

The Standard Cosmology: 10 things to know, triumphs & tests, Hubble troubles, and big questions

Michael S. Turner
UCLA and UChicago
16/17 July 2023



N3AS Summer School in
**Multi-Messenger
Astrophysics**

Intended for • advanced graduate students and beginning postdoctoral researchers interested in nuclear and particle astrophysics — theory, experiment, or observation.

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Cosmology and the Early Universe

Michael Turner, Univ. Chicago and UCLA..... The Standard Cosmology
Francois Lanusse, CEA Paris-Saclay, CNRS..... Deep Learning and Observational Cosmology
Garth Illingworth, UC Santa Cruz..... JWST: Searching for the First Galaxies

Dark Matter

Graciela Gelmini, UCLA..... Dark Matter: Theory and Laboratory Phenomenology
Ben Safdi, UC Berkeley..... Dark Matter in Astrophysics

Neutron Stars, Supernovae, Mergers, and Nucleosynthesis

Dany Page, Univ. Nacional Autónoma de México.. Neutron Stars: Structure, Evolution and Cooling
David Radice, Penn State Univ. Explosive Astrophysics: Mergers and Supernovae
Nicole Vassh, TRIUMF..... Nucleosynthesis: Connecting Nuclear Properties and Observations

Multi-Messenger Astrophysics

Glennys Farrar, New York Univ..... UHE Cosmic Rays and Multi-Messenger Astrophysics
Joshua Smith, CalState Fullerton Gravitational Wave Astronomy
Susanne Mertens, Tech. Univ. Munich..... Neutrino Properties: Masses and Mixing
George Fuller, UC San Diego..... Neutrino Astrophysics

Sponsors

The Standard Cosmology: 10 things to know, triumphs & tests, Hubble troubles, and big questions (lecture 1)

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The Standard Cosmology (Λ CDM) in plain English

... very-early accelerated expansion driven by the potential energy of a scalar field gives rise to a very-large, smooth, spatially flat patch that becomes all that we can see today. Quantum fluctuations during this inflationary phase grow into the seeds for galaxies. The conversion of potential field energy into heat produces the quark soup that evolves a baryon asymmetry and long-lived dark matter particles. The excess of quarks over antiquarks becomes neutrons and protons, later some light elements and finally atoms. The gravity of the dark matter particles drives the formation of structure from galaxies to superclusters and a mere 5 billion years ago the repulsive gravity of dark energy (Λ) again drove accelerated expansion ...

1. Not a model – a real, falsifiable theory that makes precise, testable predictions

- Matter tell space how to curve:

$$G_{\mu\nu} = 8\pi T_{\mu\nu}$$

- Space tells matter how to move/evolve:

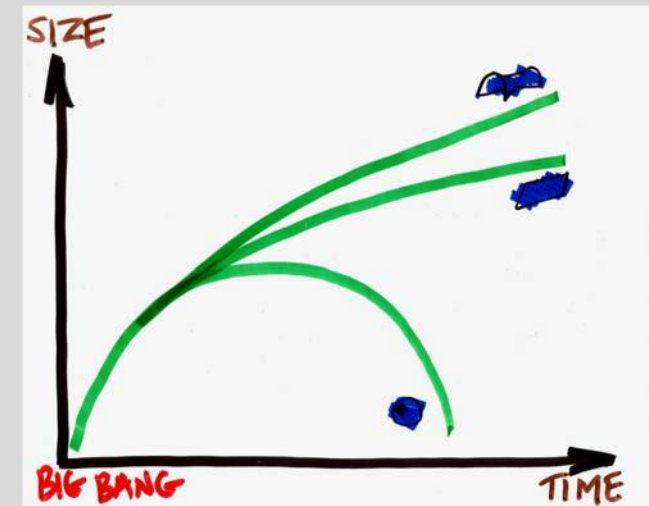
$$T^{\mu\nu}_{;\nu} = 0$$

- With the assumption of spatial flatness and isotropy/homogeneity (supported by much evidence) only the composition of the Universe needs to be fixed. Today,
 - Matter (baryons, dark matter, neutrinos): $31\% \pm 0.6\%$ (Baryons, $4.8 \pm 0.06\%$)
 - Dark Energy (Λ): $69\% \pm 0.6\%$
 - Photons (CMB): 0.005%

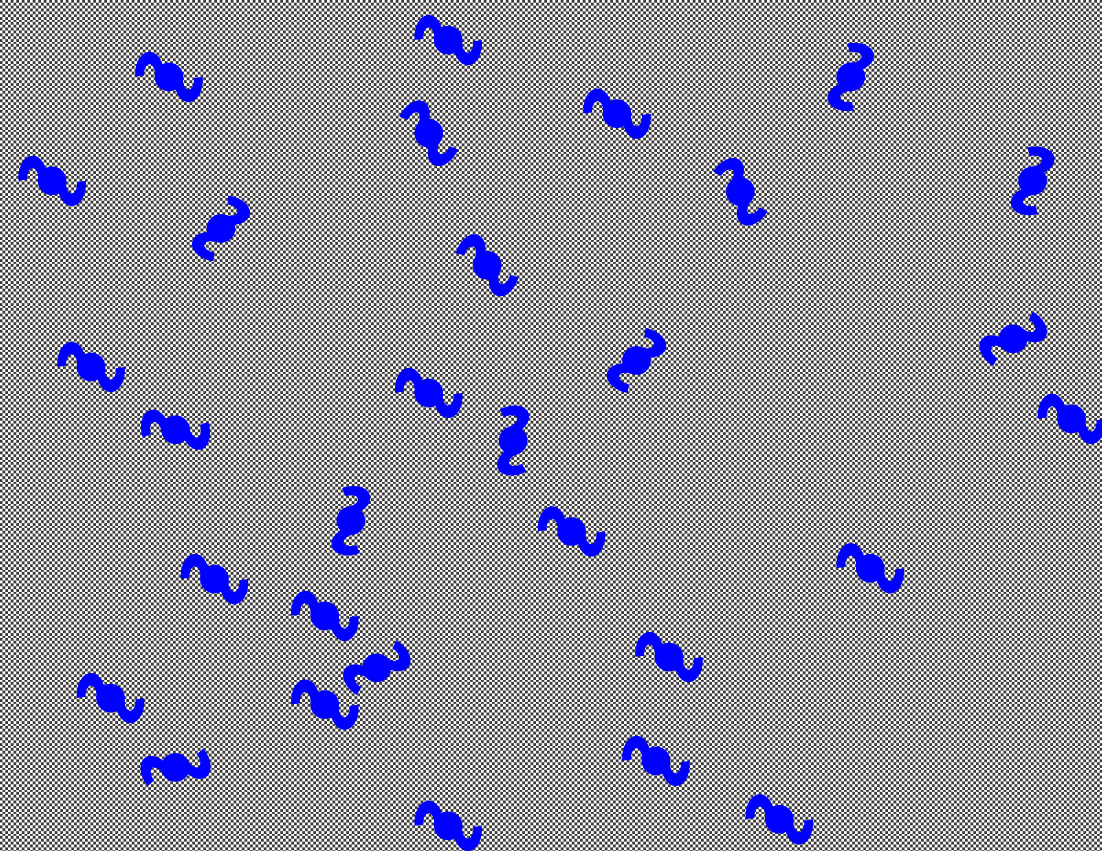
2. Expansion of space: kinematics

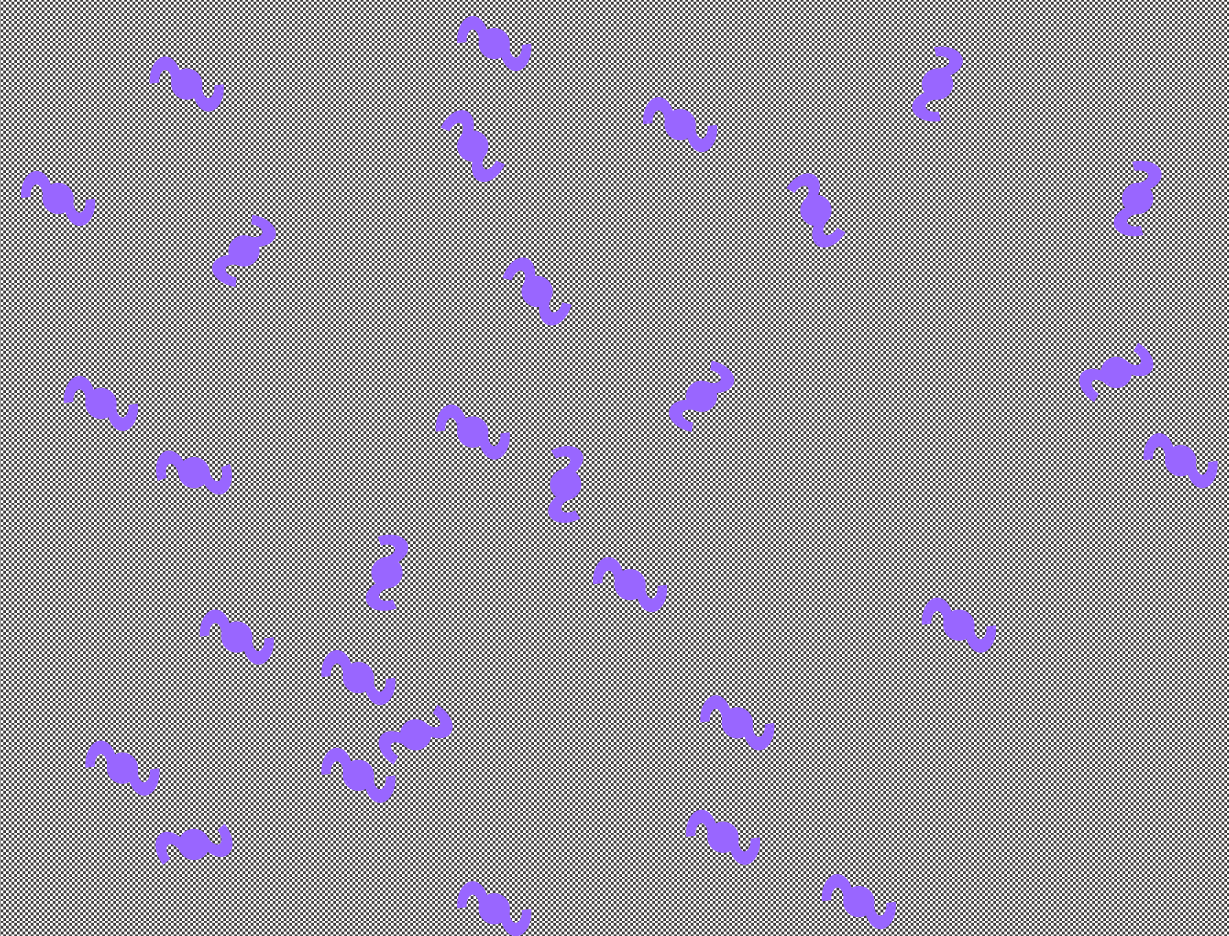
- Spacetime metric:
- One dof: the cosmic scale factor $a(t)$ or $R(t)$
- Expansion of Universe is a scale up
- Big bang is an explosion of space
- Most important number in cosmology Hubble parameter ($\sim 1/\text{age}$)

