

Nucleosynthesis: Connecting Nuclear Properties and Observations - II



Nicole Vassh

TRIUMF Theory Group

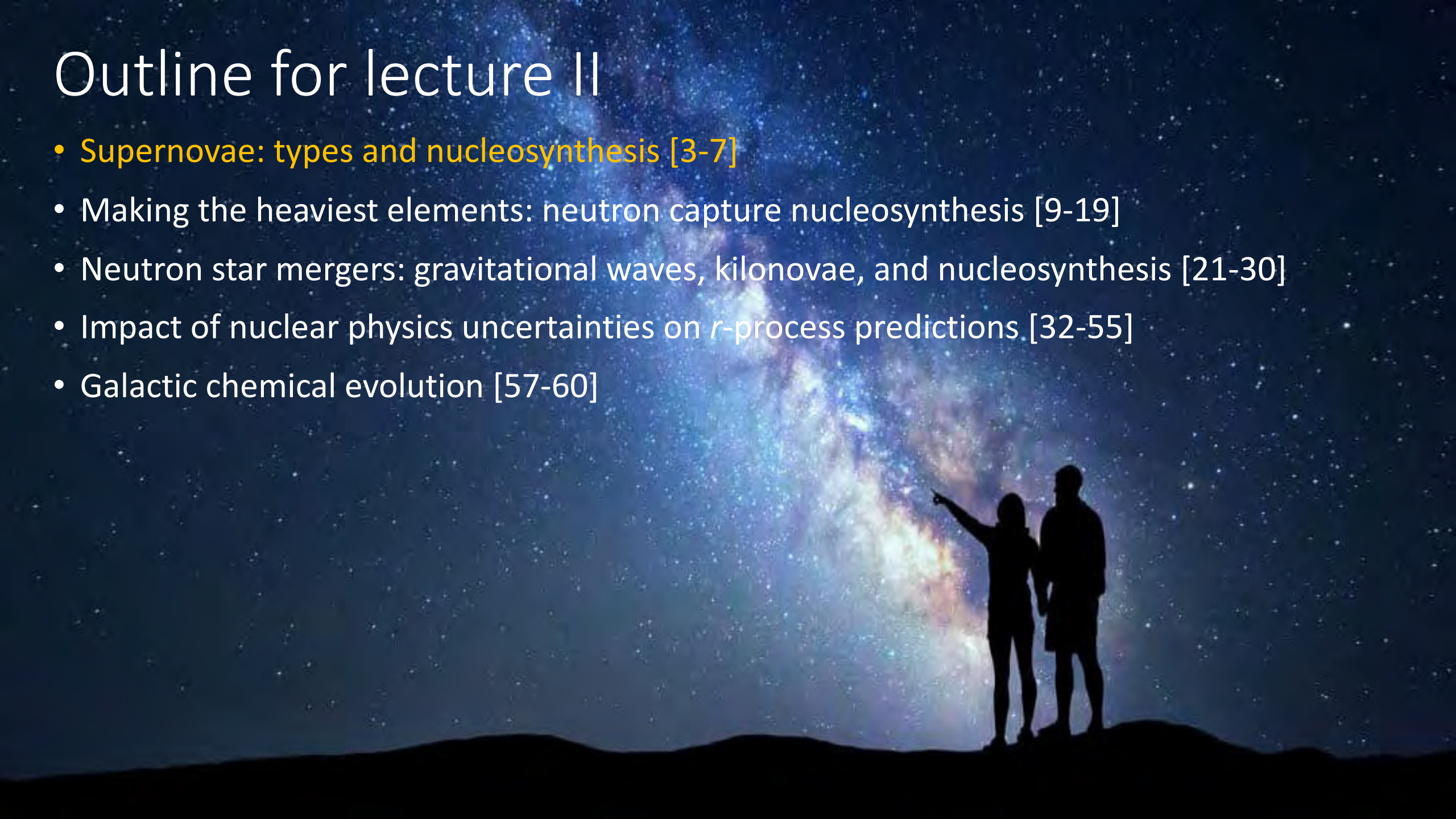
N3AS Summer School Lecture,

Santa Cruz, CA

July 20, 2023

Outline for lecture II

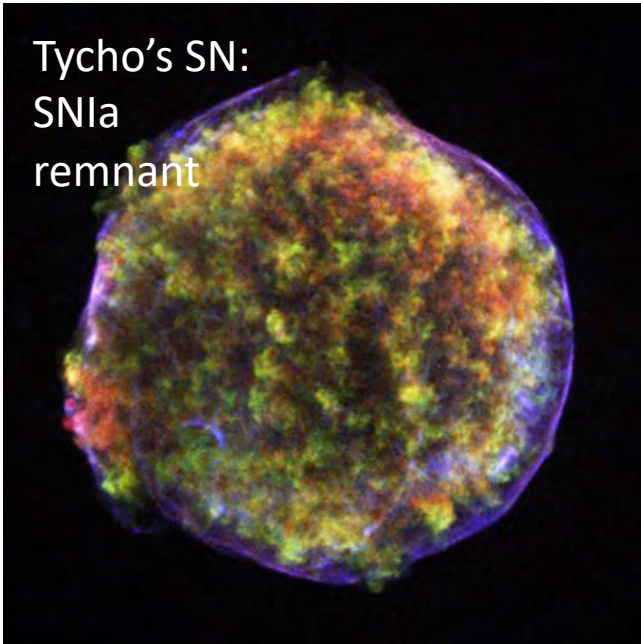
- Supernovae: types and nucleosynthesis [3-7]
- Making the heaviest elements: neutron capture nucleosynthesis [9-19]
- Neutron star mergers: gravitational waves, kilonovae, and nucleosynthesis [21-30]
- Impact of nuclear physics uncertainties on *r*-process predictions [32-55]
- Galactic chemical evolution [57-60]



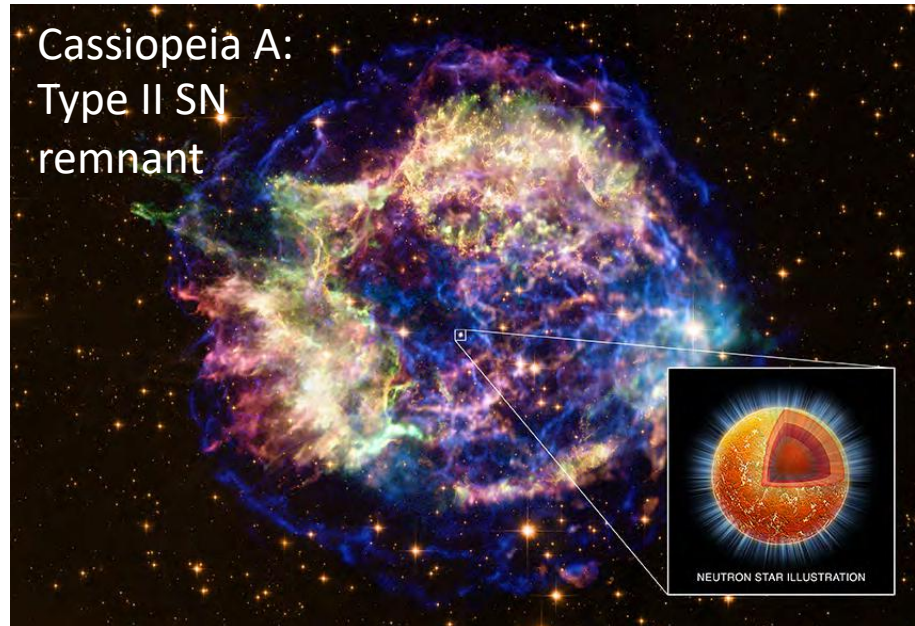
Supernova Types



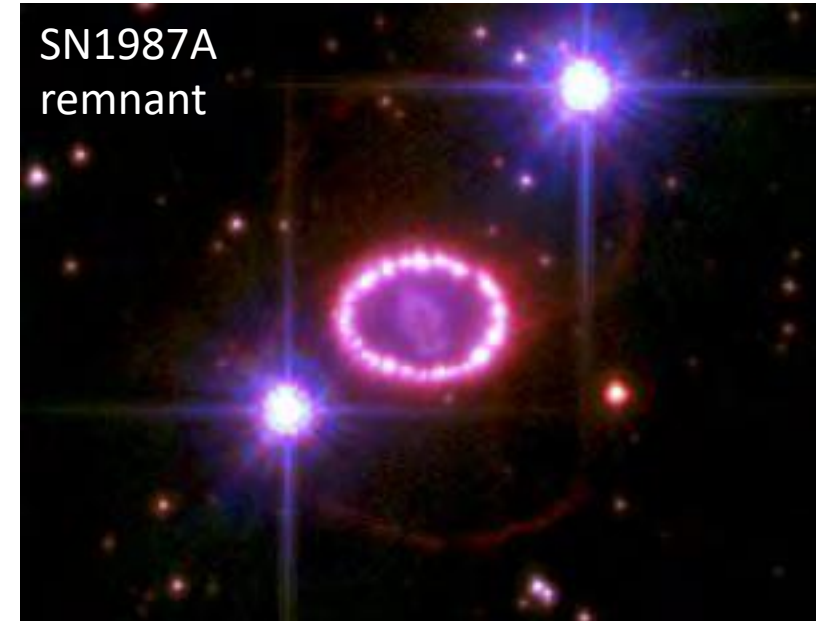
Tycho's SN:
SNIa
remnant



Cassiopeia A:
Type II SN
remnant



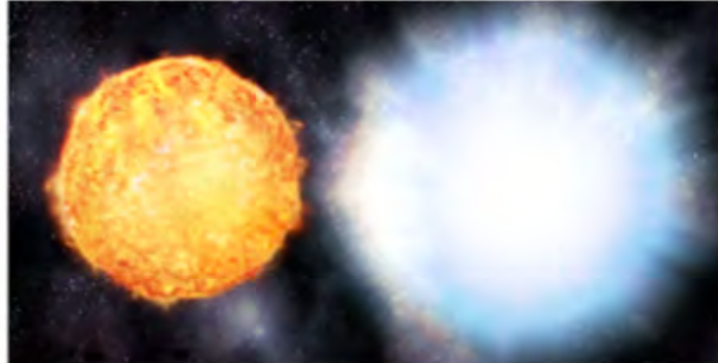
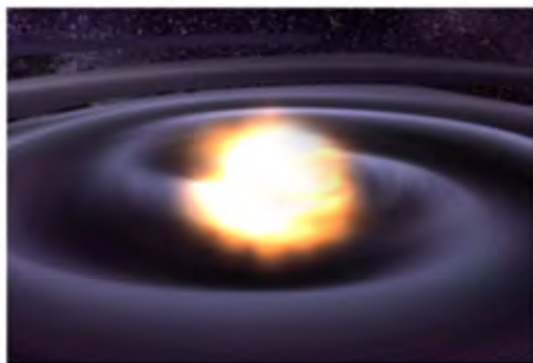
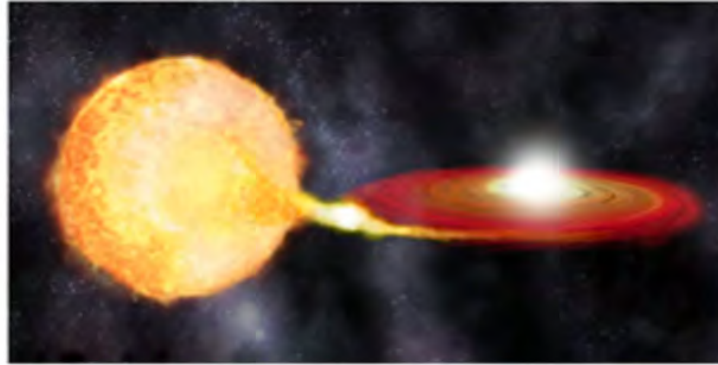
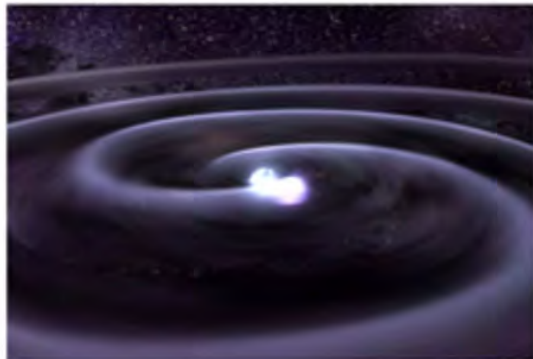
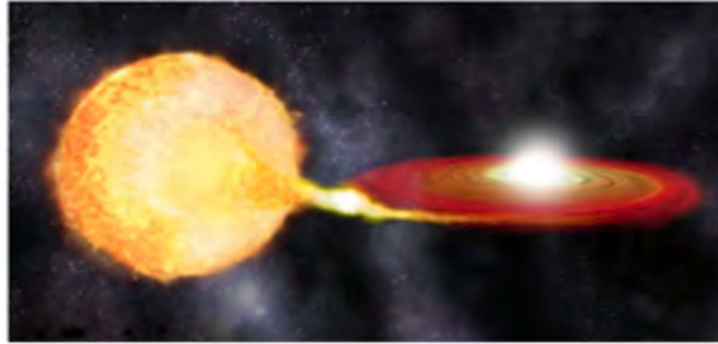
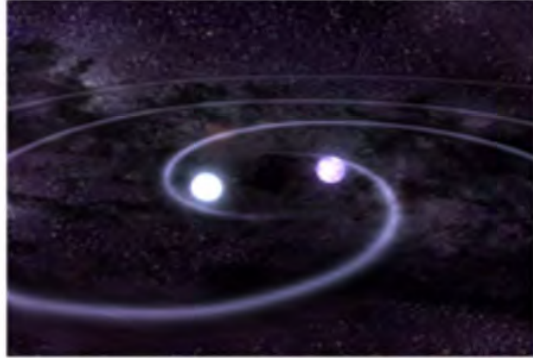
SN1987A
remnant



Thermonuclear Supernovae (Type Ia)

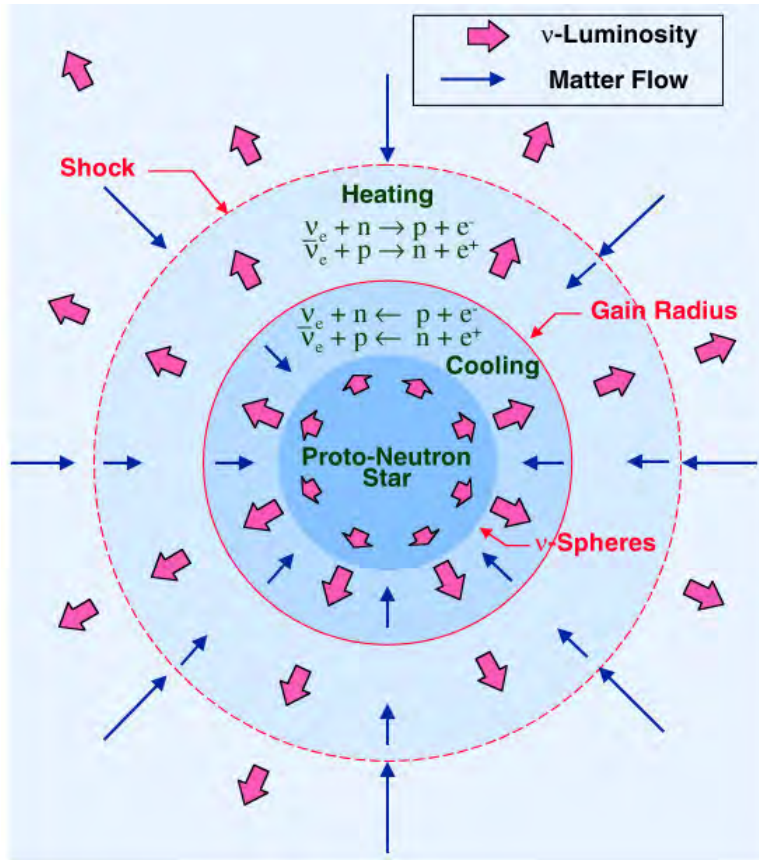
Double degenerate model

Single-degenerate model



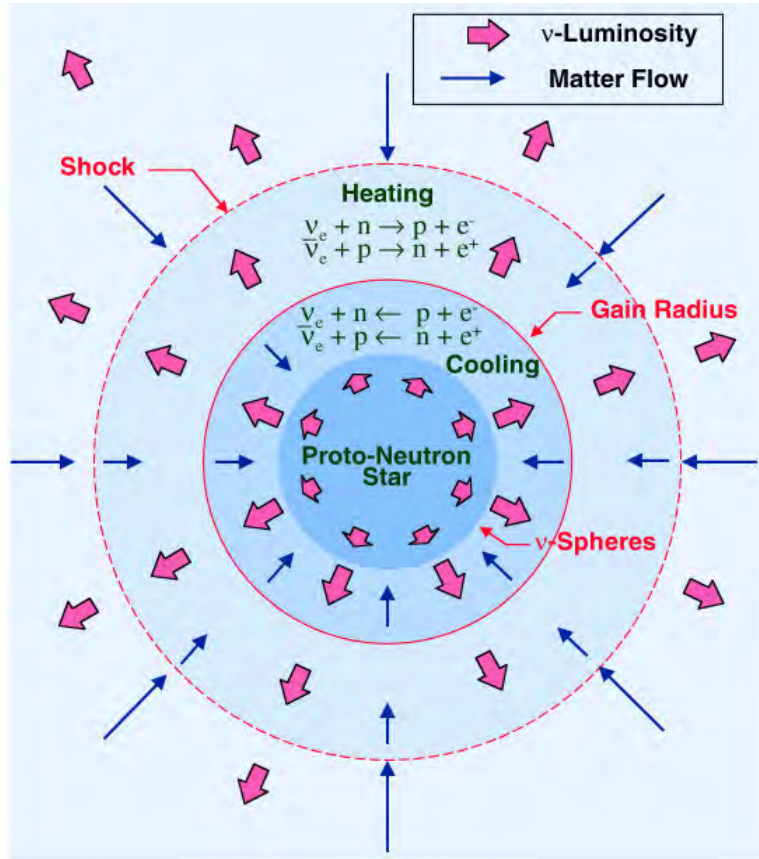
- Single-degenerate model:
 - C-O white dwarf accretes H- or He-rich matter from a companion (main sequence star, red giant, or helium star); mass of white dwarf increases until approaches Chandrasekhar limit ($1.4 M_{\odot}$), triggering explosion
 - Explains similar peak luminosity and early spectra for SN-Ia since $1.4 M_{\odot}$ implies natural limit on ^{56}Ni
 - main problem is must accrete $0.3 M_{\odot}$ to explode since max white dwarf mass is $1.1 M_{\odot}$
- Double-degenerate model:
 - Two C-O white dwarfs merge due to gravitational wave radiation, triggering explosion
 - Does not easily explain similar peak luminosity of SN-Ia due to wide range of ^{56}Ni production

Core-collapse SN and neutrino-driven winds

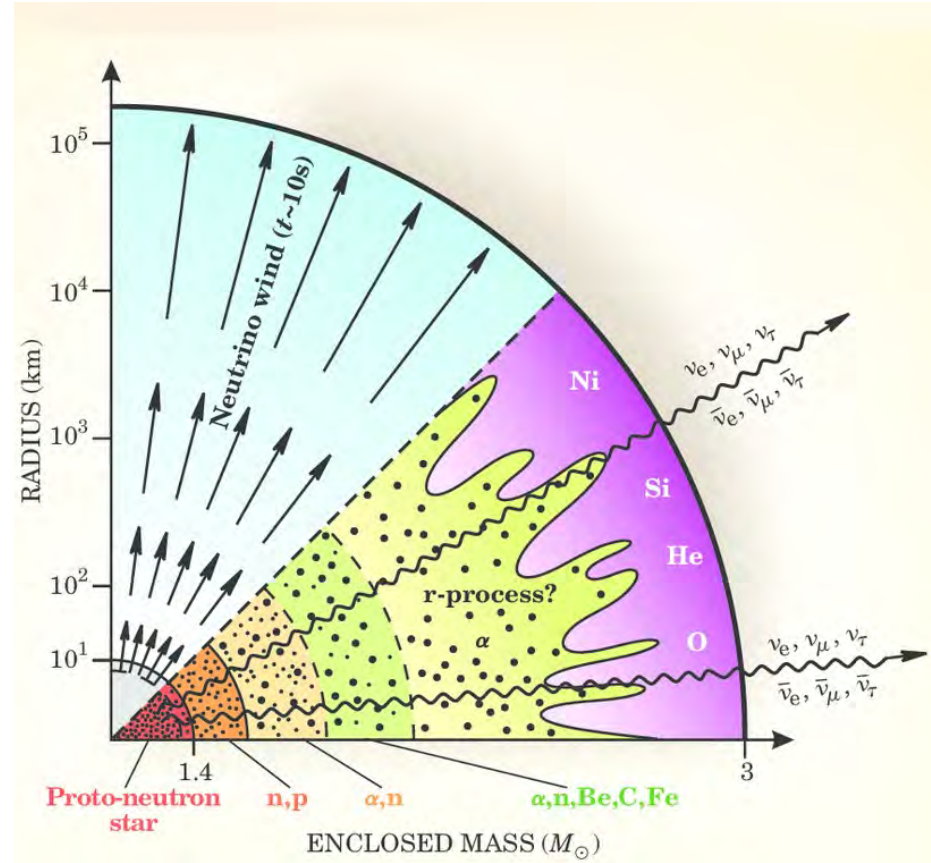


- The proto-neutron star cools through neutrino emission (99% star's binding energy released as ν)
- Cooling via *Urca processes* (lepton + baryon \rightarrow baryon + ν) as well as $e^+e^- \rightarrow \nu_l\bar{\nu}_l$ where $l = e, \mu, \tau$

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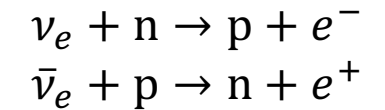


Woosley&Janka 06; see also Panov&Janka 08

Neutrinos set the neutron to proton ratio

$$Y_e = \frac{n_p}{n_p + n_n}$$

via weak interactions



and the influence of these reactions depends on the neutrino luminosities and average energies

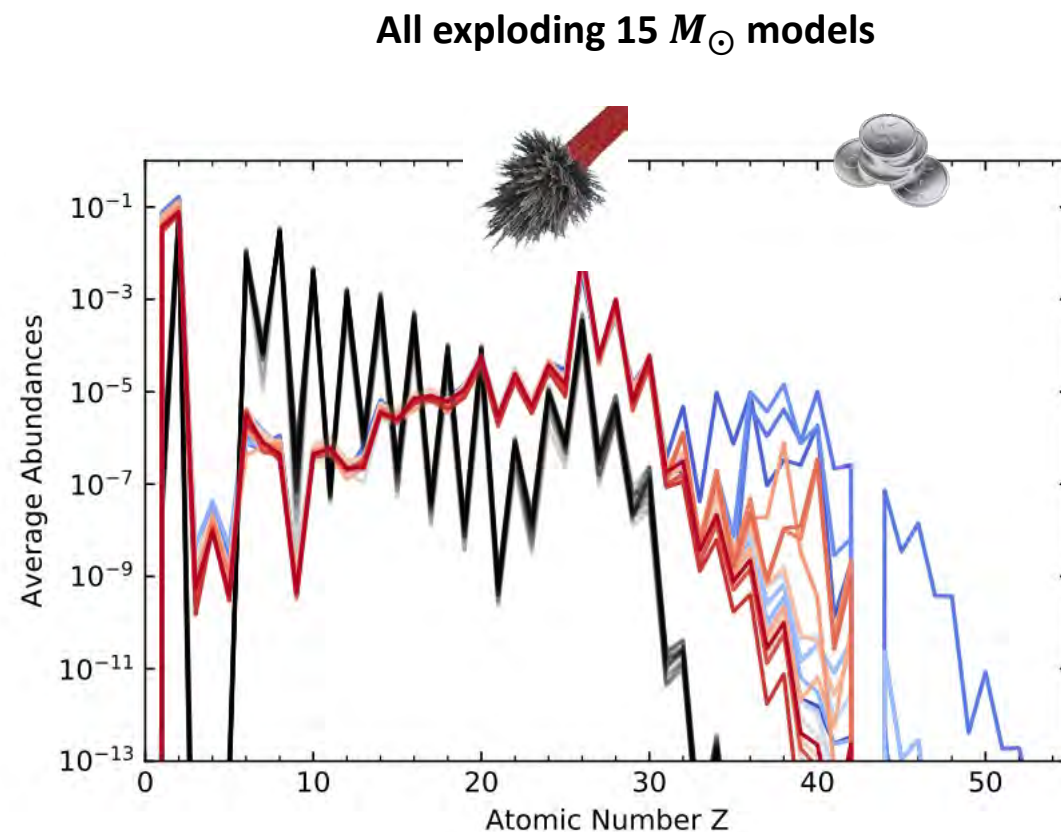
Supernovae and heavy elements?

Light heavy elements and (α, n) in core-collapse SN

Conditions which synthesize $A > 130$ are not found by most modern core-collapse SNe simulations
(e.g. Arcones+07, Wanajo+09, Fischer+10, Hudepohl+10)

In such events other processes such as (α, n) and νp process could reach up to $A \sim 100$
(e.g. Pruet+06, Fröhlich+06, Bliss+18)

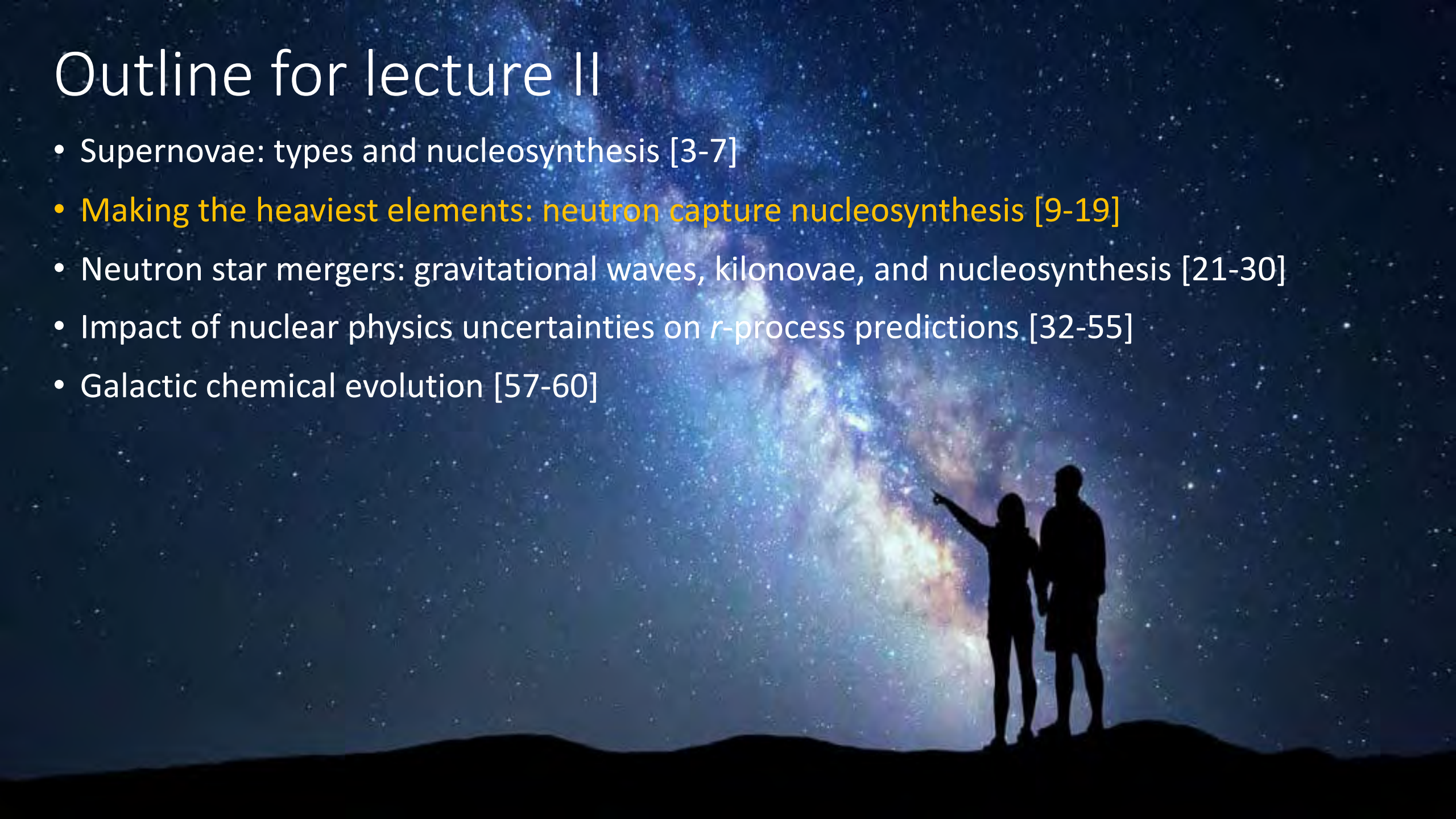
Recent simulations (right) find some cases develop neutrino driven winds but not a standard feature for successful explosions



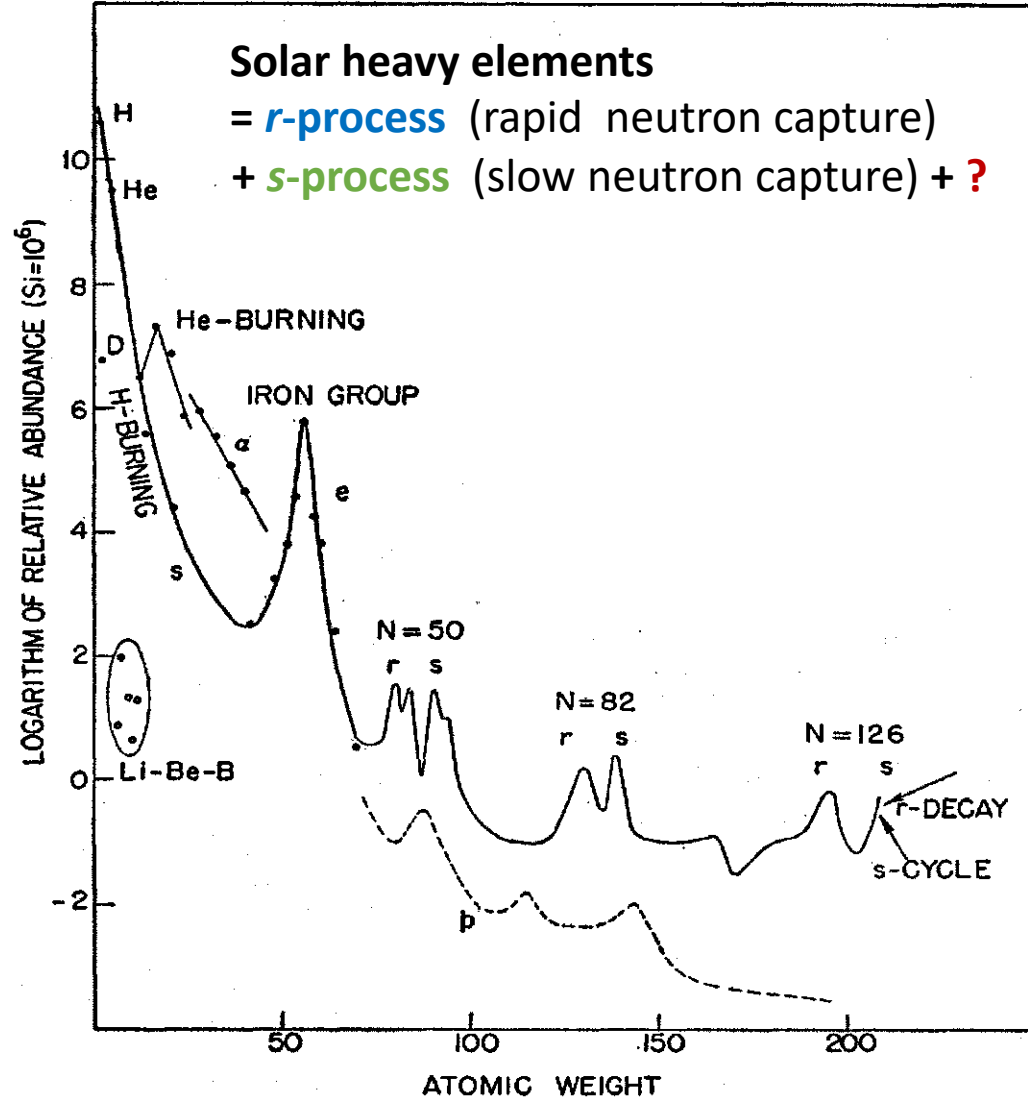
Witt+21; see also
Bliss+18

Outline for lecture II

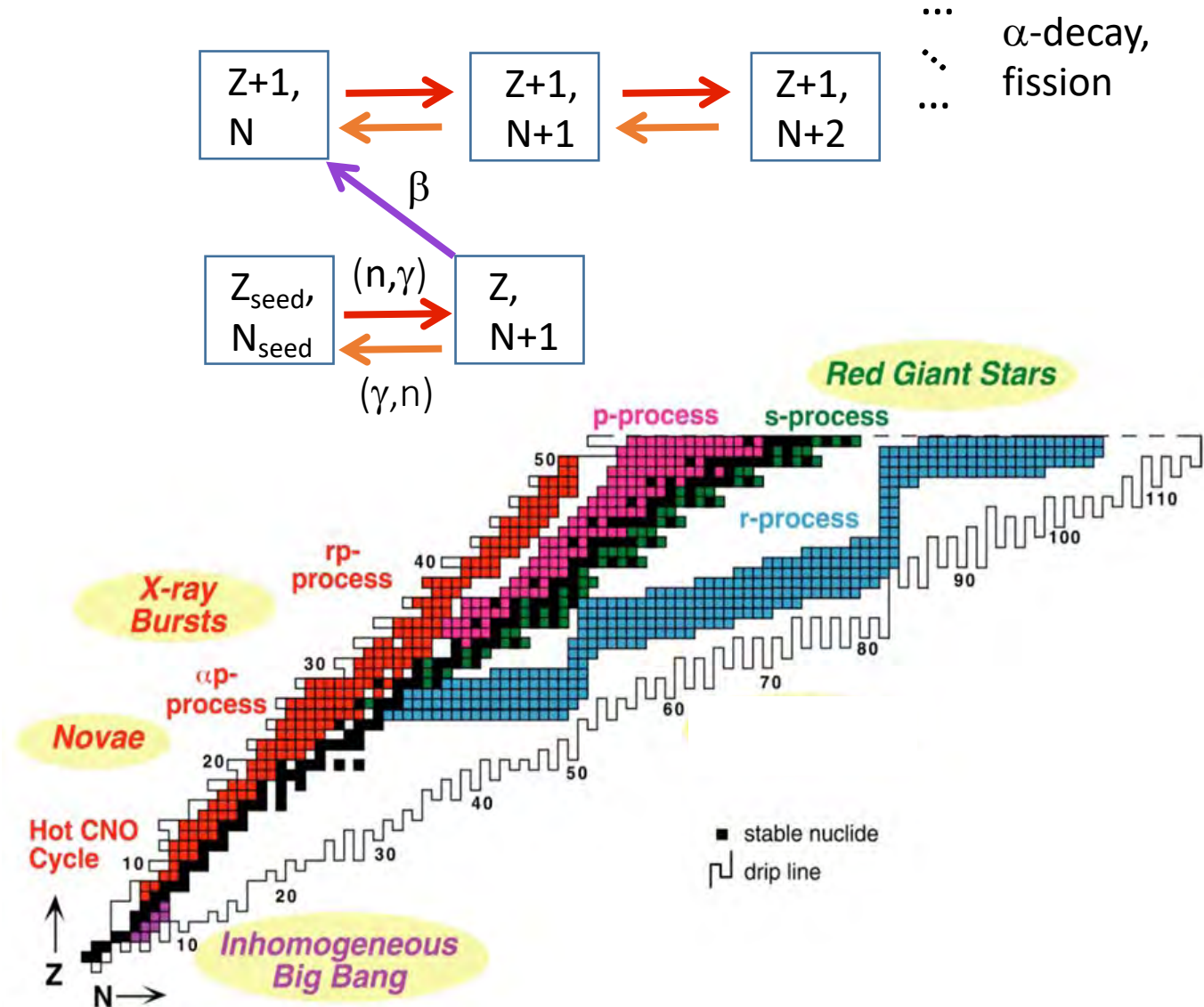
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Neutron capture processes to make the heaviest elements

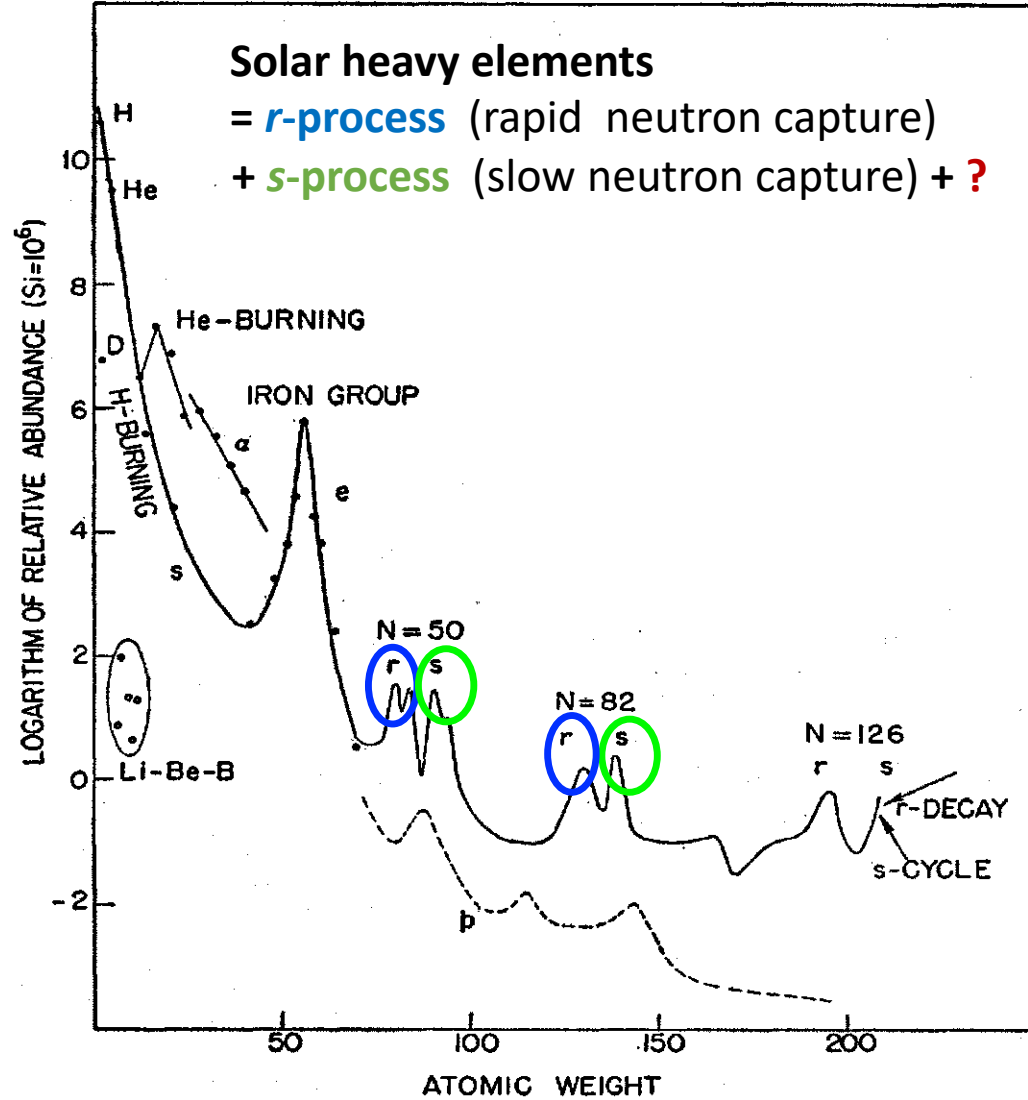


Burbidge, Burbidge, Fowler, and Hoyle (B²FH) (1957)

