with other studies



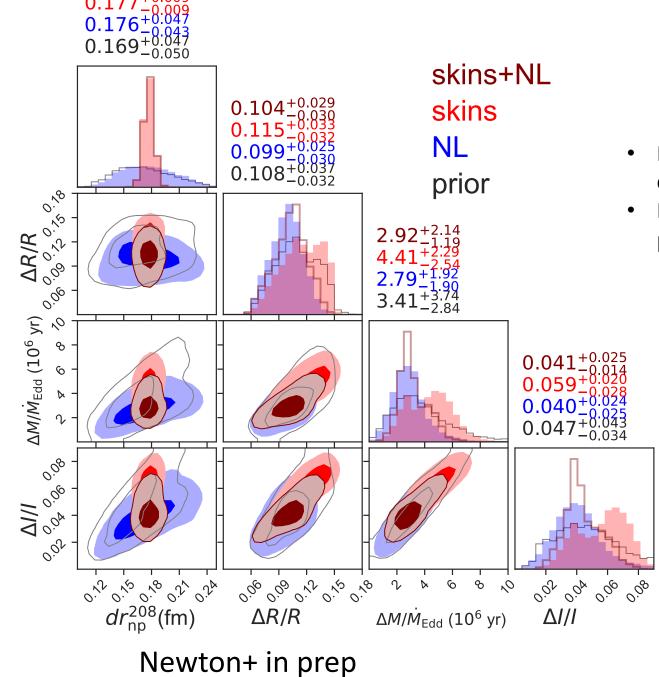
Preliminary



NL prefers stiffer *L*,*K*_{sym} neutron skins prefer softer (note: PREX alone would come to a different conclusion)

0

Preliminary

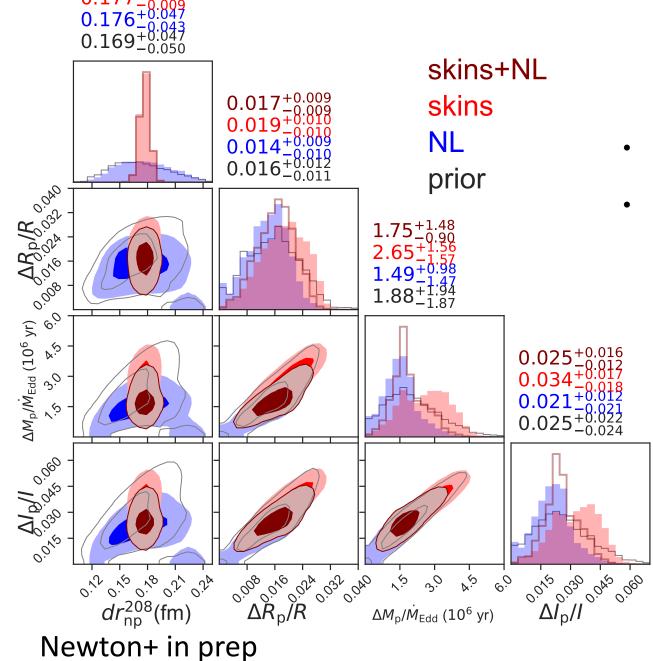


 $0.177^{+0.008}$

- NL prefers thinner, less massive crust
- Both neutron skin data and NL prove informative

Crustal glitches: 0.018 0.08 (with entrainment)

Preliminary

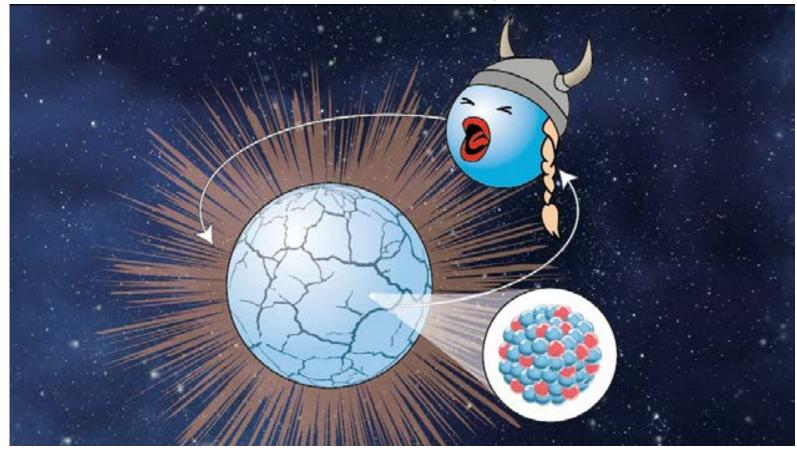


 $0.177_{-0.008}^{+0.008}$

- NL prefers thinner, less massive pasta region
- Both neutron skin data and NL prove informative