$$ds^2 = dt^2 - a(t)^2[dr^2 + r^2 d\Omega]$$

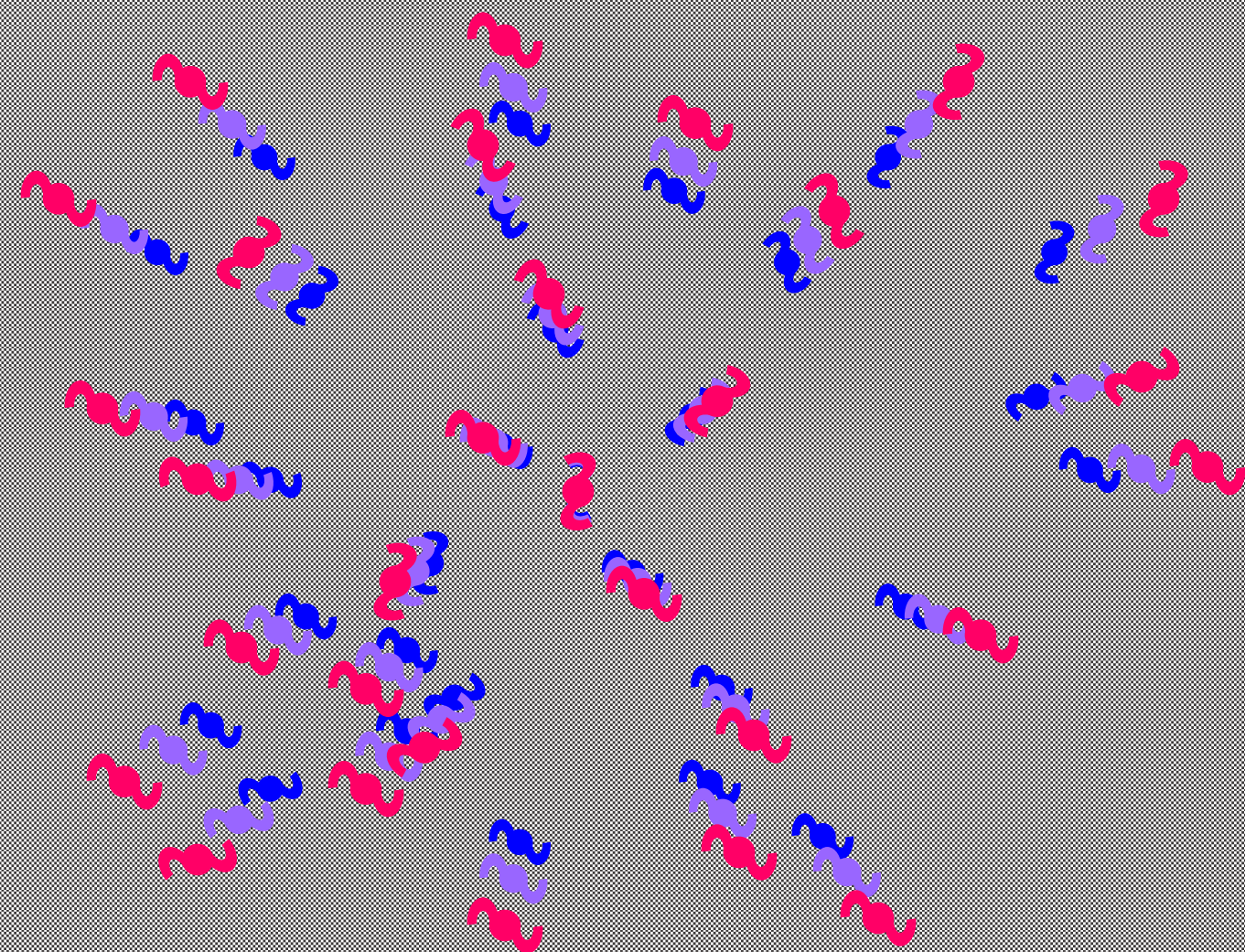


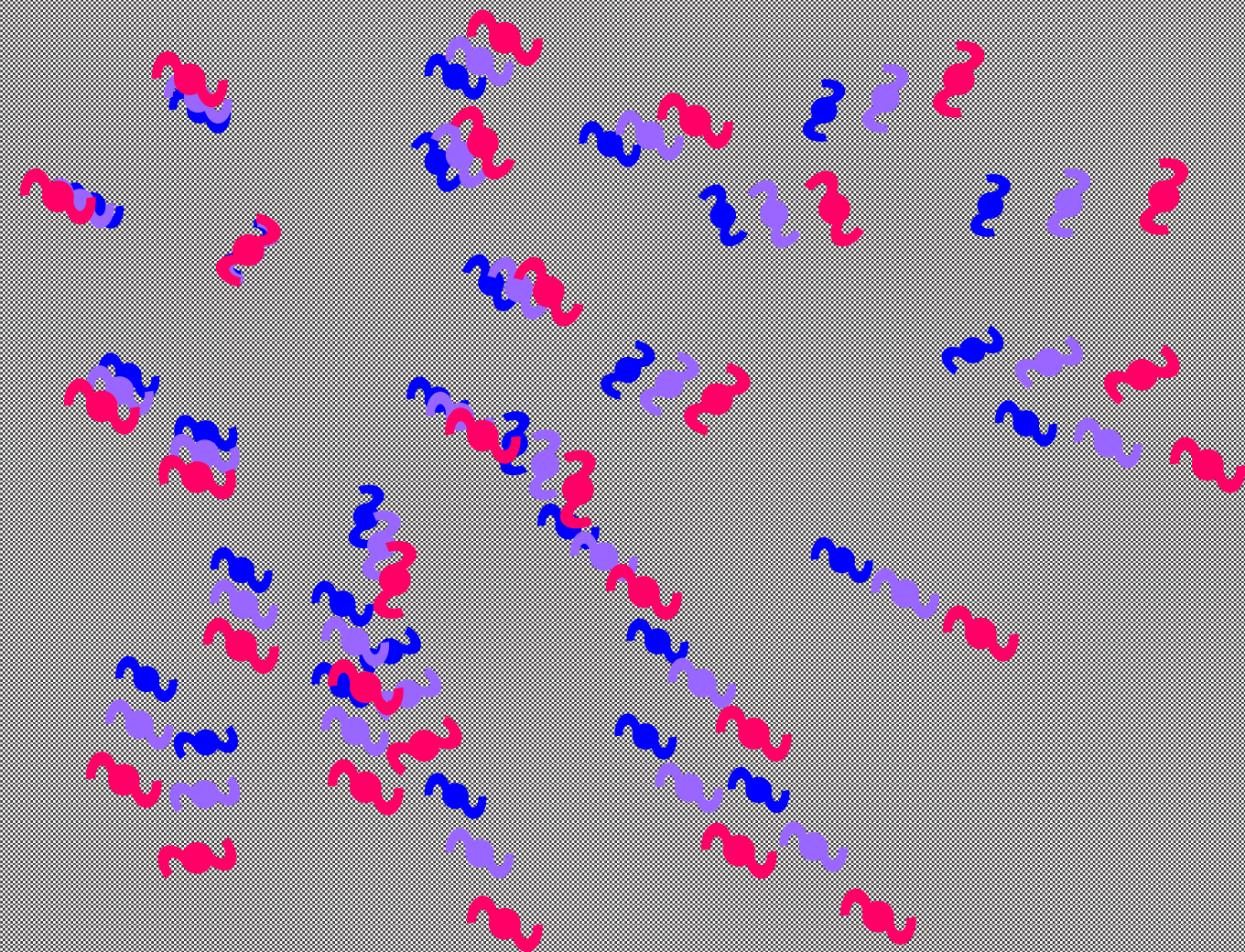
$$H \equiv \frac{\dot{a}}{a}$$

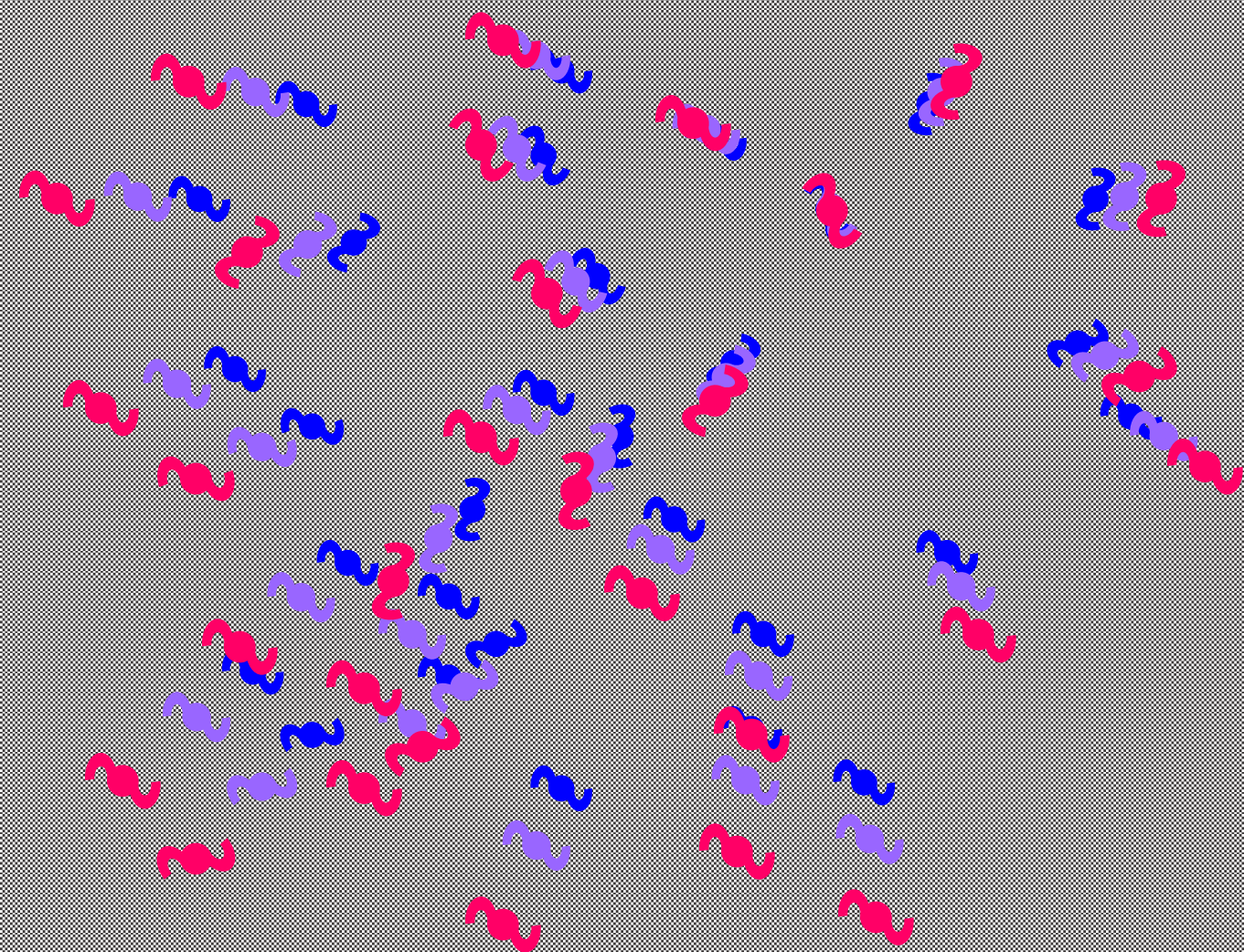






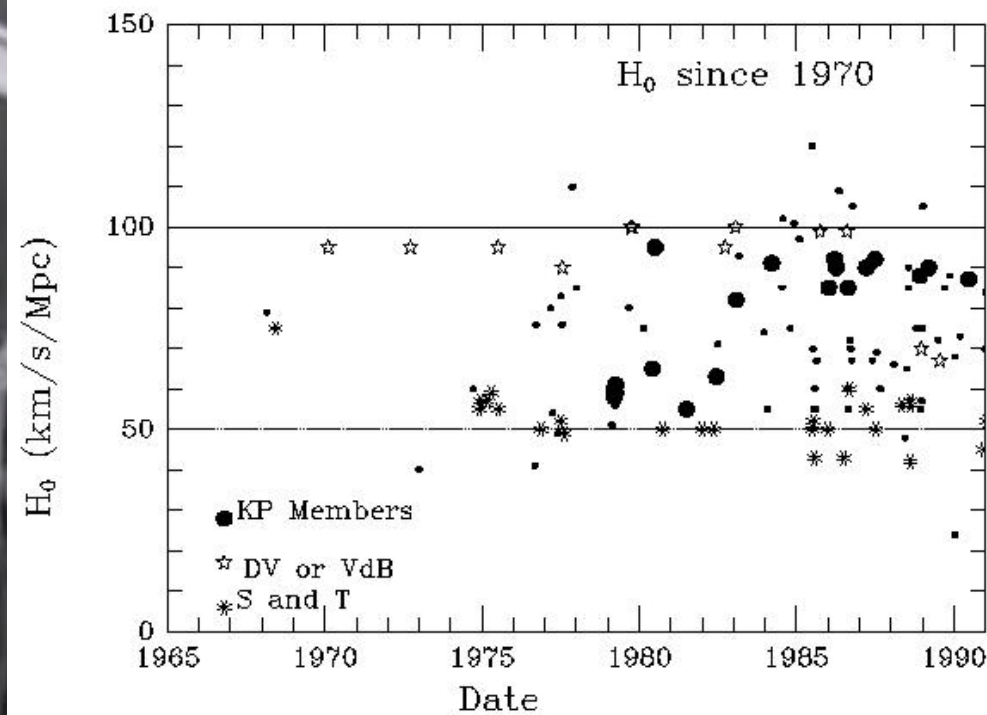
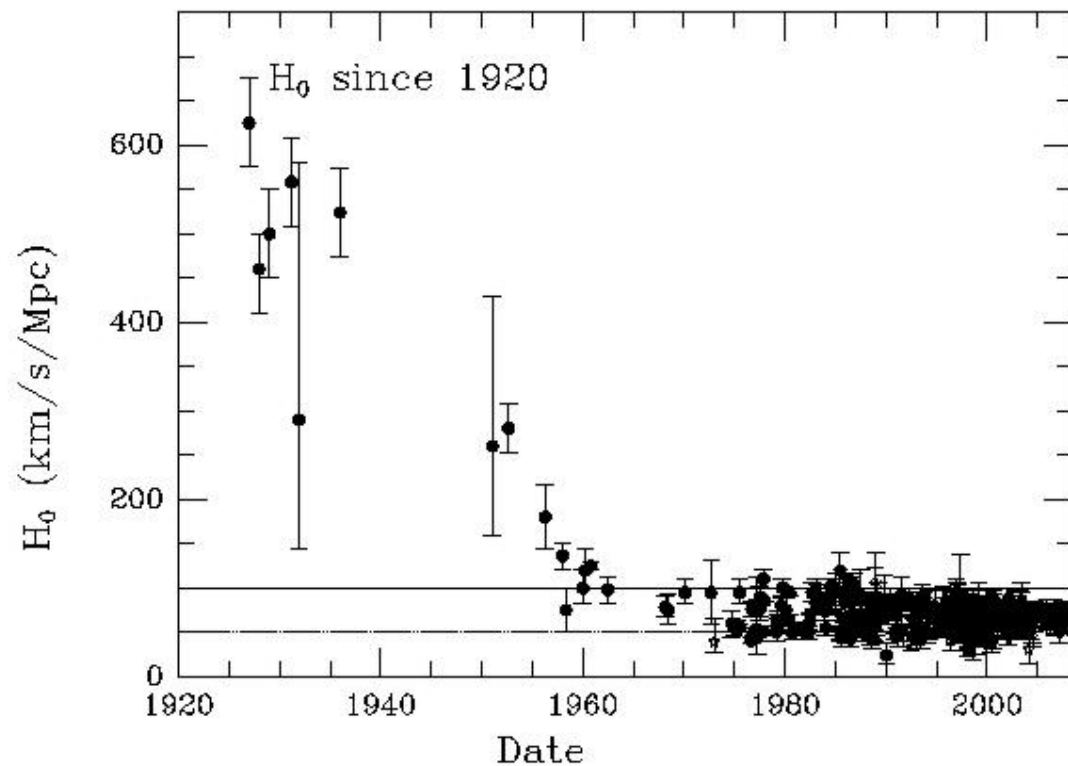






Just like Harvard, everyone is at the
center of their Universe
But really, NO CENTER

Hubble troubles, part 1



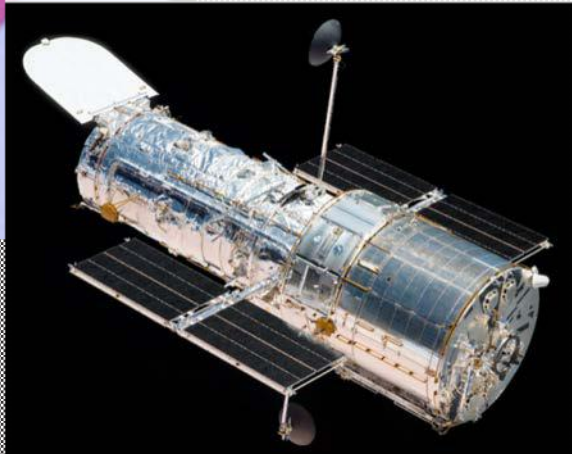
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The problem with
mirrors