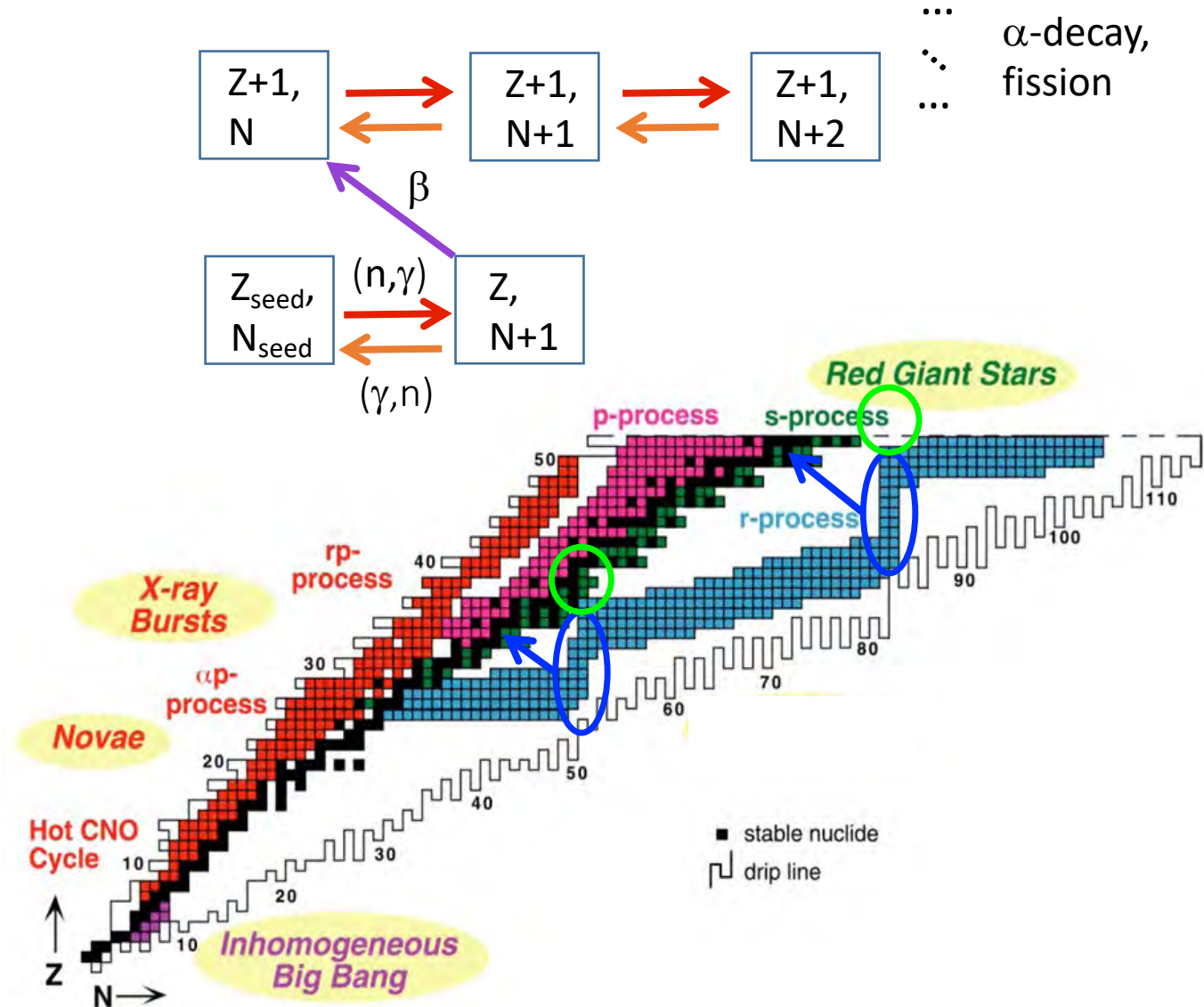


Smith&Rehm 01

Neutron capture processes to make the heaviest elements

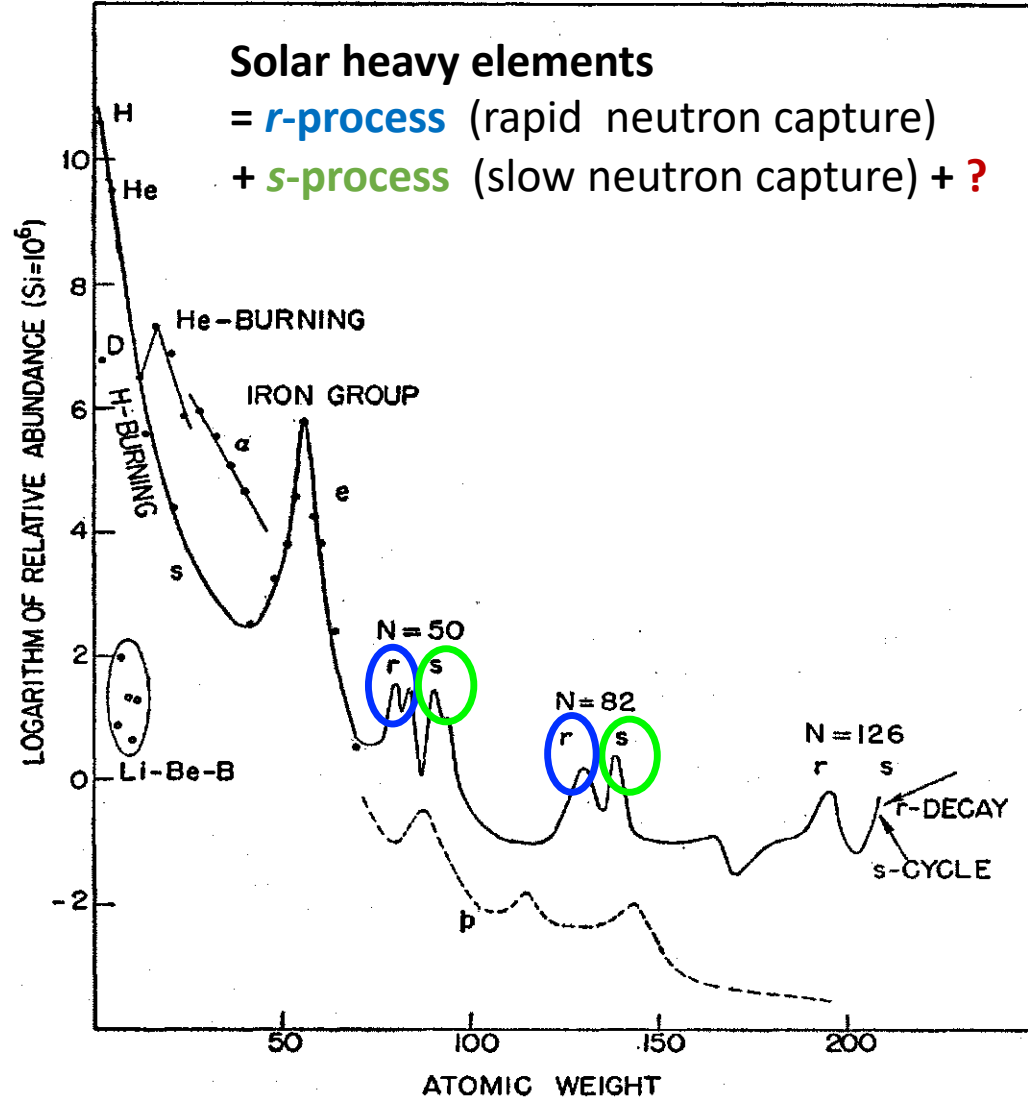


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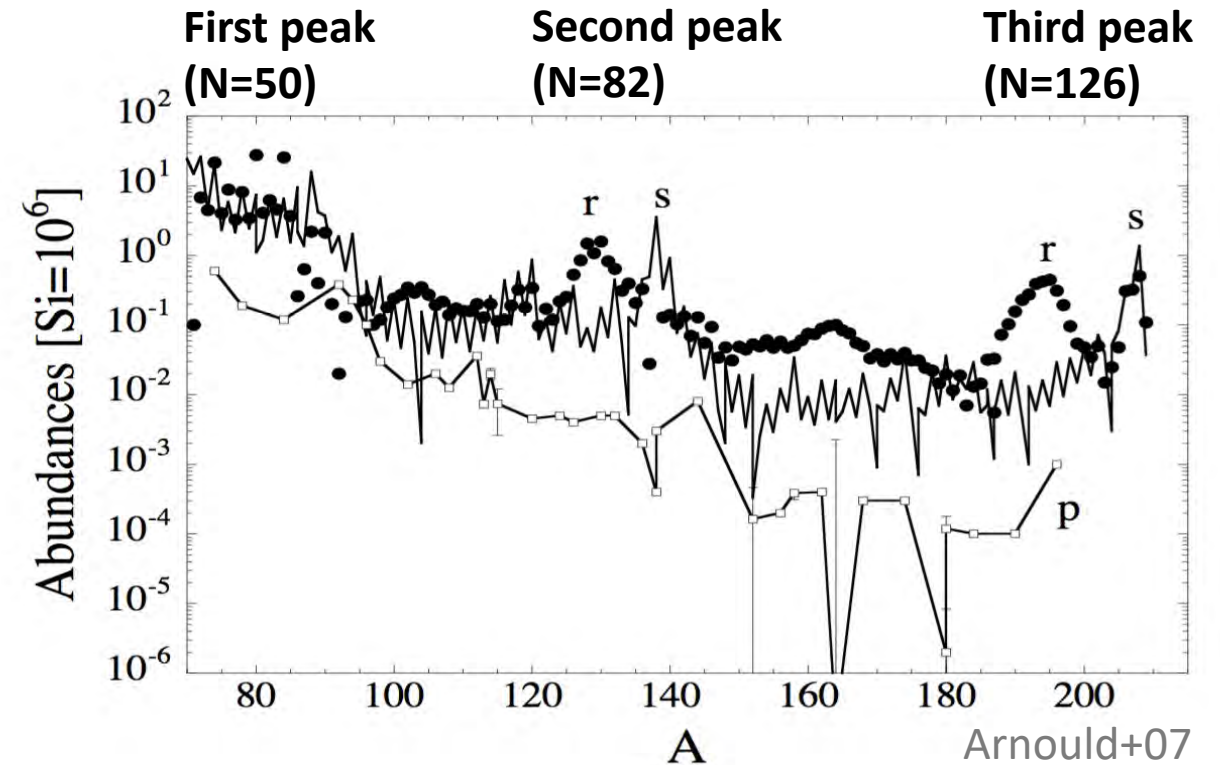
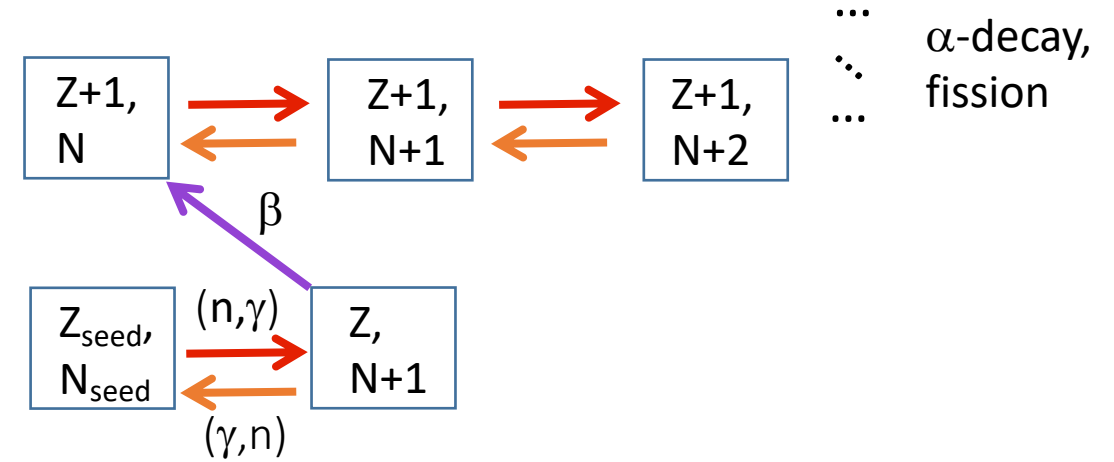


Smith&Rehm 01

Neutron capture processes to make the heaviest elements



Burbidge, Burbidge, Fowler, and Hoyle (B²FH) (1957)



Slow neutron capture (*s*-process) pathway

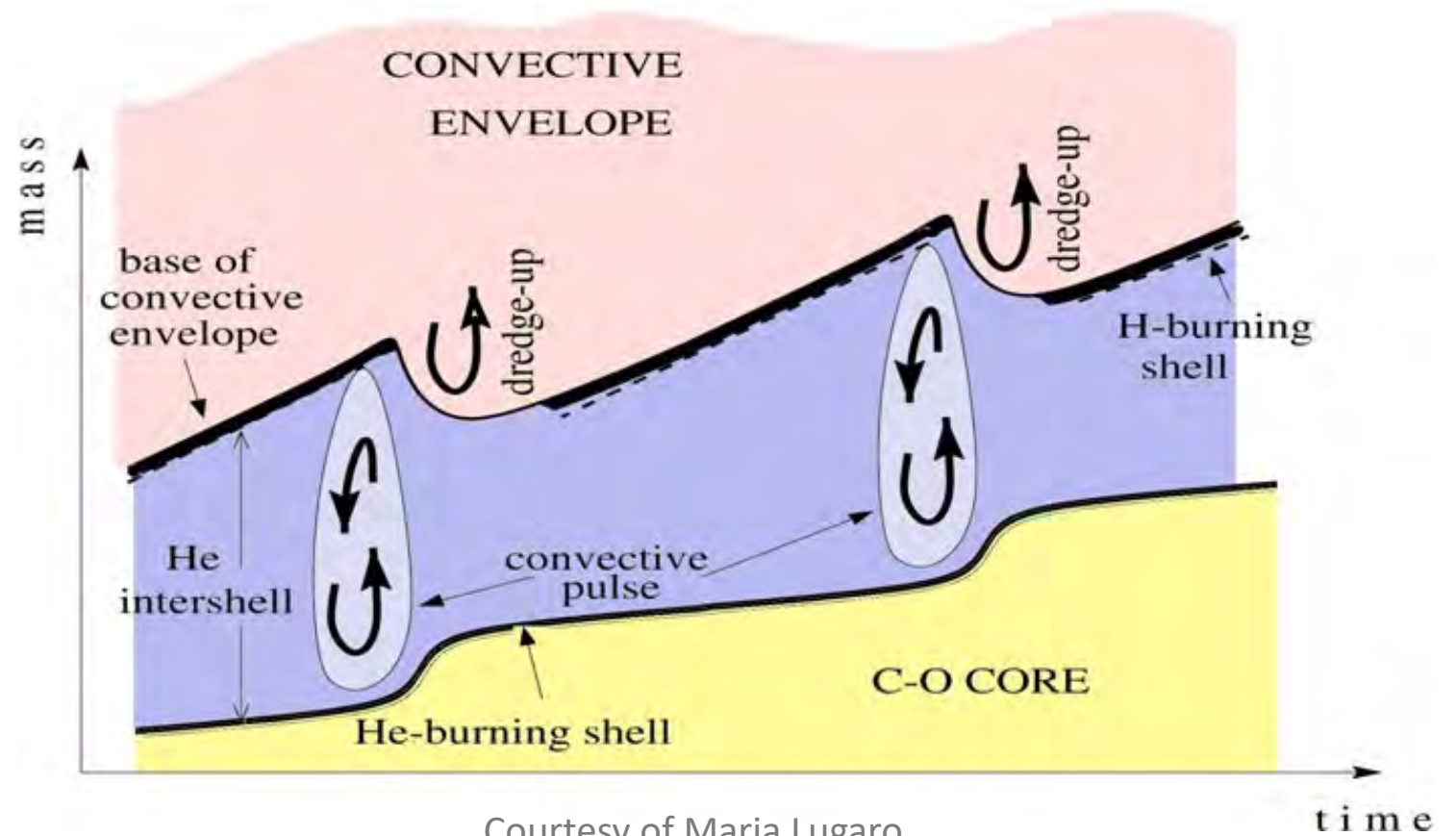
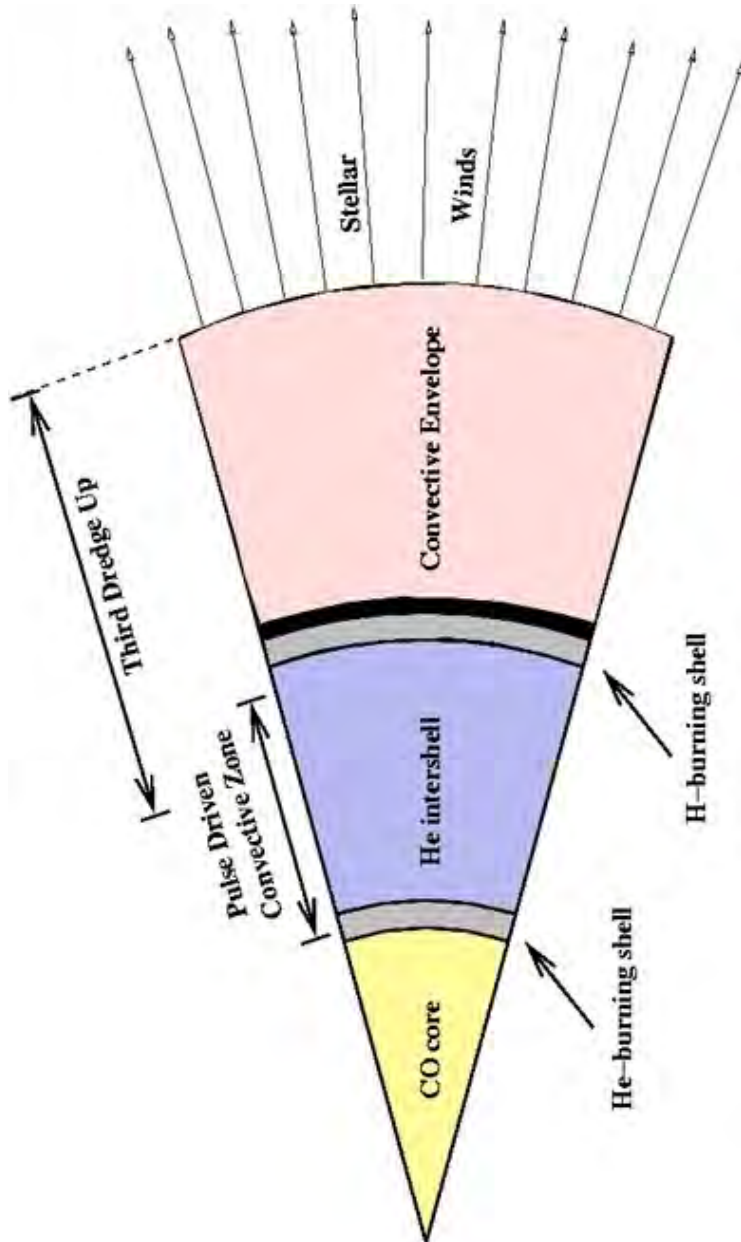
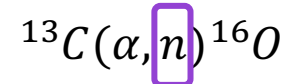
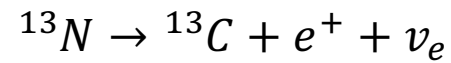
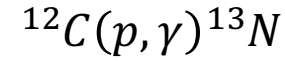


- *s*-process number density of neutrons $\sim 10^8 \text{ cm}^{-3}$ (compare to $\sim 10^{24} \text{ cm}^{-3}$ for the *r* process)
- Capture is “slow” relative to β -decay; implies a path close to stable species
- Note how the different paths imply some isotopes to be *s*-only or *r*-only
- *s*-process “seeds” are heavy nuclei such as ^{56}Fe (star enriched by past events)

Courtesy of Maria Lugaro

Slow neutron capture (s-process) in AGB stars

Where do the neutrons come from?

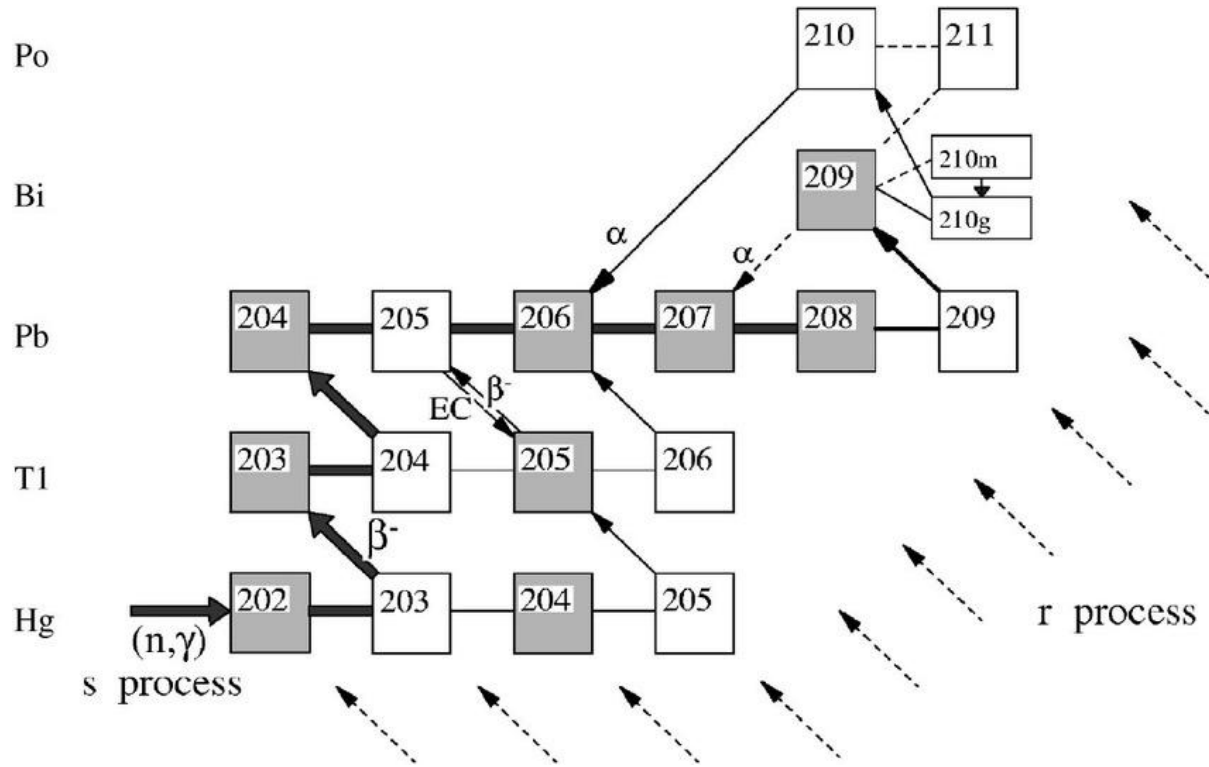


Courtesy of Maria Lugaro

How do we know there is something more than an s-process?

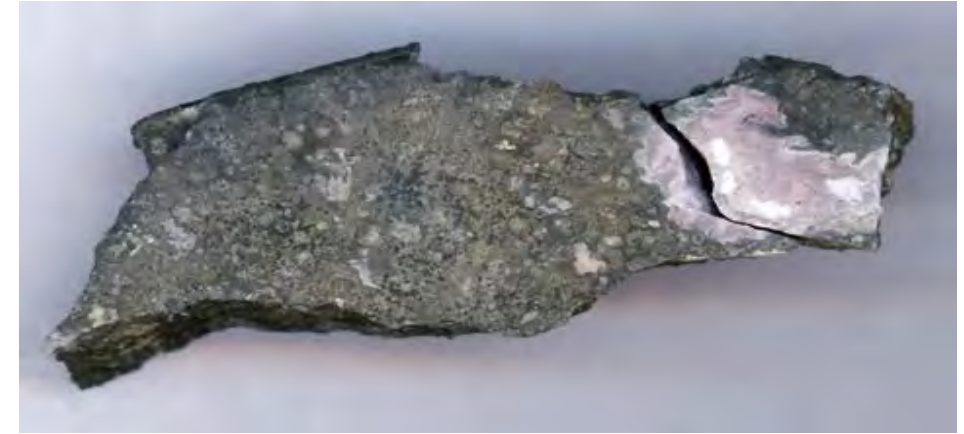
Actinides (Z=89-103)

The s-process terminates at Pb-208 (Z=82) but *we observe actinides* in meteorites, Earth ocean crusts, our Sun, and other stars

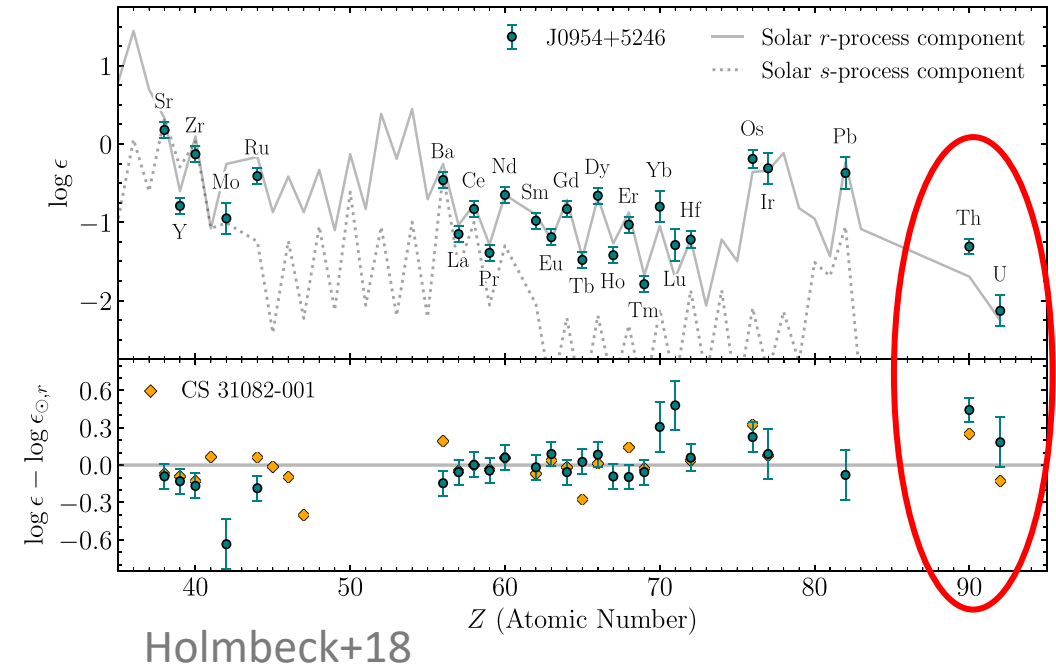


Ratzel+04

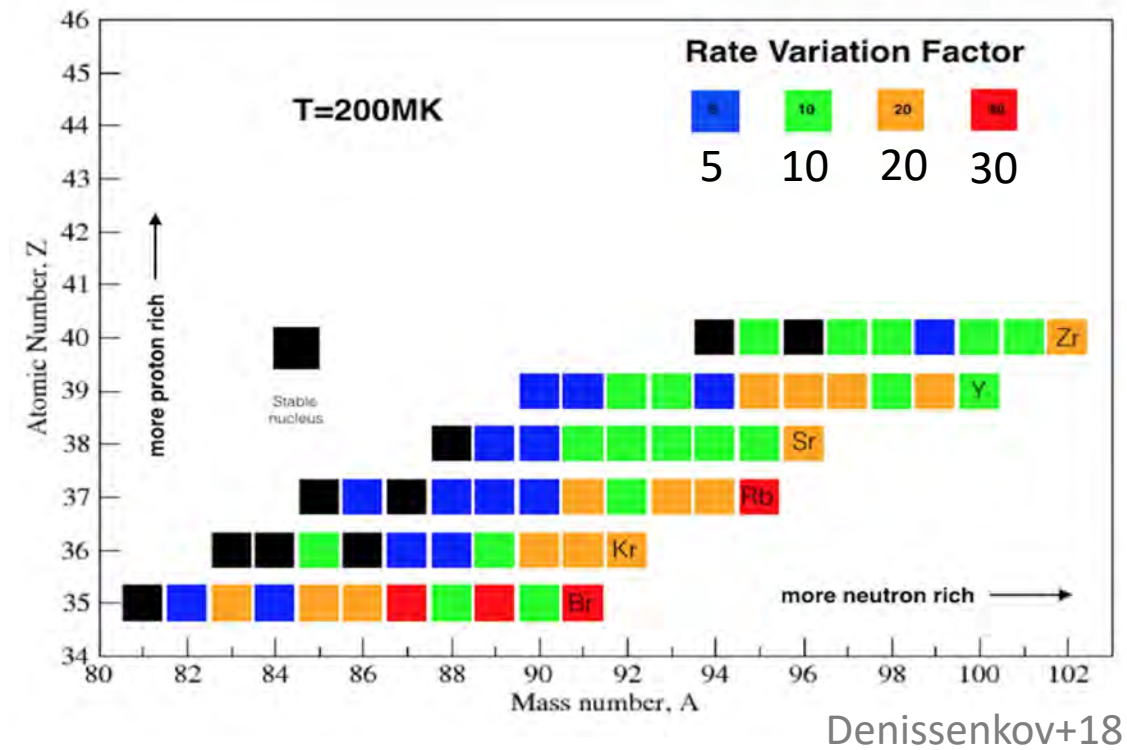
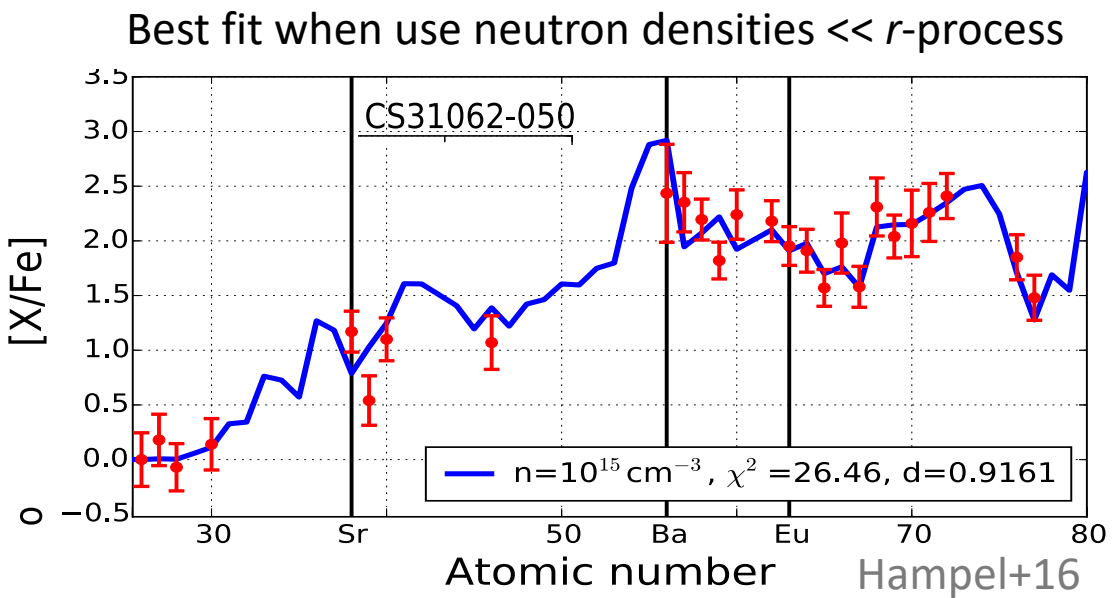
“Curious Marie” sample of Allende meteorite shows excess U-235 which is a trace of Cm-247



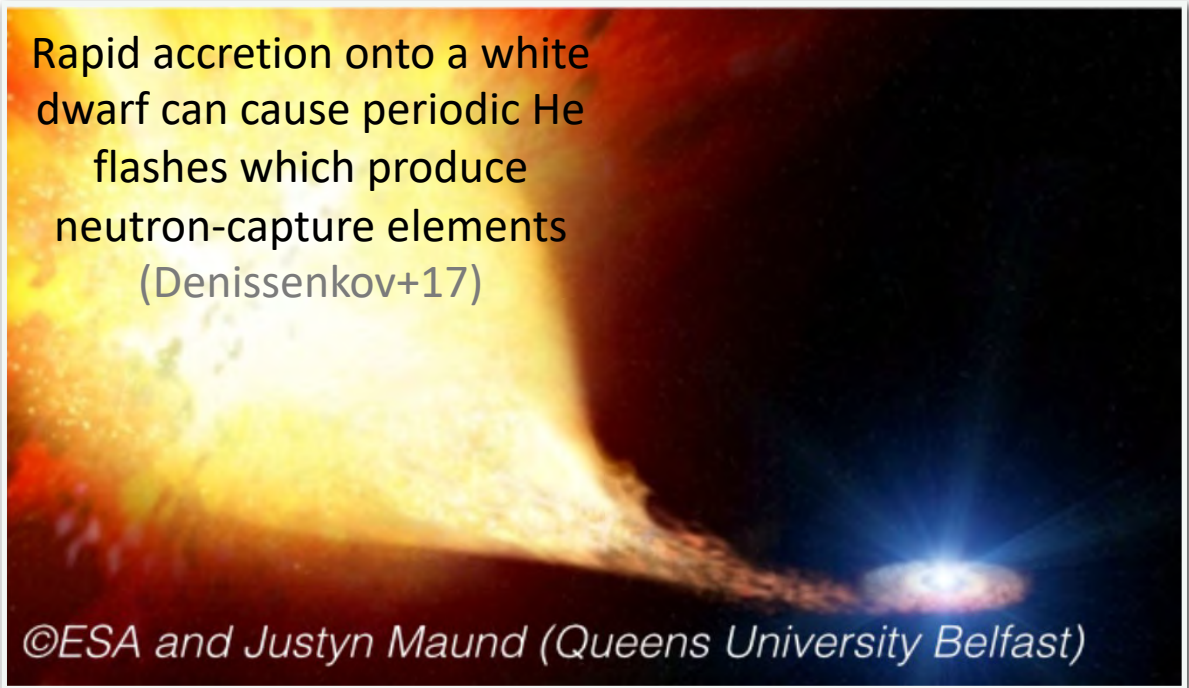
Actinide boost stars compared to solar



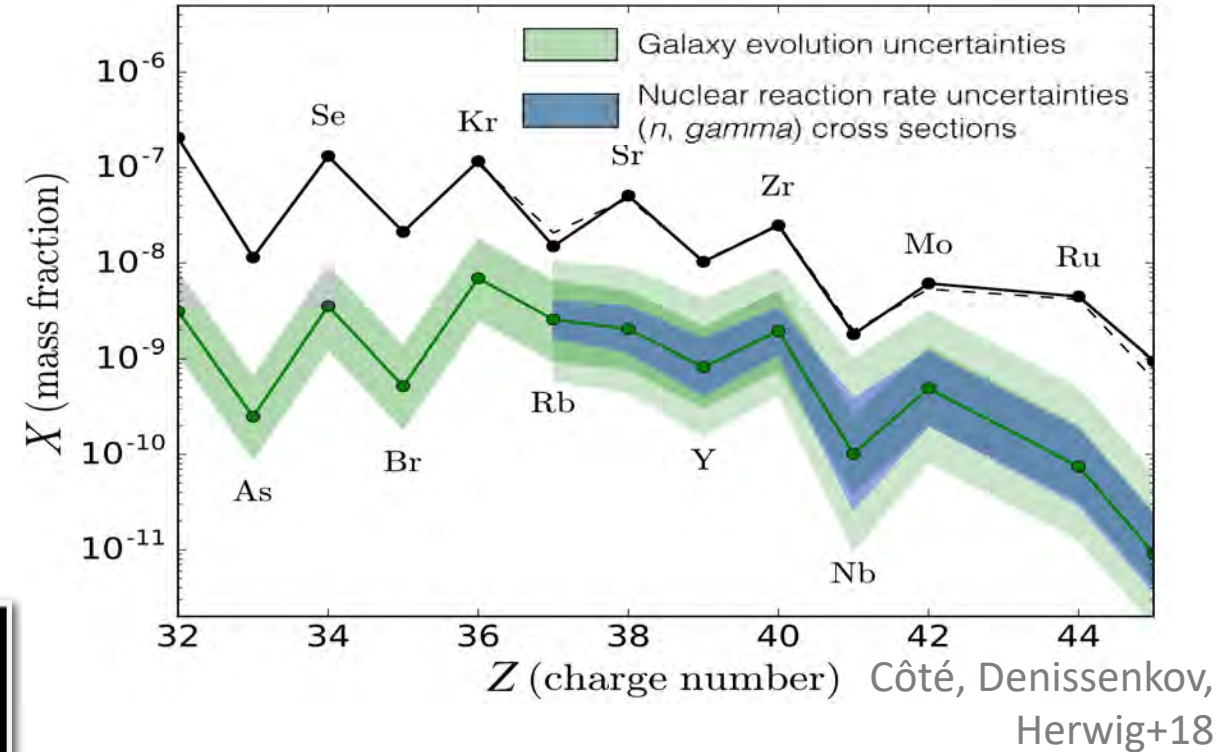
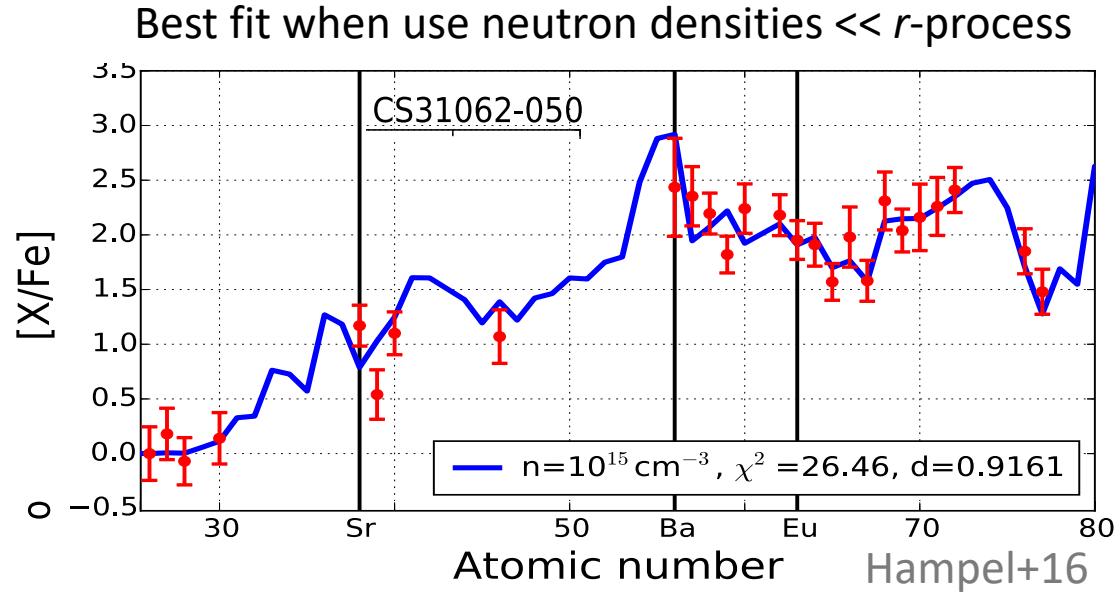
The intermediate neutron capture process (*i*-process): CEMP-i stars and (n, γ) uncertainties



(n, γ) rate variations informed by variations on theoretical **γ -strength functions** and **nuclear level densities**



The intermediate neutron capture process (*i*-process): CEMP-i stars and (*n*, γ) uncertainties



Rapid accretion onto a white dwarf can cause periodic He flashes which produce neutron-capture elements (Denissenkov+17)

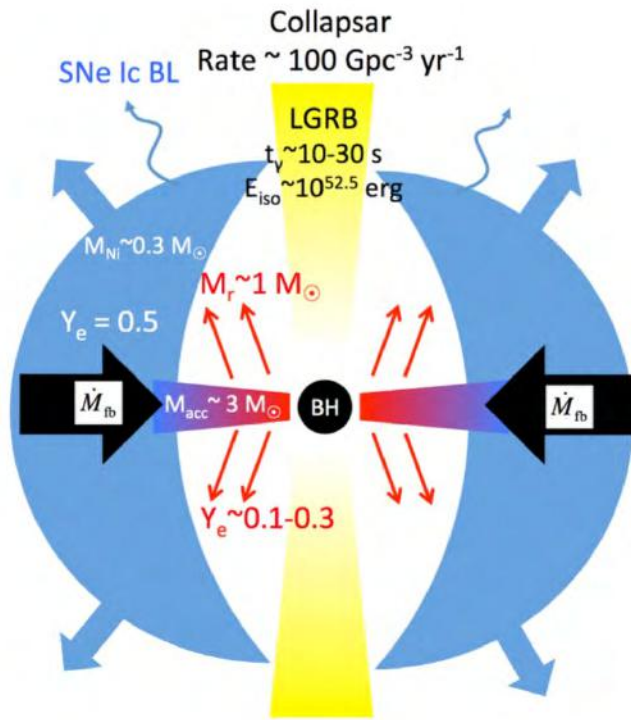
©ESA and Justyn Maund (Queens University Belfast)

(*n*, γ) rate variations informed by variations on theoretical **γ -strength functions** and **nuclear level densities**

(*n*, γ) uncertainties introduce up to \sim an order of magnitude uncertainty in the RAWD contribution to solar system heavy elements

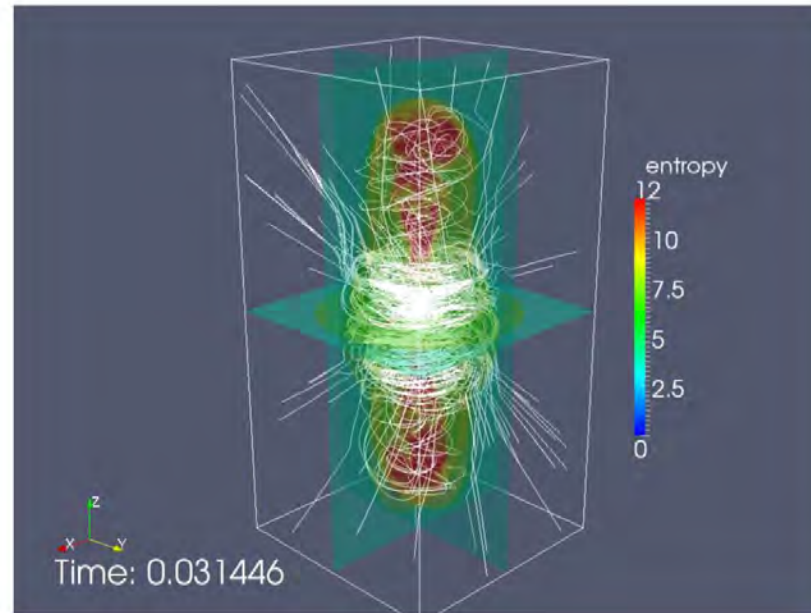
Some candidate sites for r -process element production

Collapsar disk winds



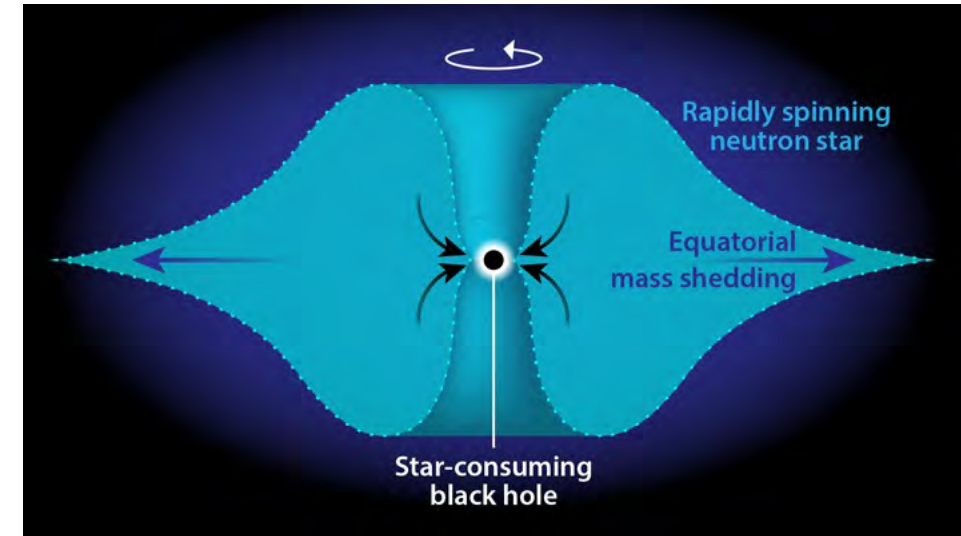
Siegel+18; see also
McLaughlin&Surman 05,
Miller+19