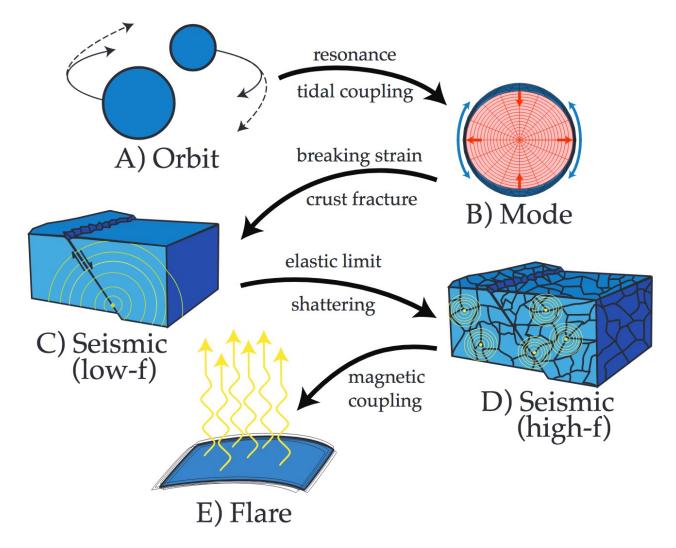
Application: resonant crust shattering flares