Distance is the
hardest thing to
measure in
cosmology



2001: $H_0 = 72 \pm 2 \pm 6$ km/s/Mpc



Final Results from the *Hubble Space Telescope* Key Project to Measure the Hubble Constant*

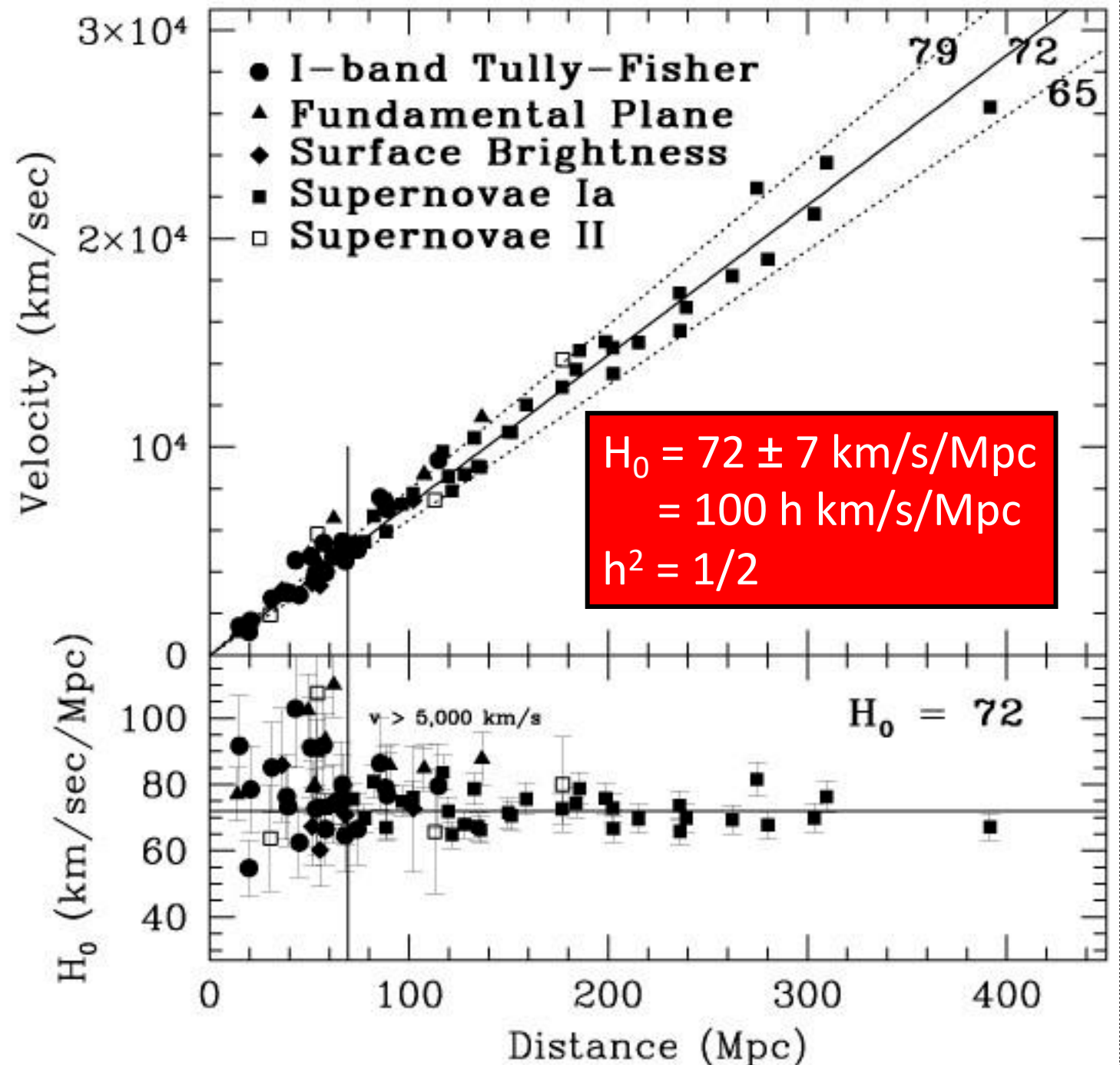
Wendy L. Freedman¹, Barry F. Madore^{1,2}, Brad K. Gibson³, Laura Ferrarese⁴, Daniel D. Kelson⁵, Shoko Sakai⁶, Jeremy R. Mould⁷, Robert C. Kennicutt, Jr.⁸, Holland C. Ford⁹, John A. Graham⁵

[+ Show full author list](#)

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[The Astrophysical Journal](#), Volume 553, Number 1

Citation Wendy L. Freedman et al 2001 ApJ 553 47



2. More kinematics of the expanding Universe

- Redshift z (IMPORTANT)

$$1 + z \equiv \frac{\lambda_{\text{rcvd}}}{\lambda_{\text{rest}}} = \frac{a_{\text{today}}}{a_{\text{emit}}} = 1/a$$

Redshift is the observable!

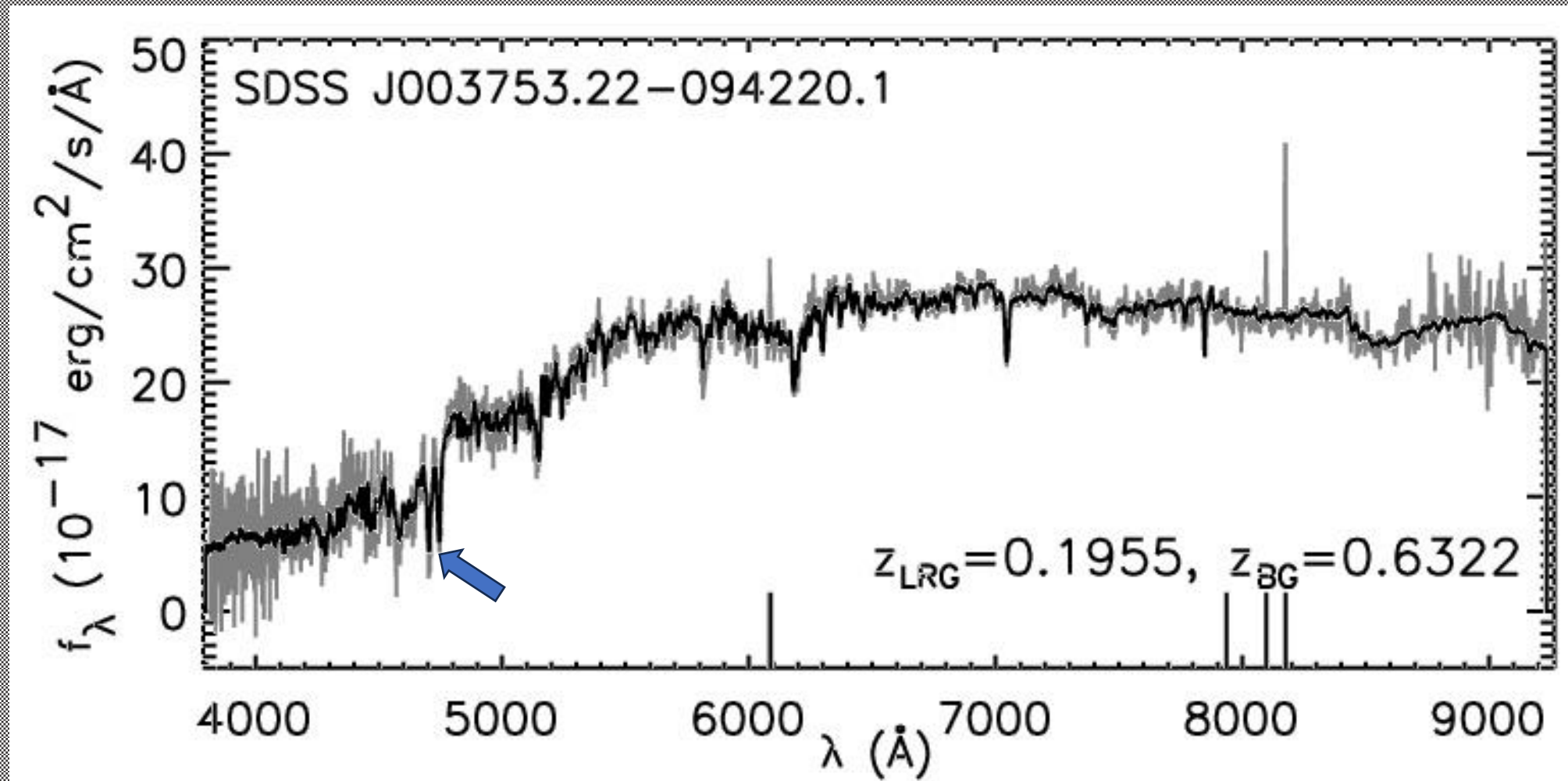
Redshift of H, K lines of Ca

Lab: 3934, 3969 Å

Observed: ~4750 Å

$1 + z = \lambda_{\text{rcvd}} / \lambda_{\text{rest}} = 4750 / 3950$

$z = 0.20$



Big redshifts: UV \rightarrow IR

$$\text{Ly-}\alpha = 1216 \text{ \AA}$$

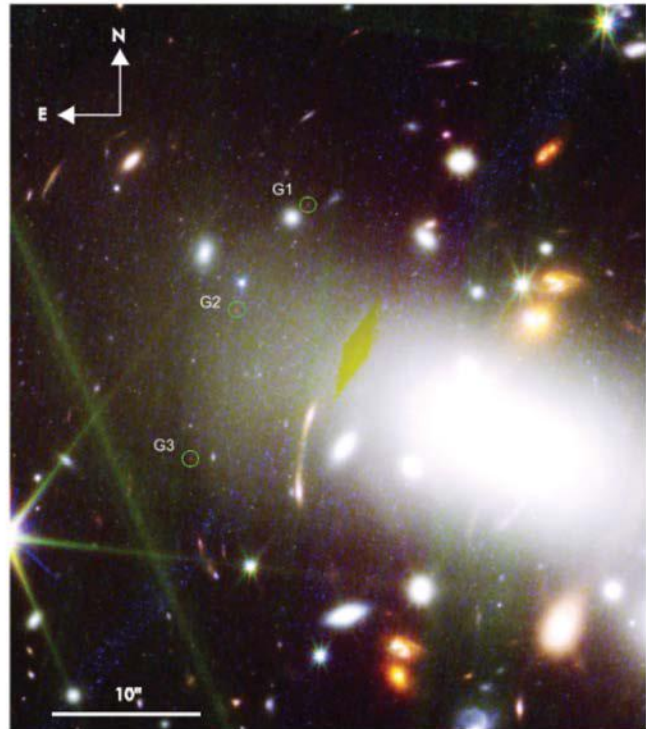


Fig. 1. Color-composite image of part of RX J2129. JWST NIRC2 + HST ACS color-composite image of galaxy cluster RX J2129, with three images of the $z = 9.51$ galaxy circled in green. We obtained spectroscopy of image G2. Filters were assigned to RGB colors as red, JWST F277W+F356W+F444W; green, JWST F115W+F150W+F200W; and blue, HST F606W + F814W. The broad blue and green bands are diffraction spikes caused by foreground stars. The yellow diamond is an artifact caused by a chip gap in the HST ACS camera. The individual red, green, and blue images are shown in Figs. S11 to S13.

JWST redshift 9.51 galaxy

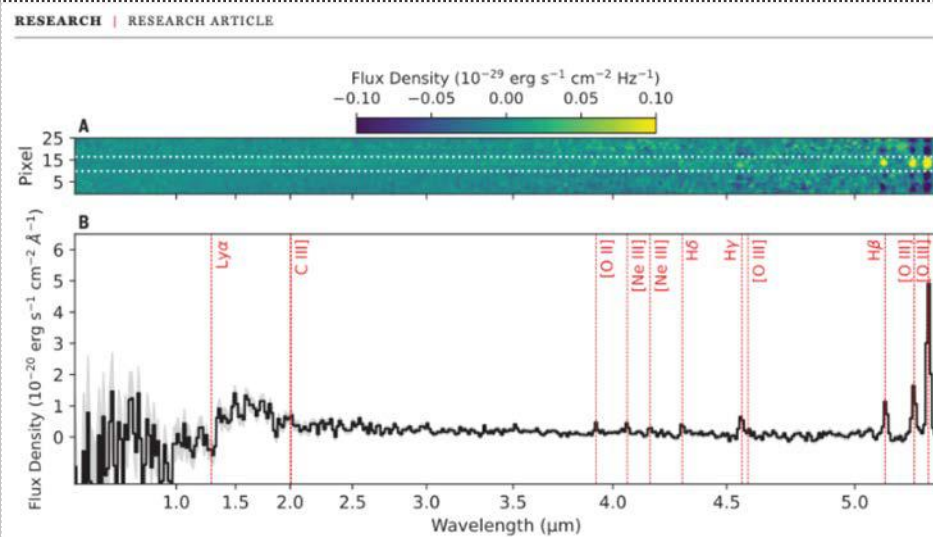
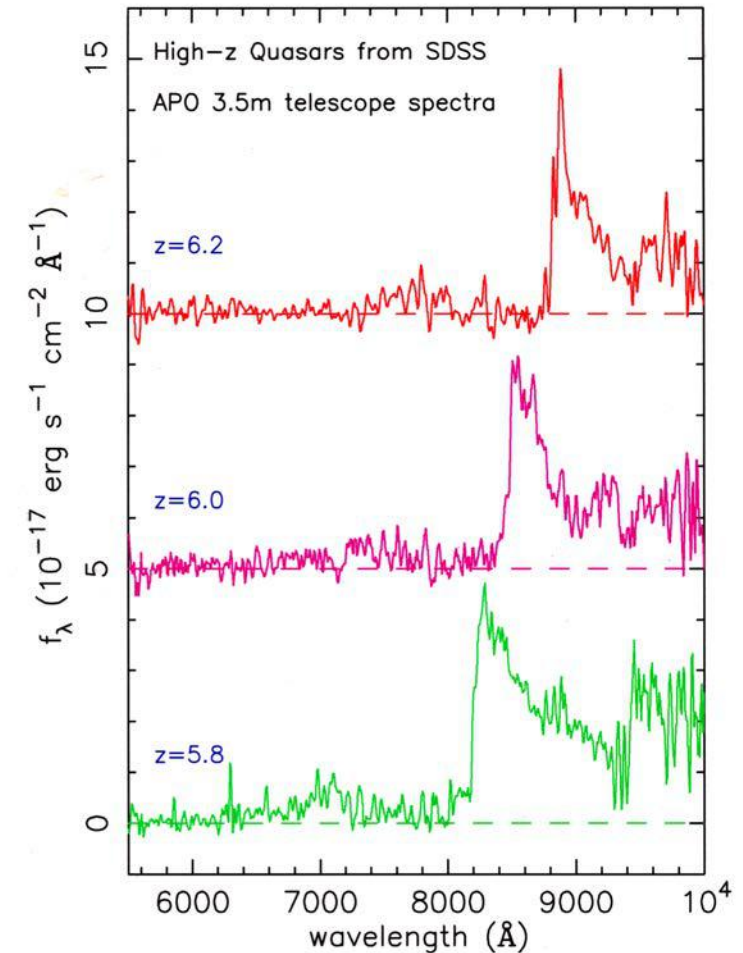