Magneto-rotationally driven (MHD) supernovae



Winteler+12; see also Mosta+17

Primordial black hole + neutron star

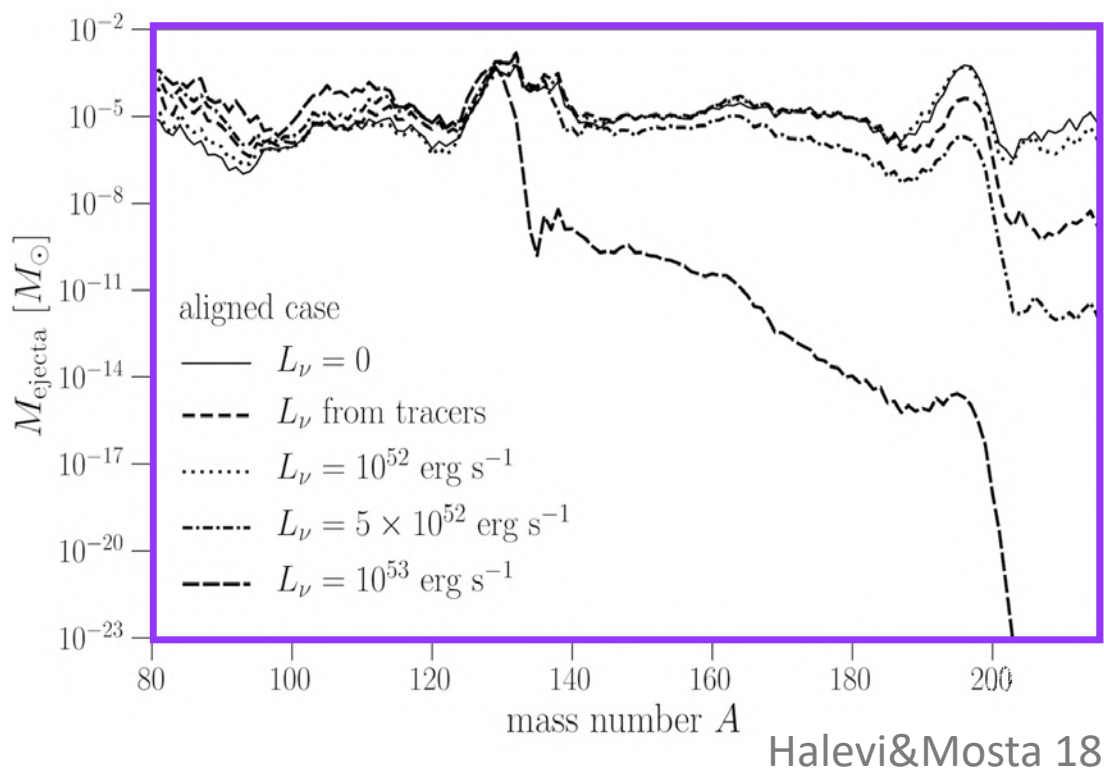


Credit: APS/Alan Stonebraker, via *Physics*

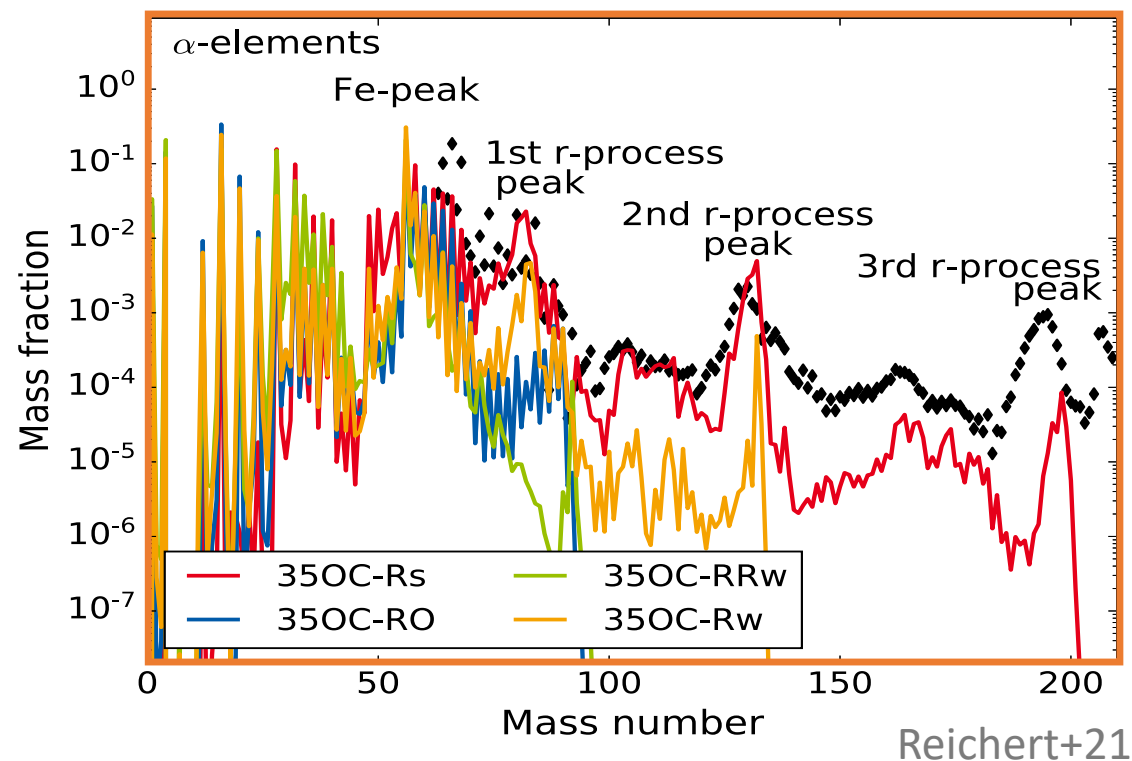
Fuller+17

Spotlight on MHD supernovae

Whether MHDs undergo only a “weak” r process reaching the second peak rather than a “main” or “strong” r process reaching the third peak or beyond depends on the **influence of neutrinos** and the **magnetic field strength**

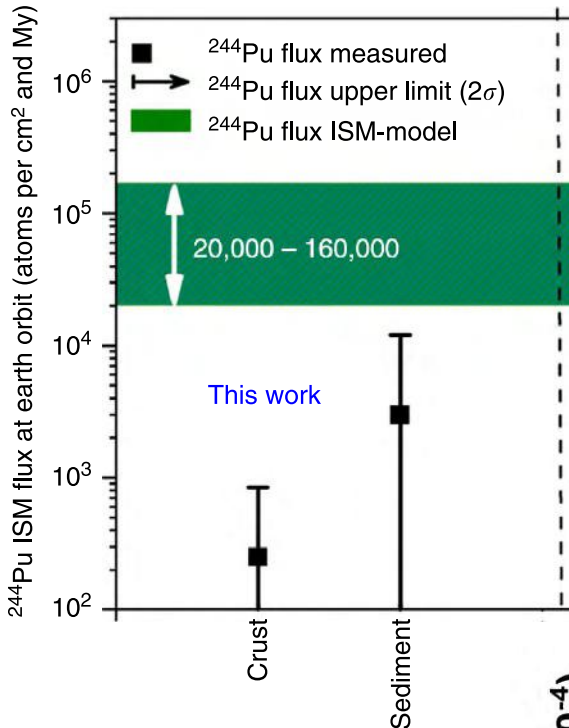


Just like in CCSNe, neutrino energies and luminosities are crucial to determine the r -process reach



Simulations with higher magnetic field strength (ex 350C-Rs $\rightarrow 10^{12} \text{ G}$) undergo a stronger r process than those with lower magnetic field strength (ex 350C-Rw $\rightarrow 10^{10} \text{ G}$)

r-process species must be from a *rare source*

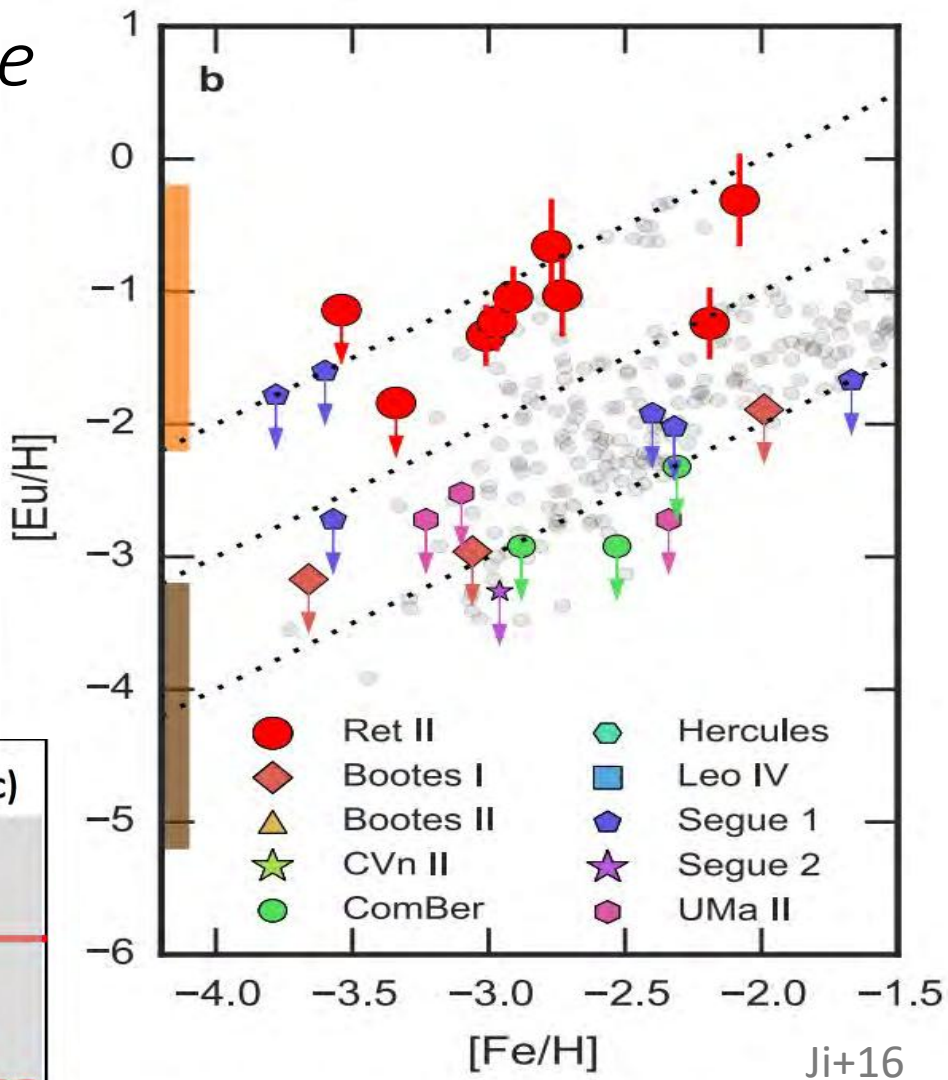
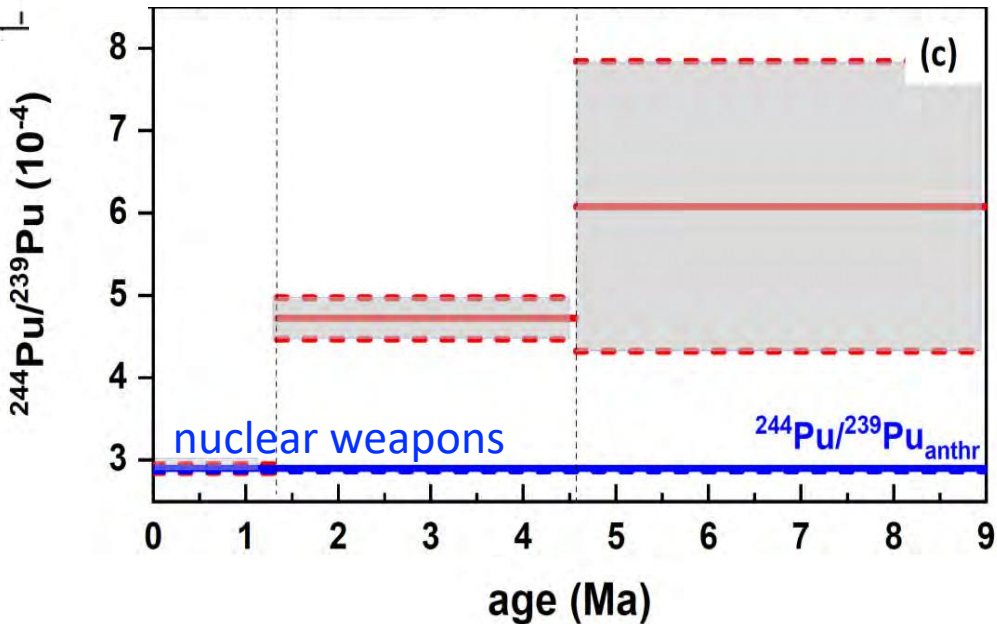


Most recent measurements are still consistent with a rare extraterrestrial source for Pu-244 (long lived compared to Pu-239)

Wallner+21

Pu-244 in deep-sea ocean crusts compared to a *model* which assumes a source as frequent as supernova

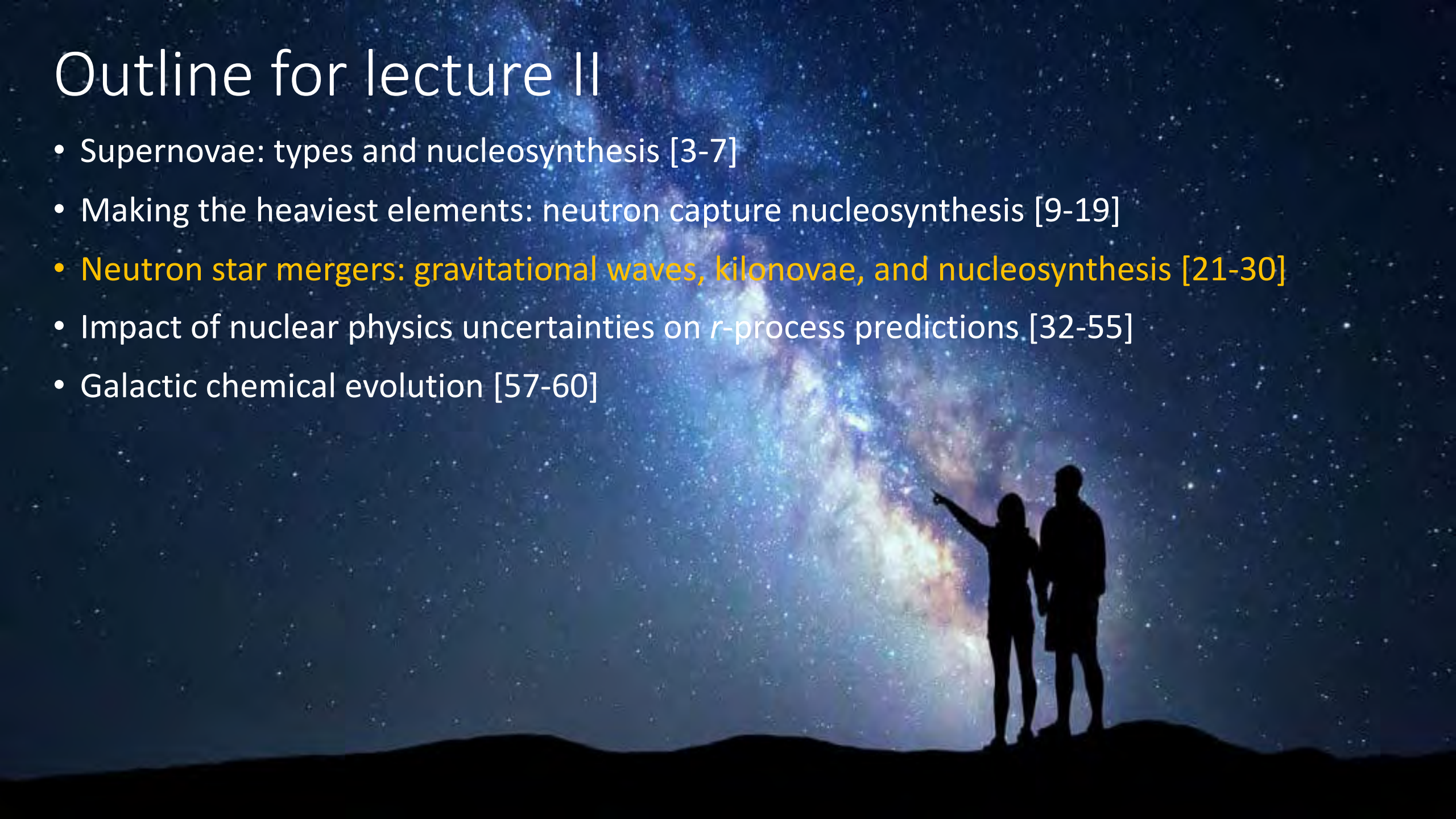
Wallner+15



Ultra faint dwarf galaxies (formed shortly after first stars) rarely show an enhancement in *r*-process elements like in *Reticulum II* (MW in grey)

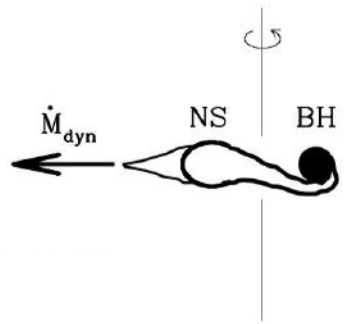
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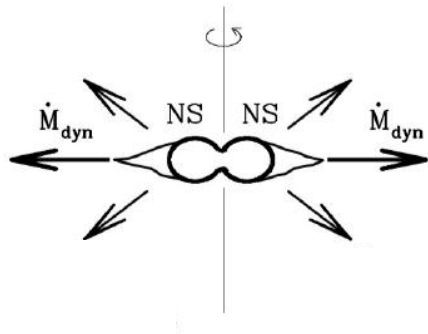


Neutron star mergers and the r process: a bit of history

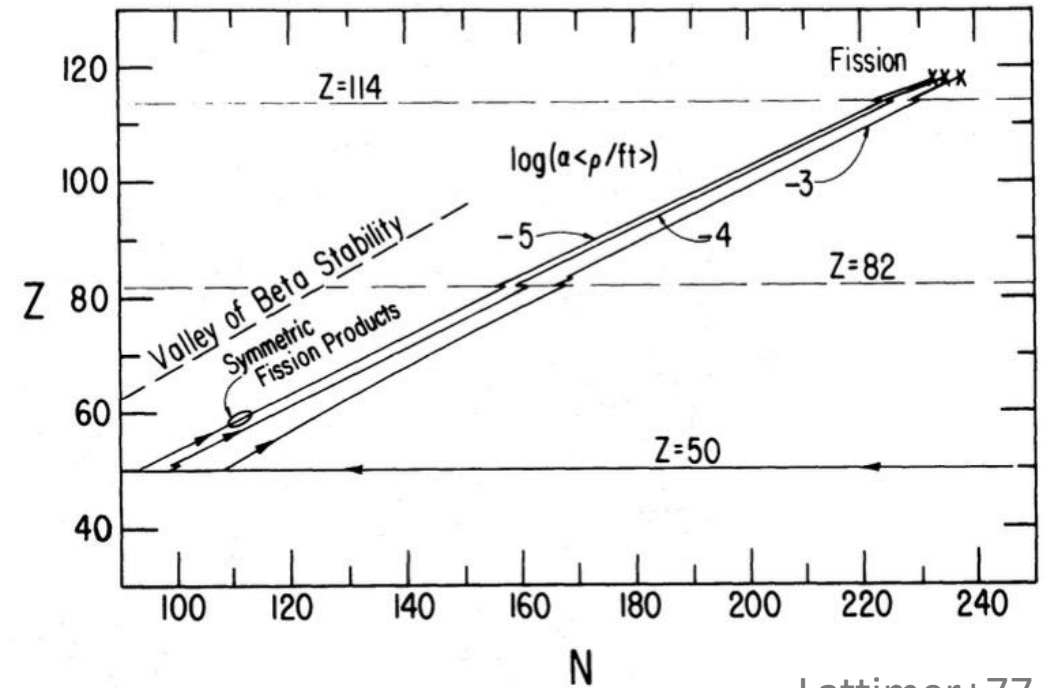
Neutron-rich ejecta from neutron stars > 40 years ago



Lattimer&Schramm (1974): ~5% of the neutron star ejected as n-rich matter



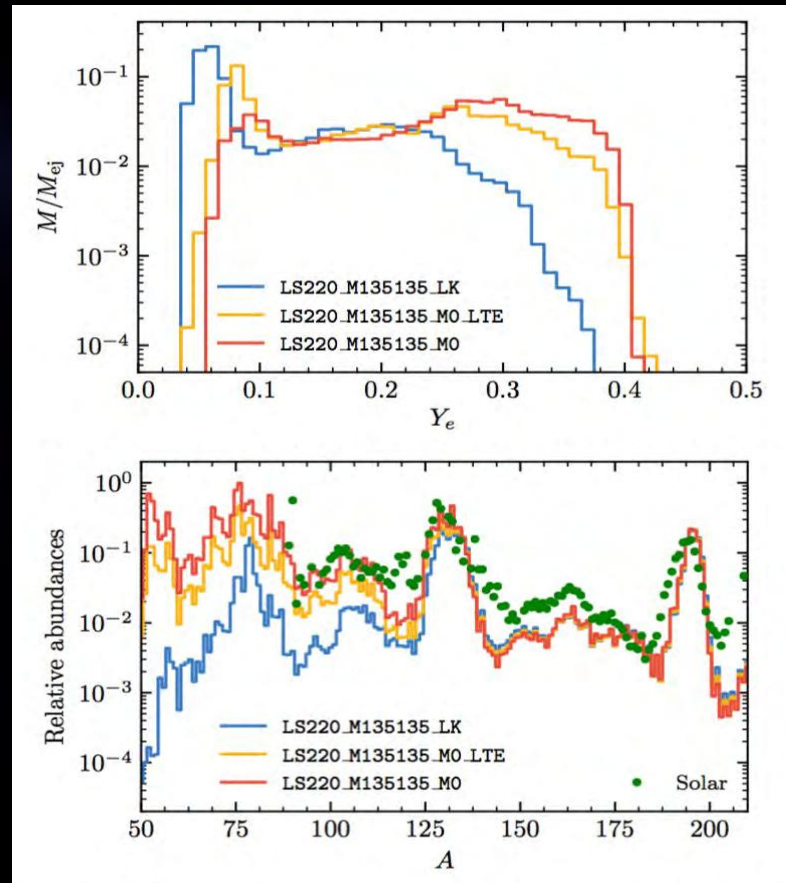
Lattimer+ (1977): initially cold, expanding neutron star matter \rightarrow fission cycling r process capable of super heavy element formation



Lattimer+77

NSM dynamical ejecta

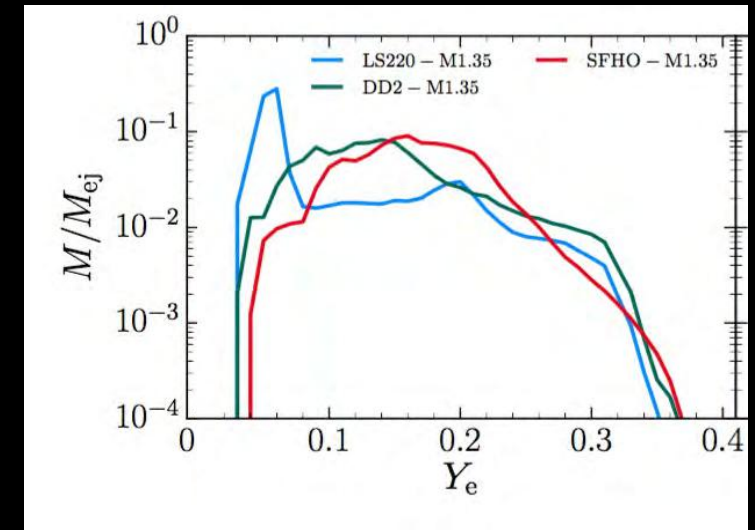
Effect of neutrinos



Rosswog+13

Radice+19; see also Perego+19

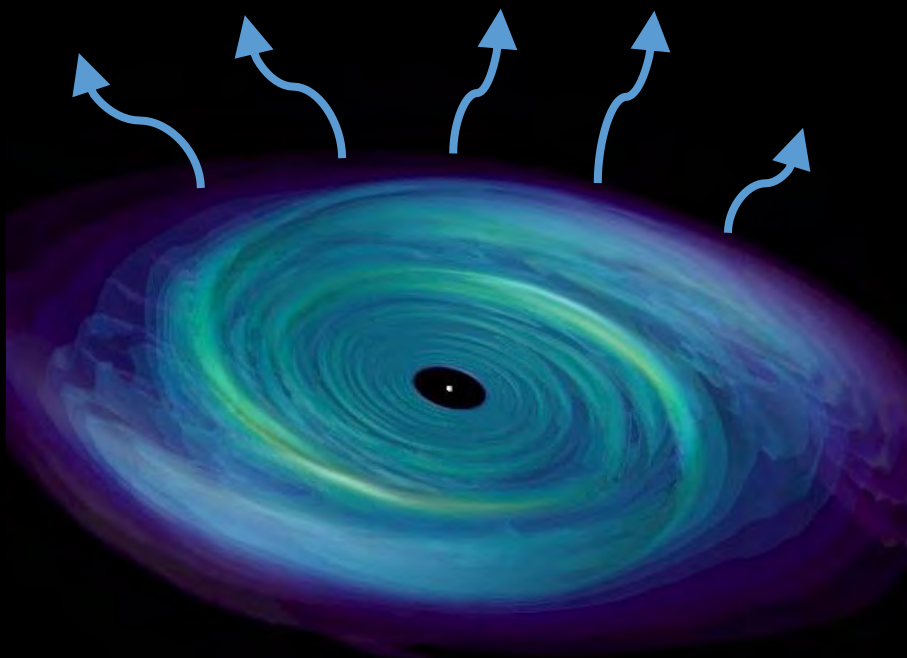
Equation of state



Bovard+17

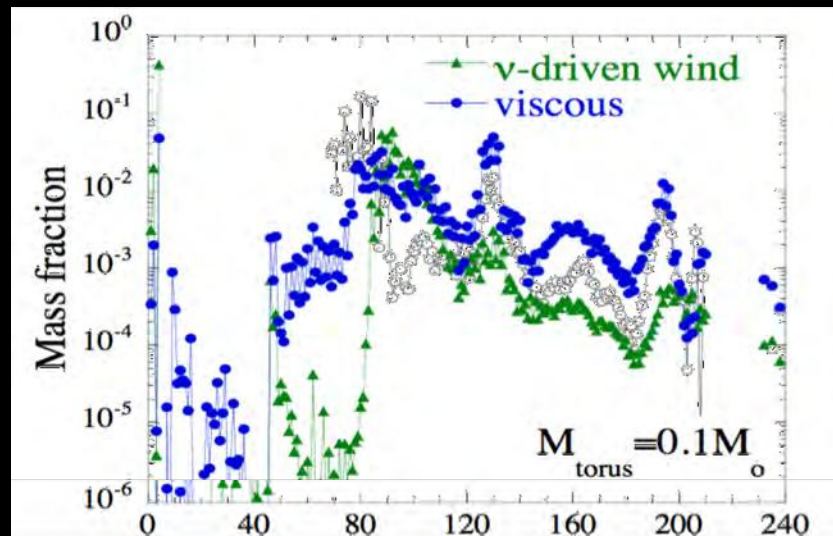
See also Wanajo+14,
Vincent+19, Foucart+20....

Post-merger disk ejecta

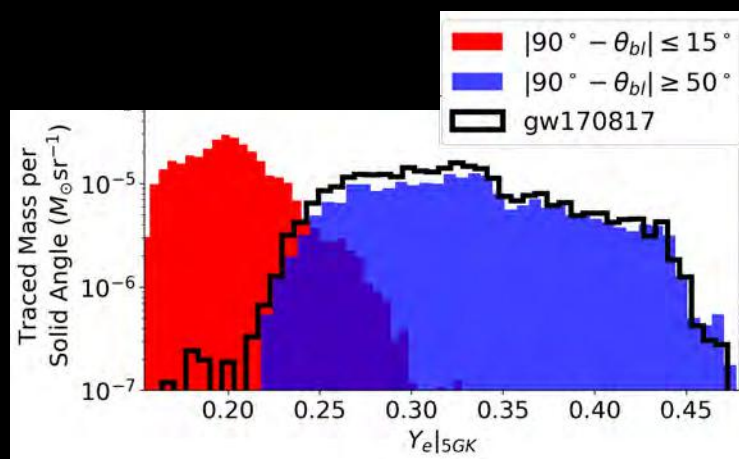


Owen&Blondin 05

Neutrino driven vs viscous

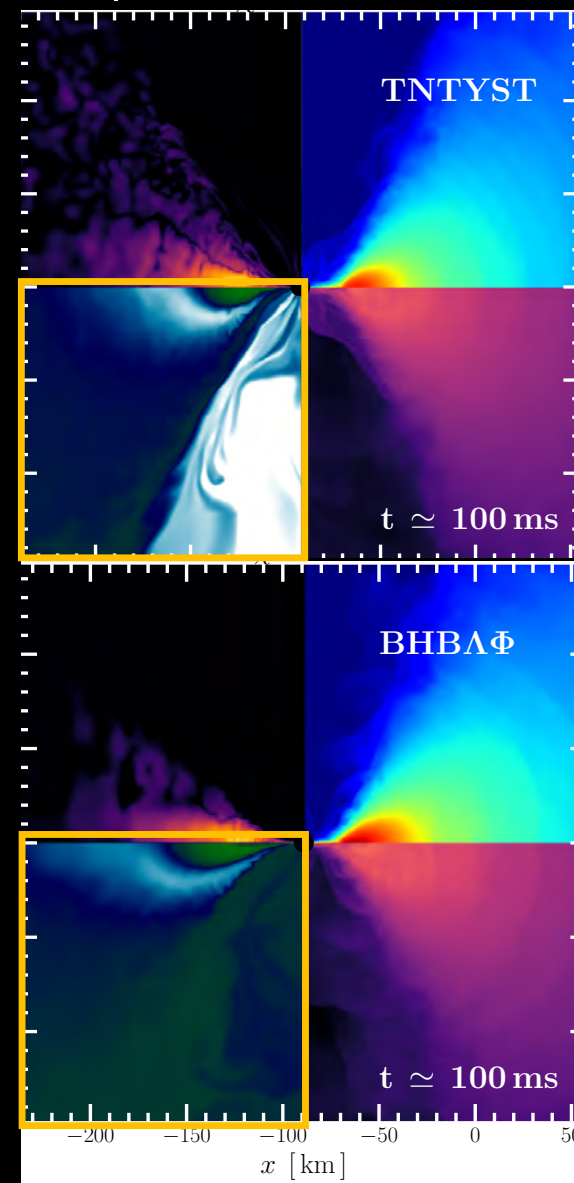


Just+16



Miller+19

Equation of state

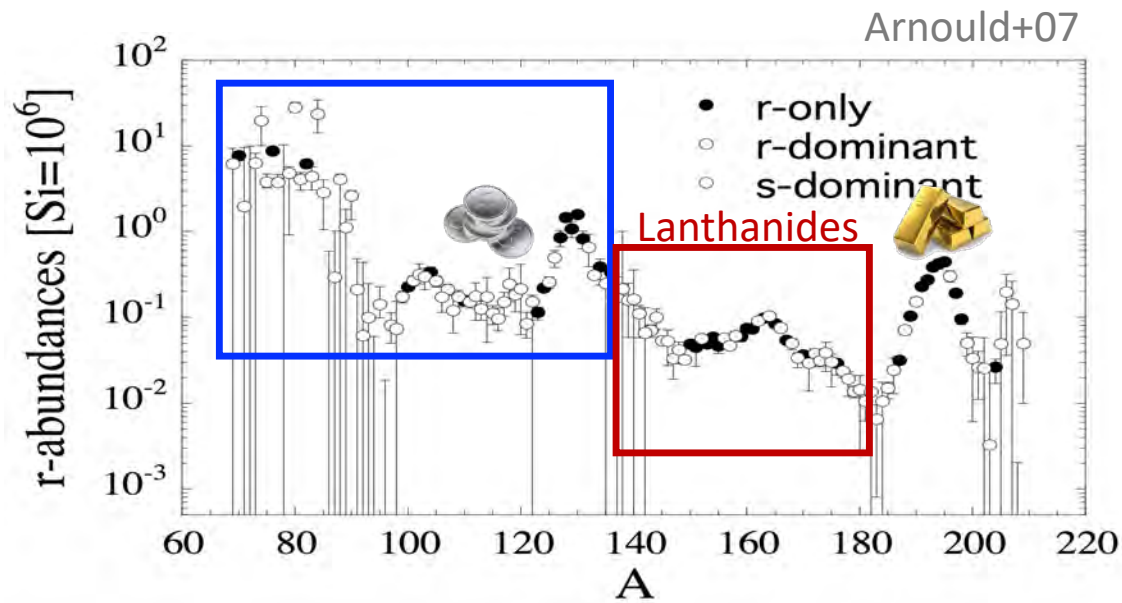


Most+21

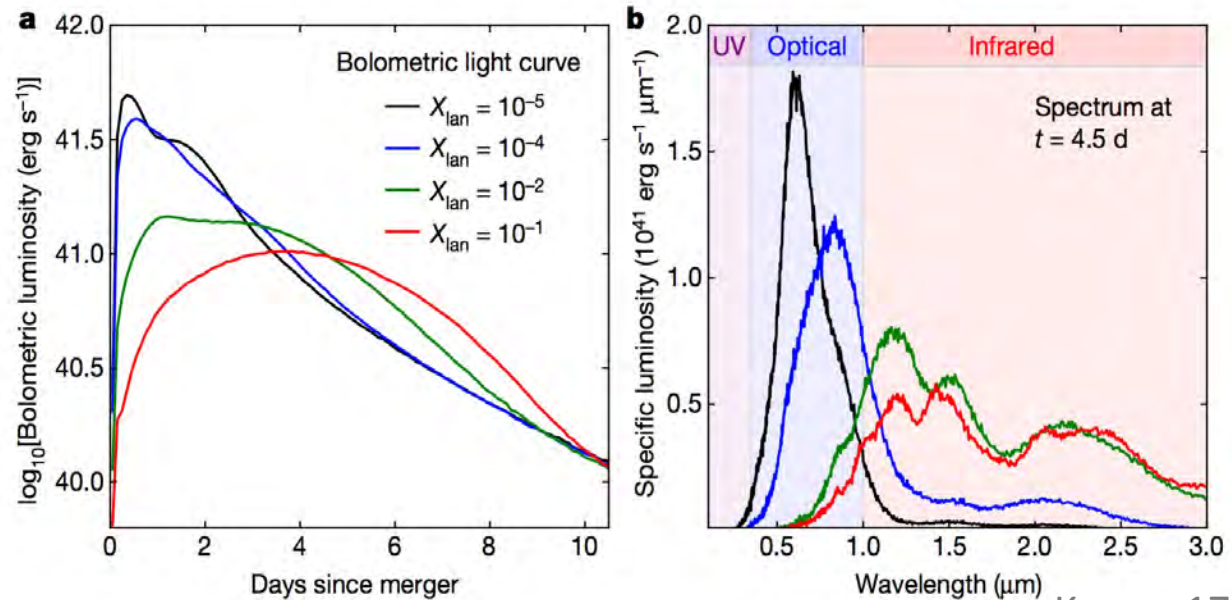
GW170817 & AT2017gfo: “red” and “blue” kilonovae

Spectra and light curves depend on the species present;
Lanthanide and/or **actinide** mass fraction \uparrow , opacity \uparrow ,
longer duration light curve shifted toward infrared

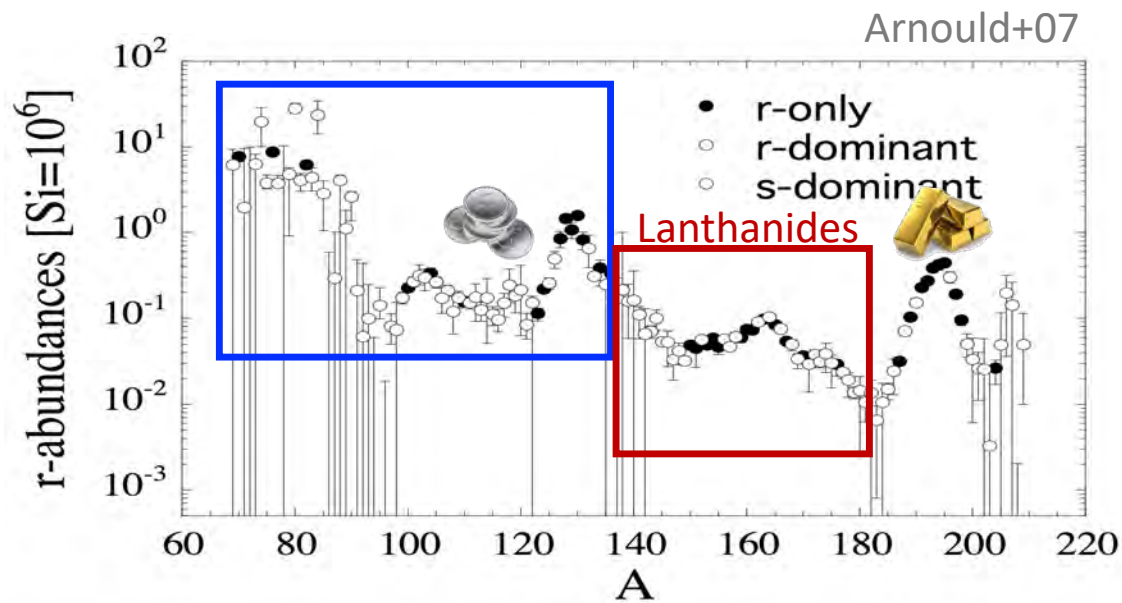
(e.g. Metzger+10, Lippuner+15, Barnes+16,21, Wanajo+18,
Watson+19, Hotokezaka+20, Korobkin+20, Zhu+18,21, Wang+20)



Model



Kasen+17

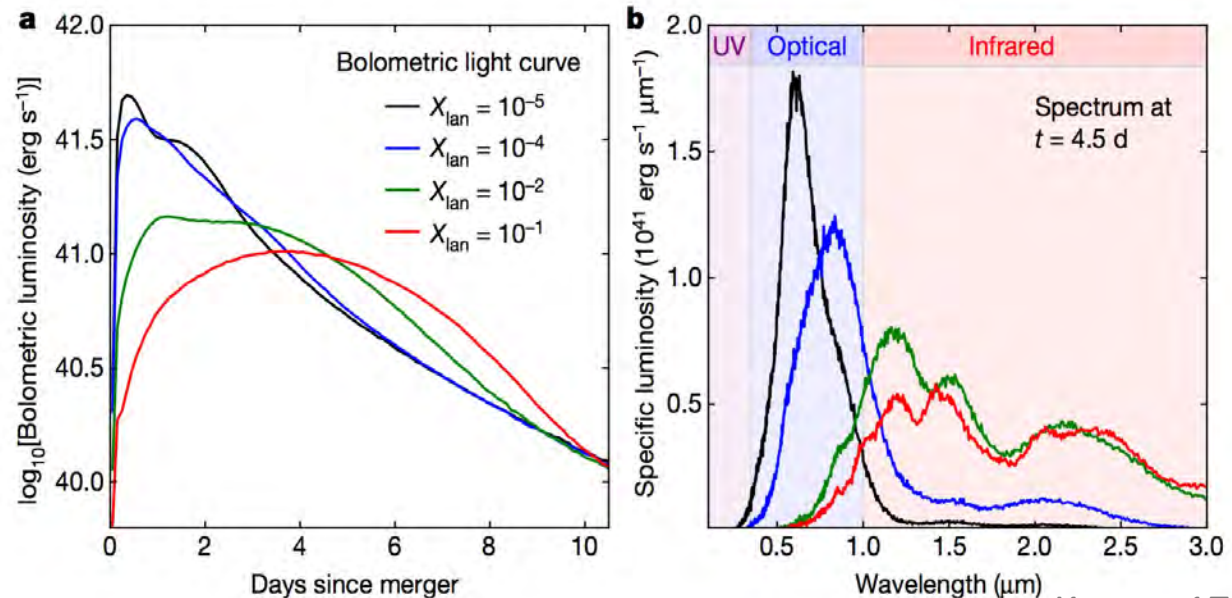


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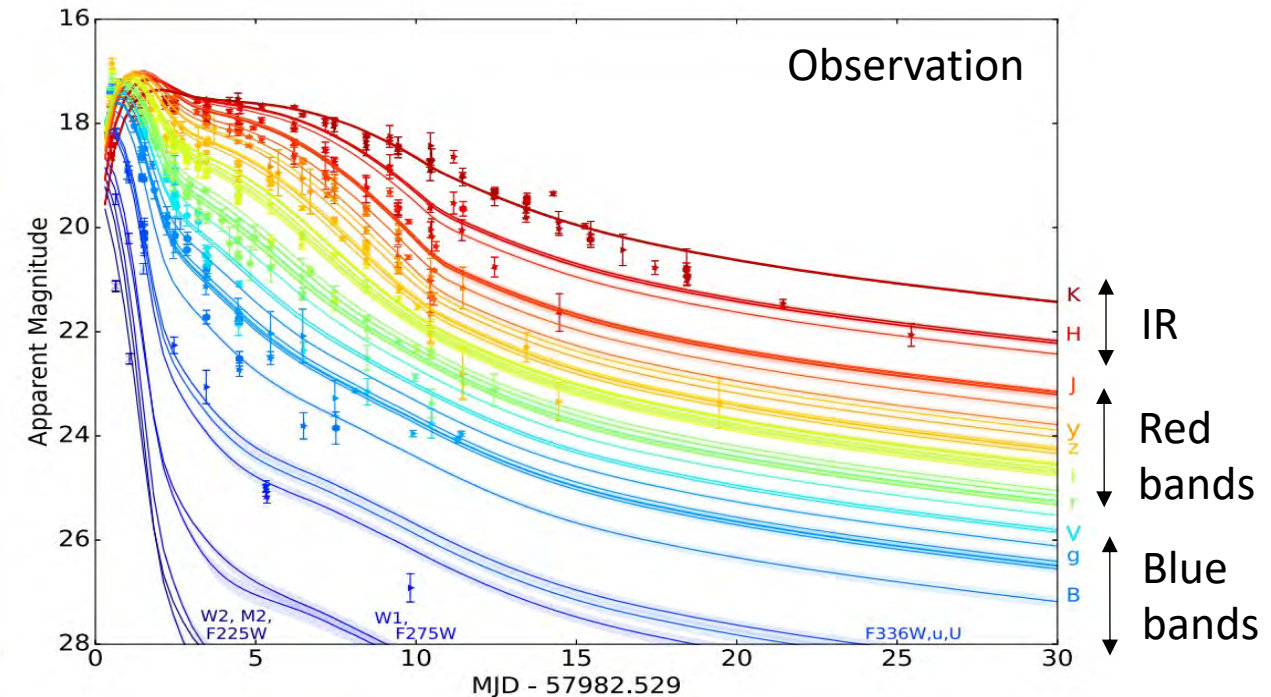
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Model