Neill, Newton & Tsang, MNRAS 504, 2021 Neill,Preston,Newton,Tsang, PRL130, 2022

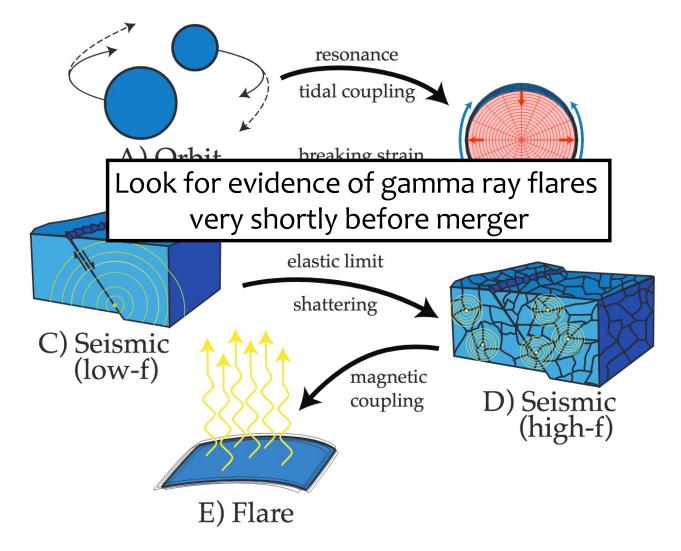


Picture: David Tsang

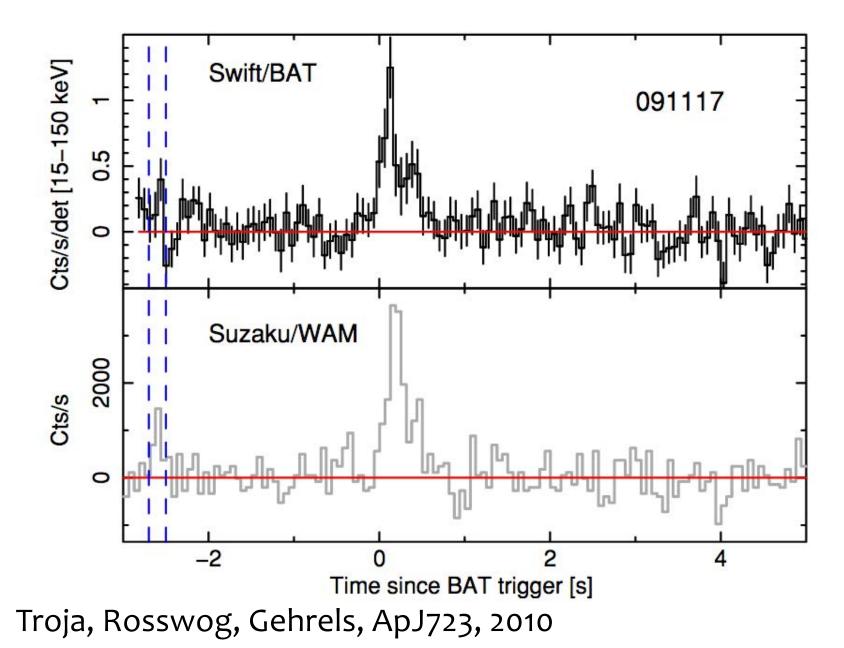
The elastic crust can be made to resonantly vibrate by the tidal field of its companion



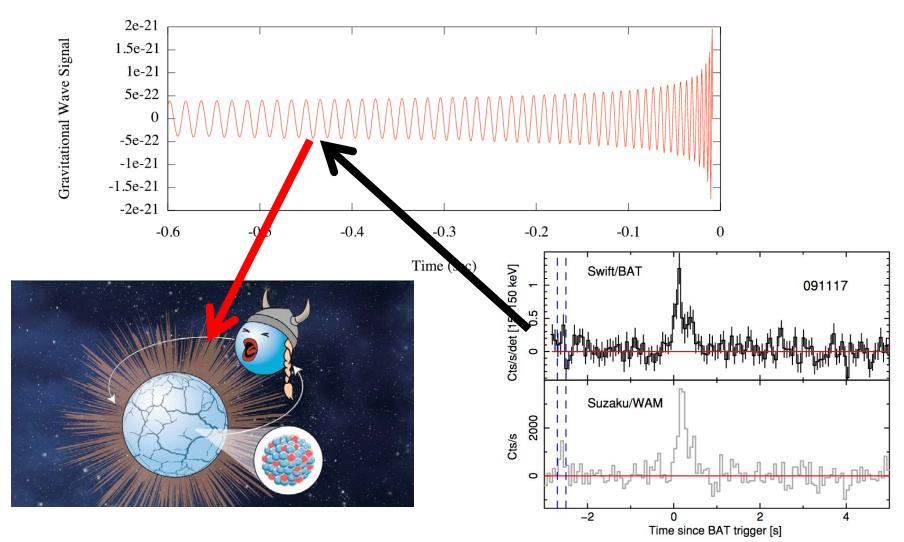
D.Tsang, Apj 777, 2013 Neill, Newton & Tsang, MNRAS 504, 2021 The elastic crust can be made to resonantly vibrate by the tidal field of its companion



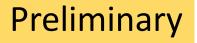
D.Tsang, Apj 777, 2013 Neill, Newton & Tsang, MNRAS 504, 2021