Fig. 2. Observed JWST spectrum of image G2. NIRSpec prism spectrum of image G2 of the $z = 9.51$ galaxy. This spectrum has not been corrected for magnification from gravitational lensing. (A) Two-dimensional spectrum, with flux densities indicated by the color bar. The apparent negative fluxes, in the background near the emission lines, are artifacts produced by the dither pattern used for the NIRSpec observations. The white dotted lines indicate the window used to extract the spectrum in (B). (B) One-dimensional spectrum. The black line is the data, with gray shading indicating its 1σ uncertainties. Red vertical lines indicate the expected wavelengths of emission lines for $z = 9.51$.

SDSS High- z quasars





A Luminous Quasar at Redshift 7.642

Feige Wang^{1,15}, Jinyi Yang^{1,16}, Xiaohui Fan¹, Joseph F. Hennawi^{2,3}, Aaron J. Barth⁴, Eduardo Banados³,
Fuyan Bian⁵, Konstantina Boutsia⁶, Thomas Connor⁷, Frederick B. Davies^{3,8}, Roberto Decarli⁹,
Anna-Christina Eilers^{10,15}, Emanuele Paolo Farina¹¹, Richard Green¹, Linhua Jiang¹², Jiang-Tao Li¹³,
Chiara Mazzucchelli⁵, Riccardo Nanni², Jan-Torge Schindler³, Bram Venemans³, Fabian Walter³,
Xue-Bing Wu^{12,14}, and Minghao Yue¹

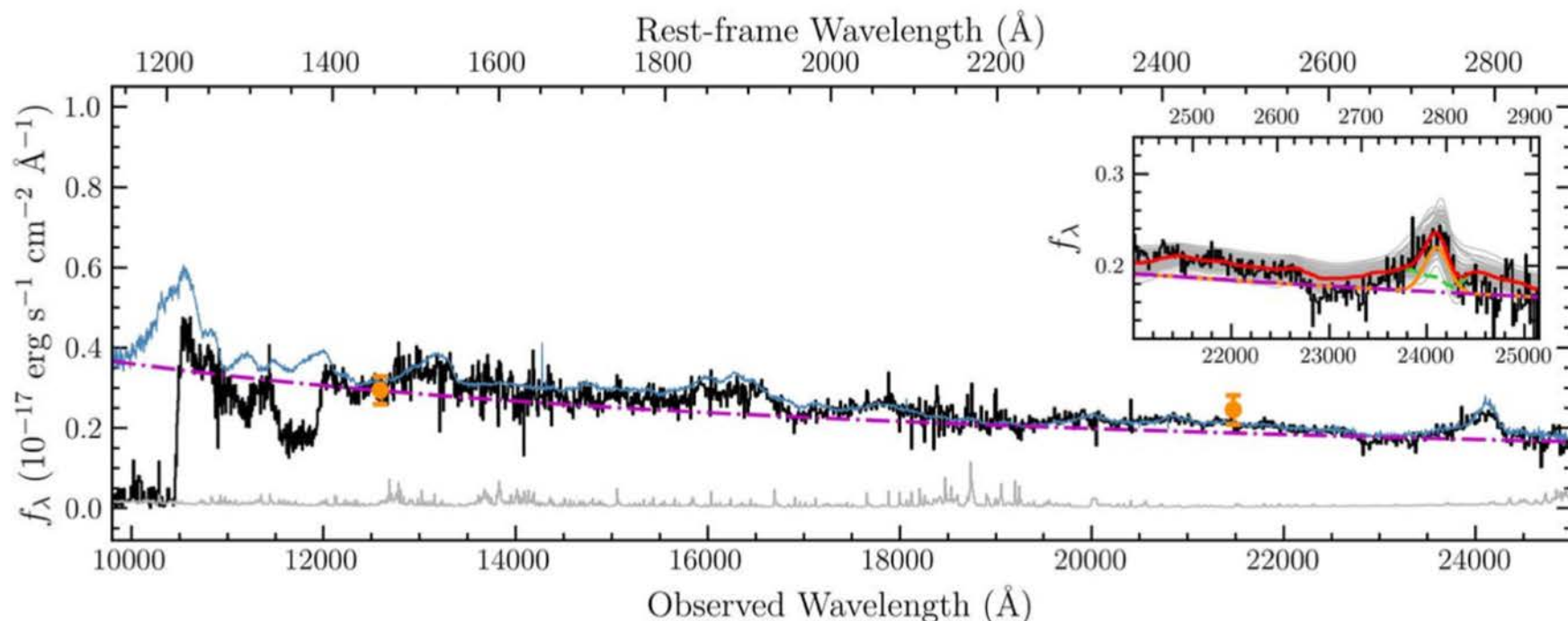


Figure 1. Top panel: image cutouts ($20'' \times 20''$, north is up and east is to the left) for J0313–1806 in PS1 z , PS1 y , DELS z , VISTA J , VISTA K_s , WISE W1, and WISE W2 bands. The photometry is given in Table 1. Bottom panel: the final stacked spectrum of J0313–1806. In the figure, we re-binned the spectrum by two spectral pixels ($\sim 173 \text{ km s}^{-1}$) for illustration purposes. The black and gray lines represent the Galactic extinction-corrected spectrum and the error vector, respectively. The blue line denotes the quasar composite spectrum constructed with Sloan Digital Sky Survey (SDSS) $z \sim 2$ quasars having similar C IV blueshifts and line strengths. The purple dashed line denotes the power-law continuum. The orange points are flux densities determined from photometry in the J - and K_s -bands. The inset panel shows the Mg II line fitting with the purple dotted-dashed line denoting the power-law continuum, the green dashed line denoting the pseudo-continuum model (the sum of power-law continuum, Fe II emission, and Balmer continuum), the orange line representing the Gaussian fitting of the Mg II line and the red line representing the total fit of pseudo-continuum and Mg II line. The thin gray lines in the inset panel represent the spectral fitting of 100 mock spectra as described in Section 3.

Just can't get enough of those high redshift galaxies: $z = 10$ to 13

Spectroscopic confirmation of four metal-poor galaxies at $z=10.3-13.2$

Emma Curtis-Lake^{1*}, Stefano Carniani^{2†}, Alex Cameron³, Stephane Charlot⁴, Peter Jakobsen^{5,6}, Roberto Maiolino^{7,8,9}, Andrew Bunker³, Joris Witstok^{7,8}, Renske Smit¹⁰, Jacopo Chevallard³, Chris Willott¹¹, Pierre Ferruit¹², Santiago Arribas¹³, Nina Bonaventura^{5,6}, Mirko Curti^{7,8}, Francesco D'Eugenio^{7,8}, Marijn Franx¹⁴, Giovanna Giardino¹⁵, Tobias J. Looser^{7,8}, Nora Lützgendorf¹⁶, Michael V. Maseda¹⁷, Tim Rawle¹⁶, Hans-Walter Rix¹⁸, Bruno Rodríguez del Pino¹³, Hannah Übler^{7,8}, Marco Sirianni¹⁶, Alan Dressler¹⁹, Eiichi Egami²⁰, Daniel J. Eisenstein²¹, Ryan Endsley²², Kevin Hainline²⁰, Ryan Hausen²³, Benjamin D. Johnson²¹, Marcia Rieke²⁰, Brant Robertson²⁴, Irene Shvaei²⁰, Daniel P. Stark²⁰, Sandro Tacchella^{7,8}, Christina C. Williams²⁵, Christopher N. A. Willmer²⁰, Rachana Bhatawdekar²⁶, Rebecca Bowler²⁷, Kristan Boyett^{28,29}, Zuyi Chen²⁰, Anna de Graaff¹⁸, Jakob M. Helton²⁰, Raphael E. Hviding²⁰, Gareth C. Jones³, Nimisha Kumari³⁰, Jianwei Lyu²⁰, Erica Nelson³¹, Michele Perna¹³, Lester Sandles^{7,8}, Aayush Saxena^{3,9}, Katherine A. Suess^{24,32}, Fengwu Sun²⁰, Michael W. Topping²⁰, Imaan E. B. Wallace³ and Lily Whitler²⁰

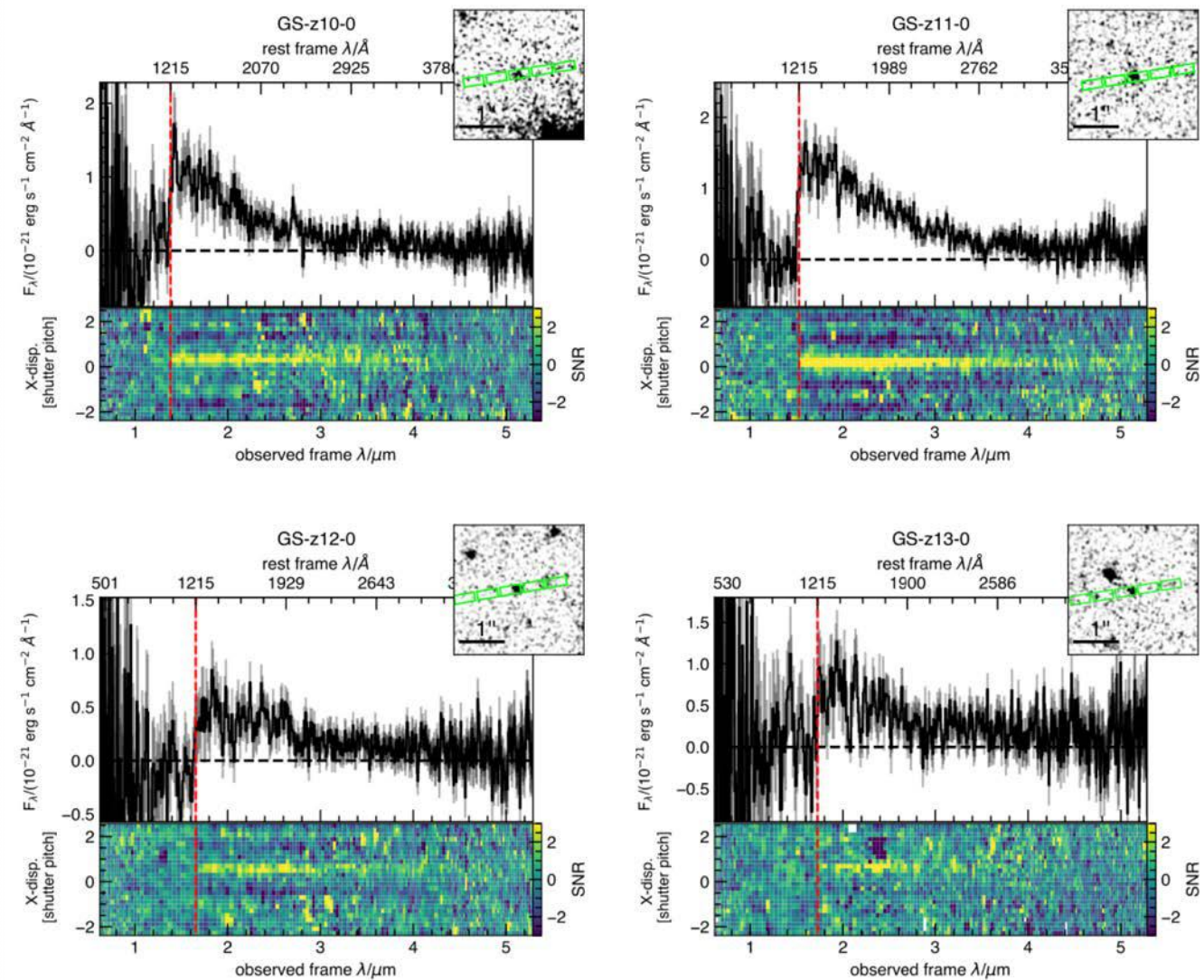


Fig. 1 NIRSpect prism $R \sim 100$ spectra for the four $z > 10$ galaxies targeted for the first deep spectroscopic pointing of the JADES survey, JADES-GS-z10-0, JADES-GS-z11-0, JADES-GS-z12-0 and JADES-GS-z13-0. For each galaxy we display the 1D spectrum and associated 1σ uncertainties (which are derived from standard error propagation through the reduction pipeline). In the bottom panel we show the 2D signal-to-noise ratio plot. The 2D plot is binned over four pixels in the wavelength direction to better show the contrast across the break. The inset panel in the top right-hand corner shows the NIRCам F444W filter image with the three nodding positions of the NIRSpect micro-shutter 3-slitlet array aperture shown in green. The red dashed line shows 1215.67\AA at the observed redshift z_{1216} .