Kasen+17



Villar+17; see also
Cowperthwaite+17

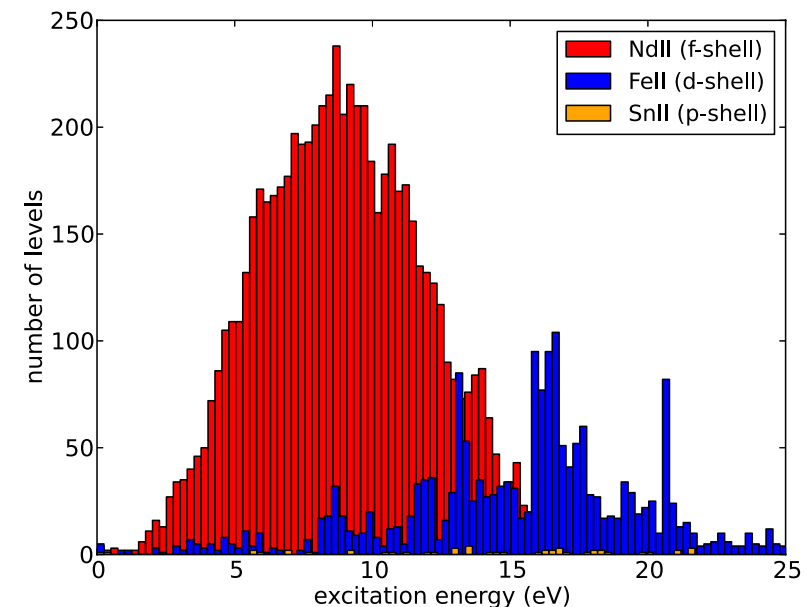
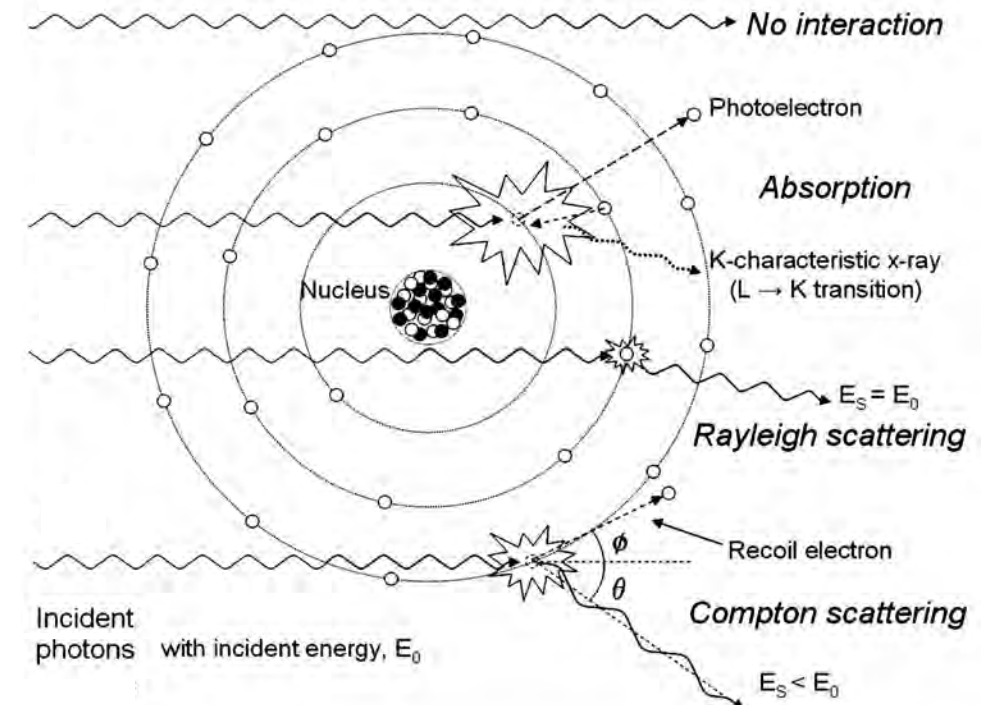
GW170817 & AT2017gfo: photon opacity

Opacity sources include (*most important in NSM ejecta):

- **bound-bound transitions*** – photoelectric absorption: photon absorbed or emitted as an electron moves between levels
- **bound-free** – photoionization: electron absorbs photon and escapes
- **free-free scattering** – bremsstrahlung: free electron passing close to ion or nucleus can emit or absorb a photon
- **electron scattering** – inelastic (Compton) scattering and elastic (Rayleigh) scattering: photons scatter off electrons

1																	2				
H																	He				
3	4															5	6	7	8	9	10
Li	Be															B	C	N	O	F	Ne
11	12															13	14	15	16	17	18
Na	Mg															Al	Si	P	S	Cl	Ar
19	20	21	22	23	24	25	26	27	28	29	30	31	32	33	34	35	36				
K	Ca	Sc	Ti	V	Cr	Mn	Fe	Co	Ni	Cu	Zn	Ga	Ge	As	Se	Br	Kr				
37	38	39	40	41	42	43	44	45	46	47	48	49	50	51	52	53	54				
Rb	Sr	Y	Zr	Nb	Mo	Tc	Ru	Rh	Pd	Ag	Cd	In	Sn	Sb	Te	I	Xe				
55	56	57-71	72	73	74	75	76	77	78	79	80	81	82	83	84	85	86				
Cs	Ba		Hf	Ta	W	Re	Os	Ir	Pt	Au	Hg	Tl	Pb	Bi	Po	At	Rn				
87	88	89-103	104	105	106	107	108	109	110	111	112	113	114	115	116	117	118				
Fr	Ra		Rf	Db	Sg	Bh	Hs	Mt	Ds	Rg	Cn	Nh	Fl	Mc	Lv	Ts	Og				

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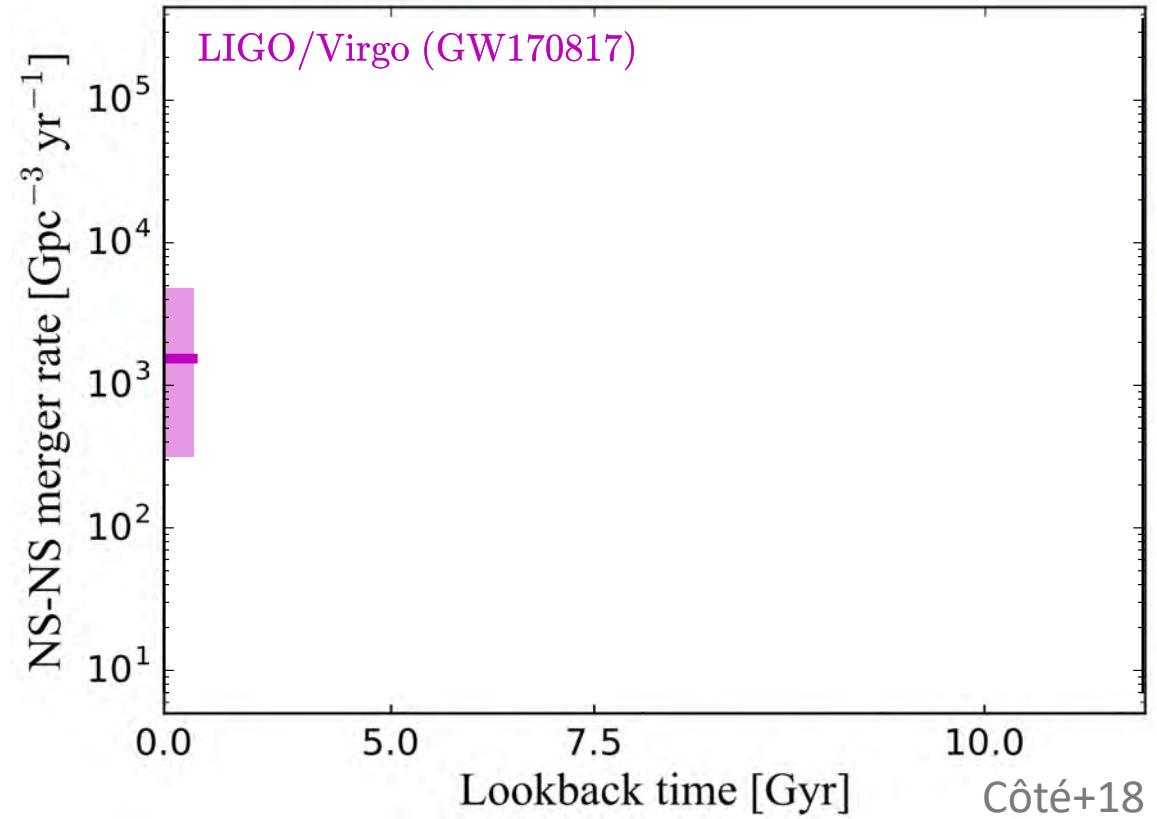
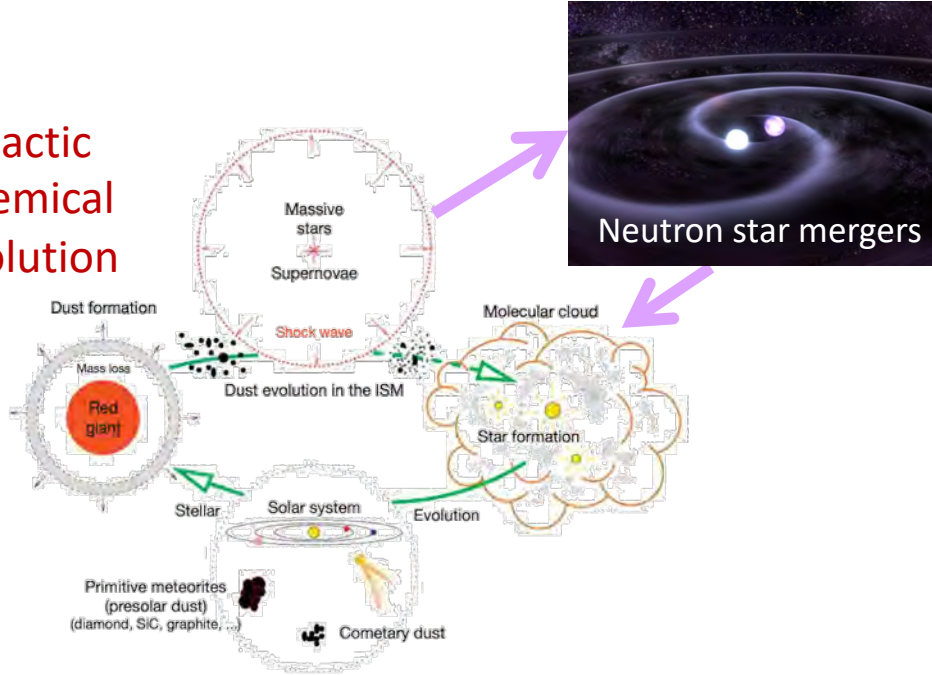


Kasen+13;
see also
Fontes+20,
Tanaka+20

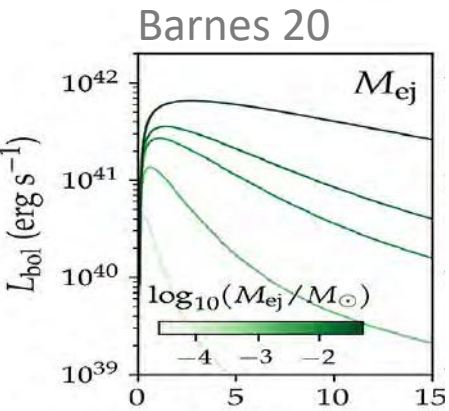
Do binary NSMs make enough heavy elements?

Galactic
Chemical
Evolution

Palm+14



Do binary NSMs make enough heavy elements?

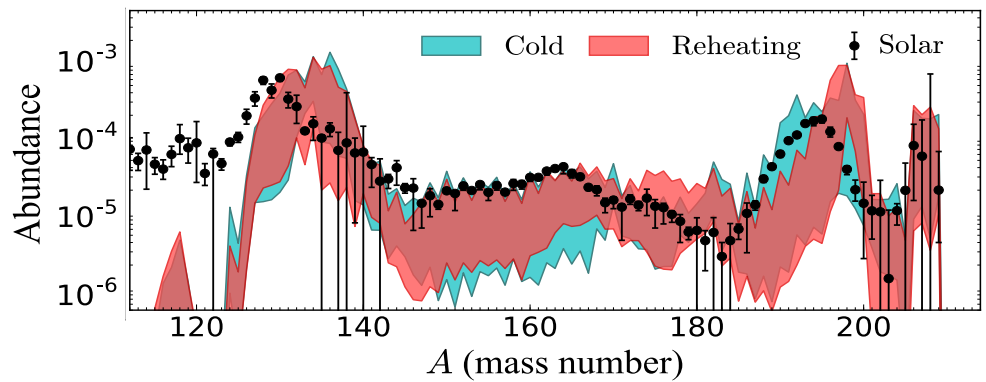
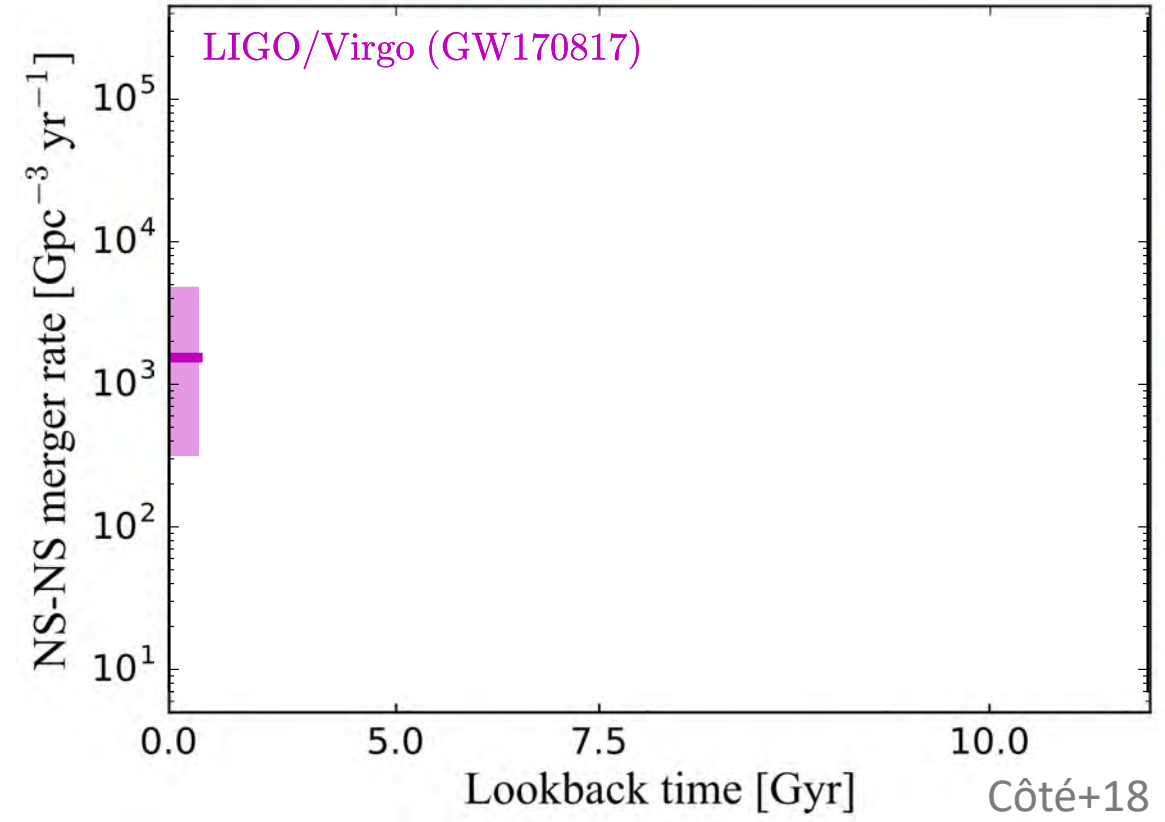
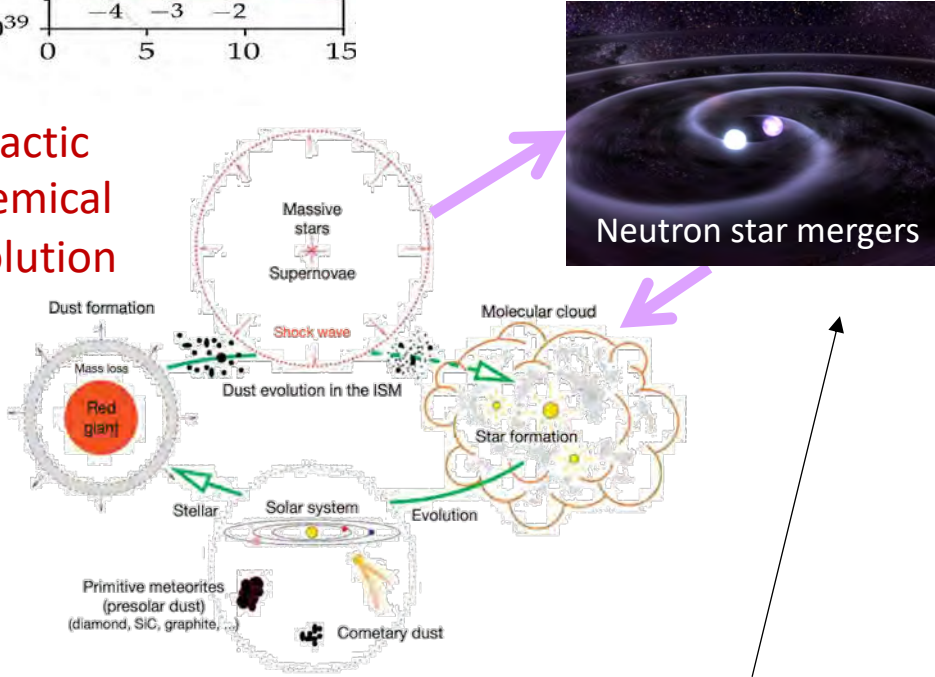


Light Curves

Take estimates for GW170817 mass ejection range from literature

Galactic Chemical Evolution

Palm+14

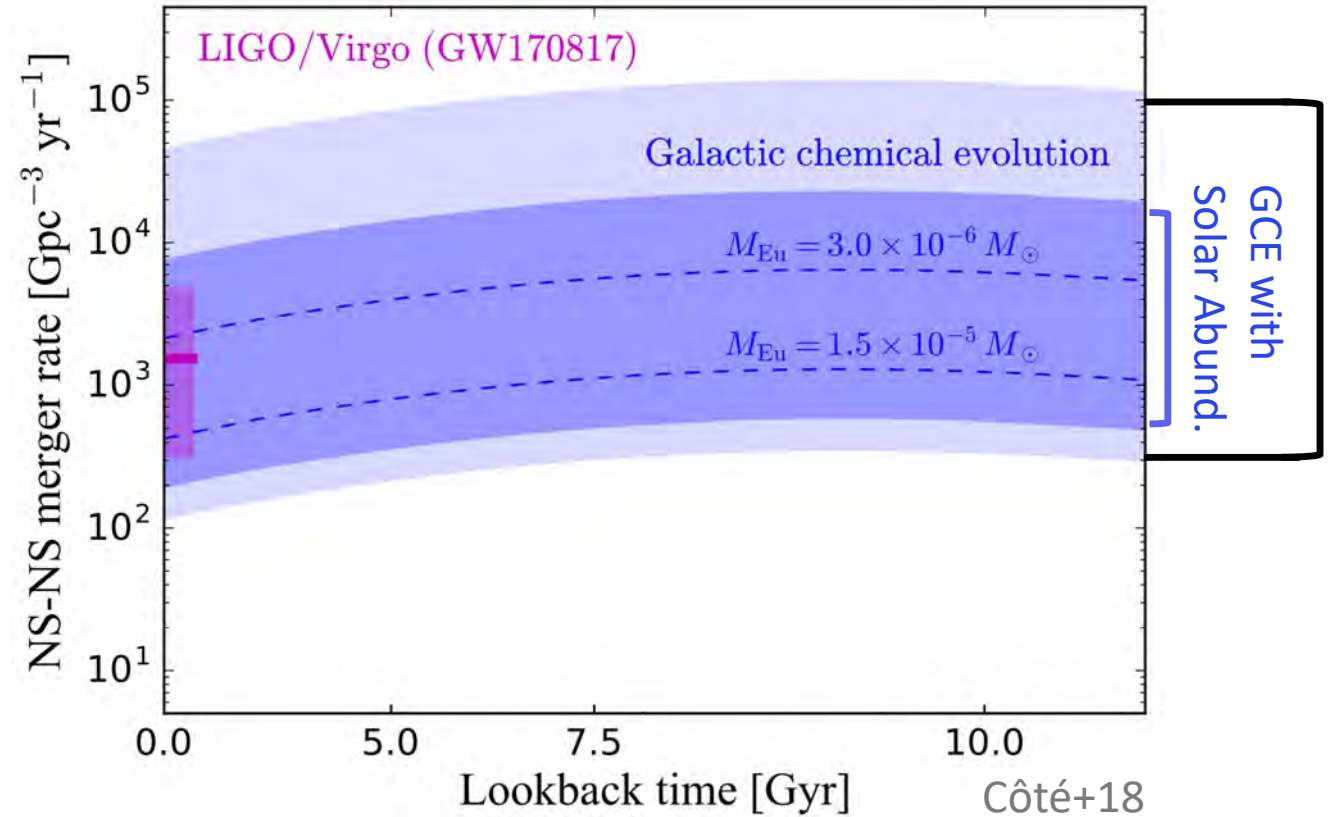


Nucleosynthesis Predictions

Abundance range of dynamical ejecta from 10 mass models

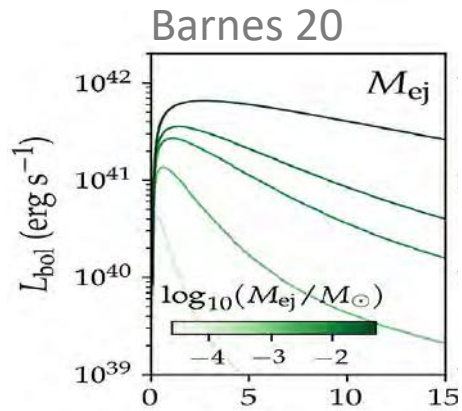
Do binary NSMs make enough heavy elements?

nuclear physics uncertainties



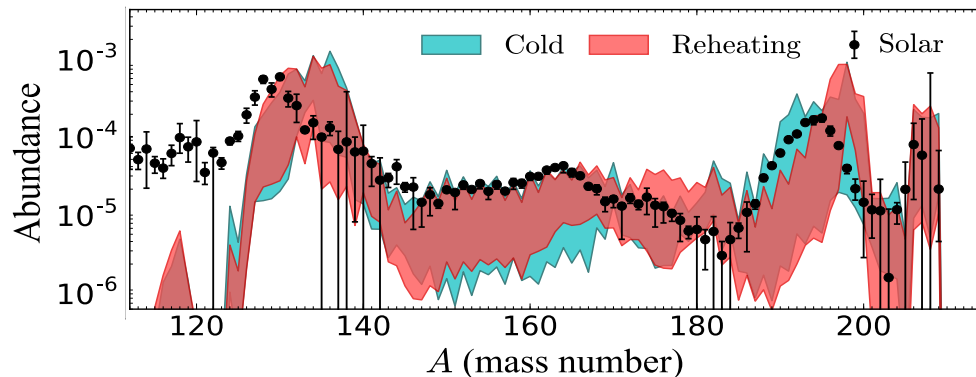
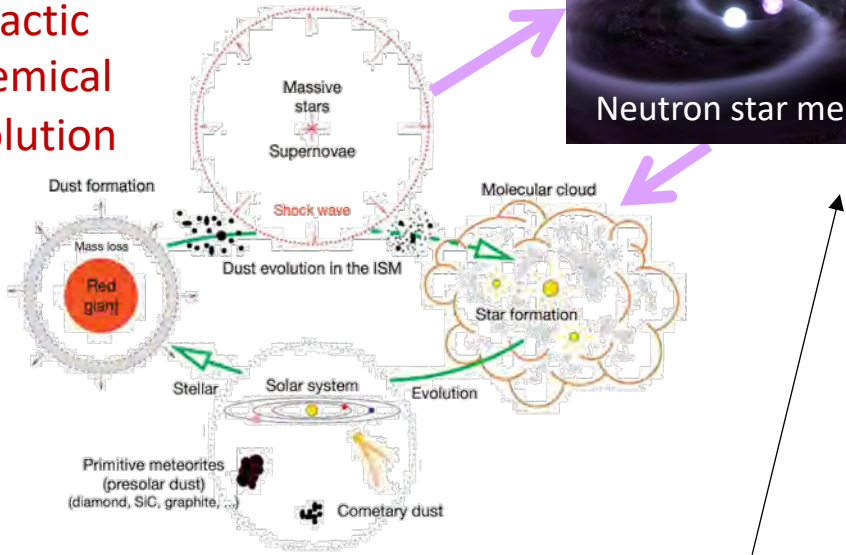
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Galactic Chemical Evolution

Palm+14

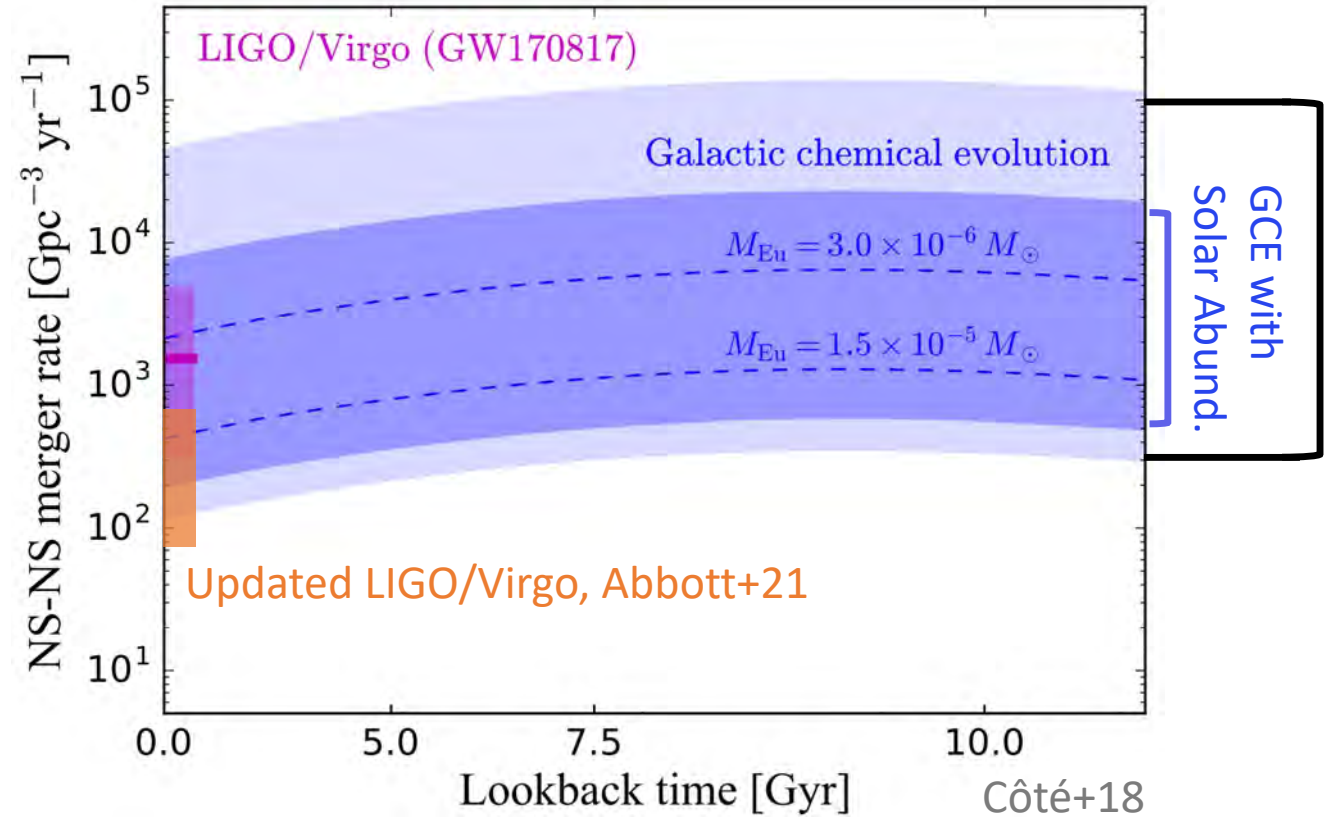


Nucleosynthesis Predictions

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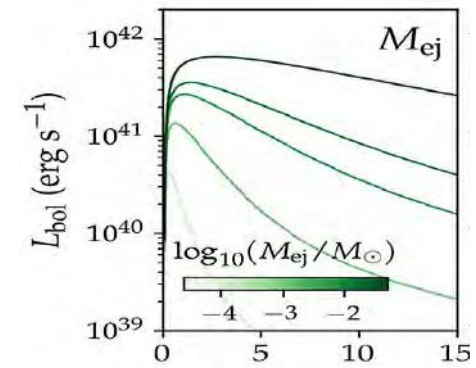
Light Curves

Take estimates for GW170817 mass ejection range from literature



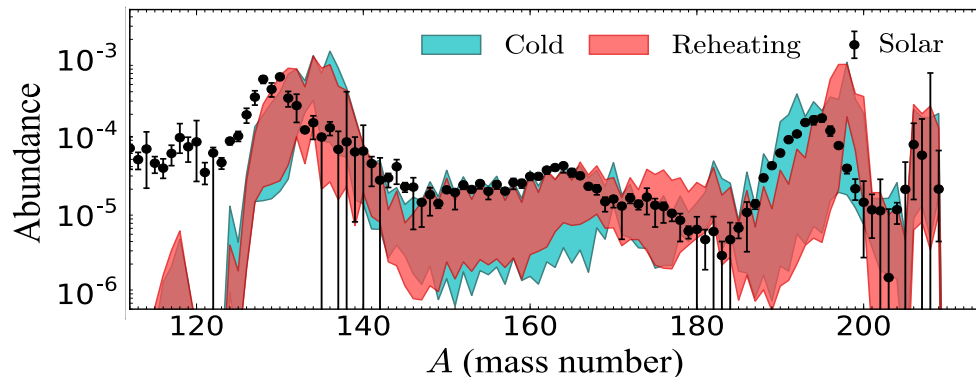
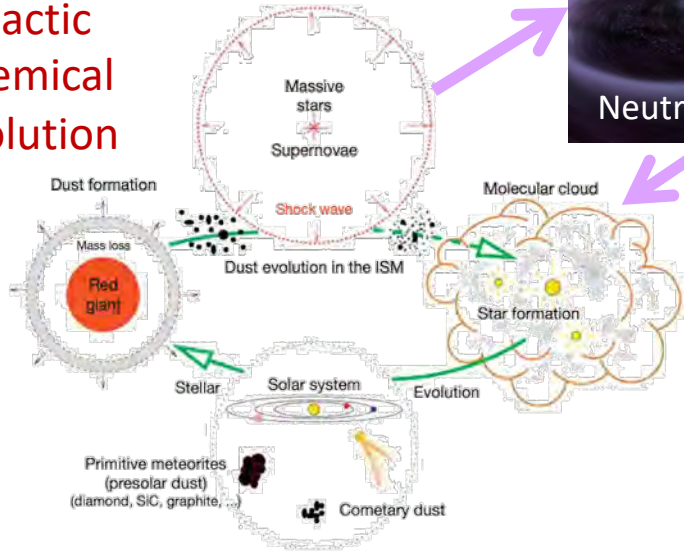
Neutron star mergers

Barnes 20



Galactic Chemical Evolution

Palm+14



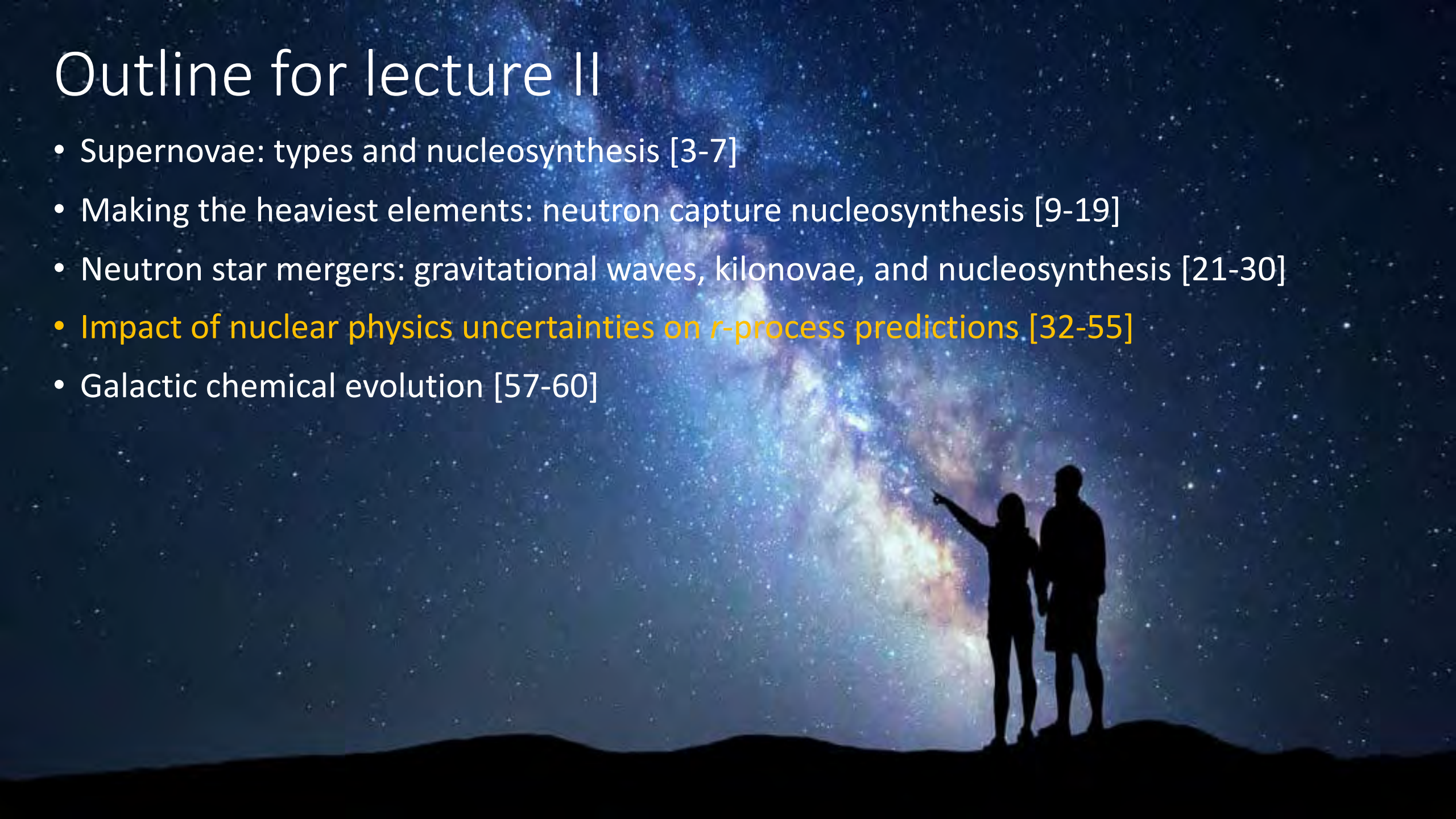
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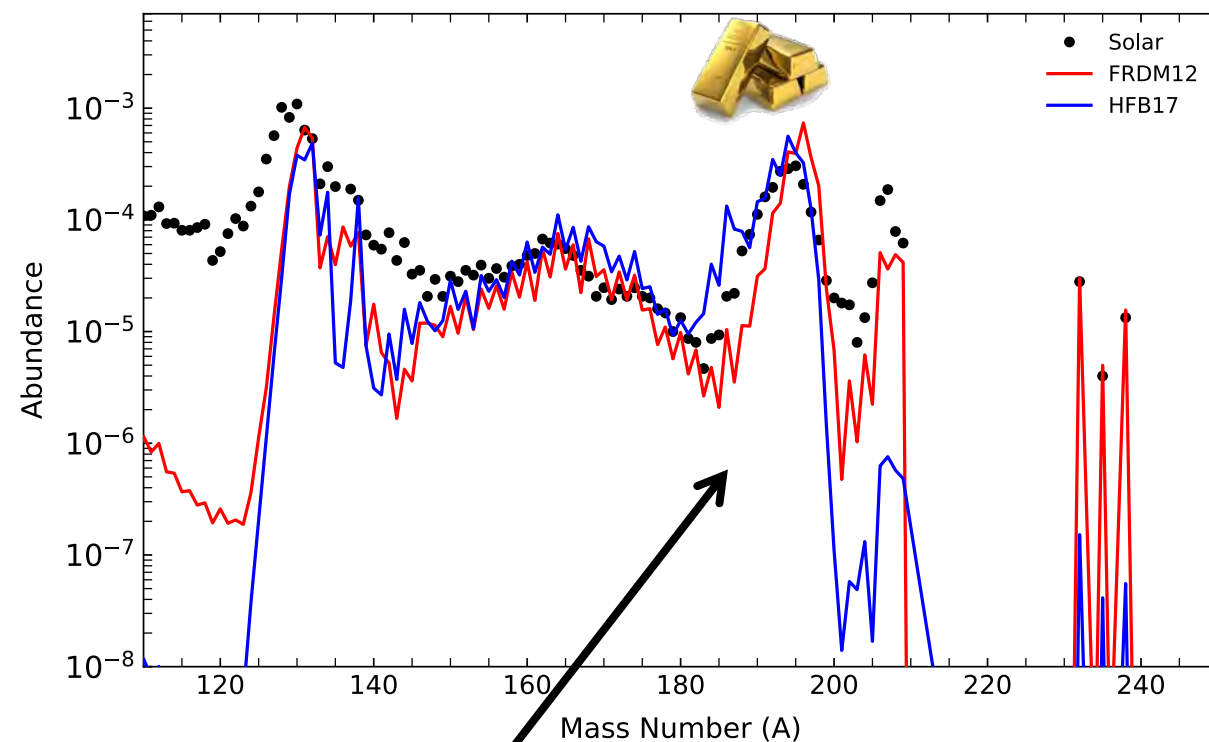
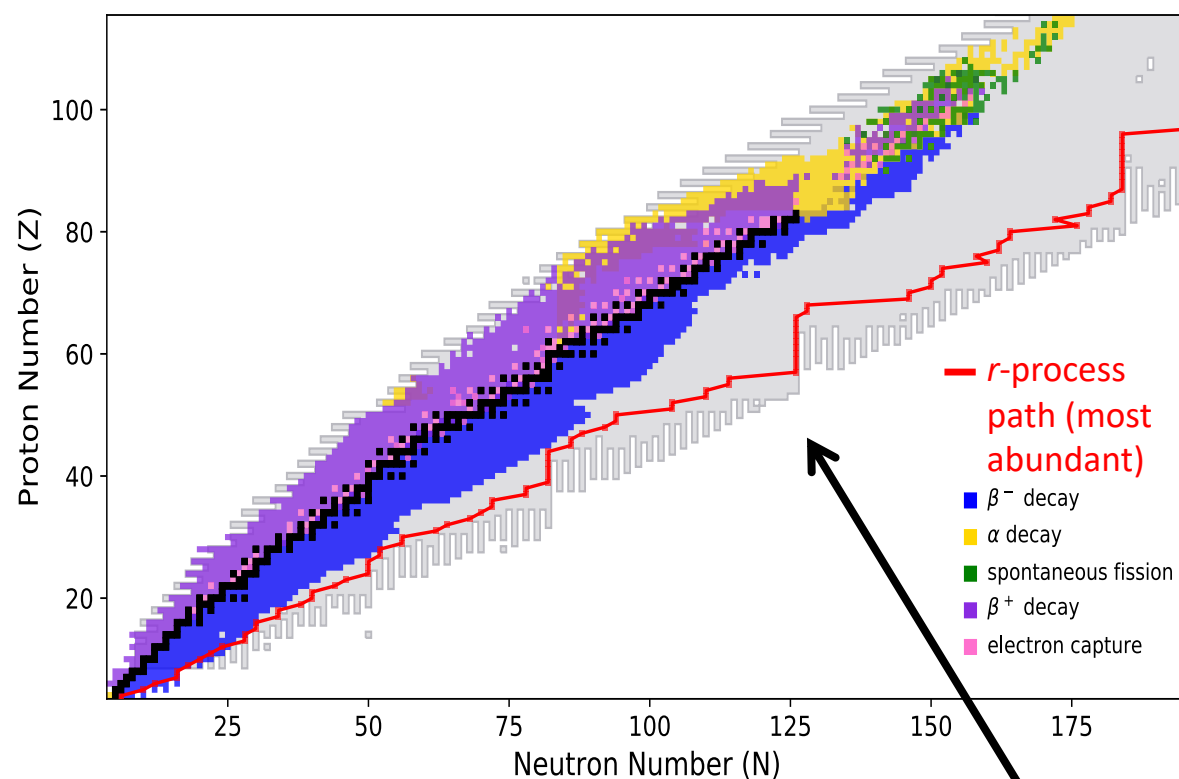
*Now another confirmed NSNS merger GW190425 as well as a June 2021 confirmation of two NSBH mergers GW200105 and 200115!