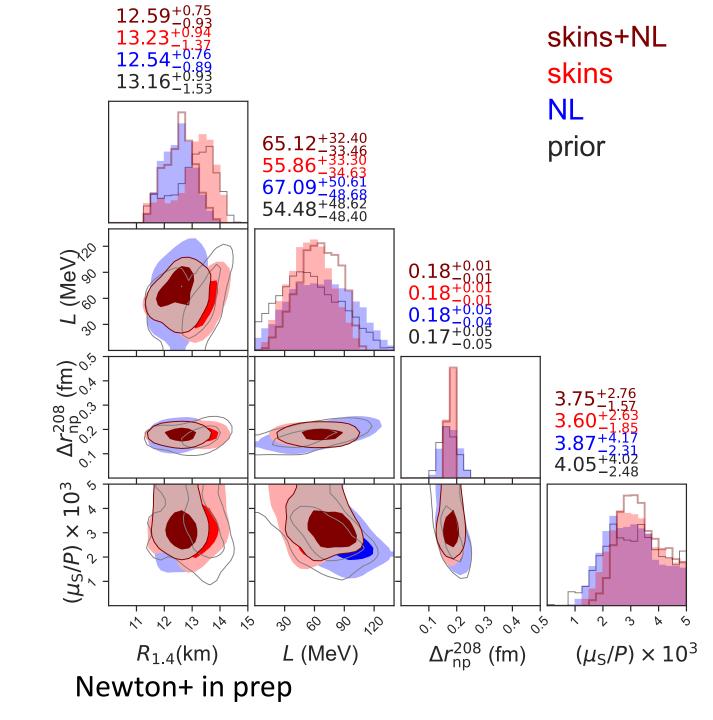


Time of gamma ray flare points out GW frequency at that time which gives the resonant frequency of the crust



Example Inspiral Gravitational Waves

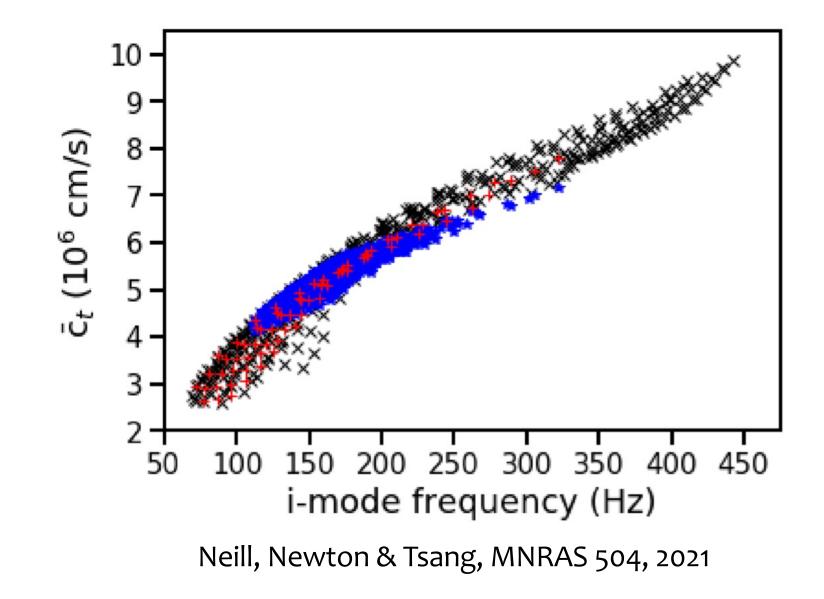




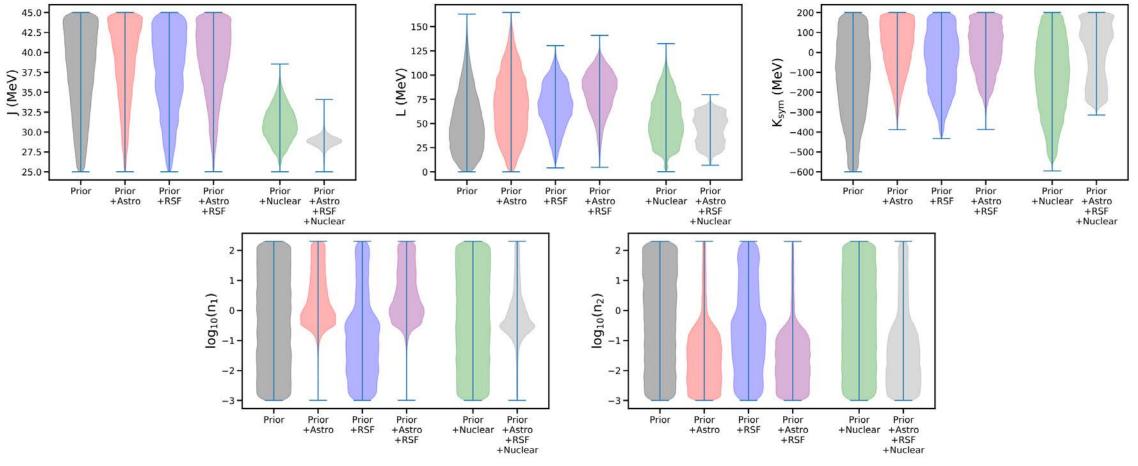
We can also obtain the posteriors on the shear modulus at the base of crust

(using form of Strohmayer+ 1991)