2. More kinematics of the expanding Universe

- Redshift z (IMPORTANT)

$$1 + z \equiv \frac{\lambda_{\text{rcvd}}}{\lambda_{\text{rest}}} = \frac{a_{\text{today}}}{a_{\text{emit}}} = 1/a$$

- Age

$$dt = -\frac{dz}{H(z)(1+z)} \Rightarrow t(z) = \int_0^z \frac{dz}{(1+z)H(z)} \quad \& \quad t_0 = \int_0^\infty \frac{dz}{(1+z)H(z)}$$

- Luminosity distance

$$F \equiv \frac{\mathcal{L}}{4\pi d_L^2} \Rightarrow d_L(z) = (1+z)r(z)$$

- Comoving distance to z

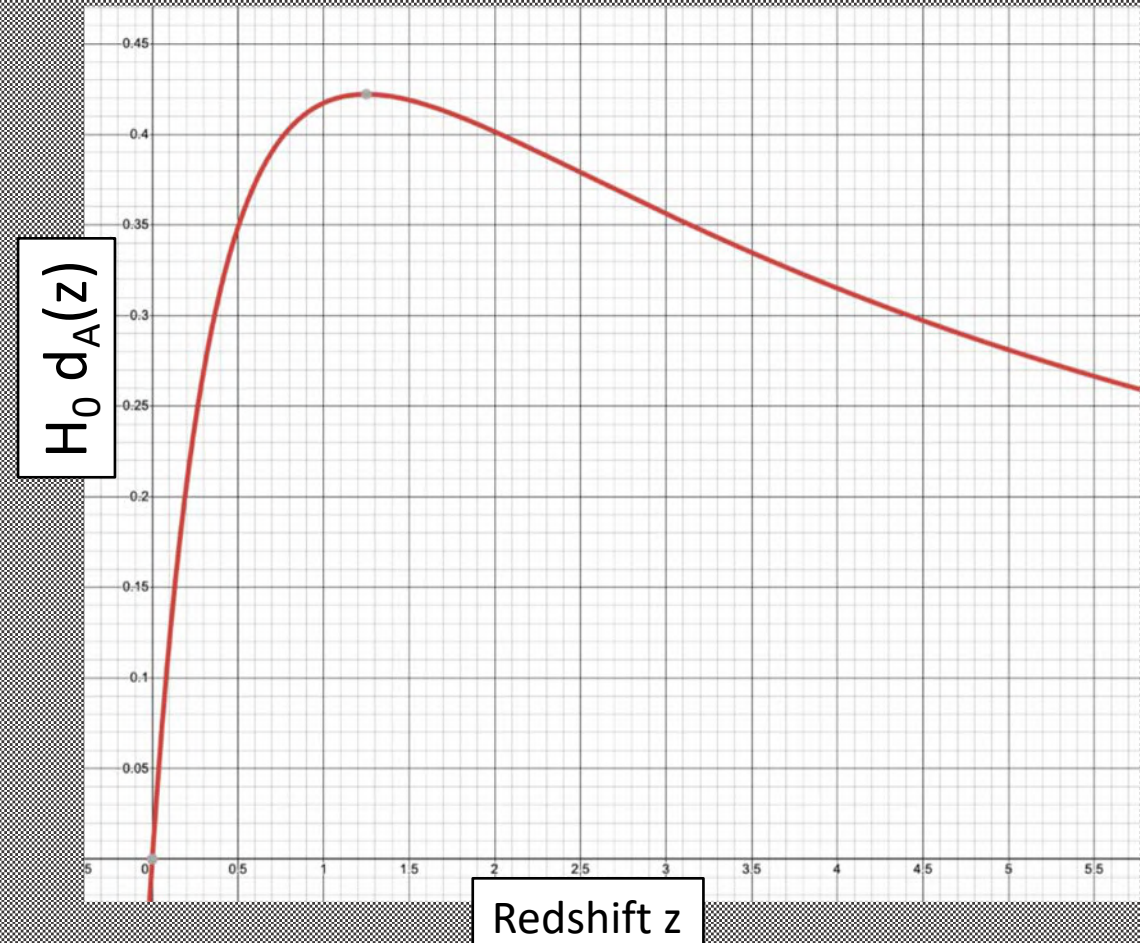
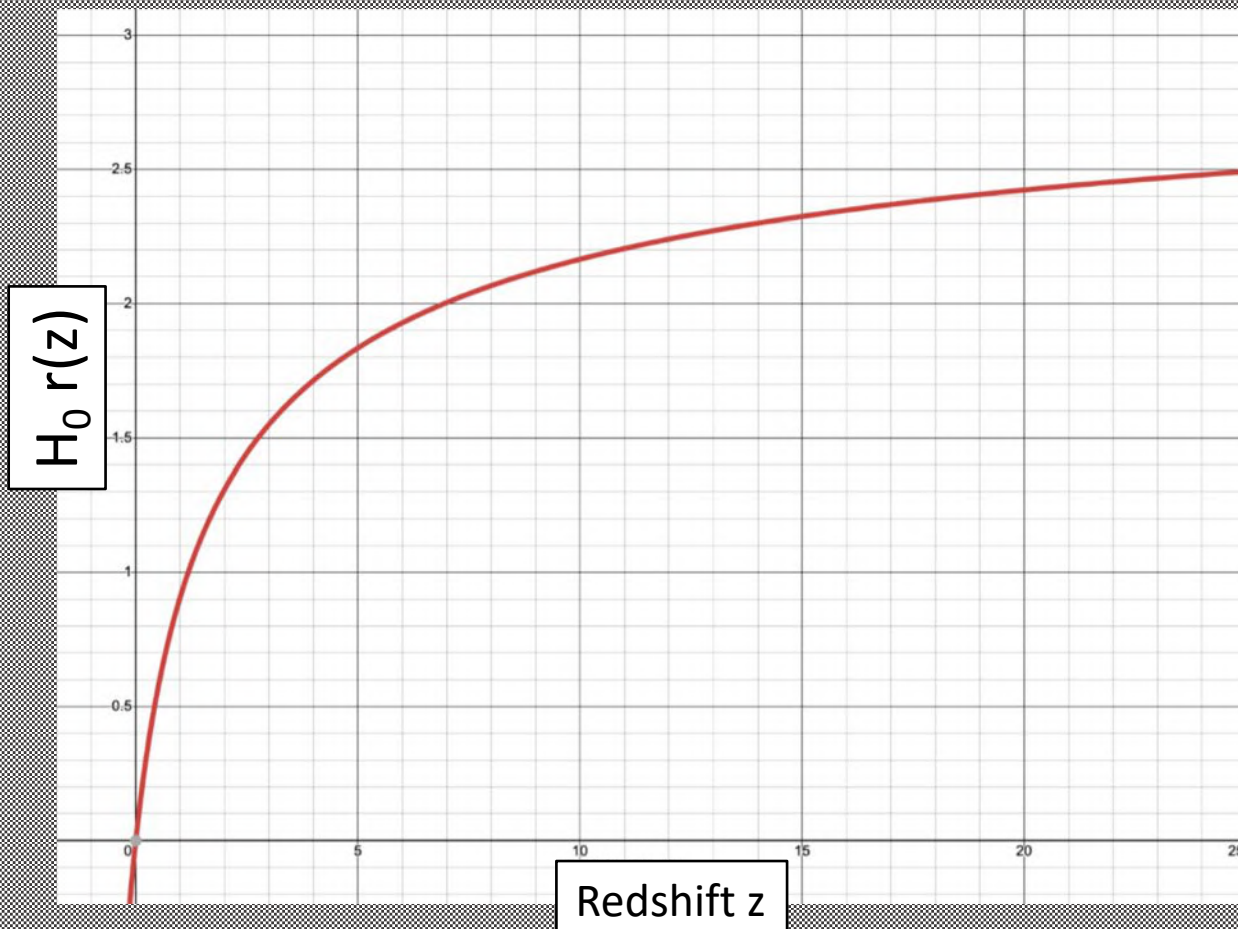
$$dr = dt/a(t) \Rightarrow r(z) = \int_0^z dz/H(z)$$

- Angular distance

$$\Delta\theta \equiv \frac{L}{d_A} \Rightarrow d_A = r(z)/(1+z) = d_L/(1+z)^2$$

Back to cosmological distances

- For small z : r , d_A , and $d_L = z H_0^{-1}$ ($H_0^{-1} = 4300 \text{ Mpc} = 14 \text{ Gyr}$)
- $r(z)_L$ asymptotes to $3.2 H_0^{-1}$ and d_A decreases at large z (rulers look bigger!)



3. Matter dynamics

- How free particles move: 3-momenta redshift as $1/a$

$$\vec{p} \propto 1/a \quad \Rightarrow \quad E \propto 1/a \quad \text{UR} \quad \text{KE} \propto 1/a^2 \quad \text{NR}$$

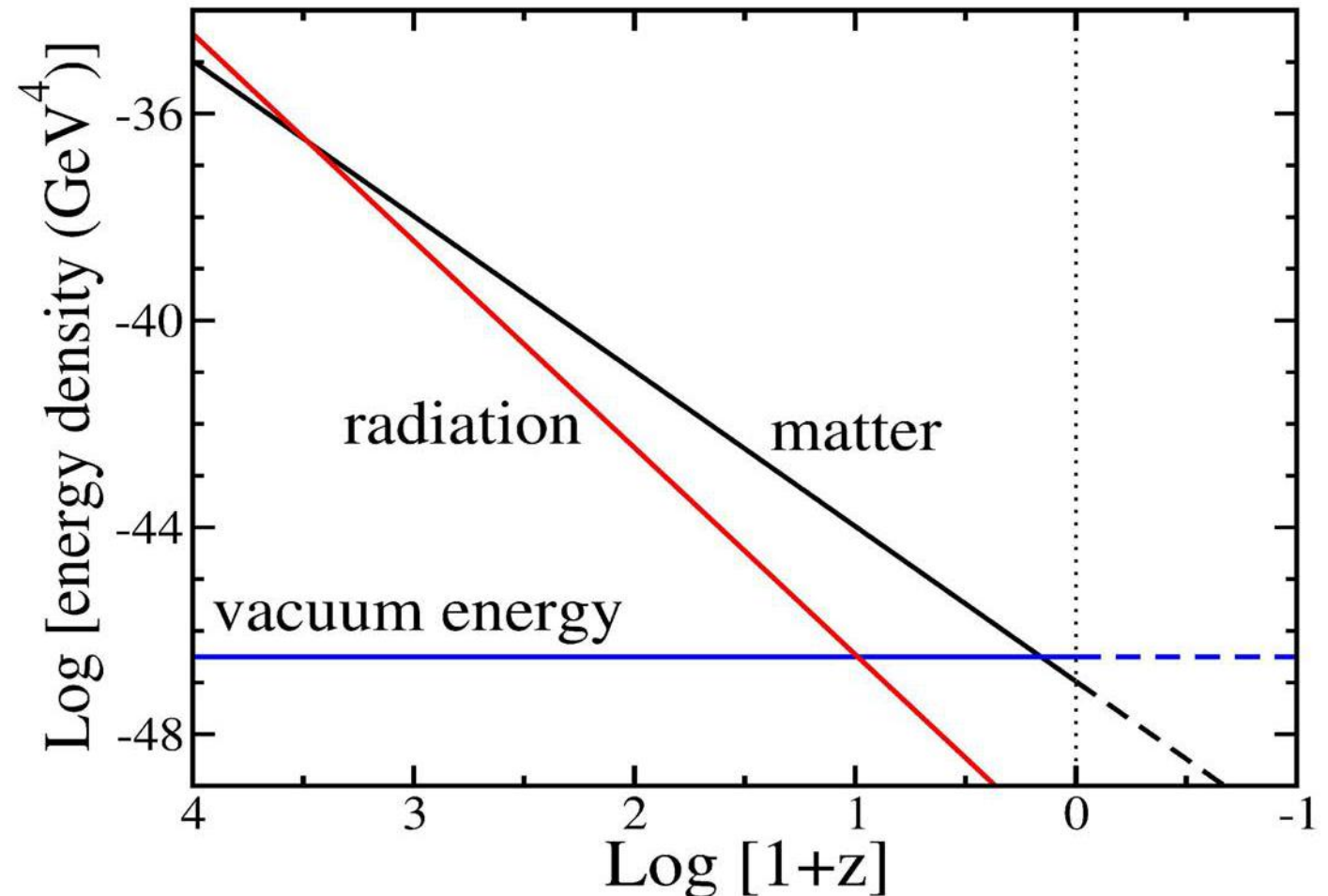
- How “fluids” evolve

$$\begin{aligned} d(\rho a^3) &= -p d(a^3) && \text{“First Law”} \\ \rho &\propto a^{-3(1+w)} && \text{for } p = w\rho \end{aligned}$$

$$\begin{aligned} \text{Matter } (w = 0): \quad \rho &\propto a^{-3} && \propto (1+z)^3 \\ \text{Radiation } (w = 1/3): \quad \rho &\propto a^{-4} && \propto (1+z)^4 \\ \text{Vacuum Energy } \Lambda \ (w = -1): &&& \rho \propto \text{const} \end{aligned}$$

Three epochs: Radiation dominated, Matter dominated, and Dark Energy dominated

- Radiation epoch
 - Thermal bath of particles (quark soup)
 - Origin of dark matter
 - Baryogenesis
 - Inflation
- Matter era
 - Growth of structure
 - Formation of CMB
- Vacuum era
 - Where we find ourselves
 - Unknown future



4. Dynamics: Matter and spacetime together

- Friedmann Equations:

$$H^2 \equiv \left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G\rho}{3}$$
$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}(\rho + 3p)$$

- Hubble parameter

$$H(z)^2 = H_0^2 \left[\Omega_\Lambda + \Omega_M(1+z)^3 + \Omega_R(1+z)^4 \right]$$