Outline for lecture II

- Supernovae: types and nucleosynthesis [3-7]
- Making the heaviest elements: neutron capture nucleosynthesis [9-19]
- Neutron star mergers: gravitational waves, kilonovae, and nucleosynthesis [21-30]
- **Impact of nuclear physics uncertainties on *r*-process predictions [32-55]**
- Galactic chemical evolution [57-60]

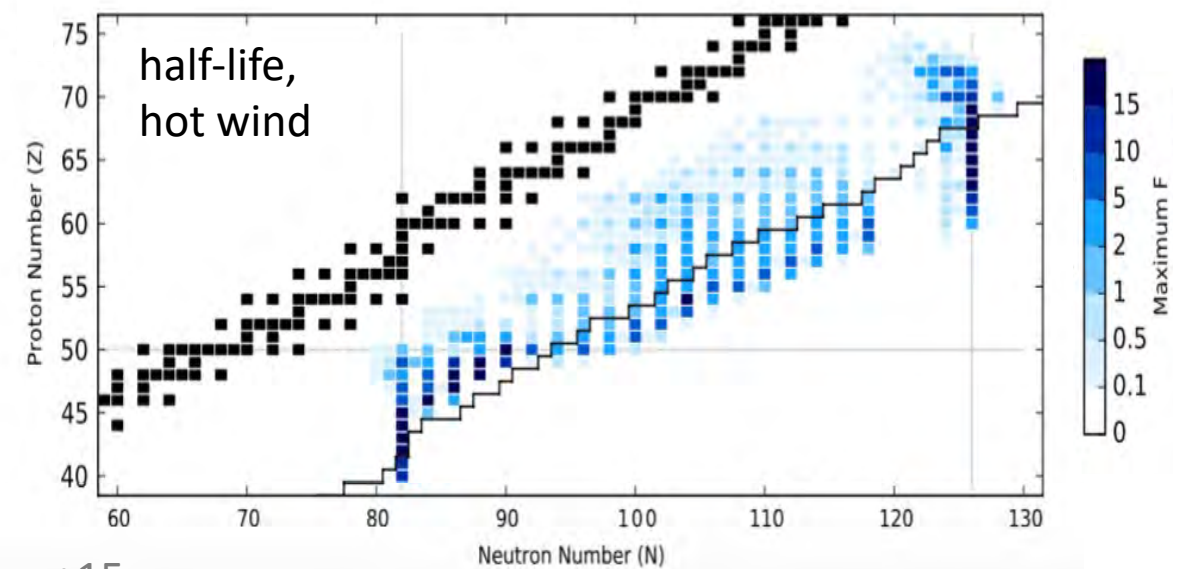
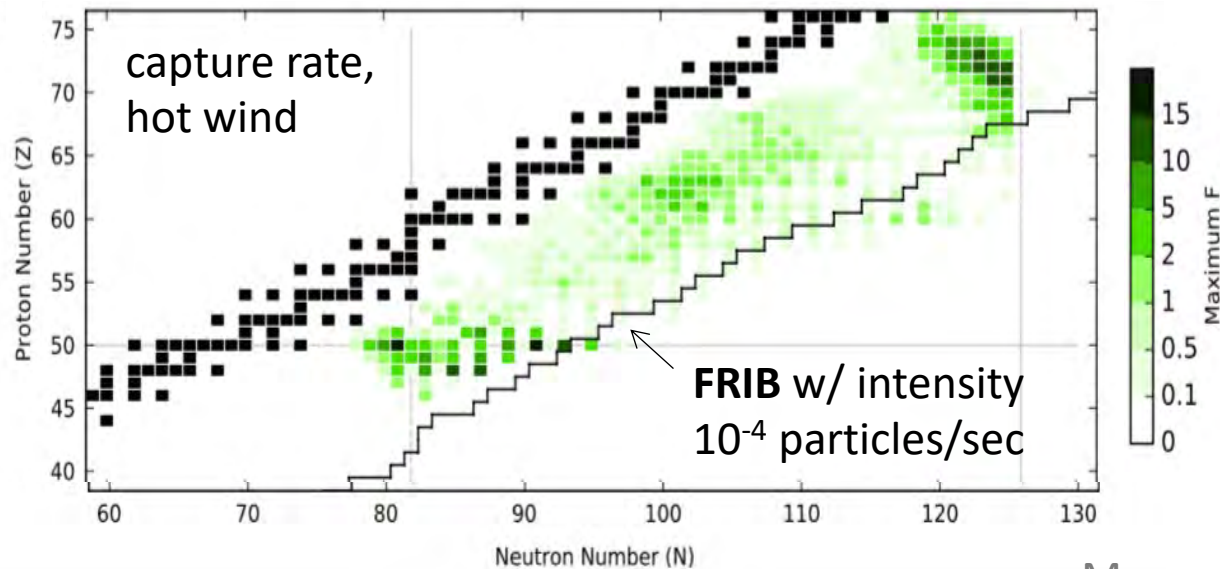
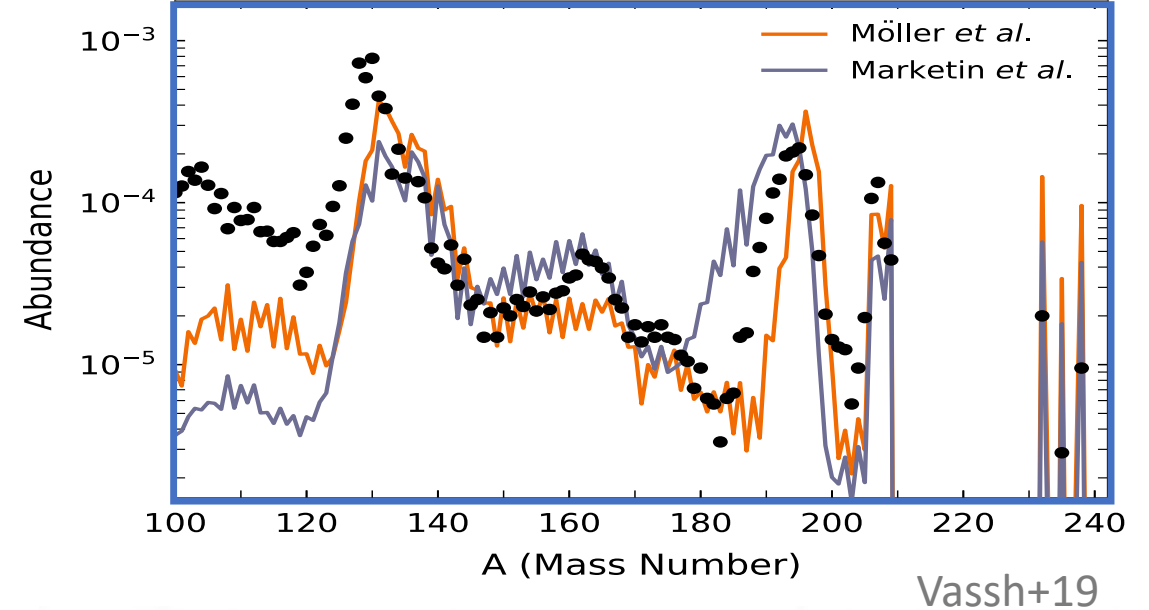
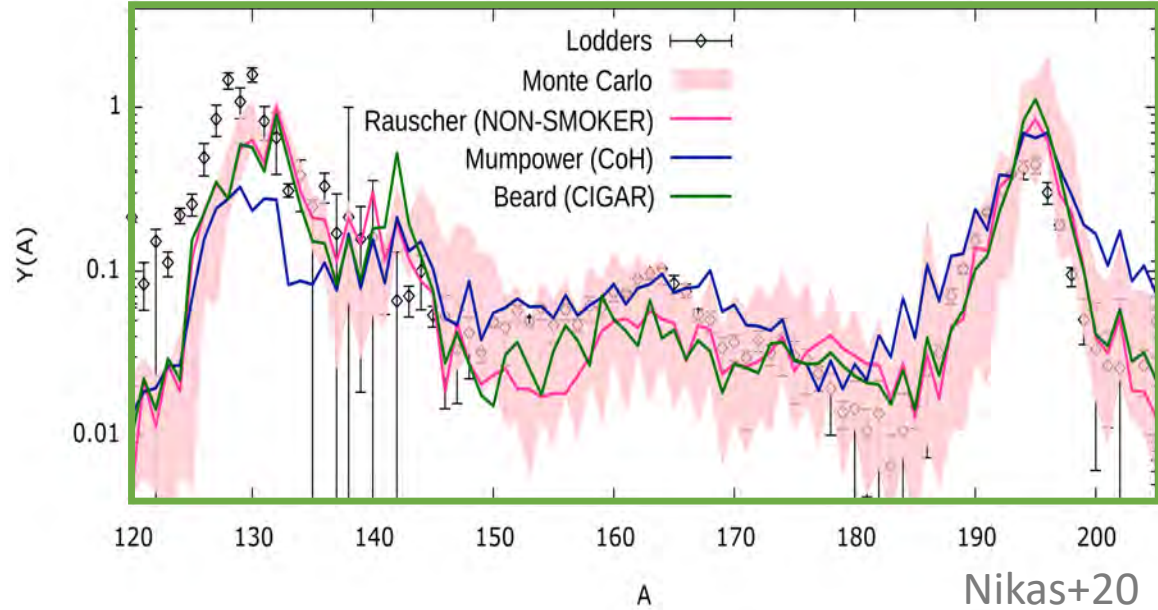


Impact of nuclear physics uncertainties: *r*-process N=126 peak example



- Little to no experimental data on the neutron-rich side at N=126; nuclear mass models predict different shell closure strengths and thus different amounts of elements like gold and platinum
- The N=126 shell closure is the “gateway” to the actinides and thus affects how strongly elements like uranium-238 are produced

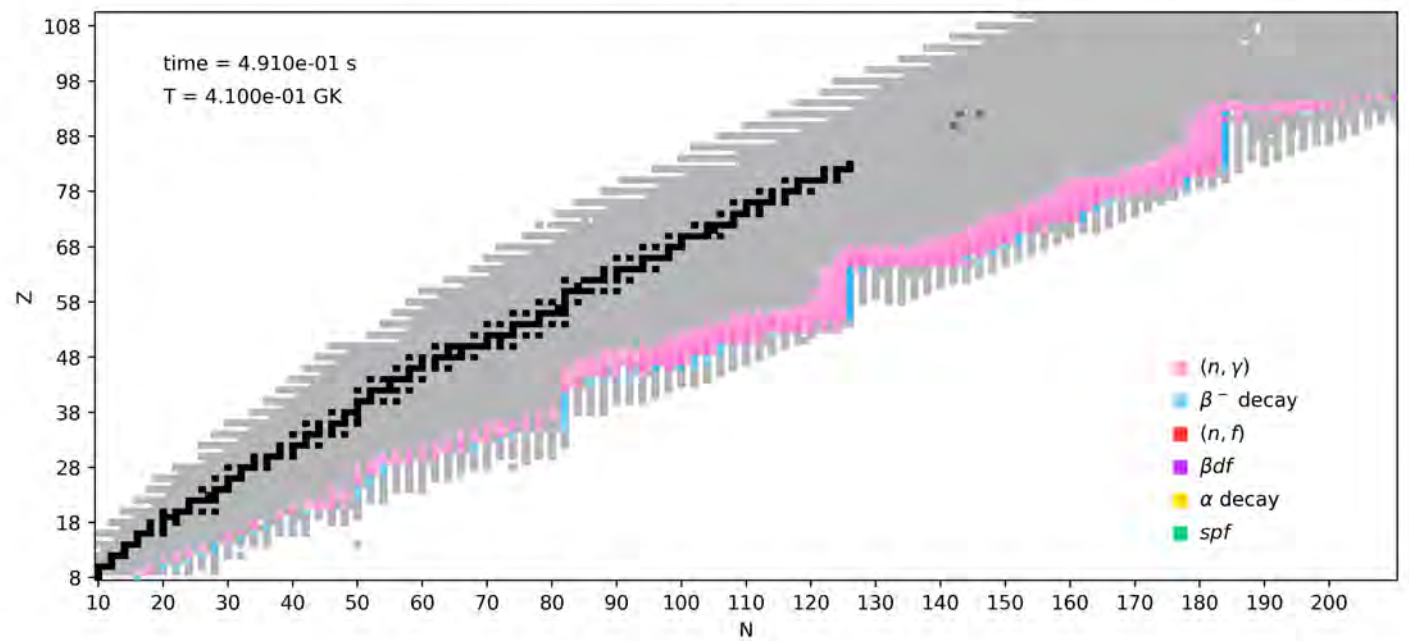
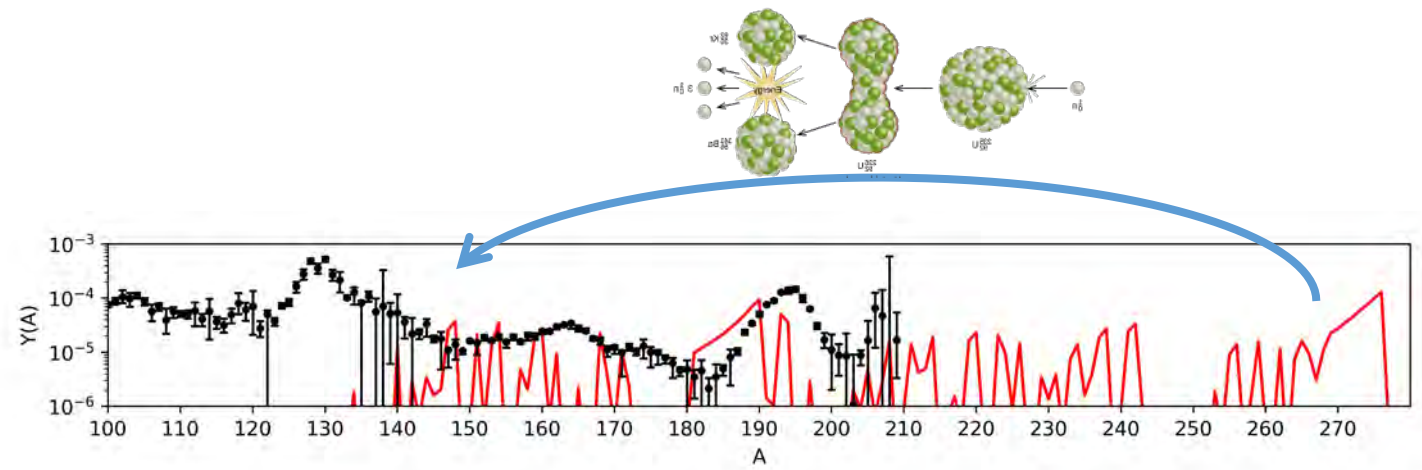
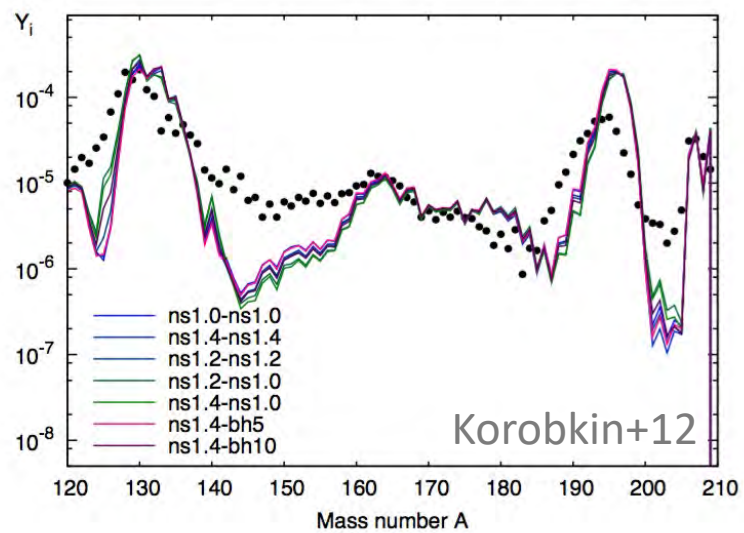
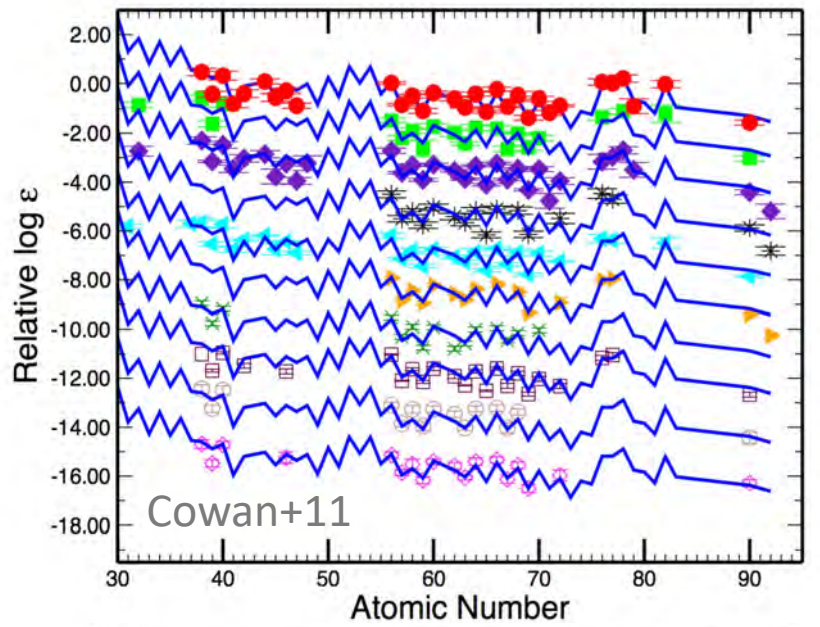
Sensitivity of r -process abundances to neutron capture and β -decay



Mumpower+15

Fission cycling to explain observed robustness of lanthanide abundances?

10 *r*-process rich halo stars compared to Solar



NSM dynamical ejecta using Rosswog+13 simulation conditions
(very neutron-rich with robust fission)

Nuclear Fission (in Astrophysics)

Incident neutron strikes

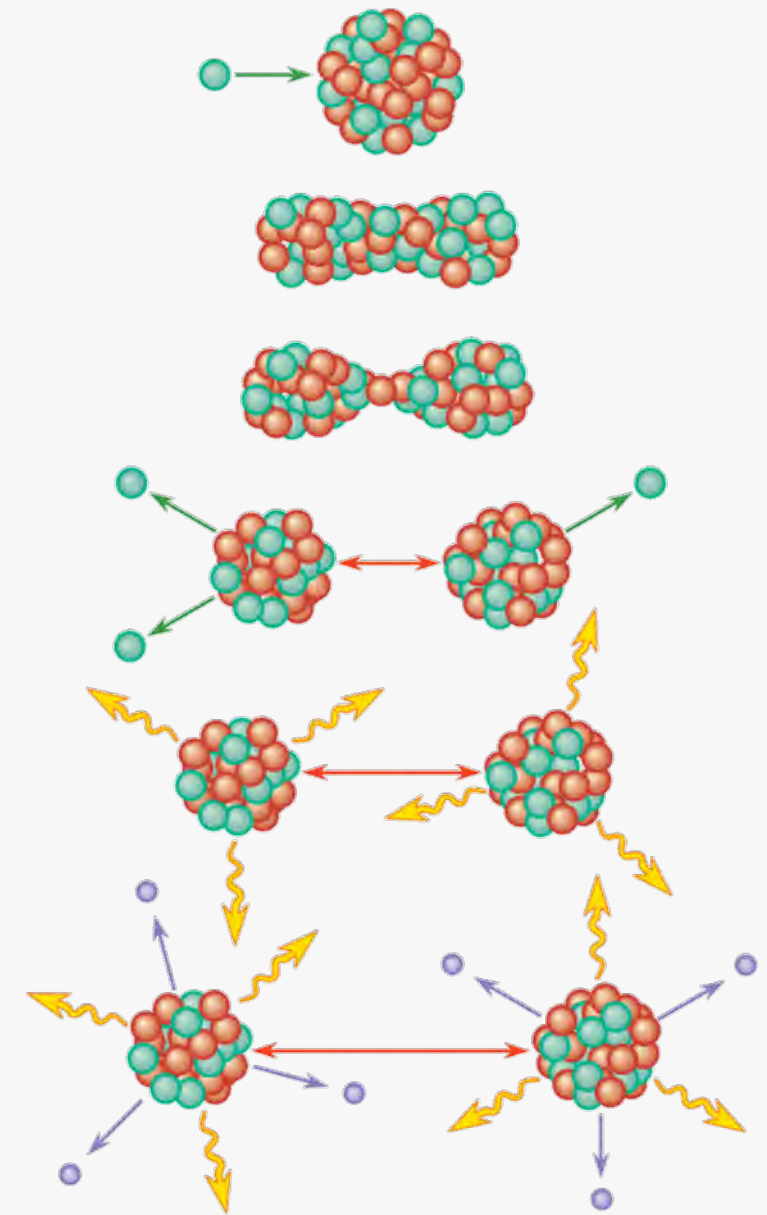
Deformation

Scission

Prompt Neutron Emission from
excited fission fragments ($\sim 2-3$)

Energy release ~ 200 MeV with
kinetic energy of fragments
 ~ 170 MeV

Delayed emission from β -decay
of neutron-rich fission products



Neutrons



Protons



Beta particles



Gamma rays

Nuclear Fission (in Astrophysics)

Incident neutron strikes

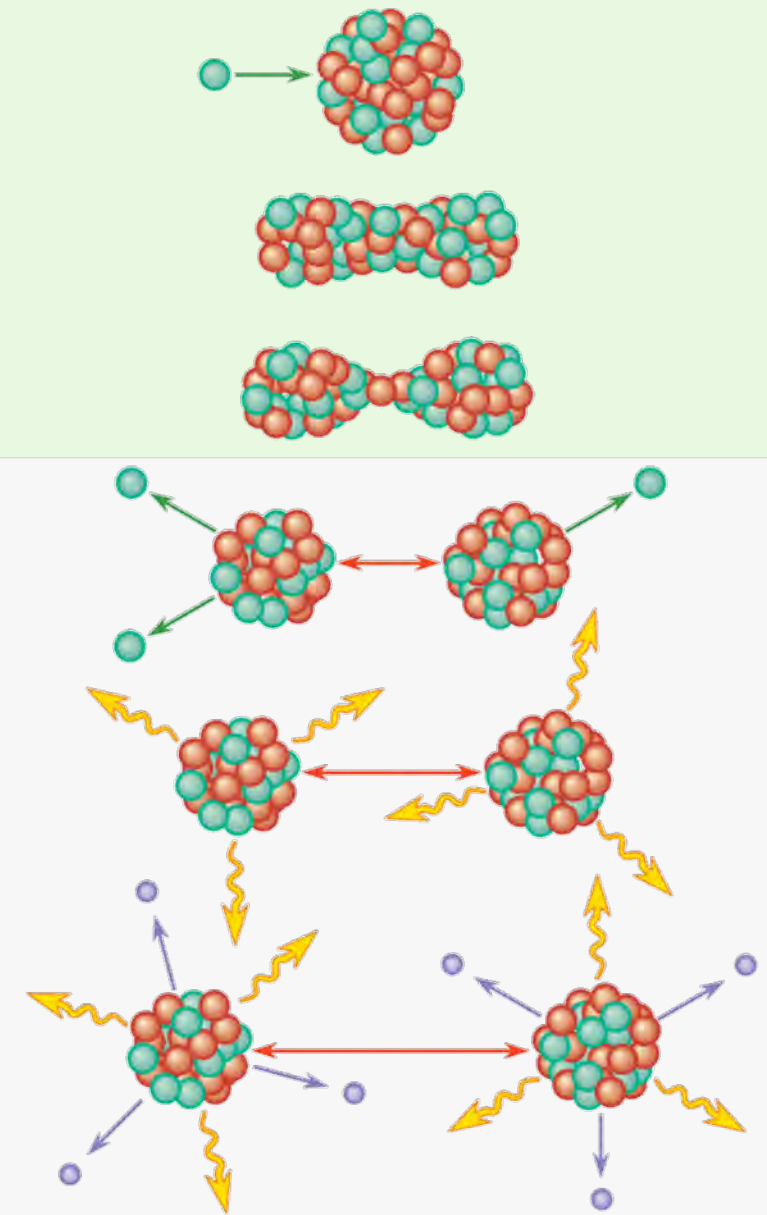
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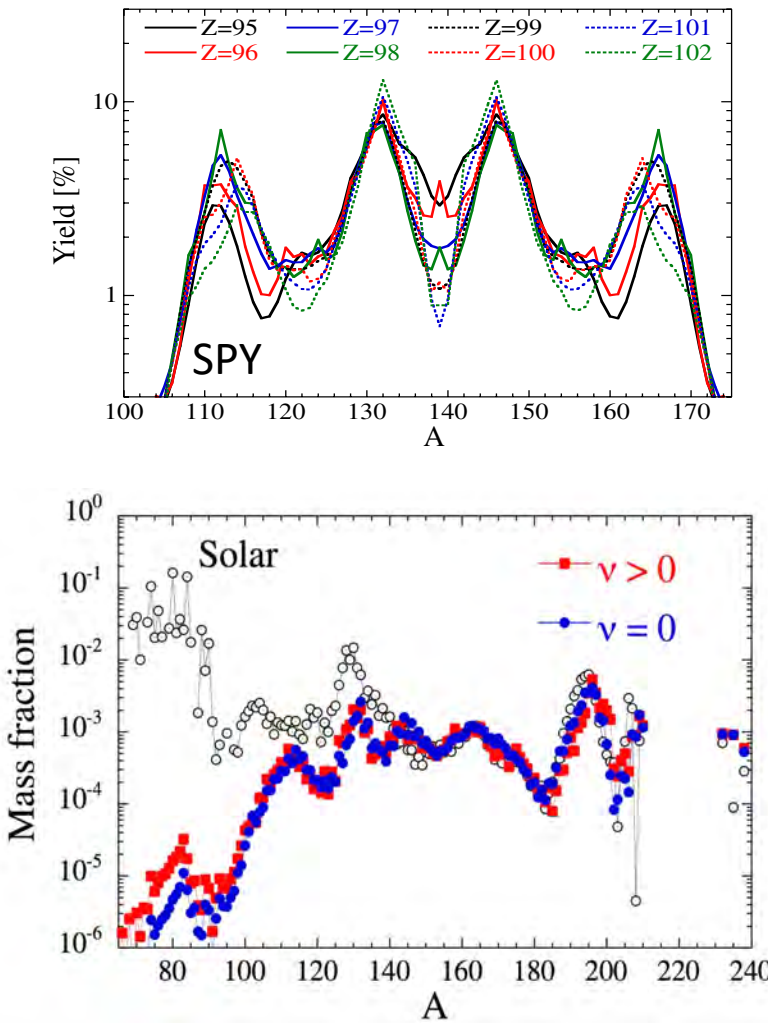
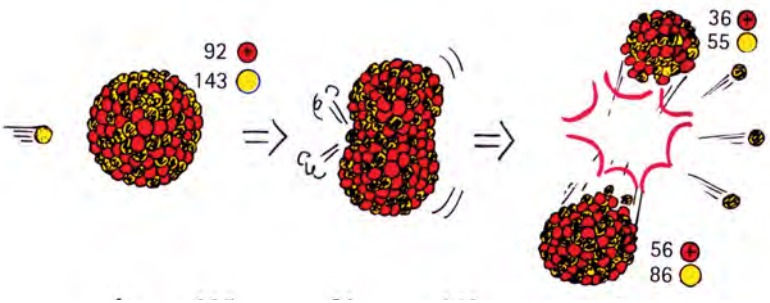


Beta particles

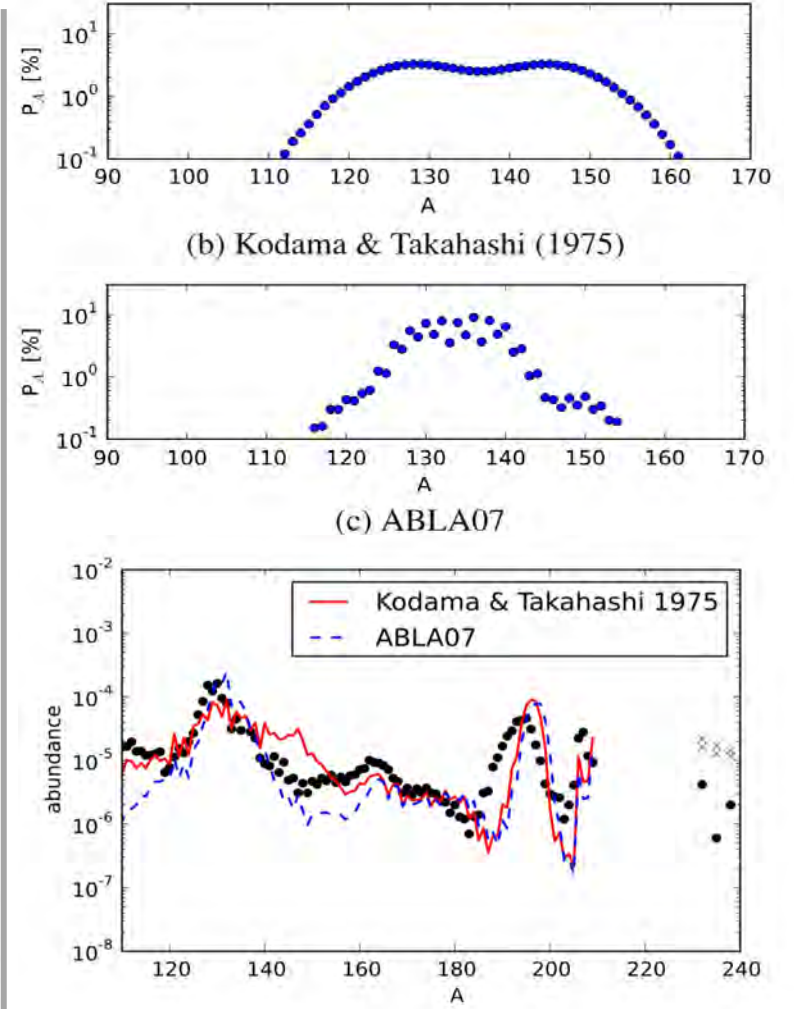


Gamma rays

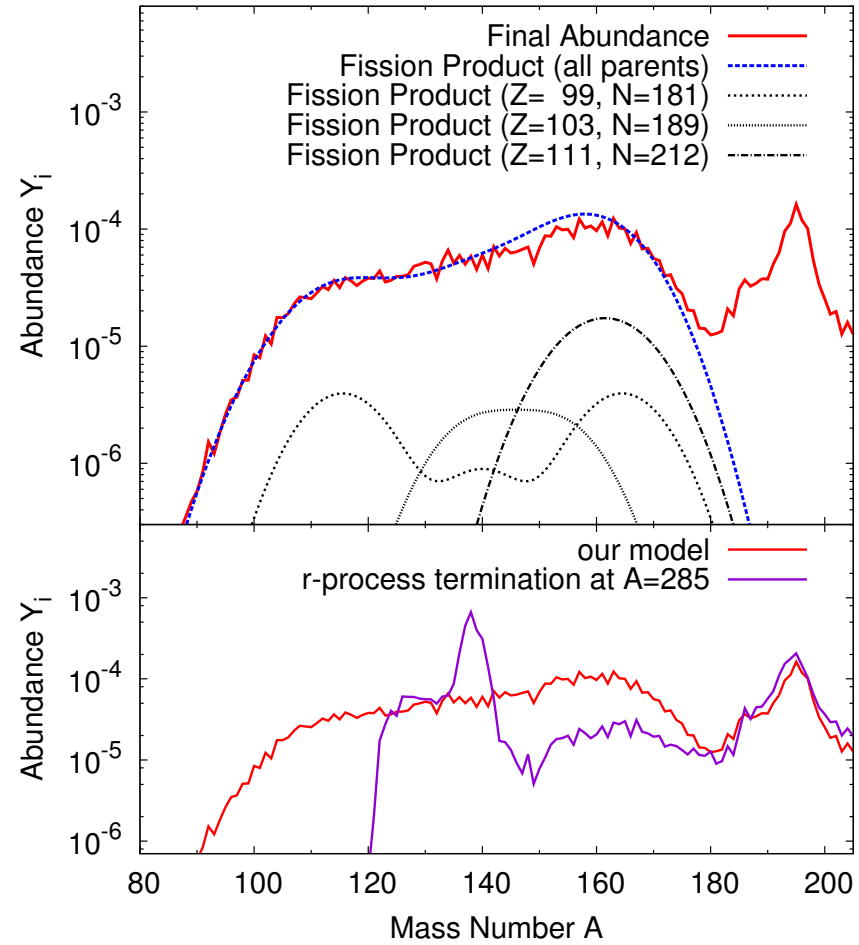
Dependence of *r*-process abundances on fission yields



Goriely+13,15; see also Lemaître+21

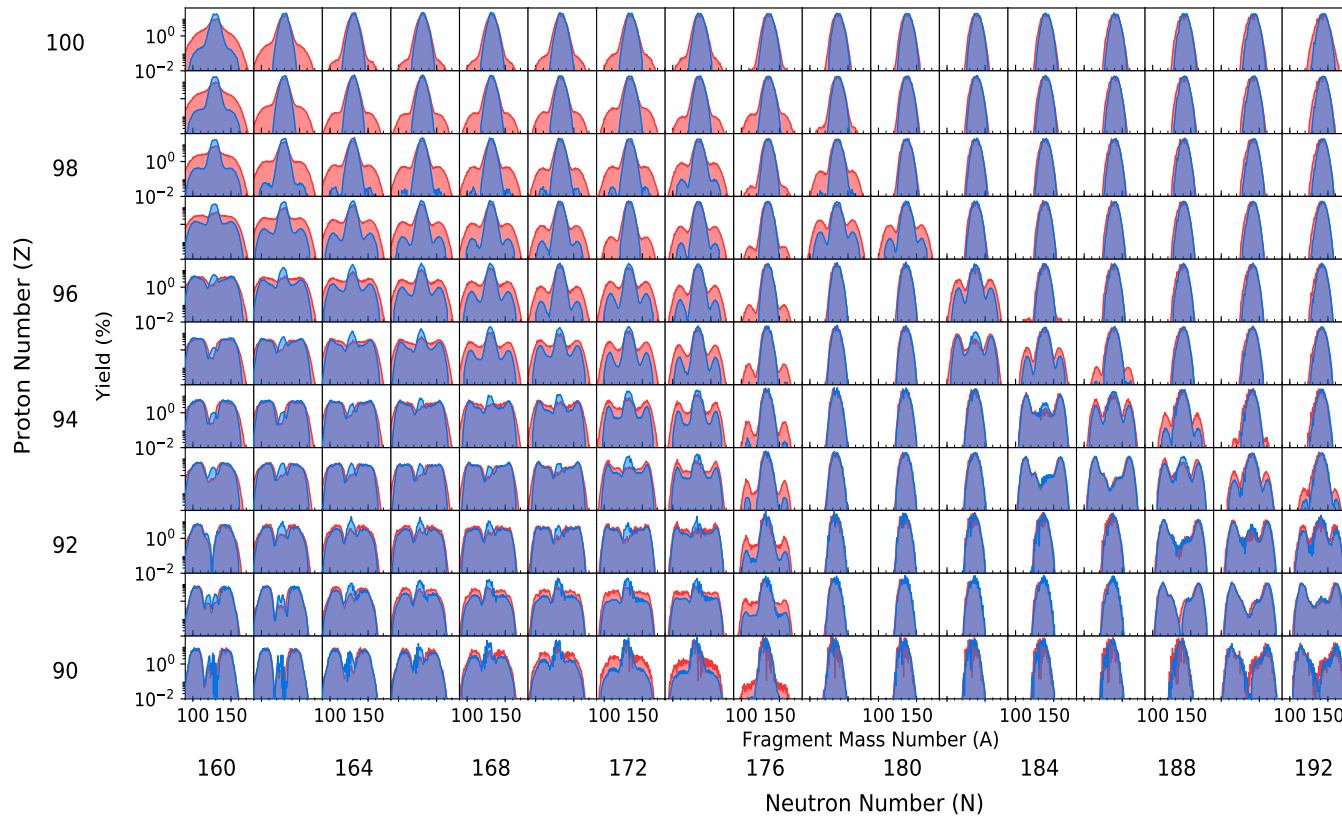


Eichler+15

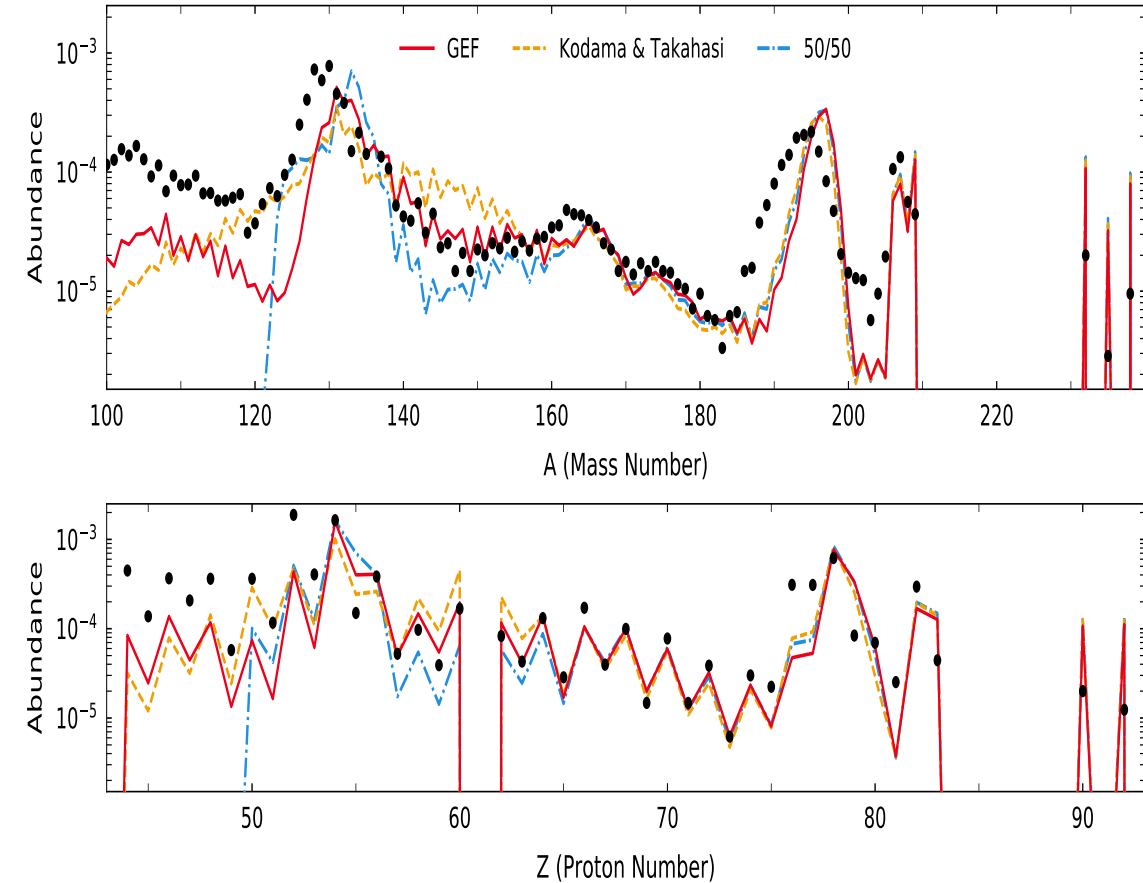


Shibagaki+16

Excitation energy dependence: distinct fission yields for neutron-induced, β -delayed, and spontaneous fission



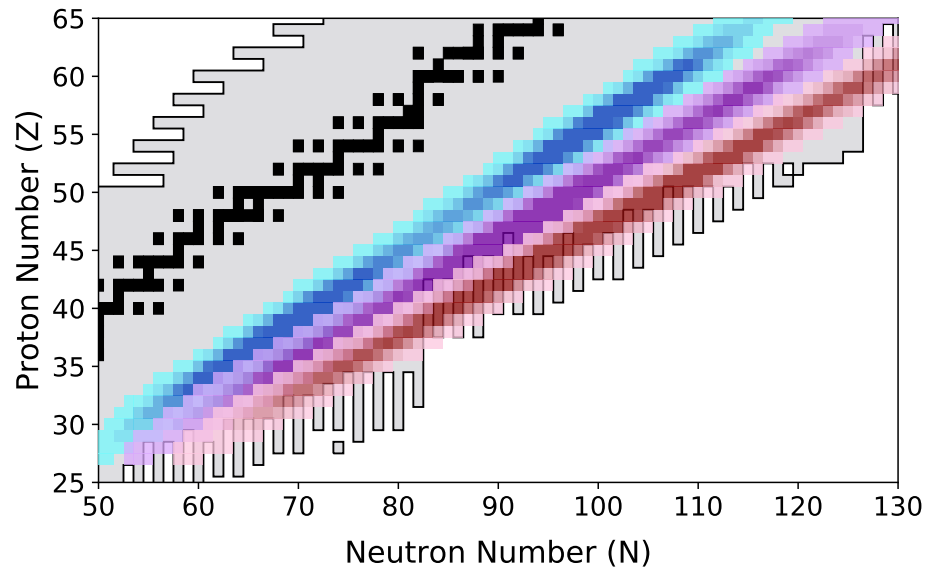
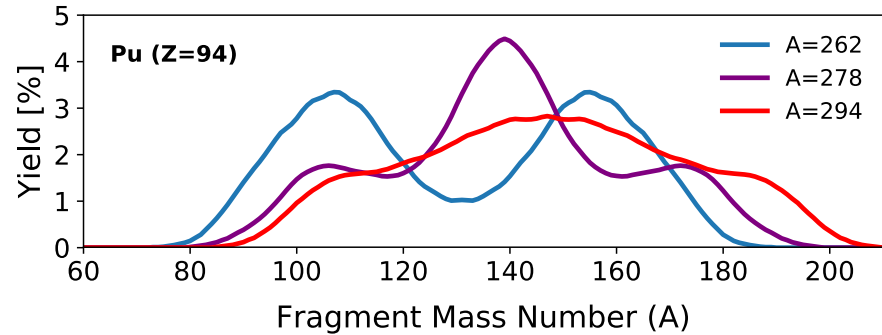
(n,f) yields with excitation energy $E_i + S_n$ differ from sf yields which have zero excitation energy (above from GEF 2016)



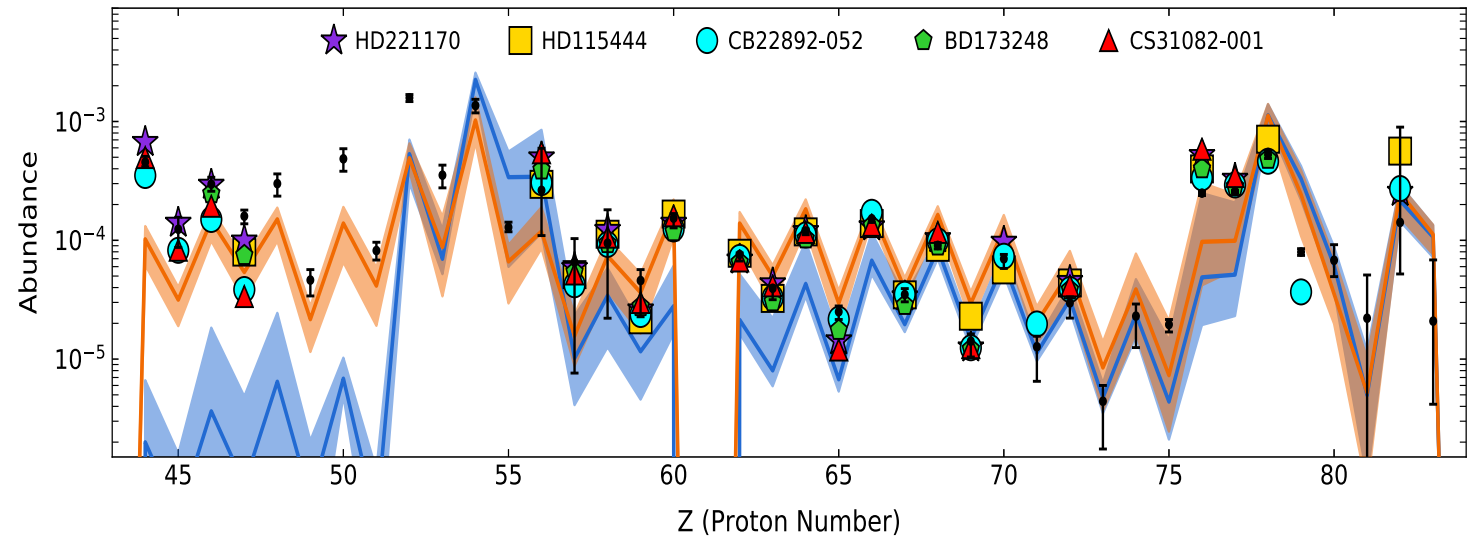
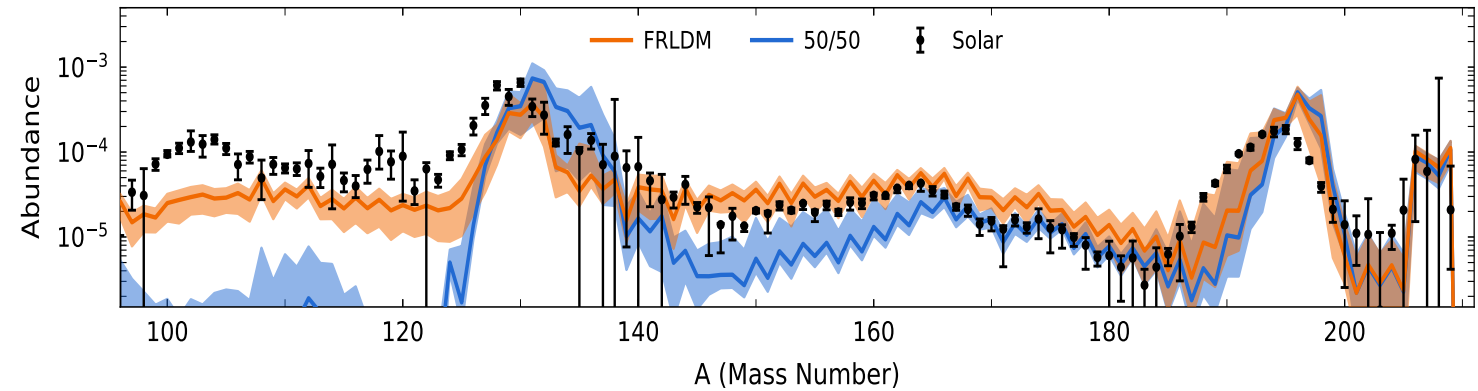
Vassh+19