Strong correlation between shear speed and *i*-mode frequency

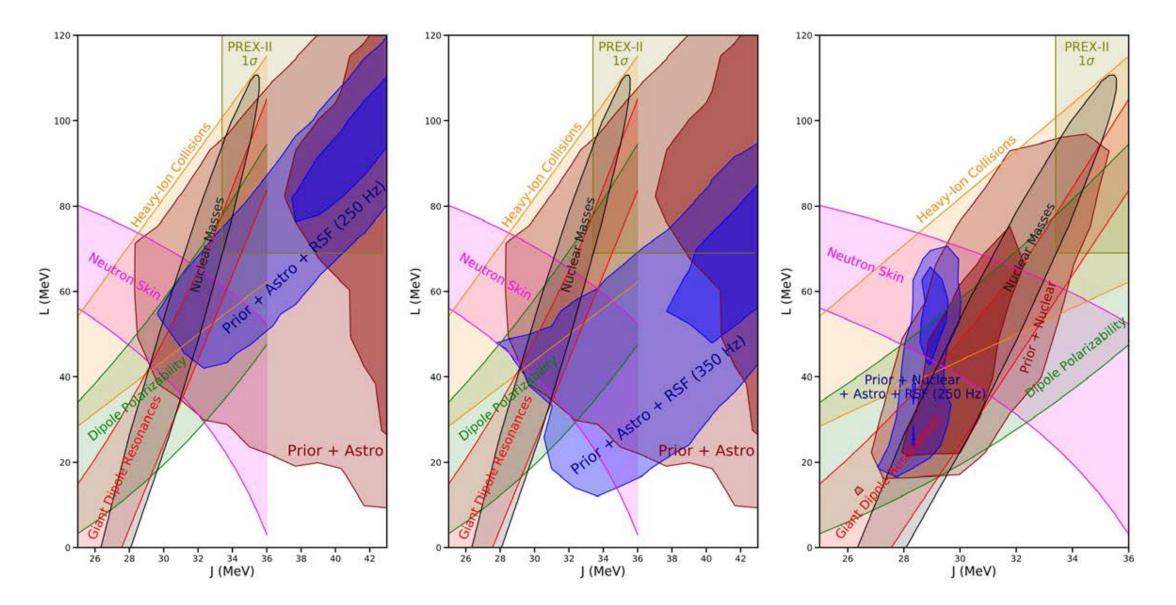


Inference using a synthetic detection of an RSF at a frequency of 250 Hz, comparison with NL and nuclear binding energy data



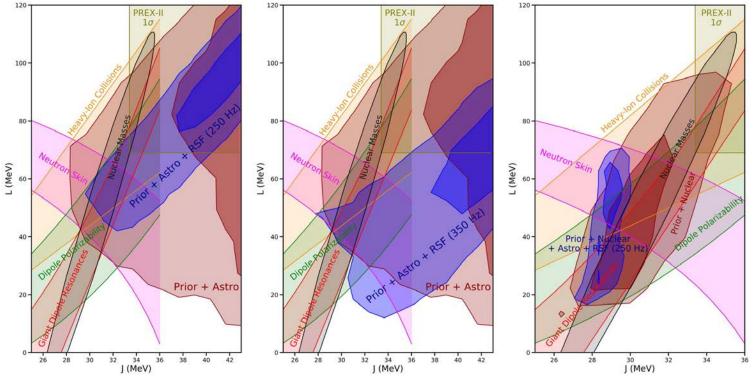
- J not constrained by RSF/NL
- Polytrope parameters not constrained by nuclear BEs
- L constrained by RSF and BE
- *K*_{sym} constrained by RSF/NL

An observation of an RSF can potentially constrain the symmetry energy



Take-aways

- We have a model that efficiently calculates crust, core and nuclear properties consistently, with polytropes added to account for uncertainty at high density
- Both nuclear and astrophysical data give us information about the neutron star crust, and consistent models are needed to take advantage of this



- During a neutron star merger, the stars may resonantly excite each other's solid crust to shattering
- A coincident detection of a flare and GW signal can measure the *i*-mode frequency and constrain the nuclear symmetry energy
- But can we know for sure that the flare is from an RSF?