- Scale factor

Radiation-dominated ($z > 3000$): $a \propto t^{1/2}$

Matter-dominated ($3000 > z > 1$): $a \propto t^{2/3}$

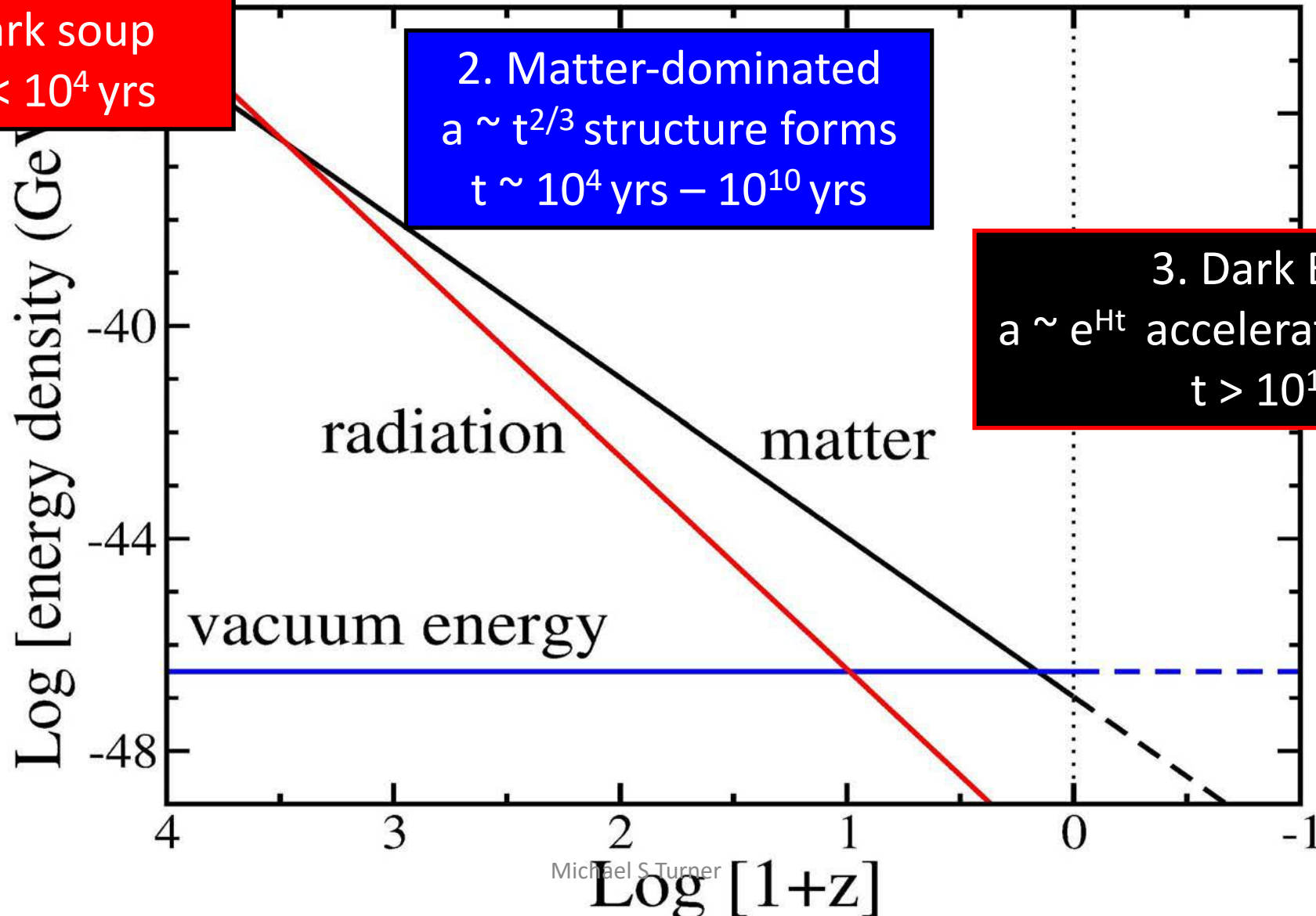
Vacuum-Energy-dominated ($z < 1$): $a \propto \exp(Ht)$

5. The 3 epochs of cosmology

1. Radiation-dominated
 $a \sim t^{1/2}$ quark soup
 $a < 10^{-4}$, $t < 10^4$ yrs

2. Matter-dominated
 $a \sim t^{2/3}$ structure forms
 $t \sim 10^4$ yrs – 10^{10} yrs

3. Dark Energy
 $a \sim e^{Ht}$ accelerated expansion
 $t > 10^{10}$ yrs



Age of the Universe

$$H = \dot{a}/a = n/t \text{ for } a \propto t^n$$

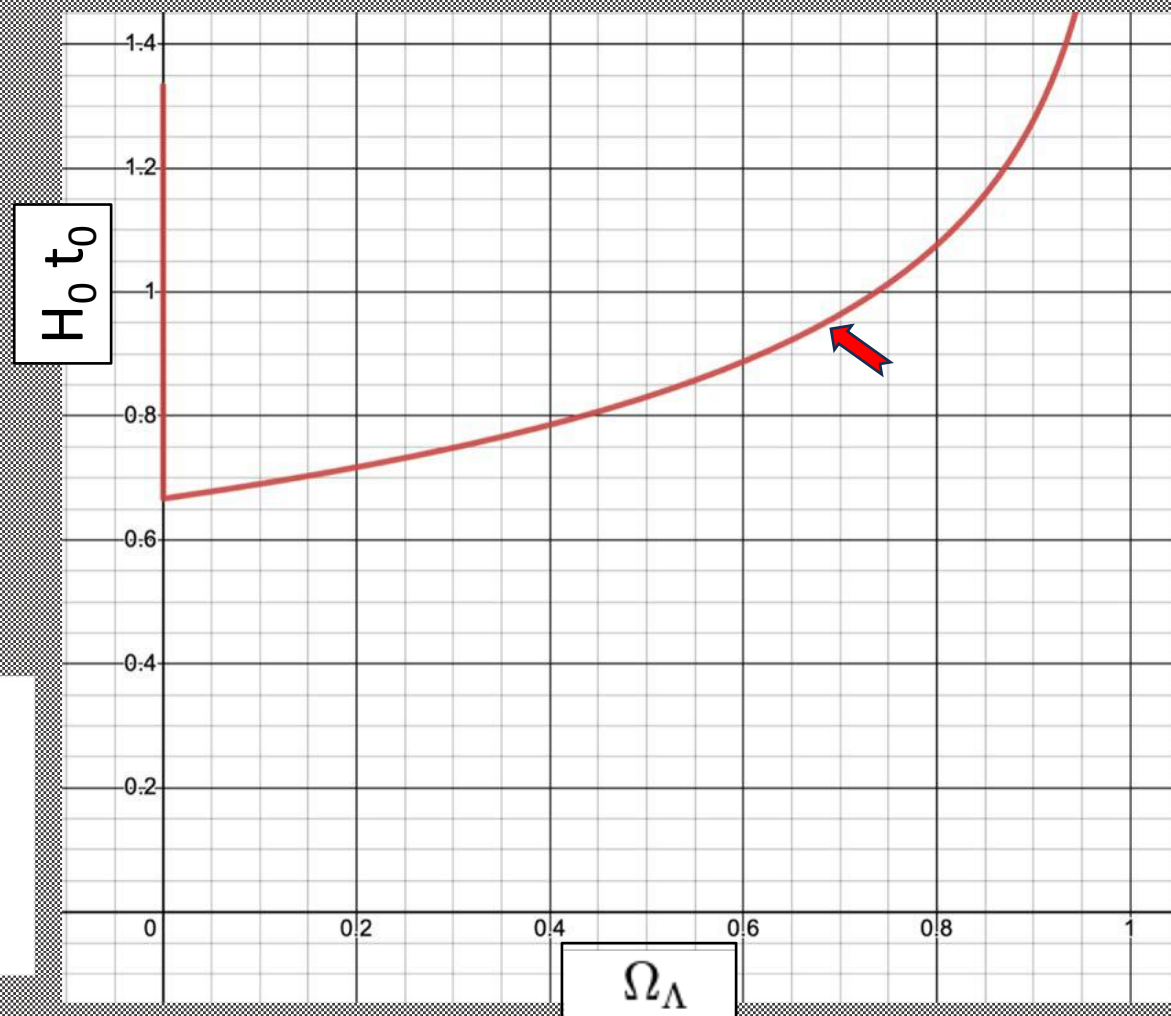
$$t = 2H^{-1}/3 \text{ MD}$$

$$t = H^{-1}/2 \text{ RD}$$

$$\text{for } \Lambda\text{CDM } t_0 \simeq H_0^{-1}$$

$$t_0 = \frac{2}{3} H_0^{-1} \Omega_{\text{VAC}}^{-1/2} \ln \left[\frac{1 + \Omega_{\text{VAC}}^{1/2}}{(1 - \Omega_{\text{VAC}})^{1/2}} \right]$$

$$t_0 = \int_0^\infty \frac{dz}{(1+z)H(z)}$$



One last thing, horizons (how far you can see on a clear day)

- Universe was smaller, but was expanding fast
- Light travels only about $c \times t$ since the beginning (mostly in the last Hubble time)

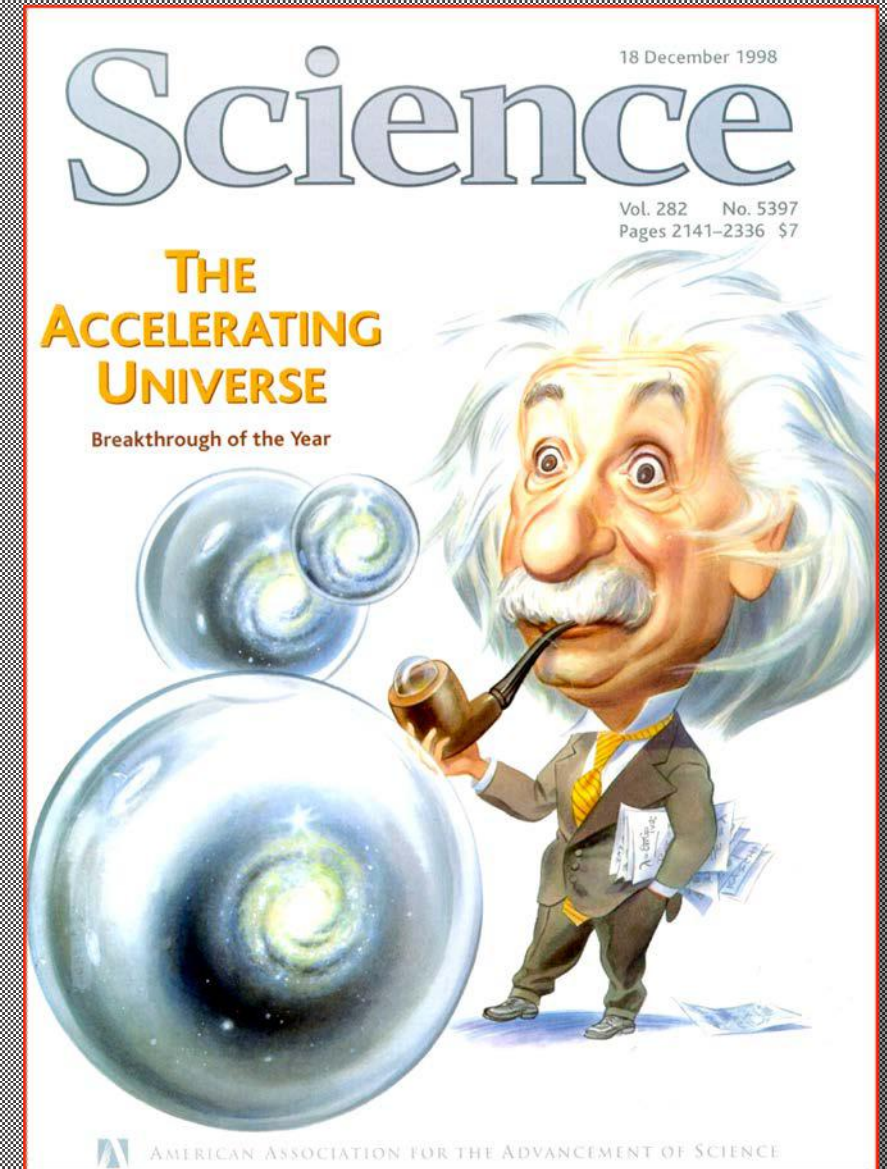
$$d_H(t) \equiv a(t) \int_0^t dr = a(t) \int_0^t \frac{dt'}{a(t')} = \frac{t}{1-n} = \frac{n}{1-n} H^{-1}$$

$$d_H(t) = 3t = 2H^{-1} \quad \text{MD}$$

$$d_H(t) = 2t = H^{-1} \quad \text{RD}$$

- This is known as the “horizon problem”: can’t smooth or create inhomogeneities on very large scales

1998: Cosmic speed up and dark energy



- Λ (vacuum energy) fits the data but why so small?
- Evidence of the rich vacua of string theory and the multiverse?
- Related to inflation (accelerated expansion) or something else?
- Describe by equation-of-state

$$p = w\rho$$

$$w = -1.0 \pm 0.04$$

REPULSIVE GRAVITY IS A FEATURE NOT A BUG! OF EINSTEIN'S THEORY

gravity is repulsive if $\rho + 3p < 0$

... but only really weird stuff has repulsive gravity

"DARK ENERGY"

QUANTUM NOTHINGNESS HAS REPULSIVE GRAVITY!



How REPULSIVE?

JUST ABOUT RIGHT -- GIVE OR TAKE 10^{55}

DARK ENERGY MAY BE THE MOST PROFOUND PROBLEM IN ALL OF SCIENCE TODAY

$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3} (\rho + 3p)$$

Gravity is sourced
by energy + 3 x
pressure \rightarrow BHs and
repulsive gravity!!!