Using fission yields and fission rates calculated with self-consistent fission barriers

FRLDM Yields from Mumpower+20

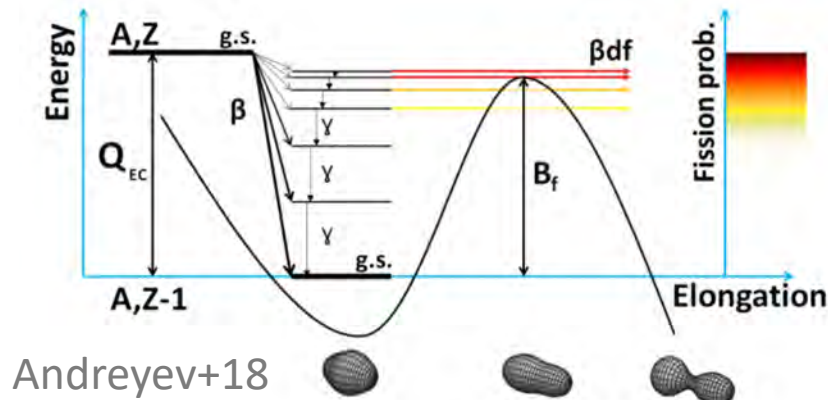


Range from Rosswog+13 NSM dynamical ejecta ($Y_e \sim 0.01-0.05$)

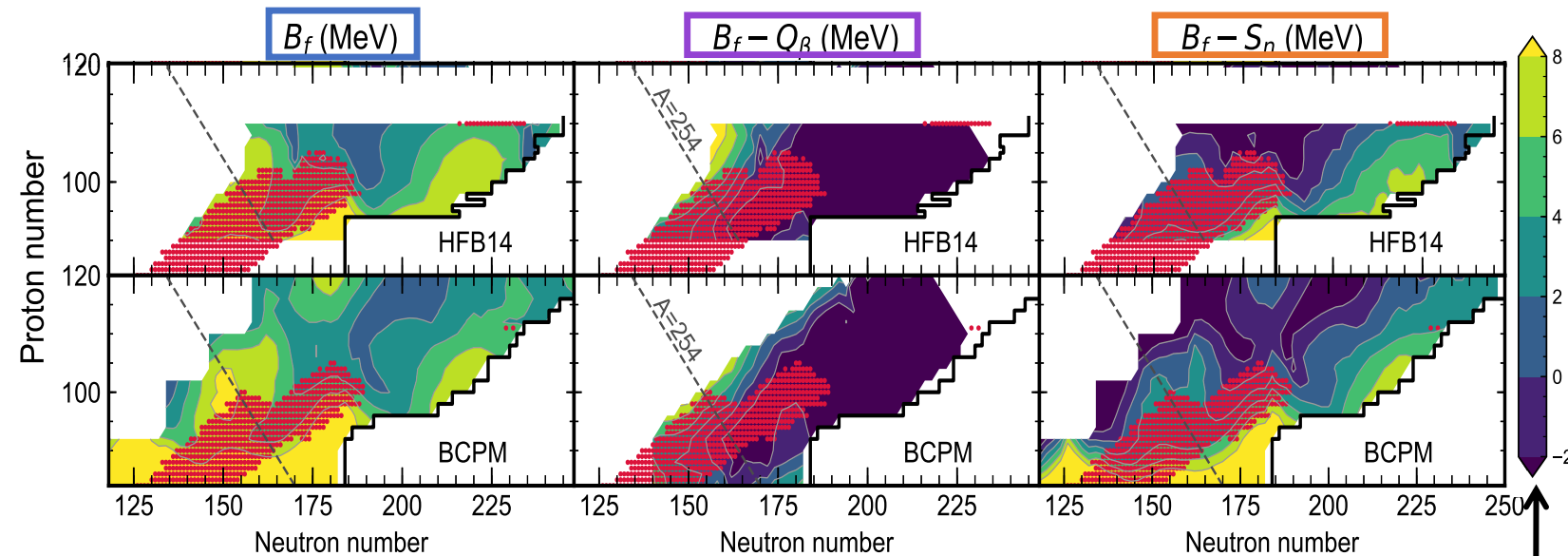


Vassh+20

Fission and the ultimate reach of the r process

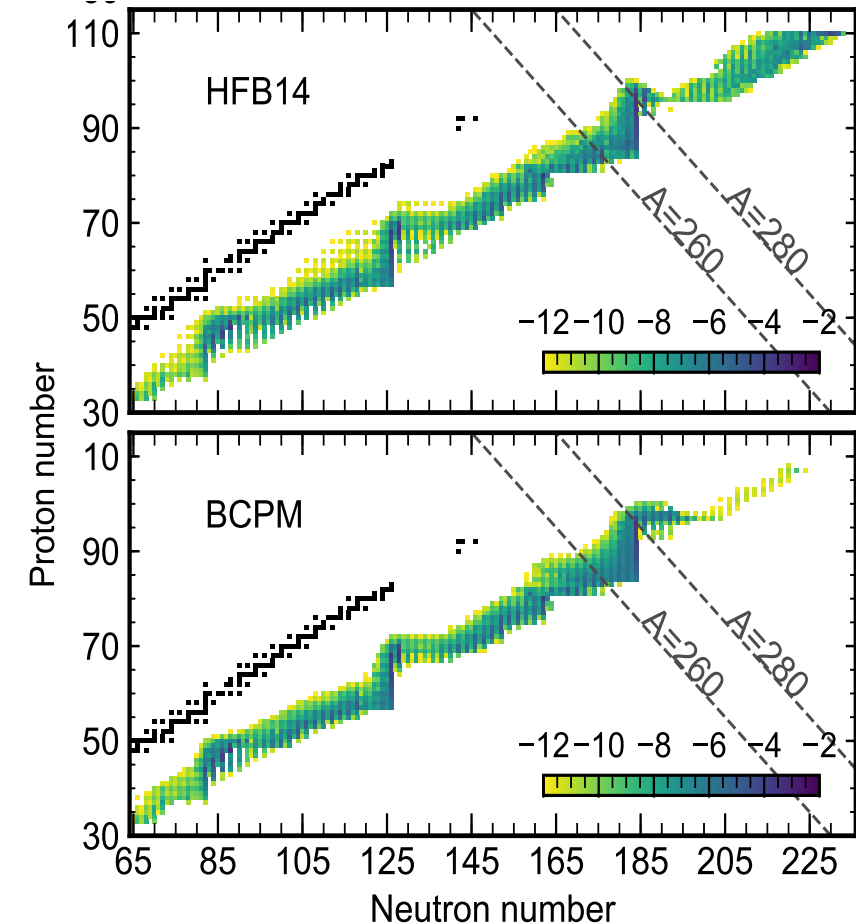


Fission barriers and masses are used to determine spontaneous, β -delayed, and neutron-induced fission rates



Giuliani+20

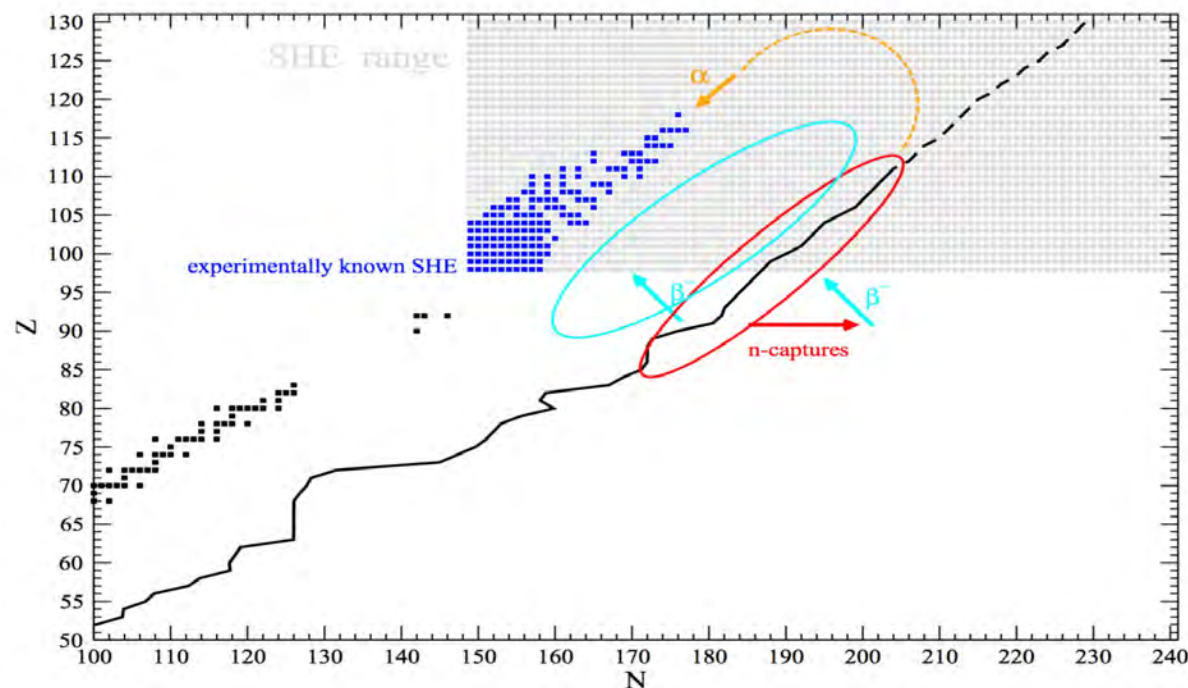
Smaller \Rightarrow larger fission probability



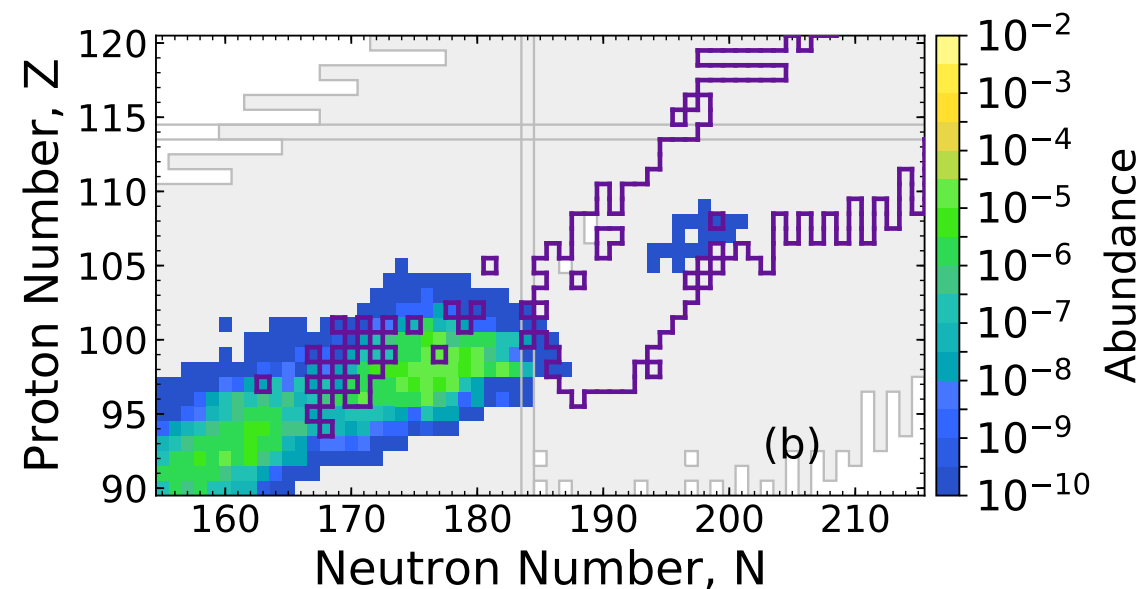
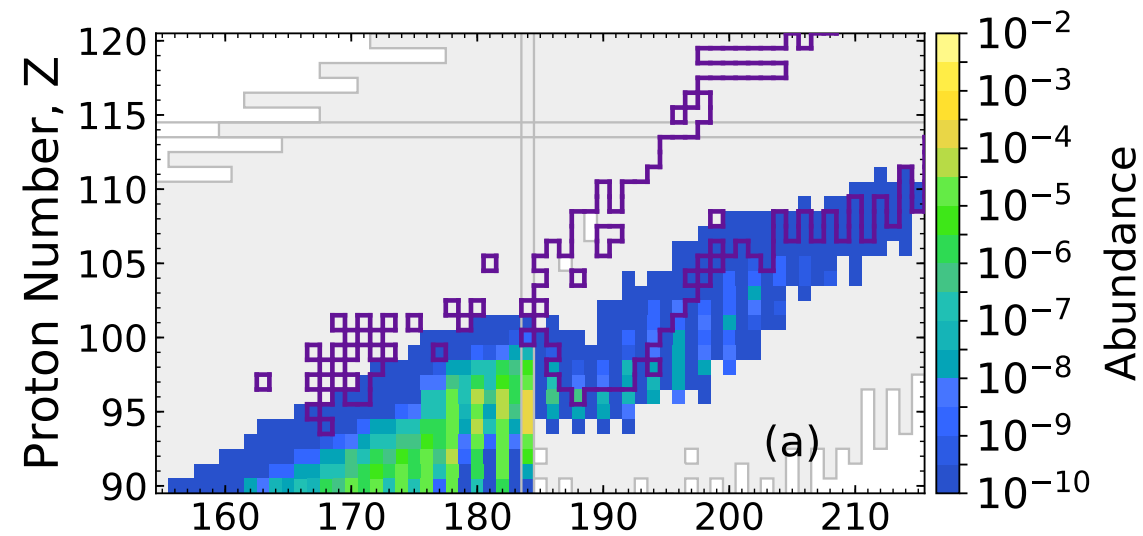
r-process production of superheavy elements?

Super heavy elements ($Z \gtrsim 103$) have been produced in laboratories and models predict an “island of stability” at $Z=114$, $N=184$, but current *r*-process calculations see fission prevent the population of such species

→ if observed in nature, fission barriers would have to differ from theory predictions



Petermann+12



(purple outline – probability of β df $\geq 90\%$)

Mumpower+18 (including Vassh)

Nuclear Fission (in Astrophysics)

Incident neutron strikes

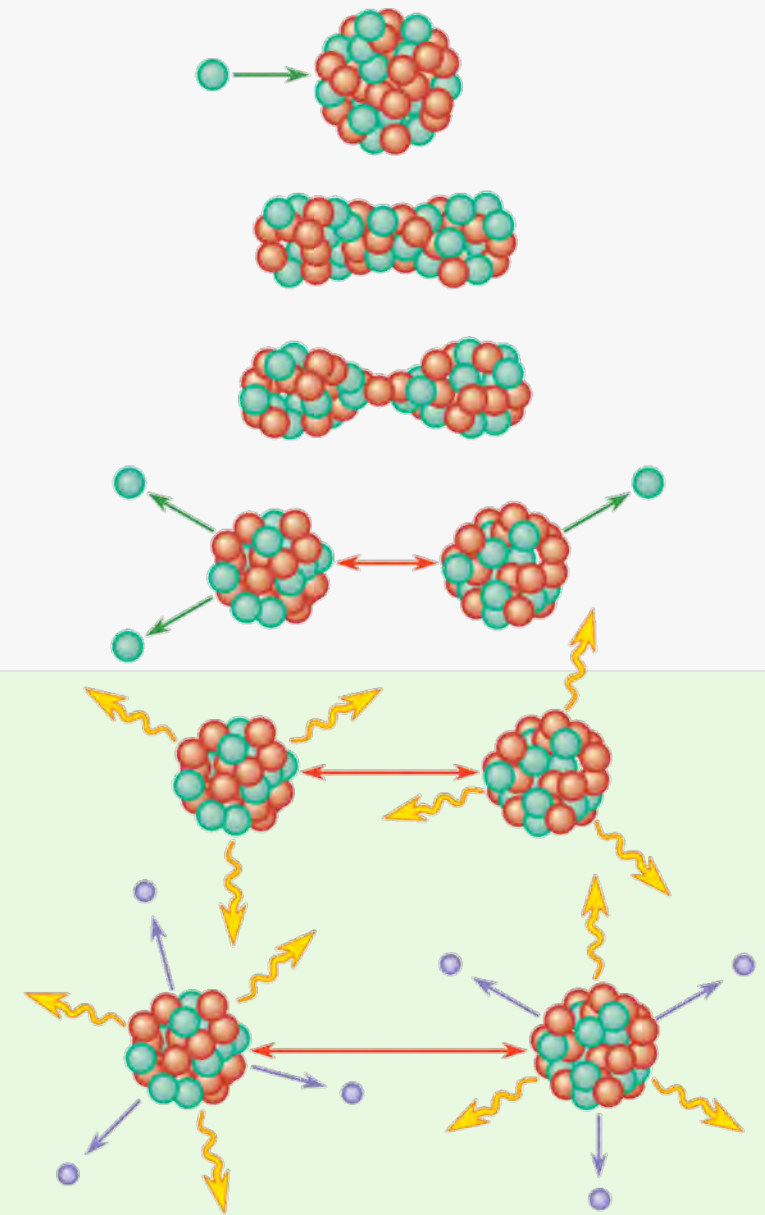
Deformation

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Energy release ~ 200 MeV with
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 ~ 170 MeV

Delayed emission from β -decay
of neutron-rich fission products



● Neutrons

● Protons

● Beta particles

~ Gamma rays

Are actinides produced in neutron star mergers?

PHYSICAL REVIEW

VOLUME 103, NUMBER 5

SEPTEMBER 1, 1956

Californium-254 and Supernovae*

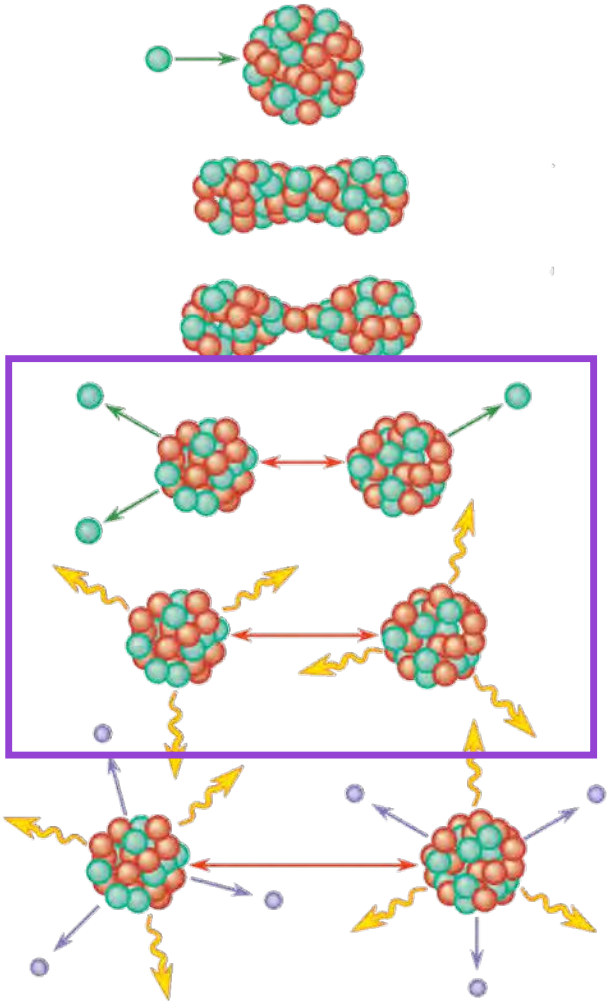
G. R. BURBIDGE AND F. HOYLE,† *Mount Wilson and Palomar Observatories, Carnegie Institution of Washington, California Institute of Technology, Pasadena, California*

AND

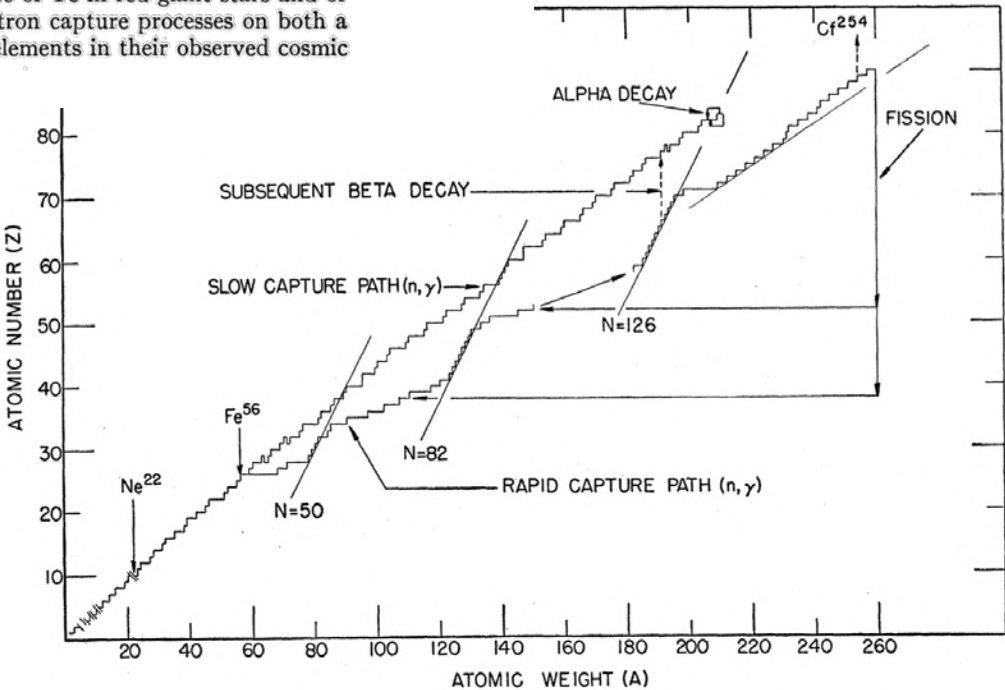
E. M. BURBIDGE, R. F. CHRISTY, AND W. A. FOWLER, *Kellogg Radiation Laboratory, California Institute of Technology, Pasadena, California*

(Received May 17, 1956)

It is suggested that the spontaneous fission of Cf^{254} with a half-life of 55 days is responsible for the form of the decay light-curves of supernovae of Type I which have an exponential form with a half-life of 55 nights. The way in which Cf^{254} may be synthesized in a supernova outburst, and reasons why the energy released by its decay may dominate all others are discussed. The presence of Tc in red giant stars and of Cf in Type I supernovae appears to be observational evidence that neutron capture processes on both a slow and a fast time-scale have been necessary to synthesize the heavy elements in their observed cosmic abundances.

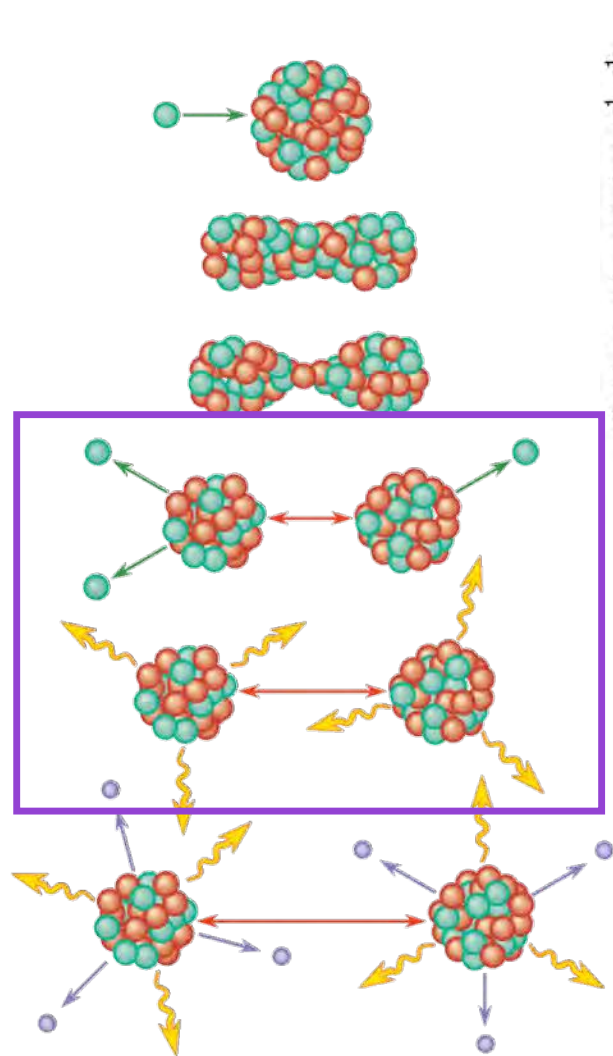


Fission heating can greatly impact light curves

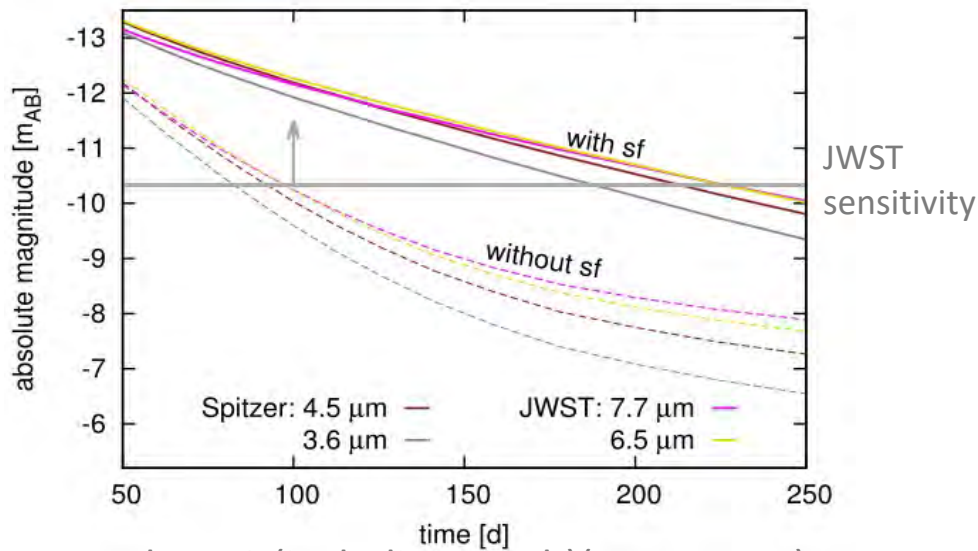
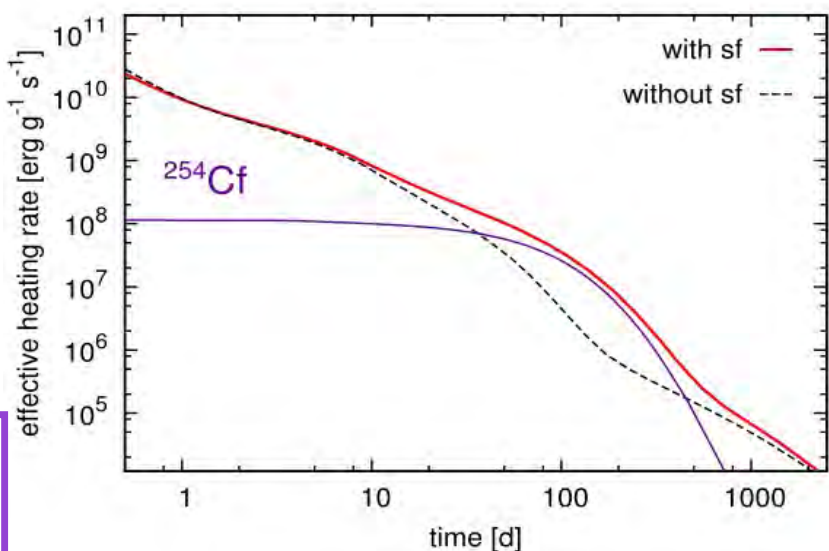


Burbidge, Burbidge, Fowler and Hoyle (1957)

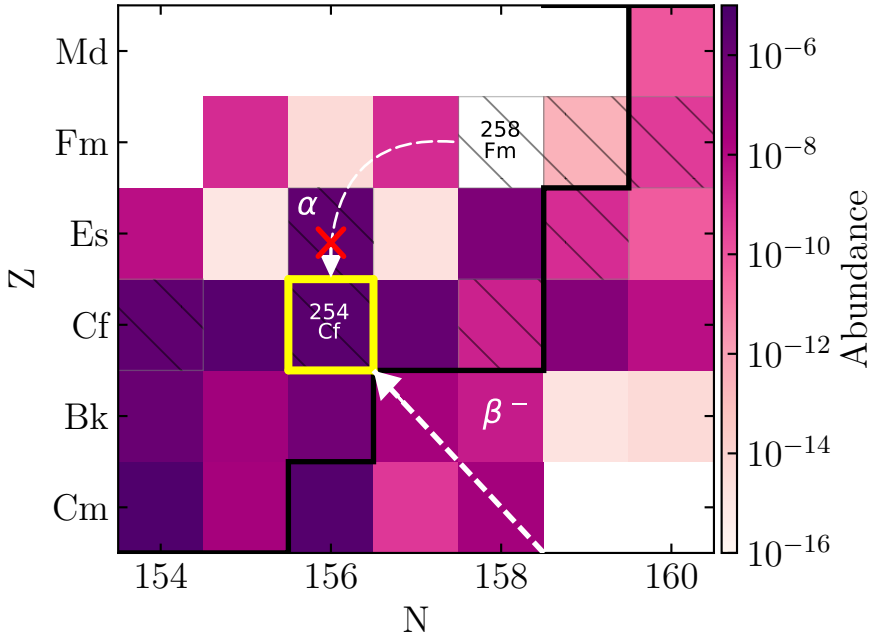
Are actinides produced in neutron star mergers?



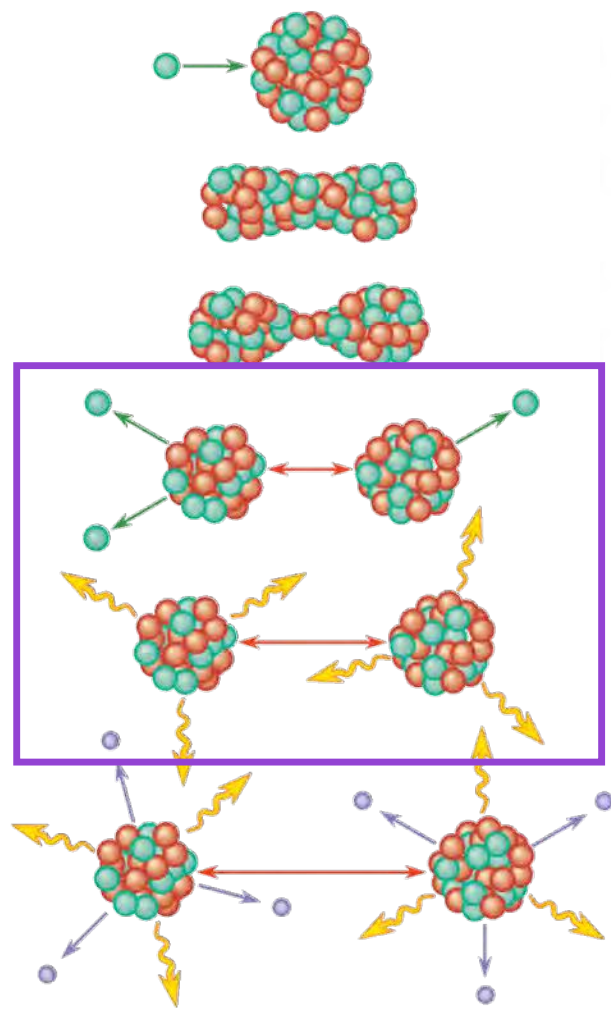
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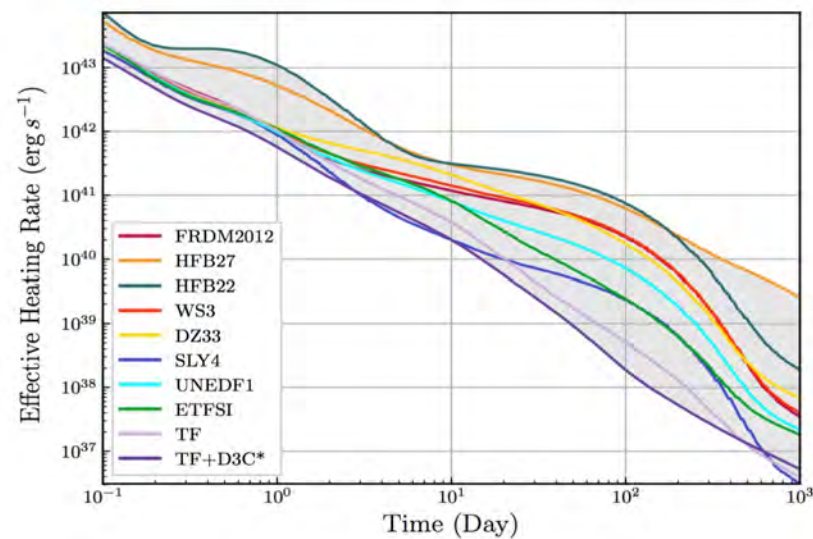
Zhu+18 (including Vassh)(ApJ Letters)



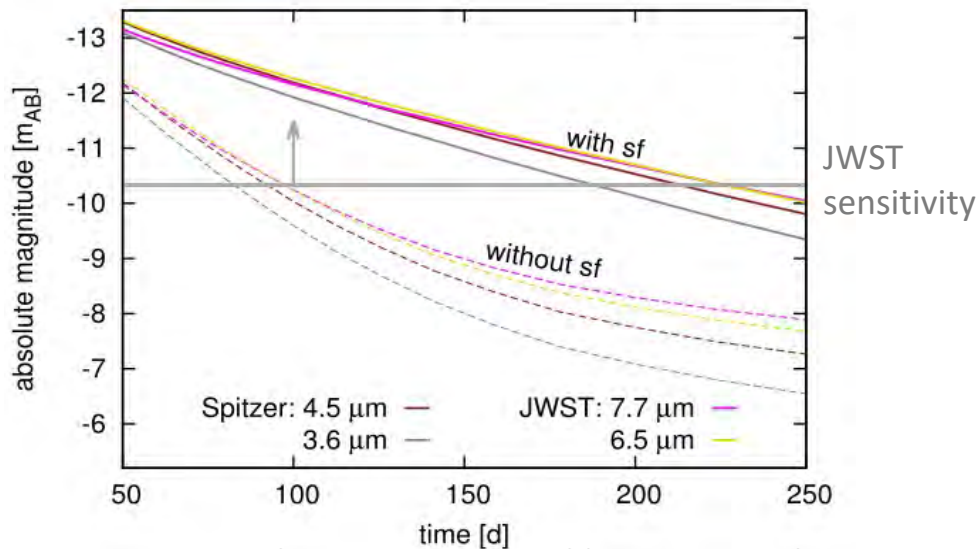
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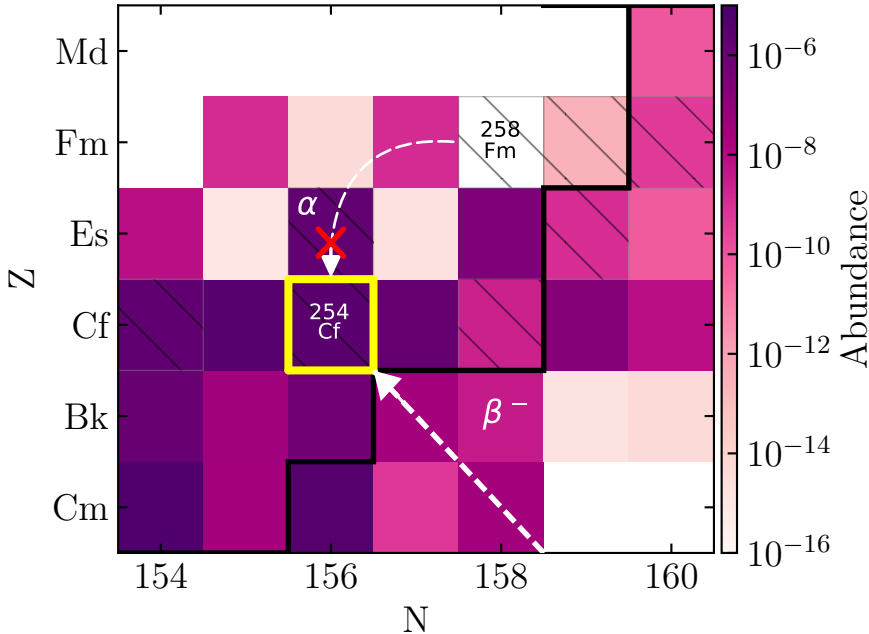
Fission heating can greatly impact light curves



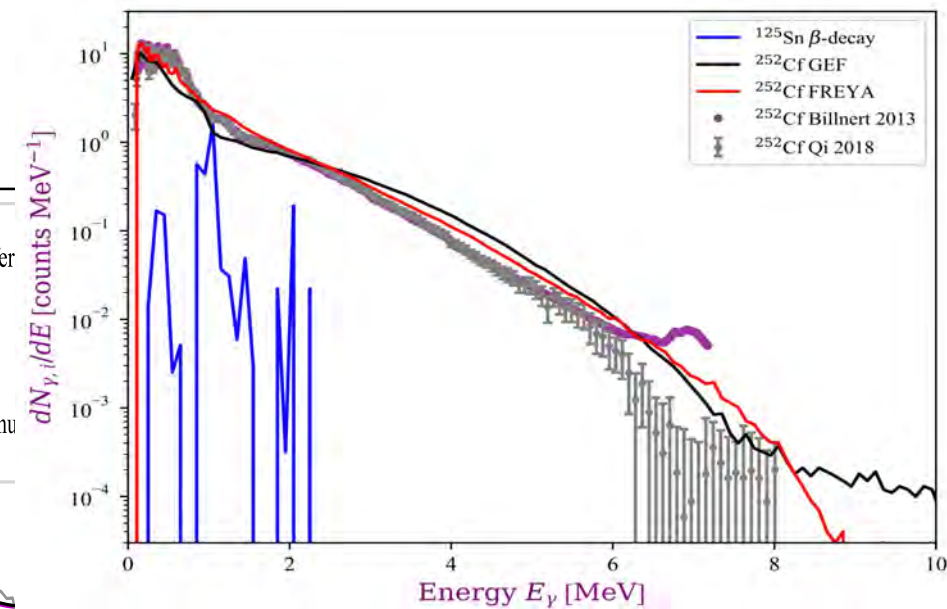
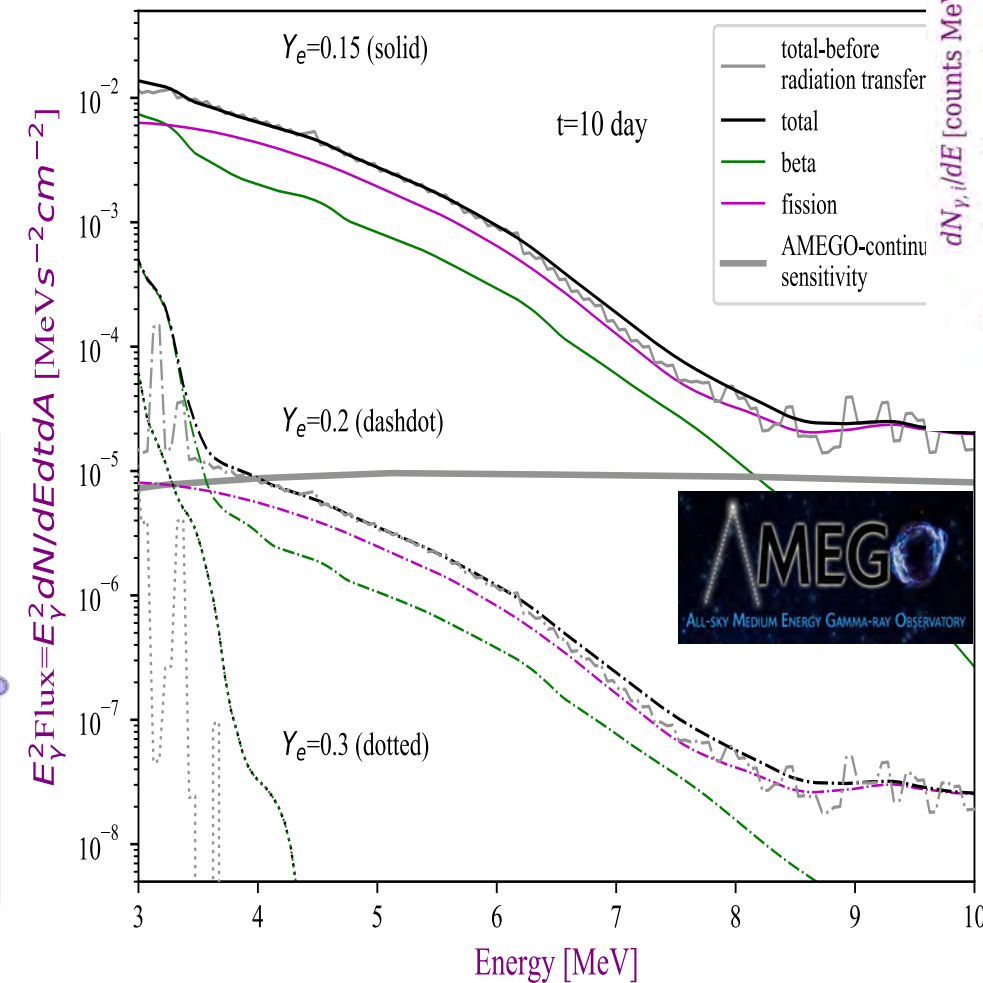
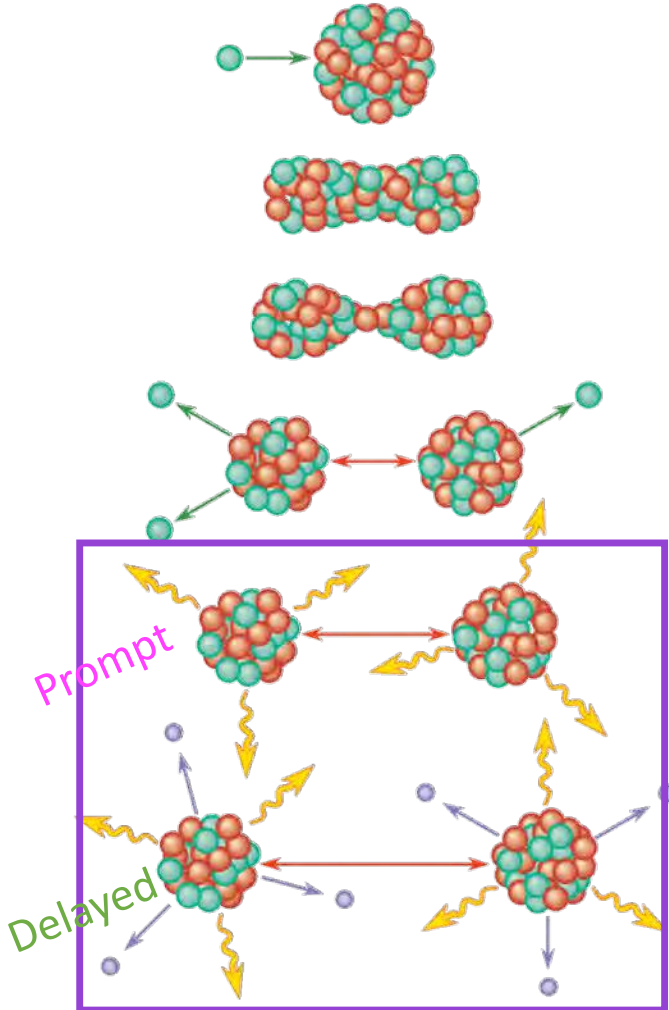
Zhu+21 (including Vassh)(ApJ)



Zhu+18 (including Vassh)(ApJ Letters)



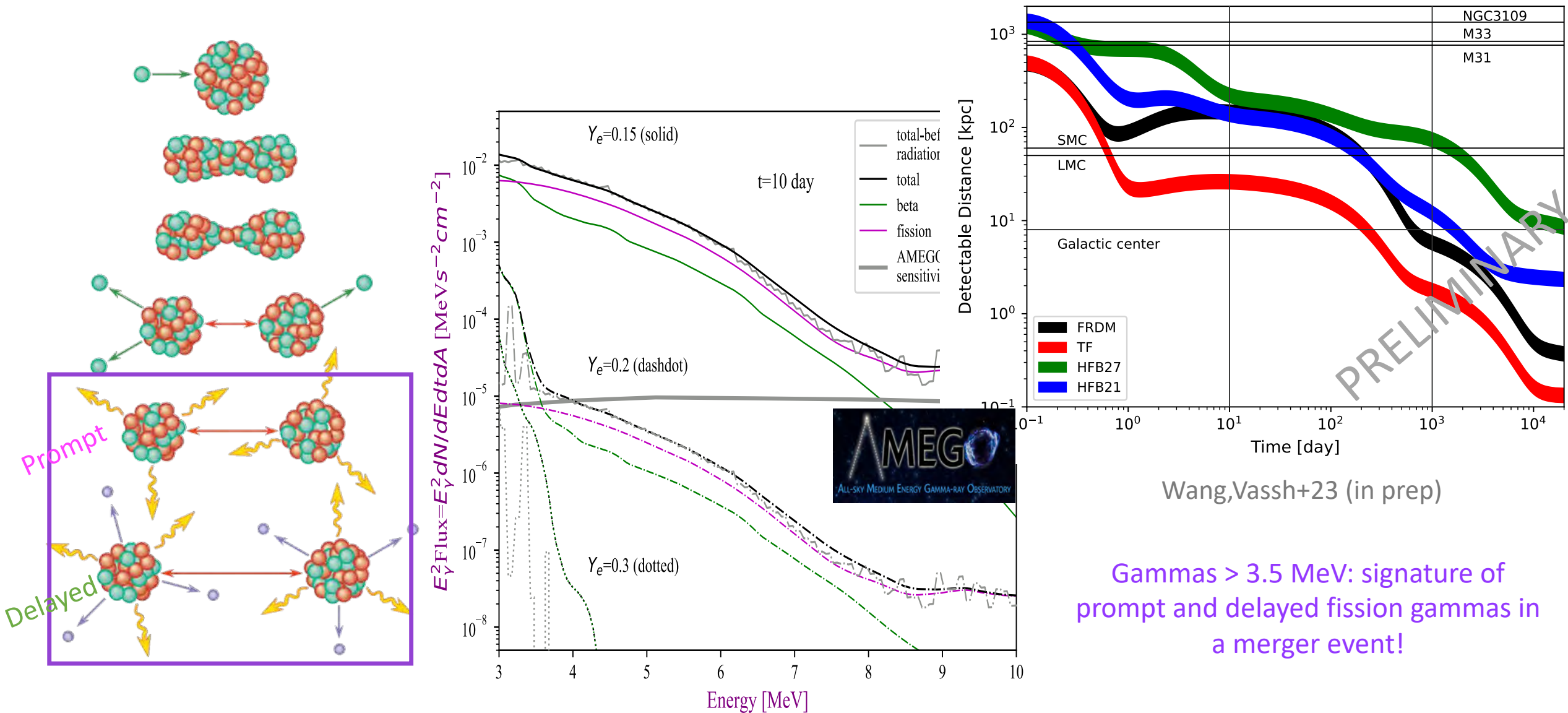
Are actinides produced in neutron star mergers?