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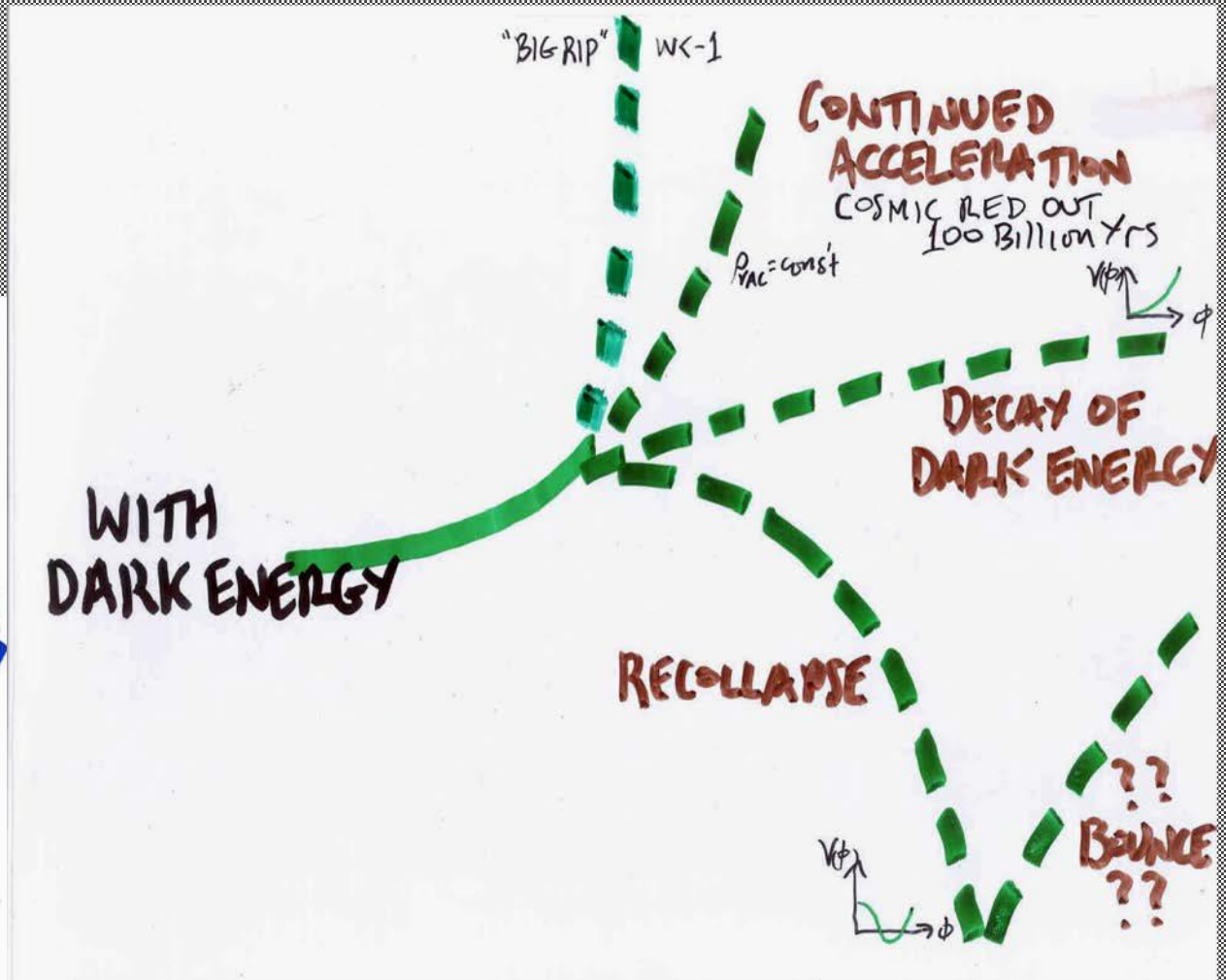
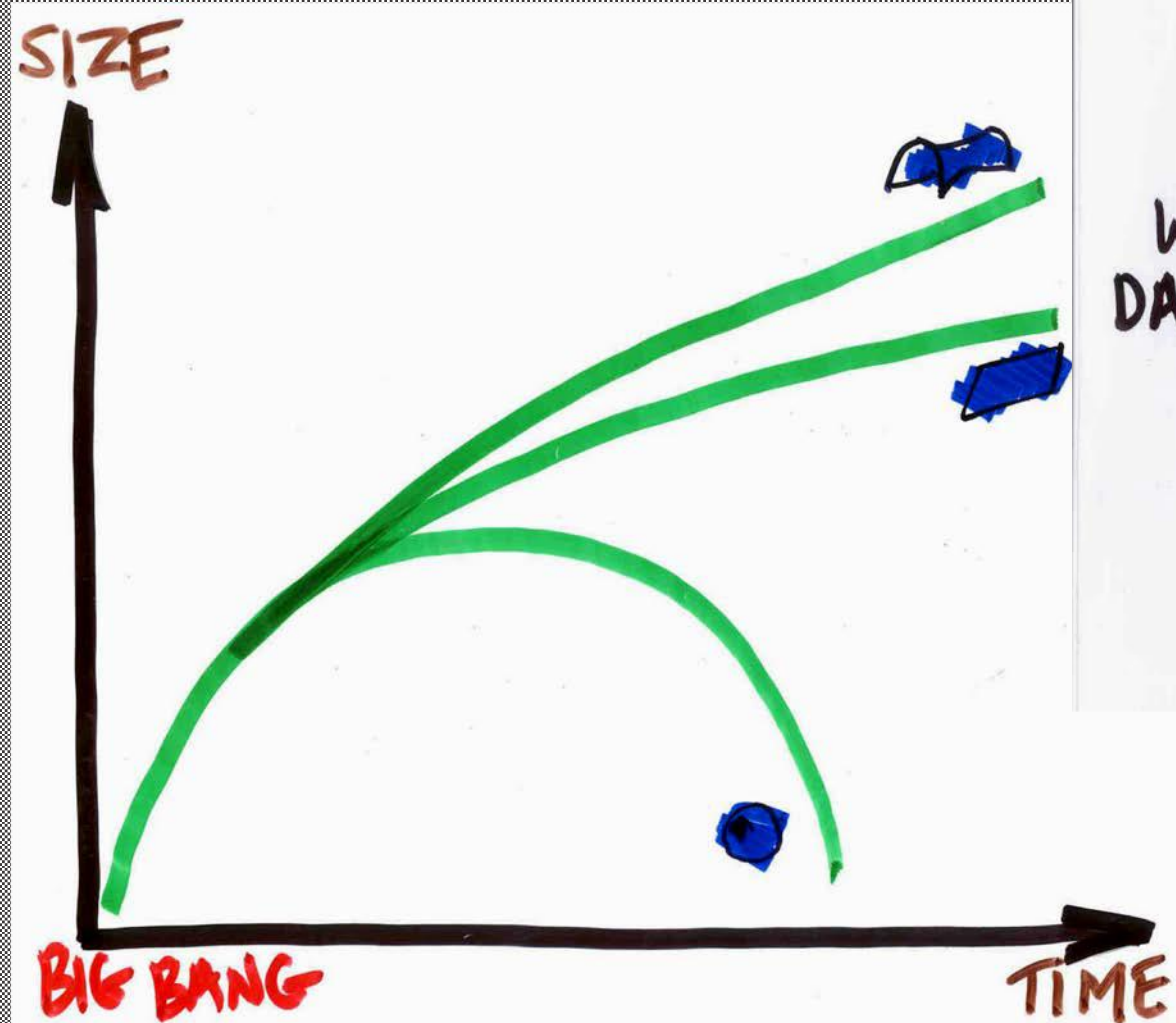
Multi-Messenger Astrophysics

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Cosmic Destiny



In the presence of dark energy, a flat Universe can expand forever, re-collapse, or even experience a big rip!

The Standard Cosmology (Λ CDM) in plain English

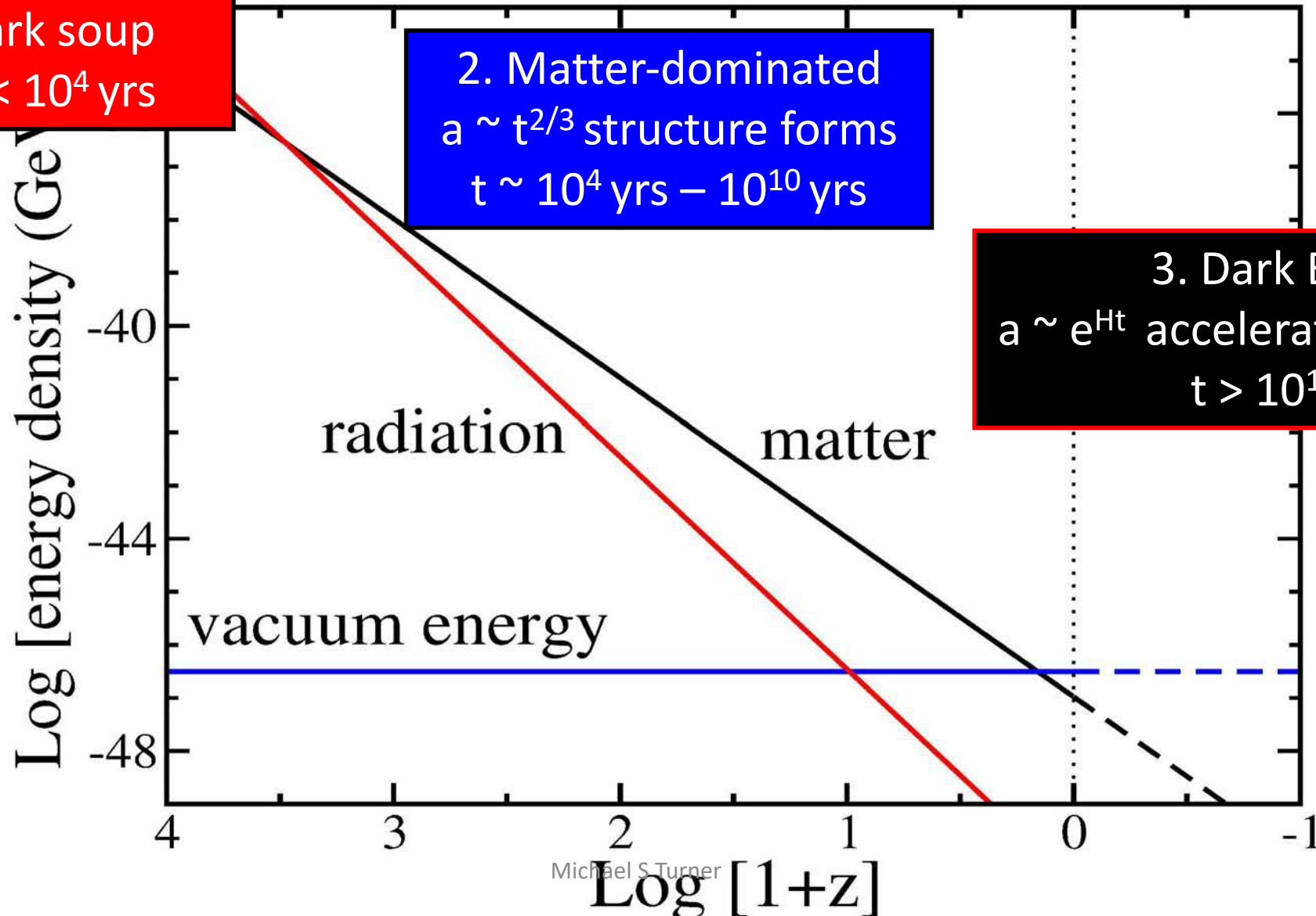
... very-early accelerated expansion driven by the potential energy of a scalar field gives rise to a very-large, smooth, spatially flat patch that becomes all that we can see today. Quantum fluctuations during this inflationary phase grow into the seeds for galaxies. The conversion of potential field energy into heat produces the quark soup that evolves a baryon asymmetry and long-lived dark matter particles. The excess of quarks over antiquarks becomes neutrons and protons, later some light elements and finally atoms. The gravity of the dark matter particles drives the formation of structure from galaxies to superclusters and a mere 5 billion years ago the repulsive gravity of dark energy (Λ) again drove accelerated expansion ...

5. The 3 epochs of cosmology

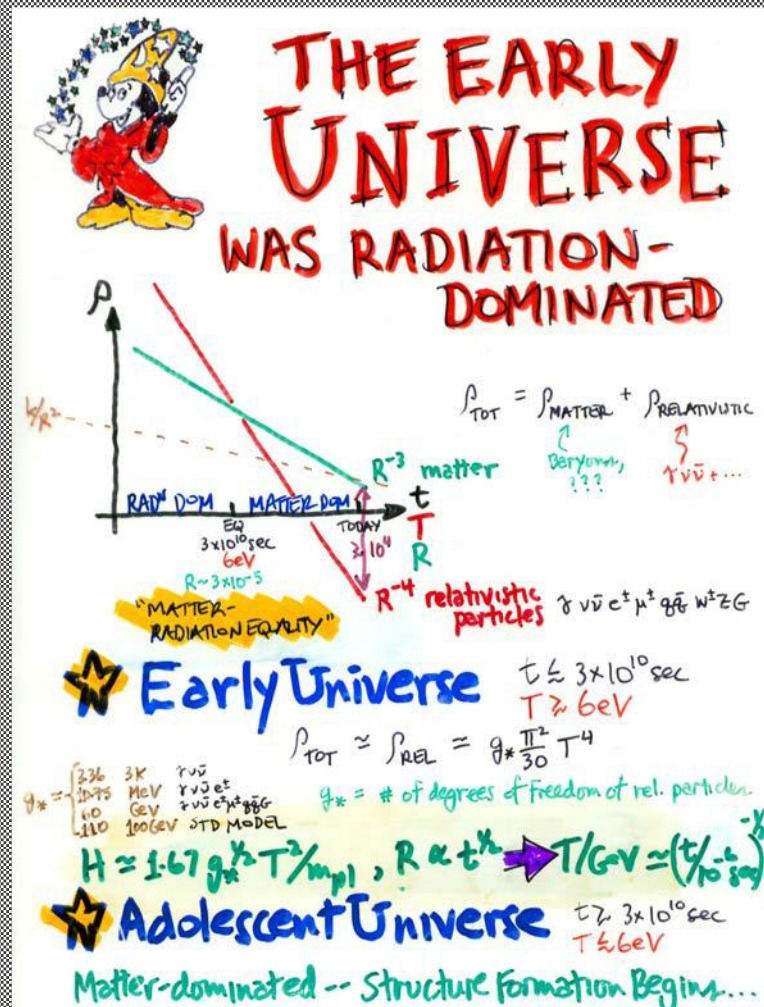
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 $a \sim e^{Ht}$ accelerated expansion
 $t > 10^{10}$ yrs