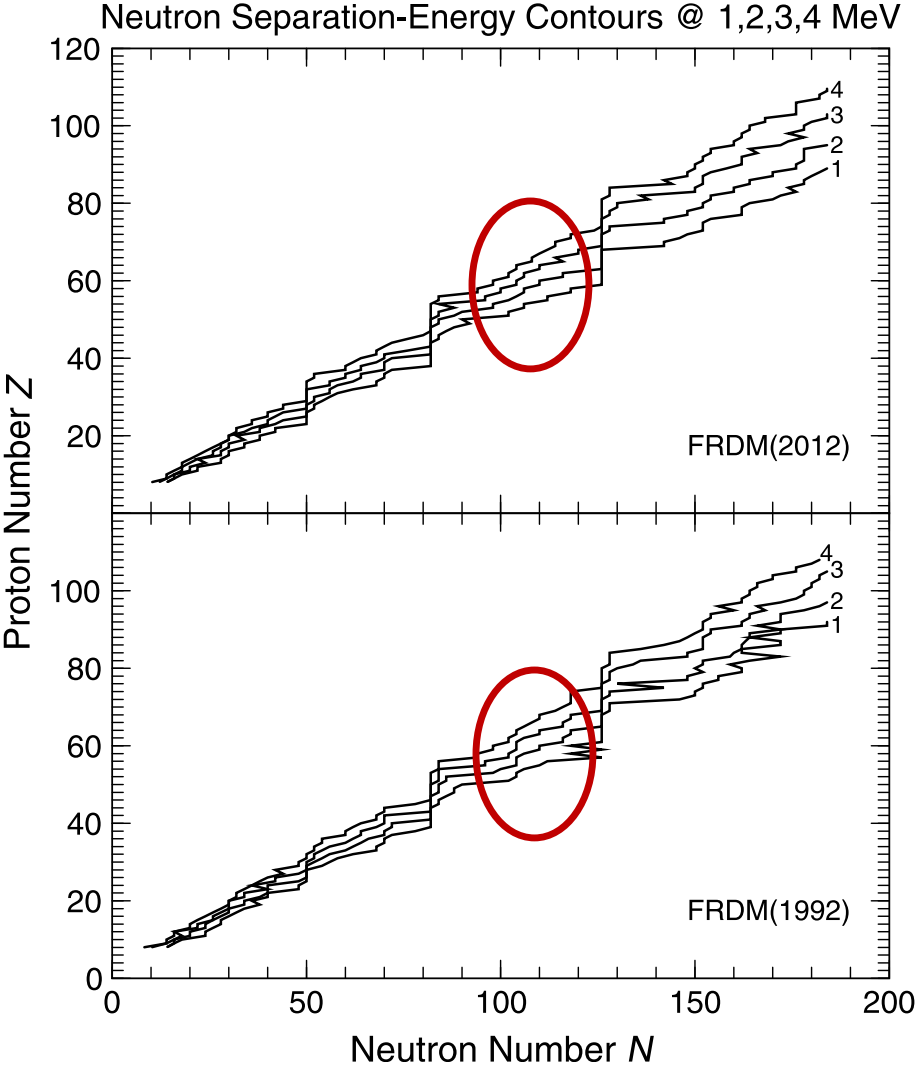
Wang,Vassh+20 (ApJ Letters)

Gammas > 3.5 MeV: signature of prompt and delayed fission gammas in a merger event!

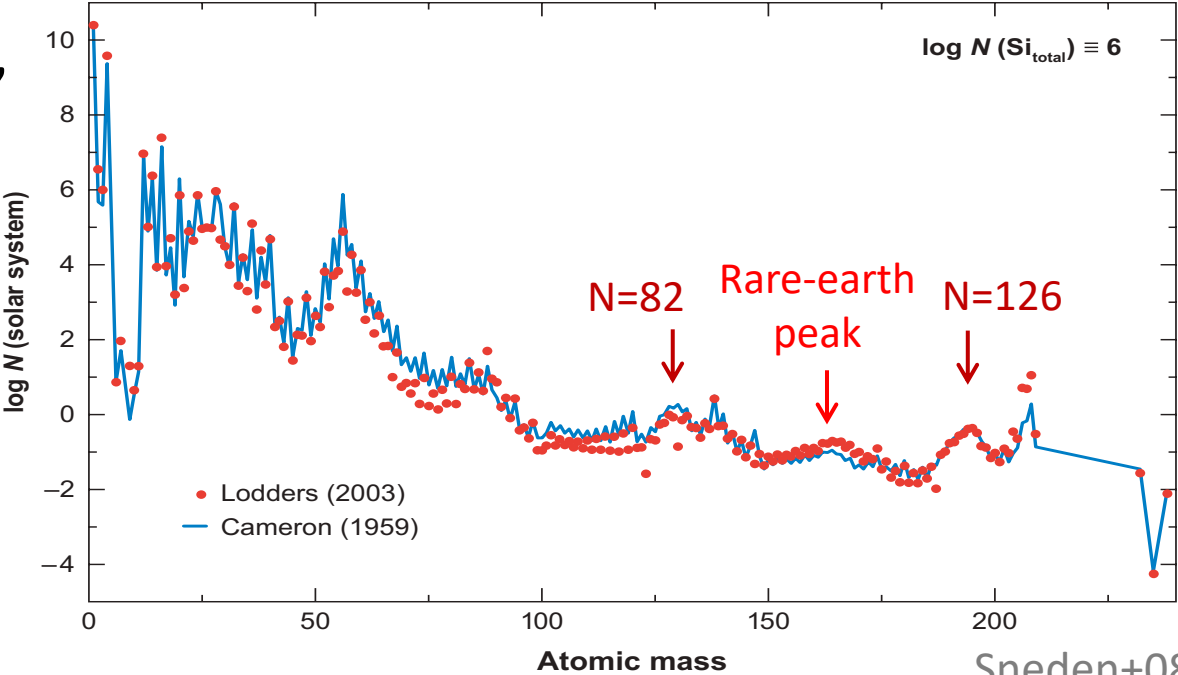
Are actinides produced in neutron star mergers?



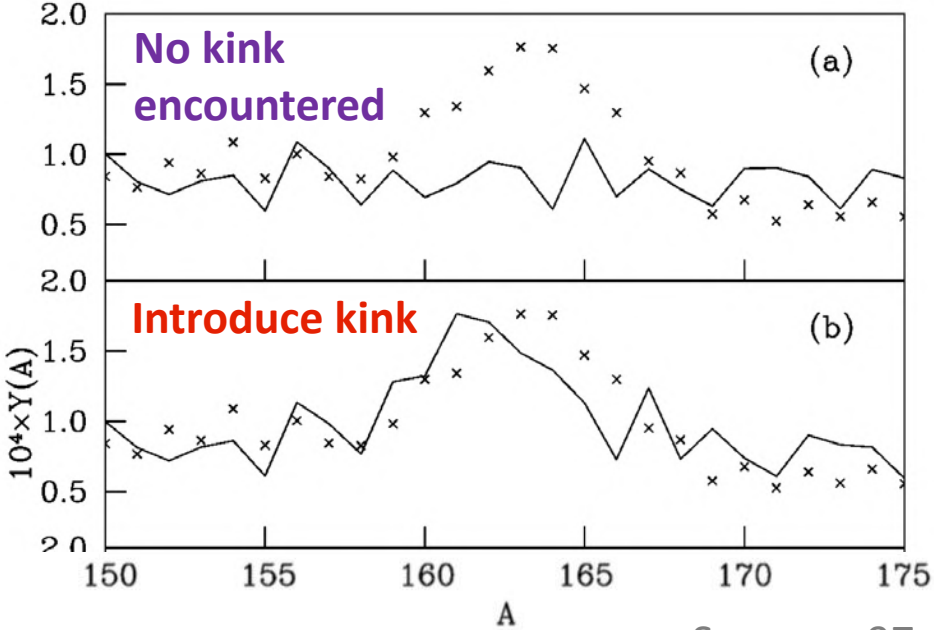
Nuclear masses, nuclear structure, and the rare-earth peak



Möller+12

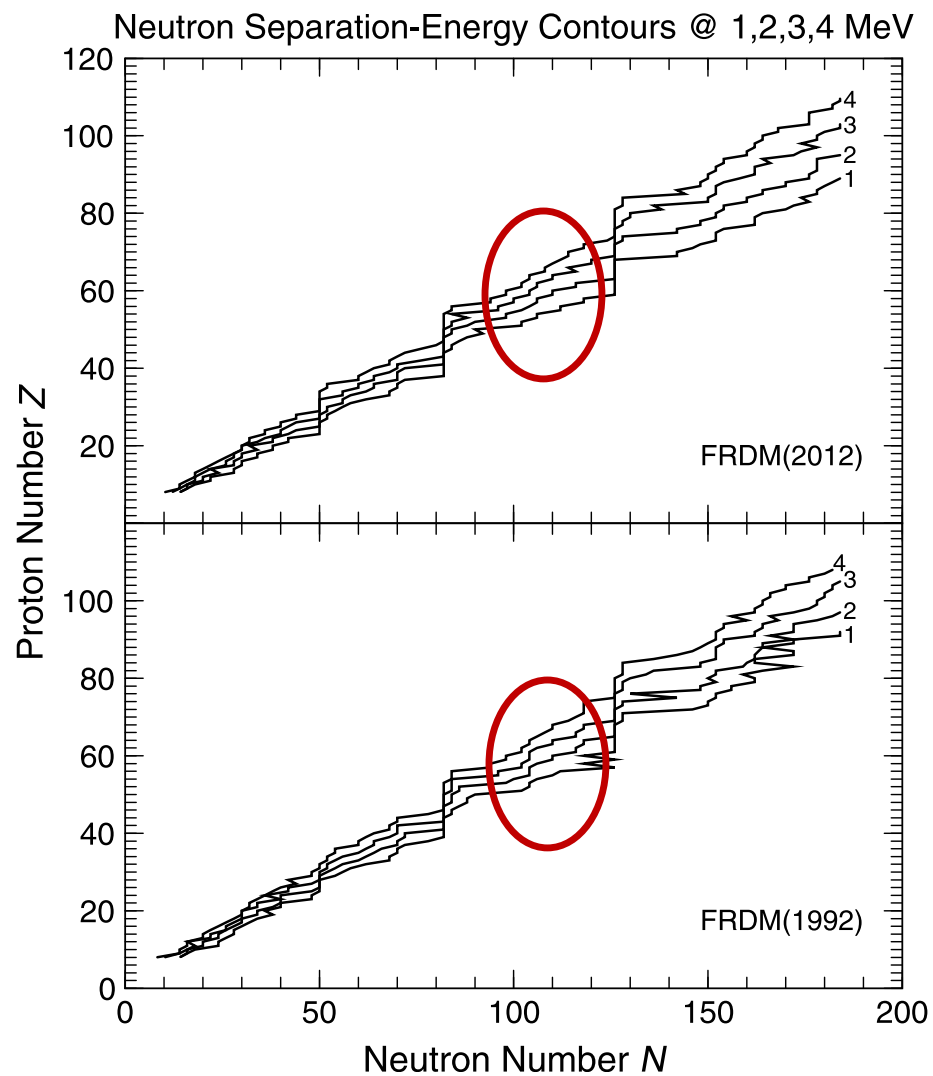


Sneden+08

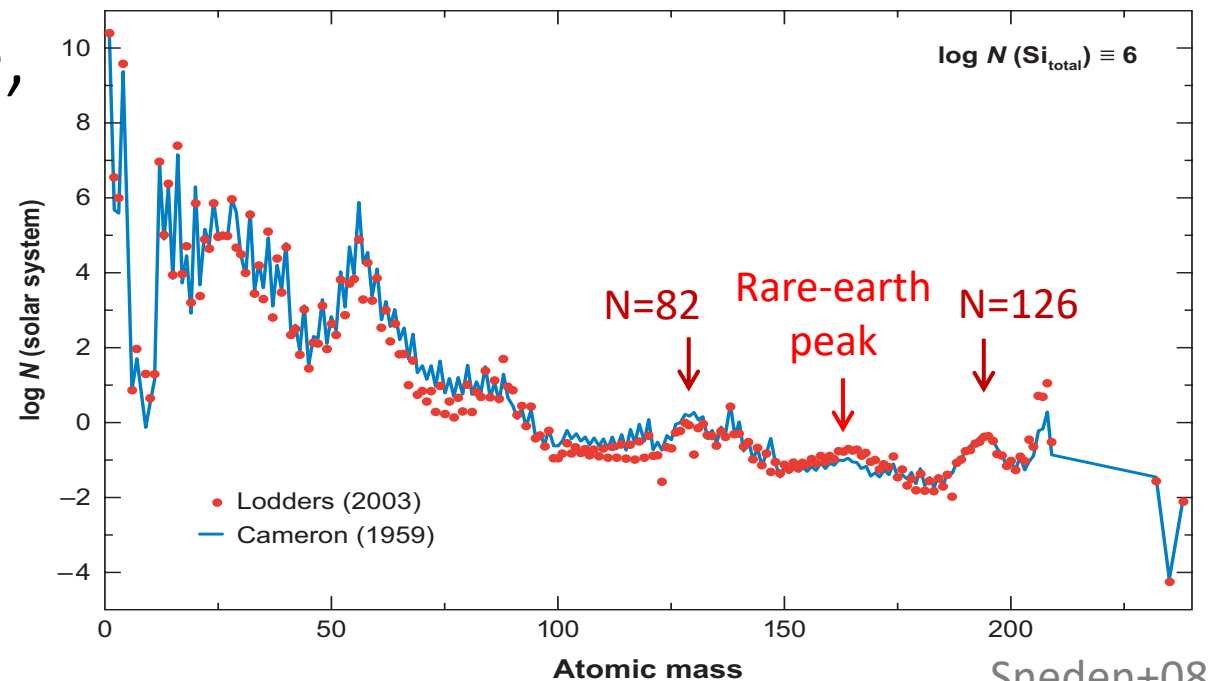


Surman+97

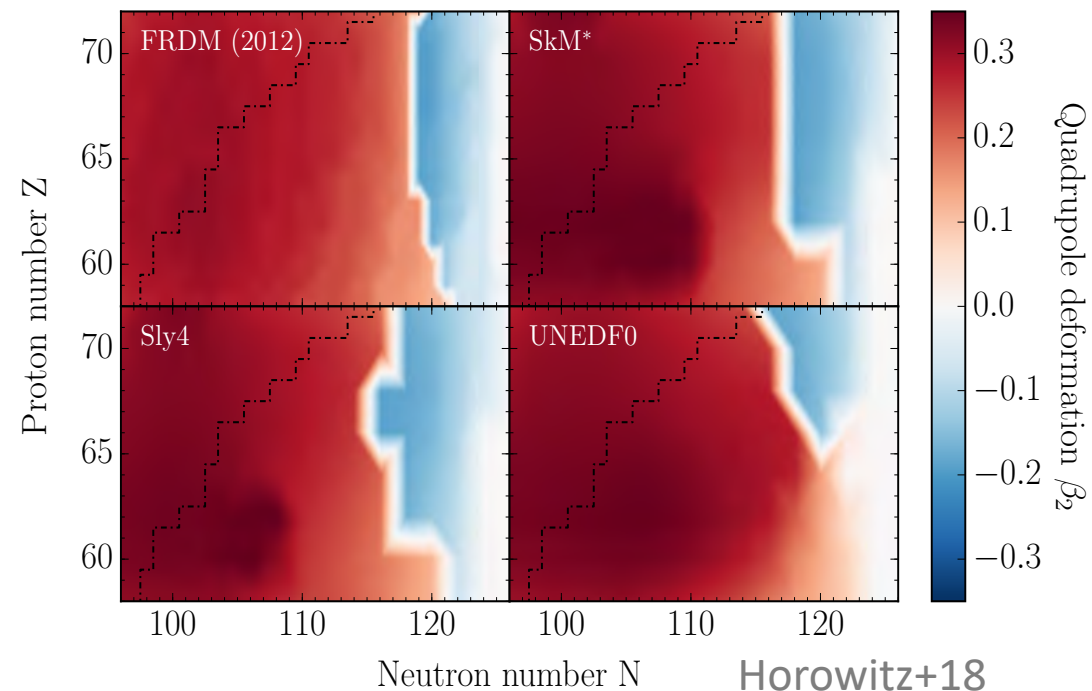
Nuclear masses, nuclear structure, and the rare-earth peak



Möller+12



Sneden+08



Horowitz+18

Markov Chain Monte Carlo (MCMC):

Uses observational data to discern nuclear properties such as masses as well as constrain the conditions present at nucleosynthesis sites

- Monte Carlo mass corrections

$$M(Z, N) = M_{DZ}(Z, N) + a_N e^{-(Z-c)^2/2f}$$

- Calculate: $\sigma_{\text{rms}}^2(M_{\text{AME12}}, M) \leq \sigma_{\text{rms}}^2(M_{\text{AME12}}, M_{DZ})$

- Calculate:

$$D_n(Z, A) = (-1)^{A-Z+1} (S_n(Z, A+1) - S_n(Z, A)) > 0$$

- Update nuclear quantities and rates

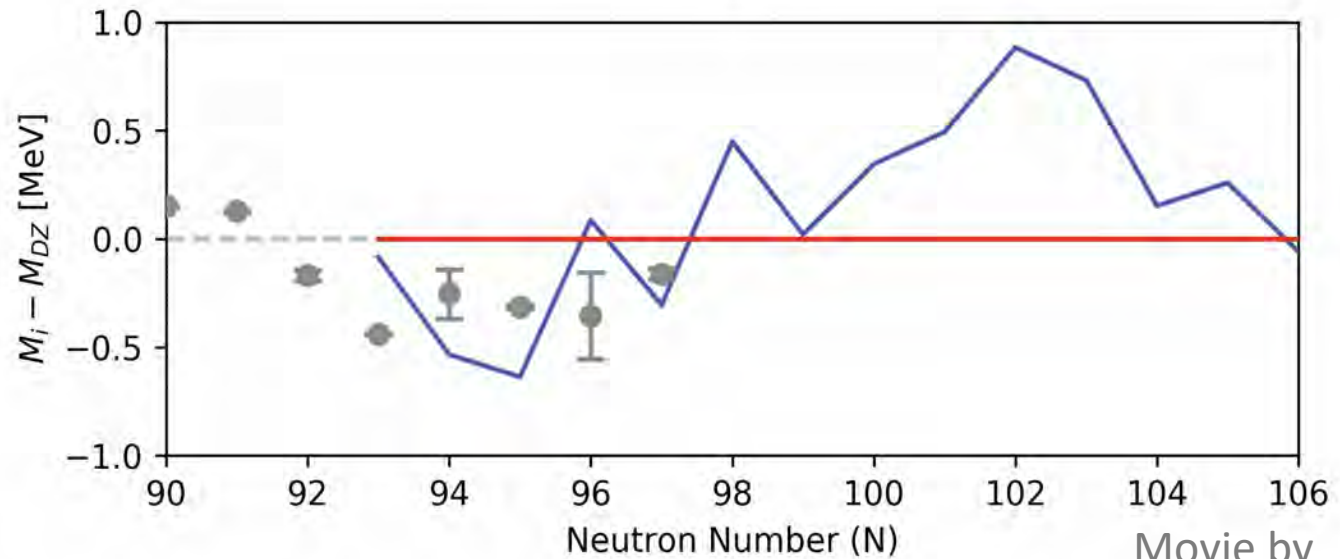
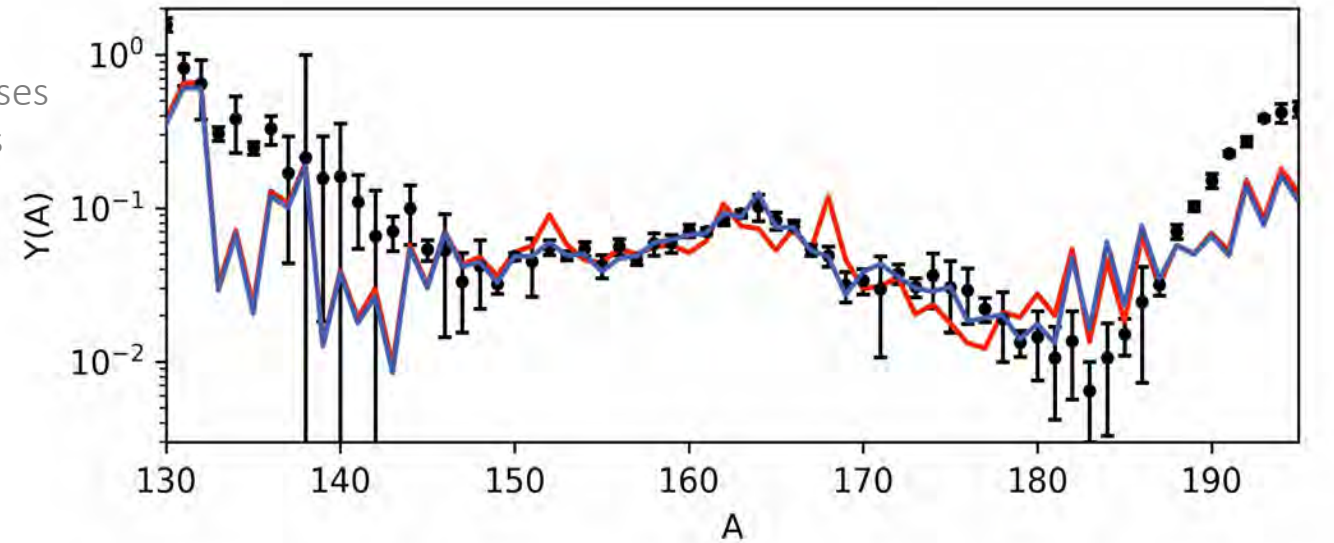
- Perform nucleosynthesis calculation

- Calculate $\chi^2 = \sum_{A=150}^{180} \frac{(Y_{\odot,r}(A) - Y(A))^2}{\Delta Y(A)^2}$

- Update parameters OR revert to last success

$$\mathcal{L}(\mathbf{m}) = \exp\left(-\frac{\chi^2(\mathbf{m})}{2}\right) \rightarrow \alpha(\mathbf{m}) = \frac{\mathcal{L}(\mathbf{m})}{\mathcal{L}(\mathbf{m}-1)}$$

See Orford, Vassh+18 (PRL), Vassh+21 (ApJ),
Orford, Vassh+22 (PRC Letters), Vassh+22 (Frontiers in Phys.)



Movie by
N. Vassh

Black – solar abundance data

Grey – AME 2012 data

Red – values at current step

Blue – best step of entire run

Markov Chain Monte Carlo (MCMC):

Uses observational data to discern nuclear properties such as masses as well as constrain the conditions present at nucleosynthesis sites

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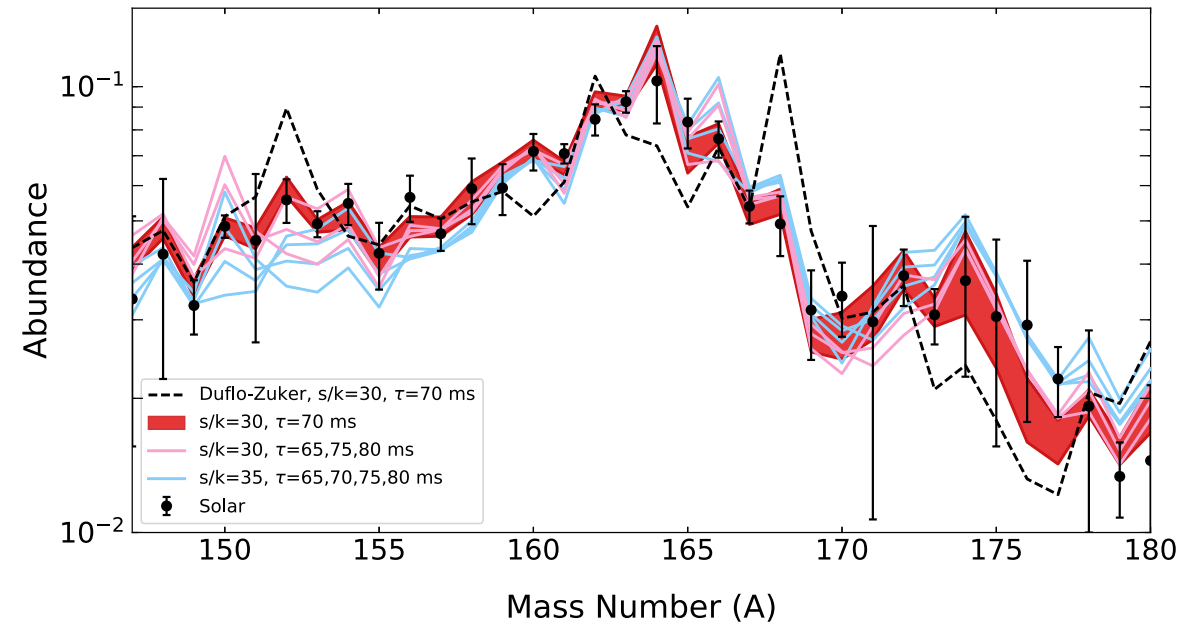
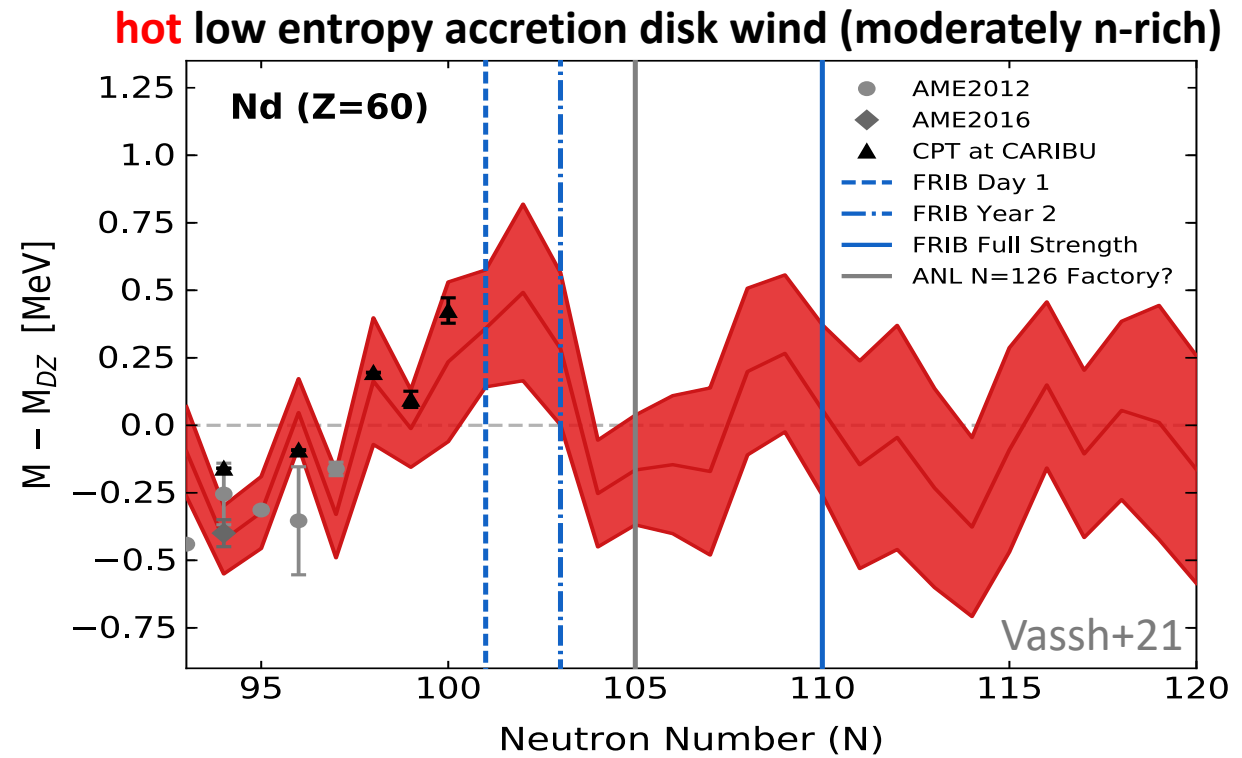
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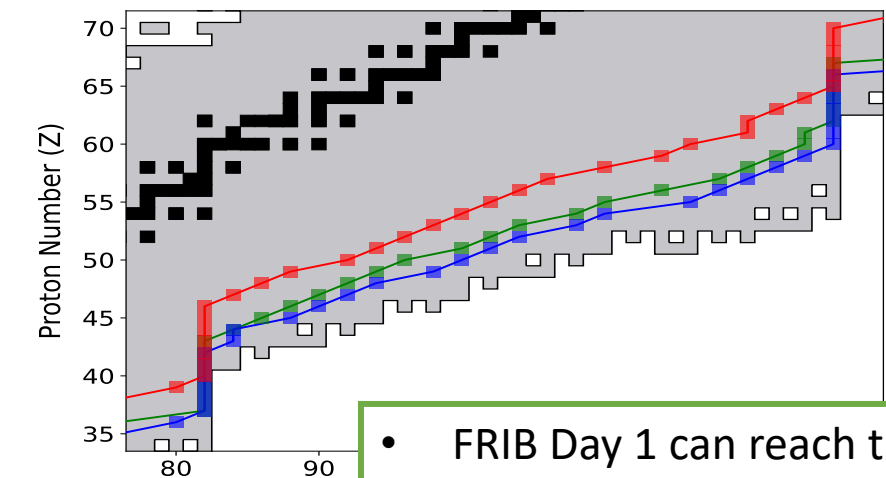
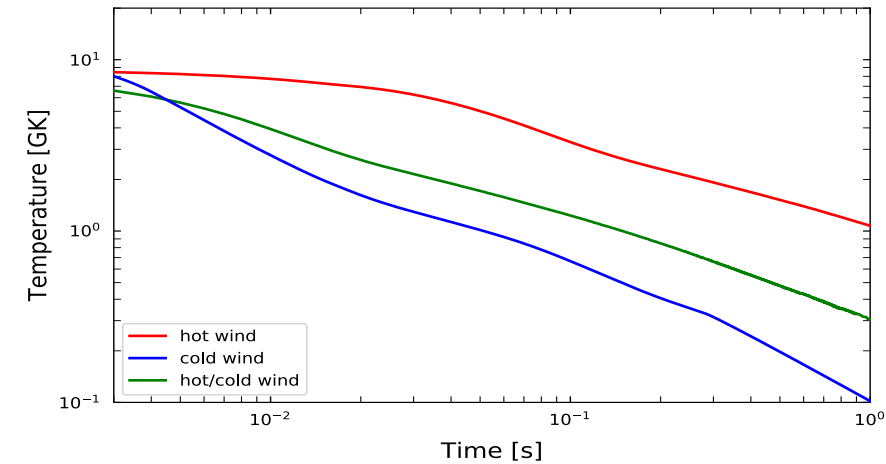
$$\mathcal{L}(\mathbf{m}) = \exp\left(-\frac{\chi^2(\mathbf{m})}{2}\right) \rightarrow \alpha(\mathbf{m}) = \frac{\mathcal{L}(\mathbf{m})}{\mathcal{L}(\mathbf{m}-1)}$$

See Orford, Vassh+18 (PRL), Vassh+21 (ApJ),
Orford, Vassh+22 (PRC Letters), Vassh+22 (Frontiers in Phys.)

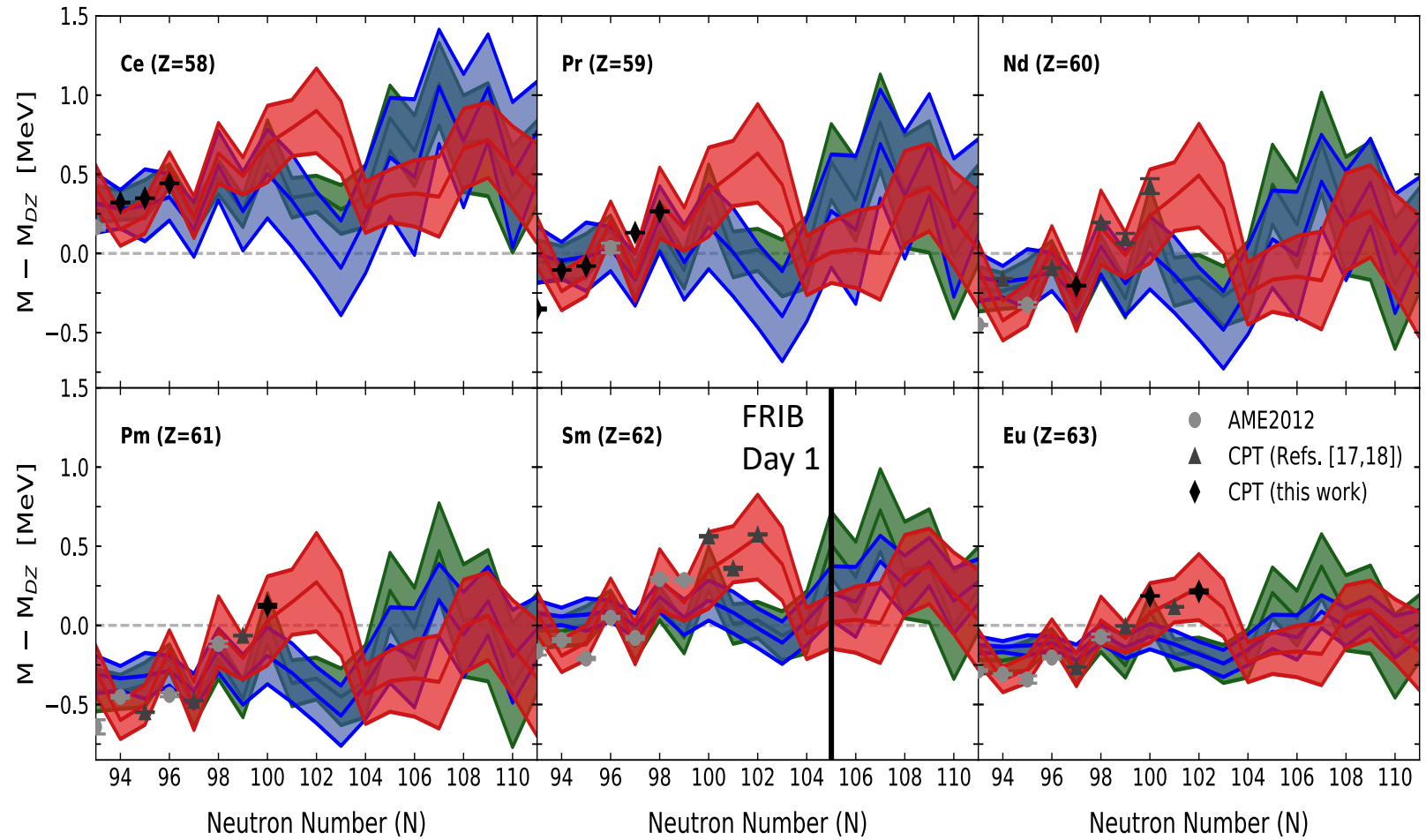


Ejecta Outflow Parameters

Outflow Type	Entropy (s/k_B)	Timescale (ms)	Y_e
Hot	30	70	0.2
Hot/cold	20	10	0.2
Cold	10	3	0.2



hot vs cold low entropy accretion disk winds (moderately n-rich)

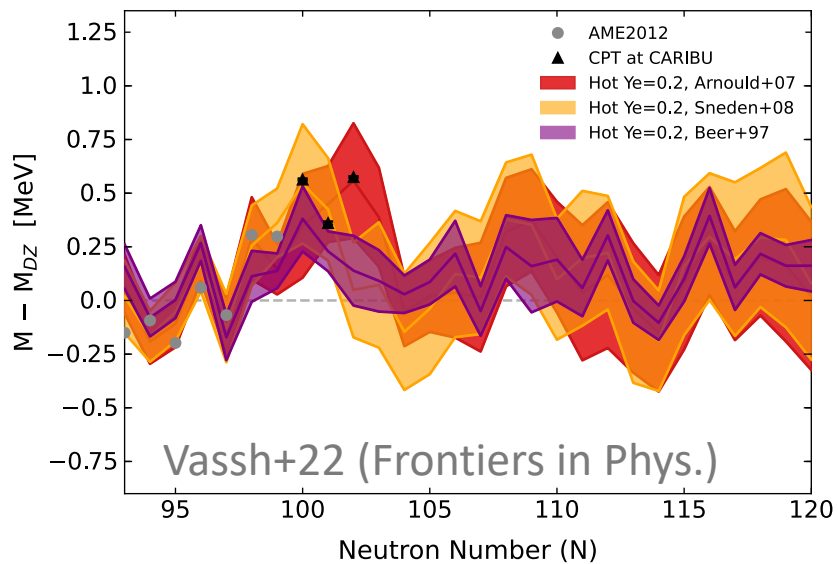
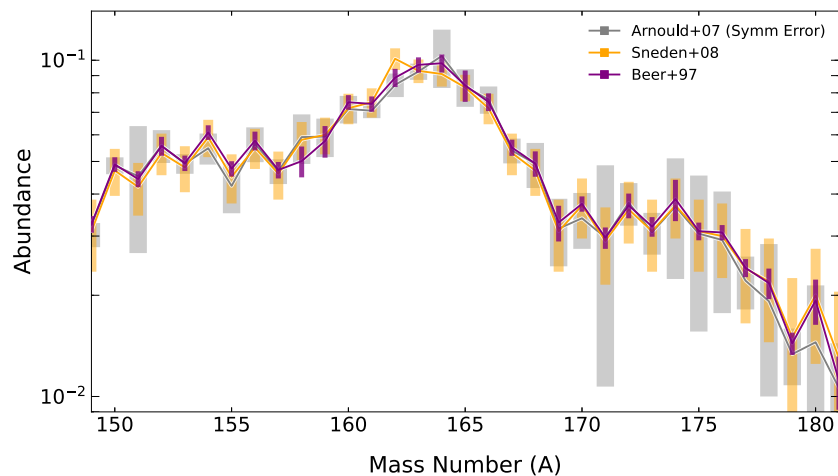


Orford,Vassh+22 (PRC Letters)

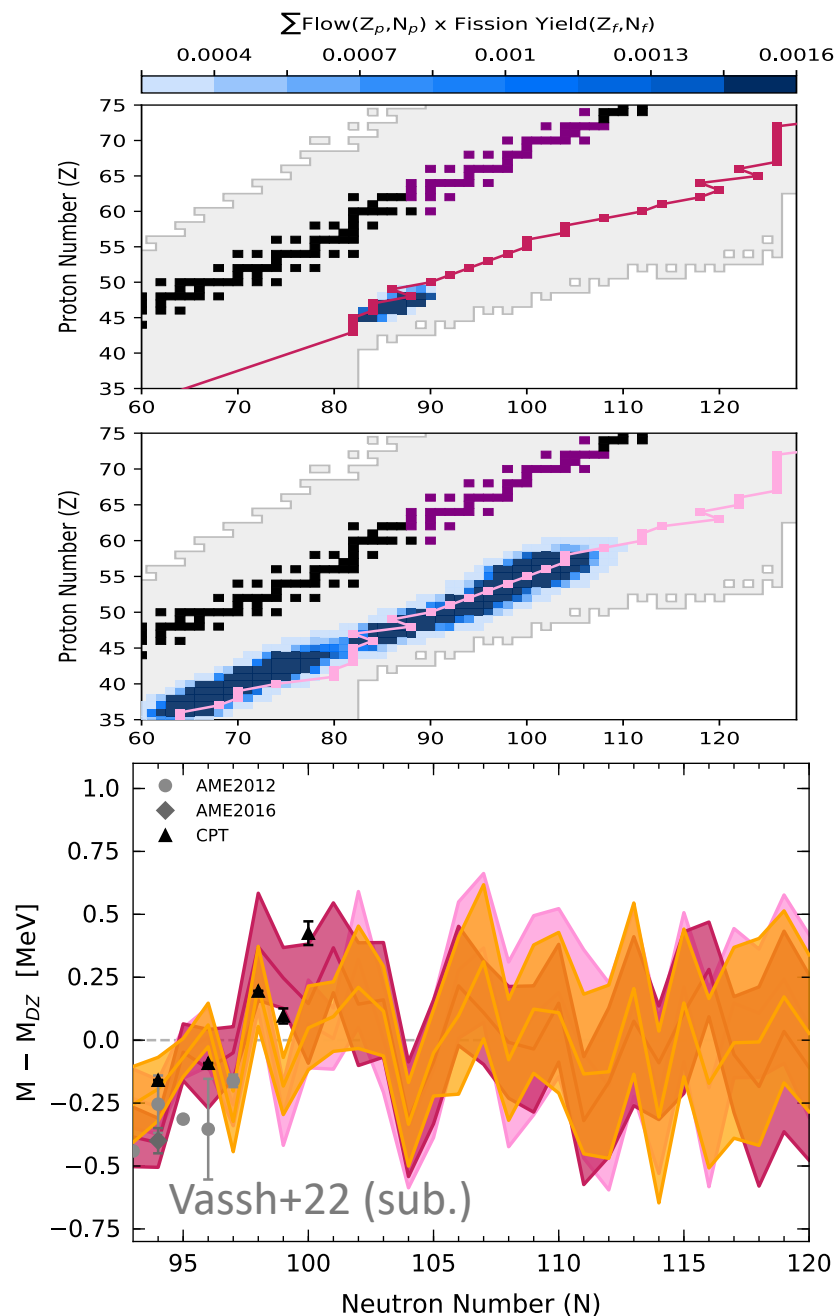
- FRIB Day 1 can reach the **N=104** feature forming the peak in **hot** conditions
- Future FRIB reach will cover the **N=108** and **N=106** features utilized with **cold** and **in between** dynamics

Recent work:

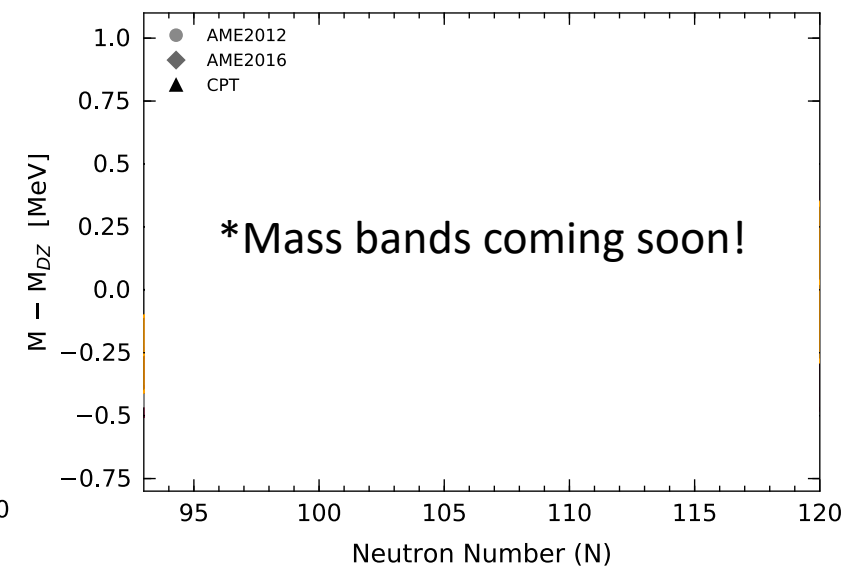
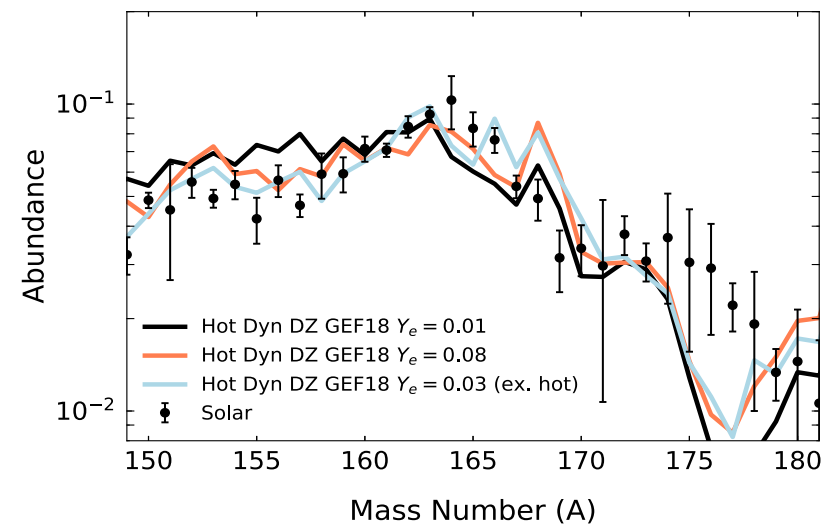
Solar abundance variations



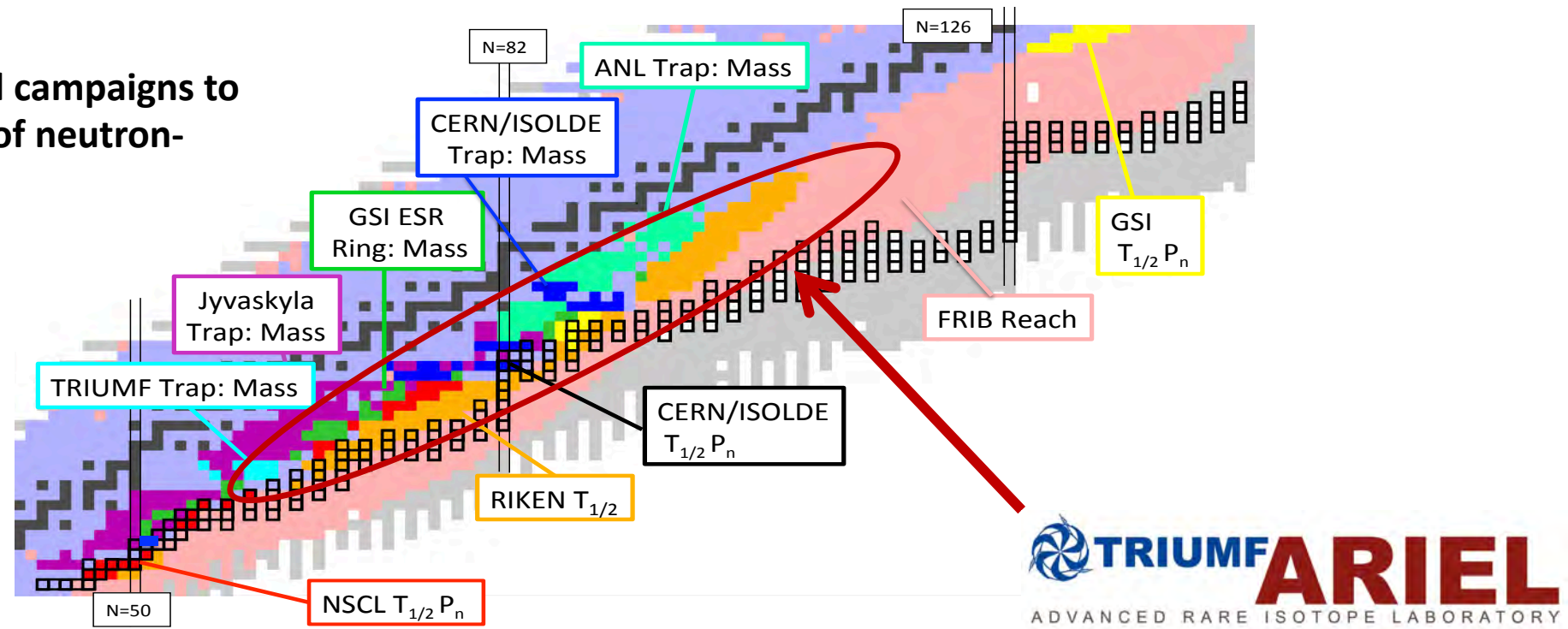
Very n-rich conditions with fission deposition



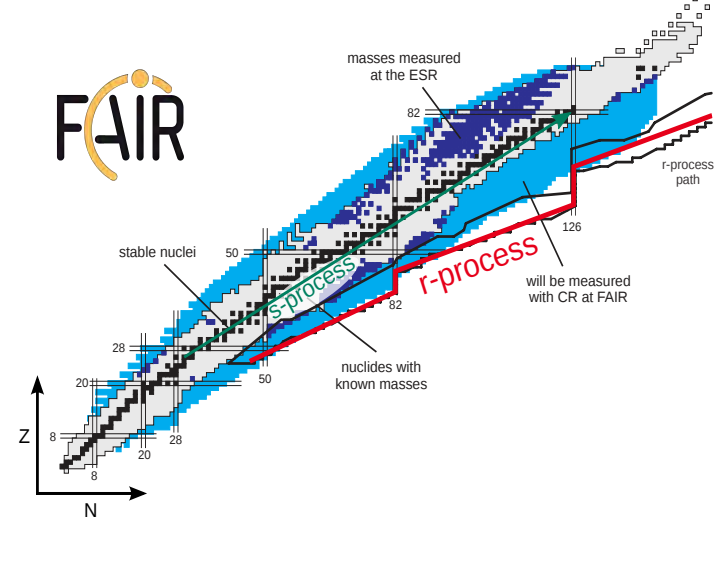
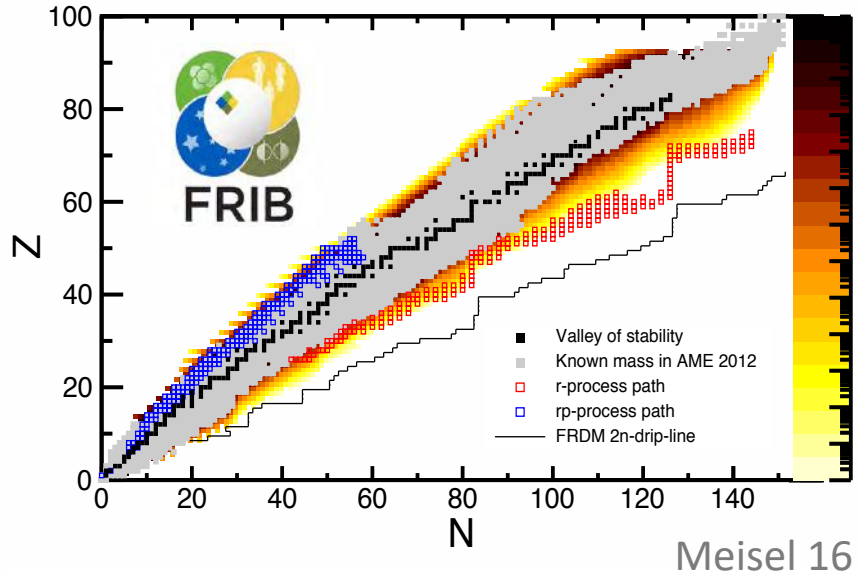
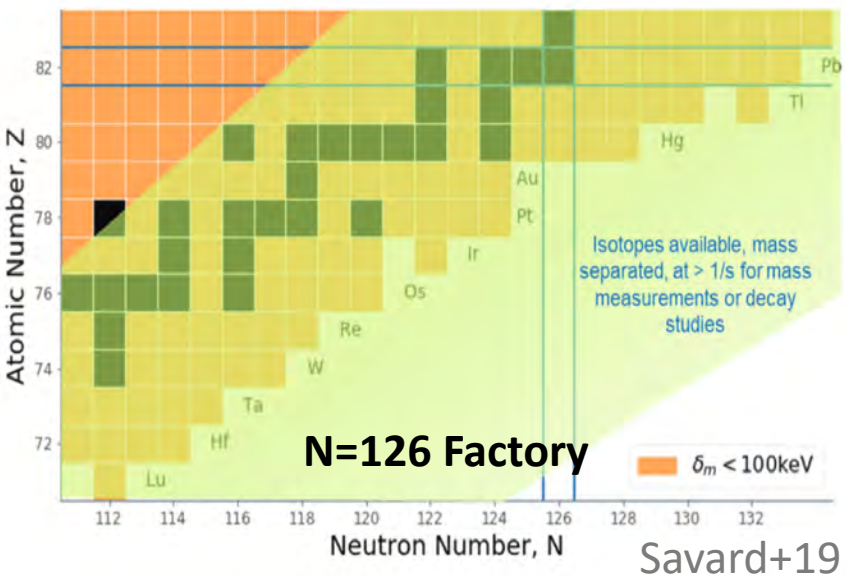
A wider variety of n-richness



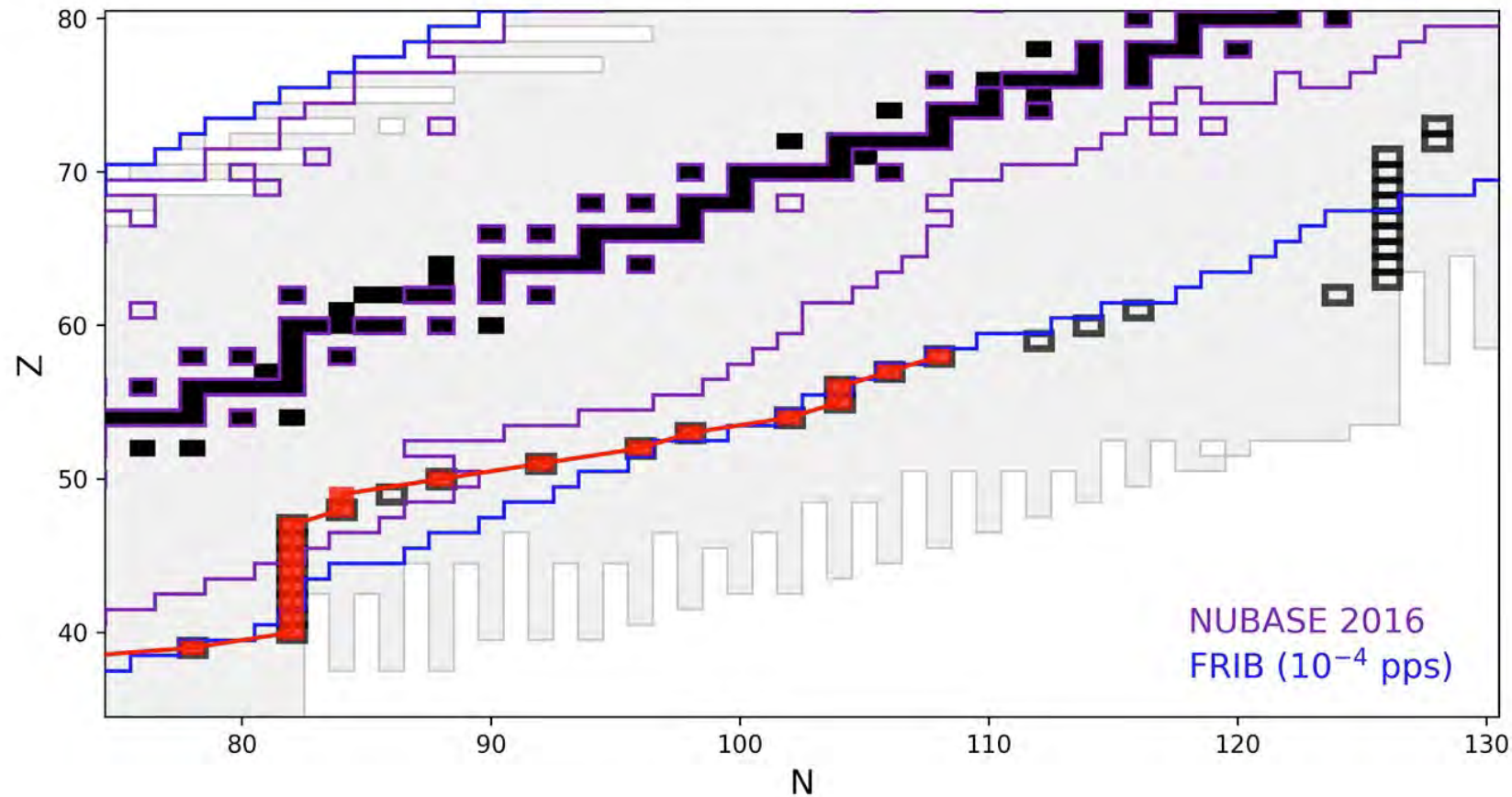
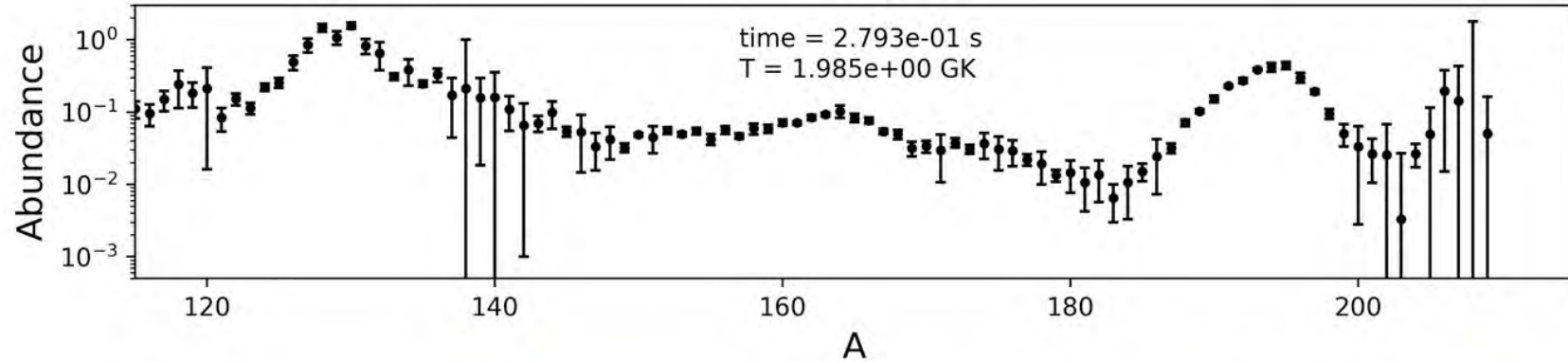
Worldwide experimental campaigns to measure the properties of neutron-rich nuclei:
 masses, half-lives,
 reaction rates...



Horowitz+18



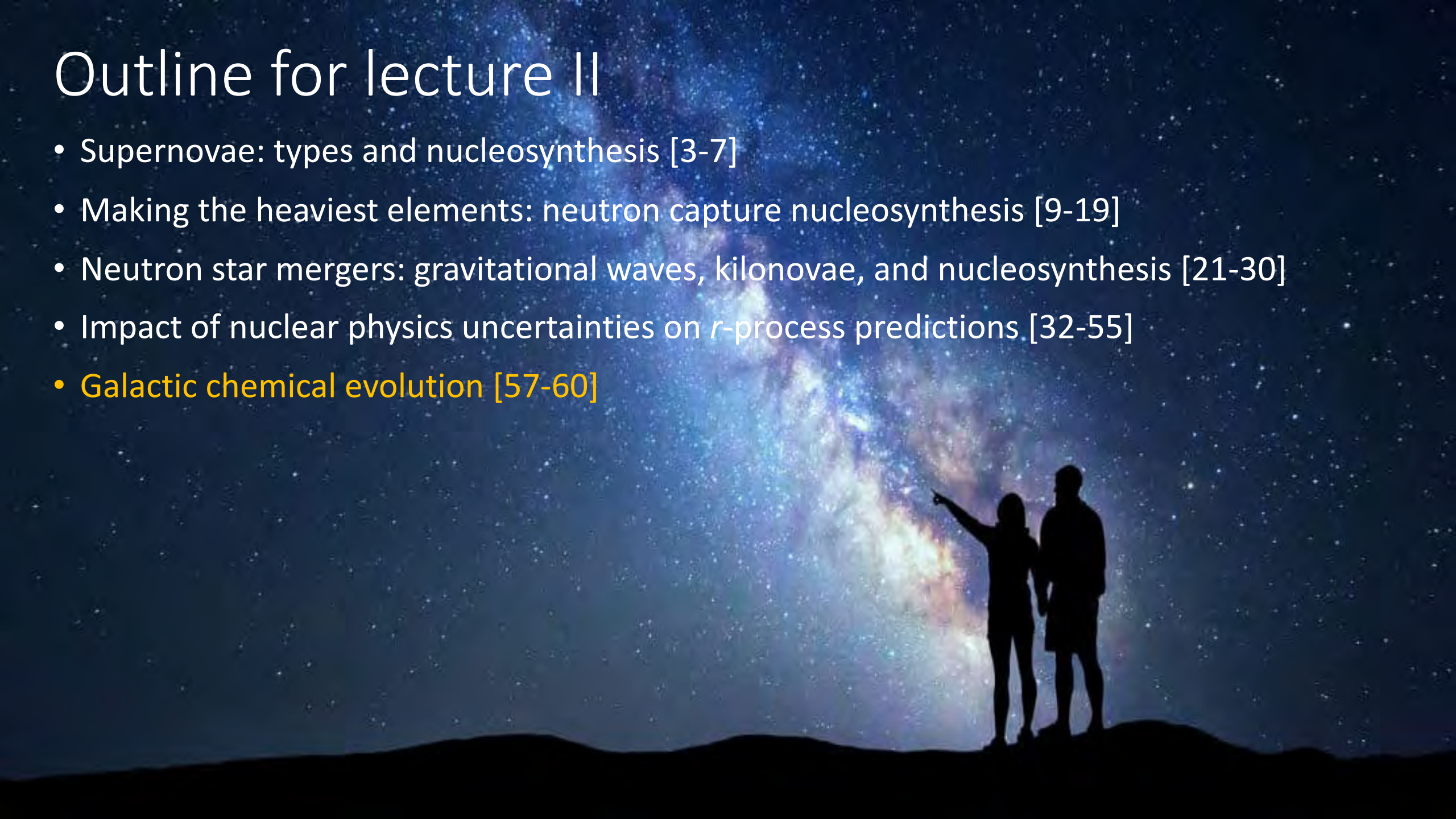
Future experiment meets the r -process path



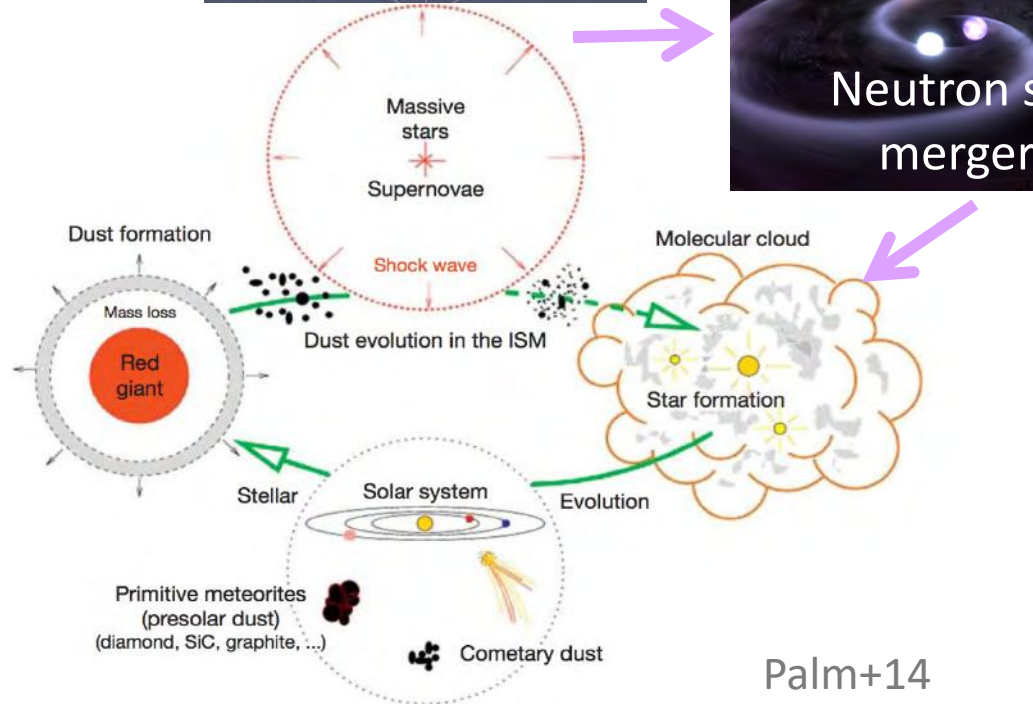
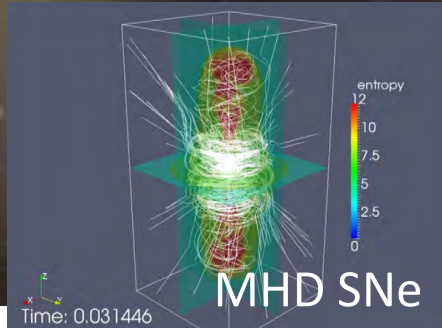
*reach of future experiment in key regions impacting the evolution of abundances (note moderately n-rich conditions used here)

Outline for lecture II

- Supernovae: types and nucleosynthesis [3-7]
- Making the heaviest elements: neutron capture nucleosynthesis [9-19]
- Neutron star mergers: gravitational waves, kilonovae, and nucleosynthesis [21-30]
- Impact of nuclear physics uncertainties on *r*-process predictions [32-55]
- Galactic chemical evolution [57-60]



Where *and when* were the heavy elements we see in stars produced?

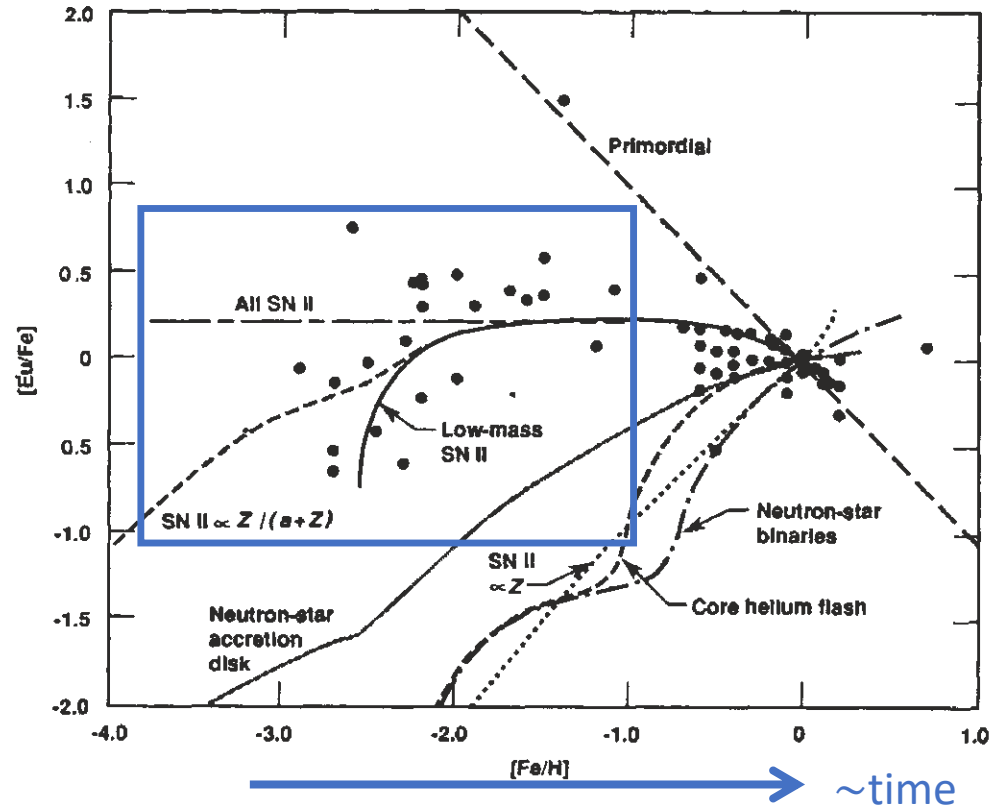


For our Milky Way (present day):

- CCSNe ~ 2 per century
- SNe Ia ~ 0.4 per century
- RAWDs ~ 0.05 per century
- NSMs ~ 0.004 per century

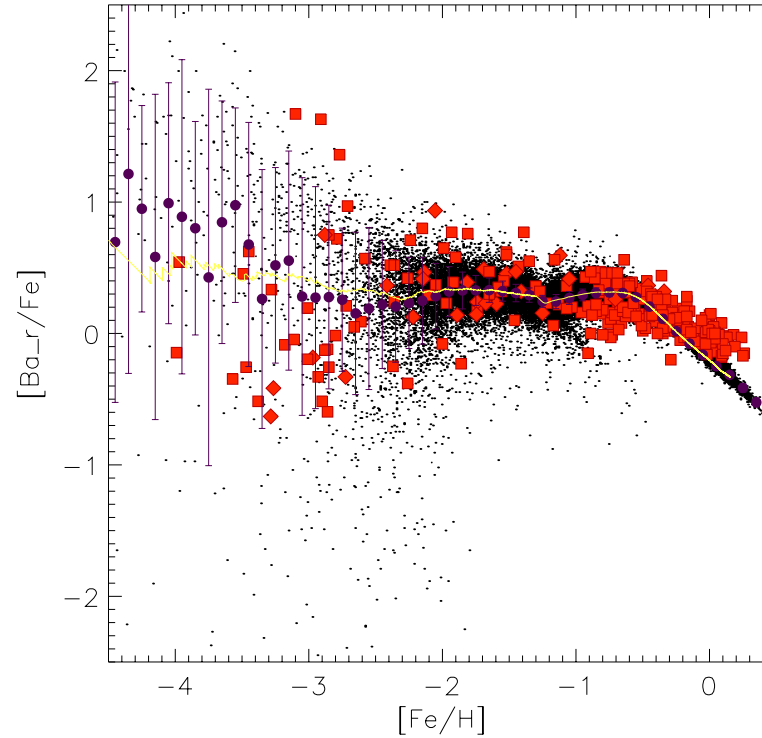
Supernovae as the *r*-process source?

Galactic chemical evolution (GCE) and low metallicity stars

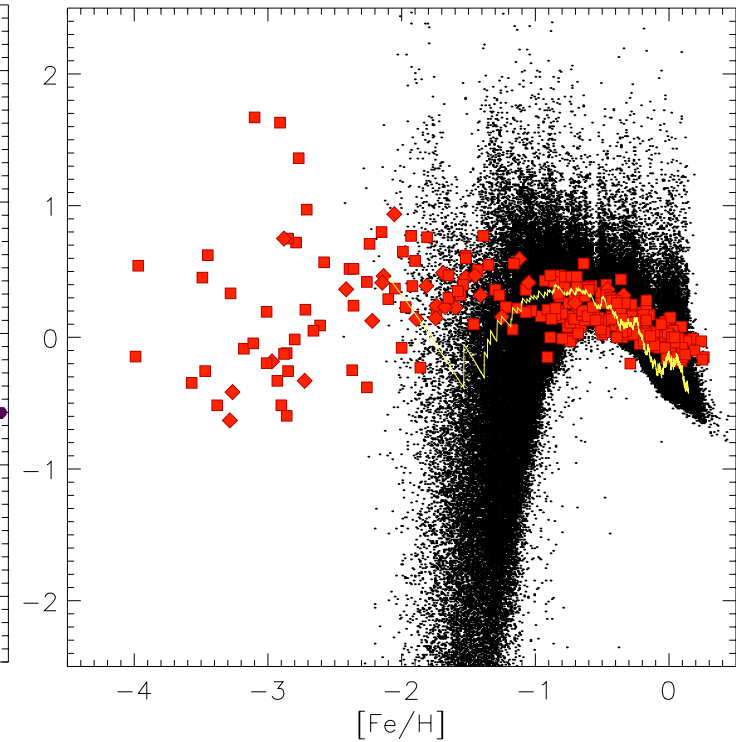


*The stars in the box seem to be more consistent with supernovae since neutron stars take time to merge

Type II Supernova



Neutron Star Mergers

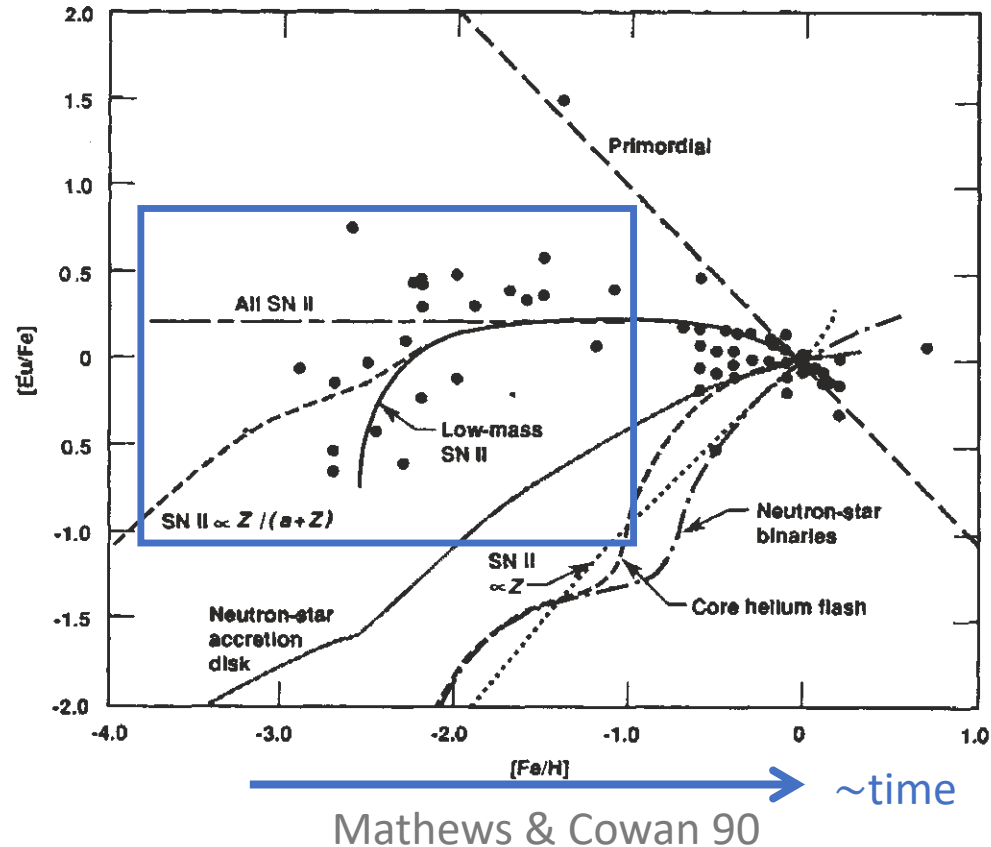


Argast+04

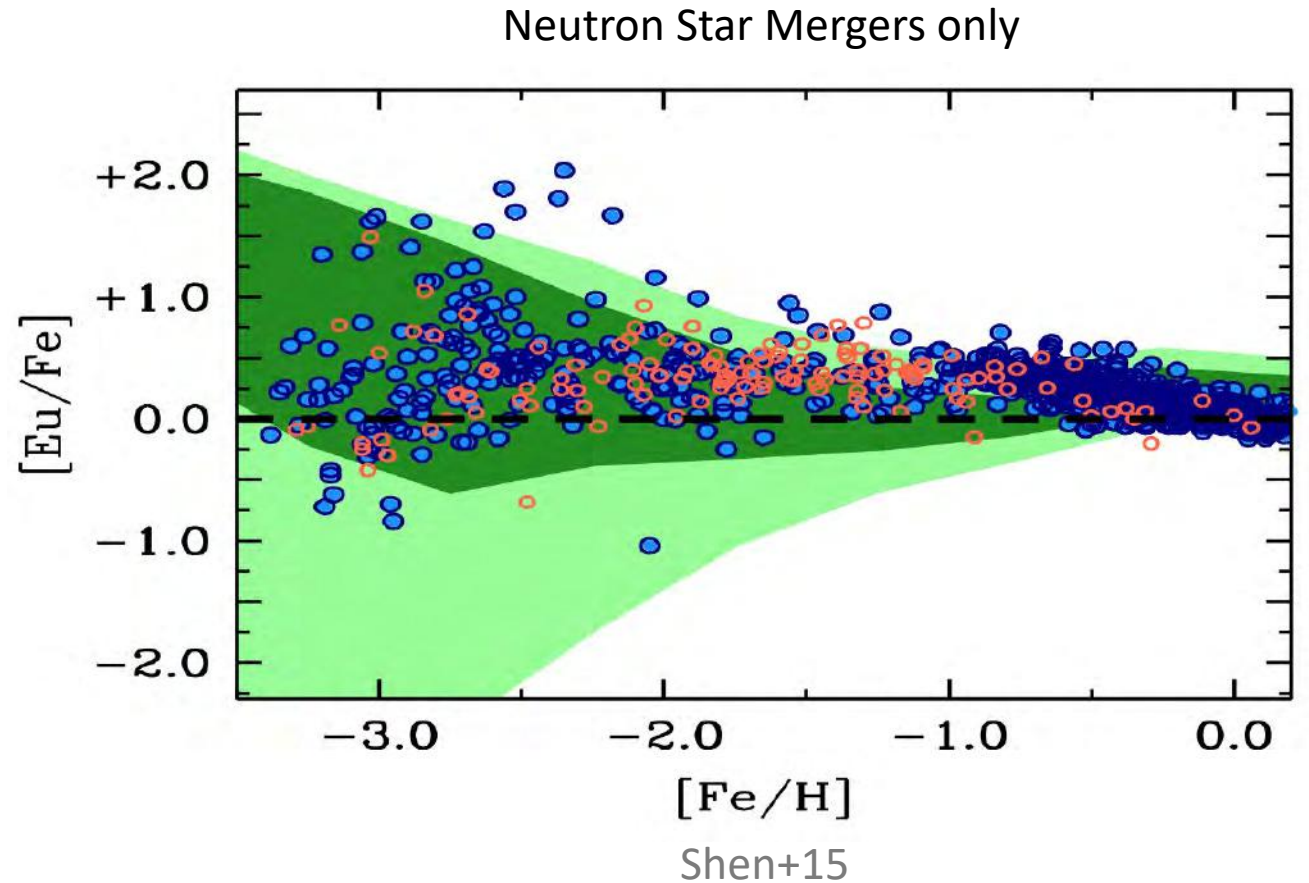
◆ ■ observations
● model stars
~ average ISM abundances

Supernovae as the r -process source?

Galactic chemical evolution (GCE) and low metallicity stars

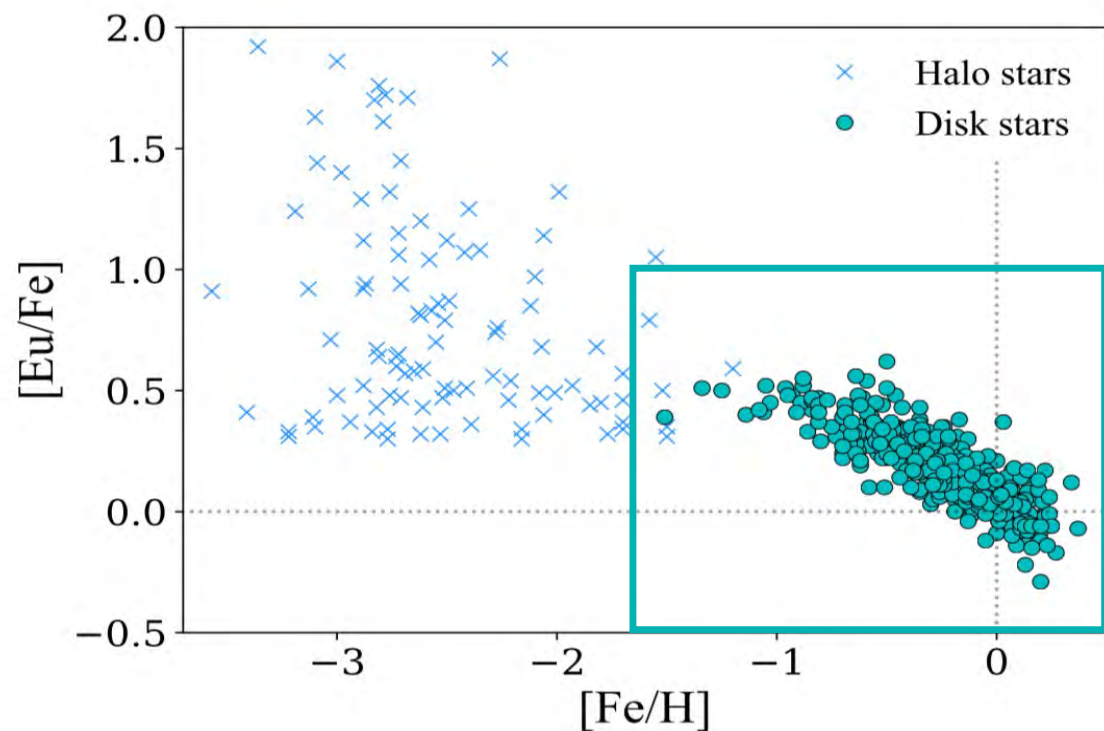


*The stars in the box seem to be more consistent with supernovae since neutron stars take time to merge



Hydrodynamic mixing accounting for inhomogeneities in the interstellar medium could explain how r -process elements find their way to low metallicity regions

Could NSMs be the only r -process source?
 Consider $[\text{Eu}/\text{Fe}]$ again but now stars in
 the Galactic disk



Eu production rate must reach equilibrium before onset
 of SNe Ia in order to reproduce $[\text{Eu}/\text{Fe}]$ of disk stars

NSM with delay times $\sim t^{-1}$ don't reproduce this behavior:
 earlier sources?