6. Thermodynamics in the early Universe



- Thermal bath of particles
- For $kT > mc^2$ particle/antiparticle pairs as abundant as photons

$$\rho_\gamma = \frac{\pi^2 T^4}{15} \quad n_\gamma = \frac{2\zeta(3)T^3}{\pi^2}$$

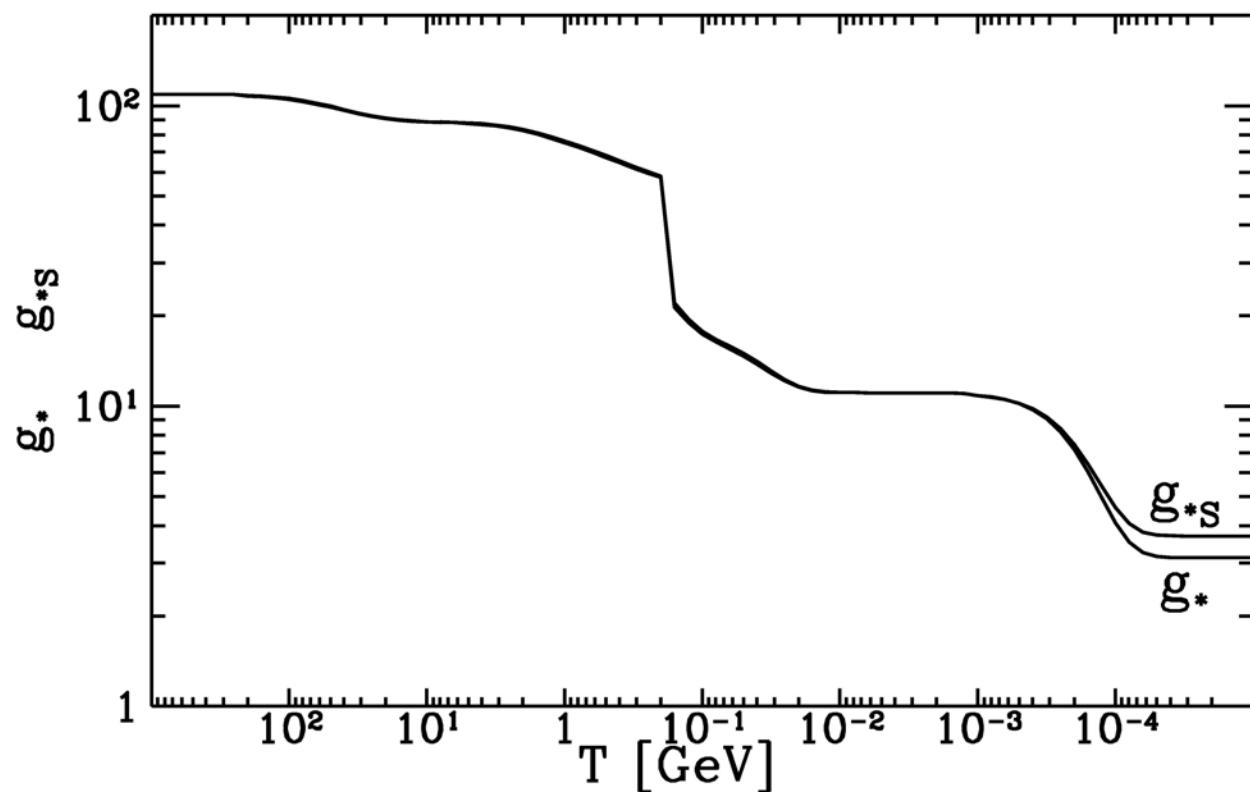
- For most of its early history: thermal equilibrium
- But, departures are very important – not all Fe today



$$\rho_R = g_* \frac{\pi^2}{30} T^4 \quad g_*(T) = \sum_{\text{bosons}} g_i + \frac{7}{8} \sum_{\text{fermions}} g_i$$

$$\Rightarrow H \simeq 1.67 g_*^{1/2} T^2 / m_{\text{pl}}$$

$$a(t) \propto t^{1/2} \quad T/\text{GeV} \sim (t/10^{-6} \text{ sec})^{-1/2}$$



UNITS

BART SAYS:

$$\hbar = k_B = c = 1$$

$$1 \text{ GeV} = (2 \times 10^{-14} \text{ cm})^{-1} = (6.6 \times 10^{-25} \text{ sec})^{-1} \\ = 1.6 \times 10^3 \text{ erg} = 1.2 \times 10^{13} \text{ K}$$

$$1 \text{ GeV}^2 = (0.4 \text{ mb})^{-1}$$

$$1 \text{ GeV}^3 = 1.3 \times 10^{41} \text{ cm}^{-3}$$

$$1 \text{ GeV}^4 = 2.32 \times 10^{17} \text{ g cm}^{-3}$$

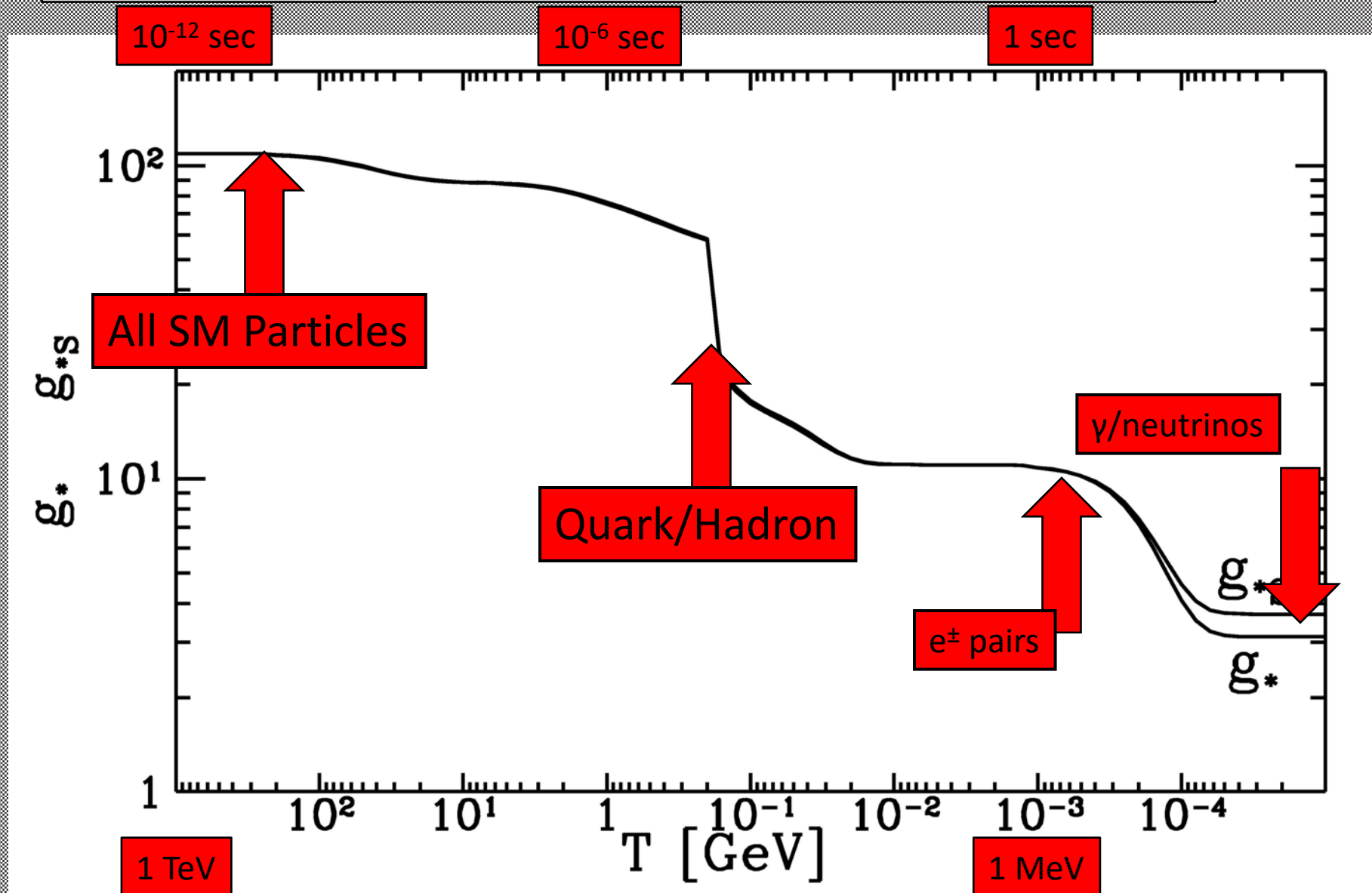
$$1 \text{ GeV}^5 = 3.2 \times 10^{62} \text{ erg cm}^{-3} \text{ s}^{-1}$$

$$G \equiv m_{\text{pl}}^{-2} \quad m_{\text{pl}} = 1.22 \times 10^{19} \text{ GeV}$$

$$1 \text{ Mpc} = 3.1 \times 10^{24} \text{ cm} = 1.6 \times 10^{38} \text{ GeV}^{-1}$$

$$H_0 = 2.1 h \times 10^{-42} \text{ GeV}$$

Relativistic Degrees of Freedom



Basic thermodynamics review

- Thermal phase space density

$$f_i = \exp [\mu_i - E_i/T]$$

- With zero chemical potentials

- Ultrarelativistic limit

$$n_i \sim T^3$$

- Non-relativistic limit

$$n_i \sim (mT)^{3/2} \exp[-m_i/T]$$

- Chemical equilibrium establishing by $i + j + \dots \leftrightarrow a + b + \dots$
occurring “rapidly”

$$\mu_i + \mu_j + \dots = \mu_a + \mu_b + \dots$$

Entropy conservation (in the absence of departures from equilibrium and entropy production)

$$s = \frac{\rho + p}{T} = g_* \frac{2\pi^2}{45} T^3$$

$$S \equiv a^3 s \propto g_* a^3 T^3 = \text{const}$$

$$\Rightarrow a^3 \propto 1/s$$

$$\Rightarrow n_X/s \propto \text{number of } X'\text{'s per comoving volume}$$

$$\Rightarrow T \propto g_*^{-1/3} a^{-1}$$

High entropy/small baryon number

- When g_* is constant, e.g., since $t = 1$ sec, $s \simeq 7n_\gamma \simeq 3000 \text{ cm}^{-3}$
- n_B/s = baryon number per comoving volume = n_b/s (few or no anti-baryons) = (baryon-to-photon ratio $\eta = 6 \times 10^{-10}$)/7 = 10^{-11}
- Note $7/\eta = 10^{11}$ is the entropy per baryon – VERY HIGH – meaning lots of photons per baryon! Cf, newly born neutron star entropy per baryon is a few per baryon. BIG consequences in cosmology!
- The old question: where did all the entropy come from?
- The other way to look at: Where did the small net baryon number $n_B/s = 10^{-11}$ come from? Baryogenesis!

Thermal equilibrium requires rapid interaction rates relative to expansion rate

- Interaction rate
= number density of targets x cross section x v_{relative}

$$\Gamma = n\sigma v$$

- Expansion rate

$$H \simeq T^2/m_{\text{pl}}$$

- Expansion time ($= 1/H$) x interaction rate = number of interactions per “expansion time” (doubling of scale factor, halving of temperature)
- Rapid interaction rate (thermal equilibrium):
- Slow interaction rate (“frozen out” reactions):

$$\Gamma/H \gg 1$$

$$\Gamma/H \ll 1$$

Worked example: decoupling of neutrinos

- Interactions

$$\bar{\nu}\nu \leftrightarrow e^+e^-, \nu e \leftrightarrow \nu e, \text{ etc.}$$

- Interaction rate

$$\sigma \simeq G_F^2 T^2$$

$$\Gamma_{int} = n\sigma|v| \simeq G_F^2 T^5$$

- Decoupled at $T \sim 1 \text{ MeV}$

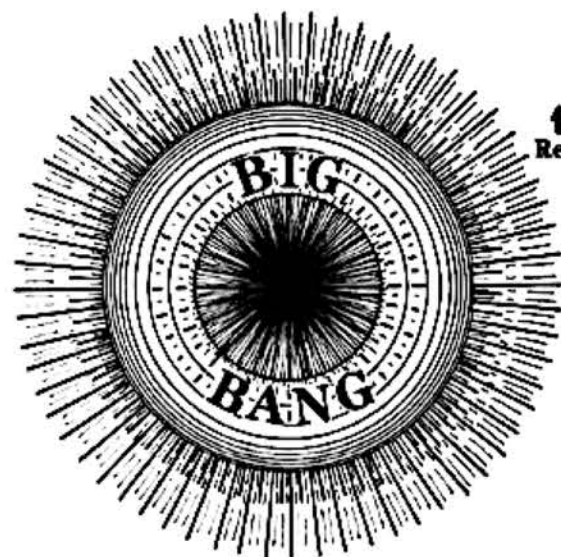
$$\frac{\Gamma_{int}}{H} \simeq \frac{G_F^2 T^5}{T^2/m_{Pl}} \simeq \left(\frac{T}{1 \text{ MeV}} \right)^3$$

Worked example: neutrino to photon temperature

- Universe at 1 sec/1 MeV: photons ($g = 2$) and electrons/positrons ($g = 4 \times 7/8 = 3.5$)
- Neutrinos decouple and evolve adiabatically
- $T < 1$ MeV, entropy from electron/positron pairs resides in photons only

$$T_\nu \propto \frac{1}{a} \quad T_\gamma \propto \frac{g_*^{-1/3}}{a}$$
$$g_* = 11/2 \rightarrow 2 \Rightarrow g_*^{-1/3} = (11/2)^{-1/3} \rightarrow (2)^{-1/3}$$
$$\Rightarrow T_\nu/T_\gamma = \left(\frac{4}{11}\right)^{1/3}$$

PS: Neutrinos participate in a small part of the e^+/e^- entropy transfer (about 1%)



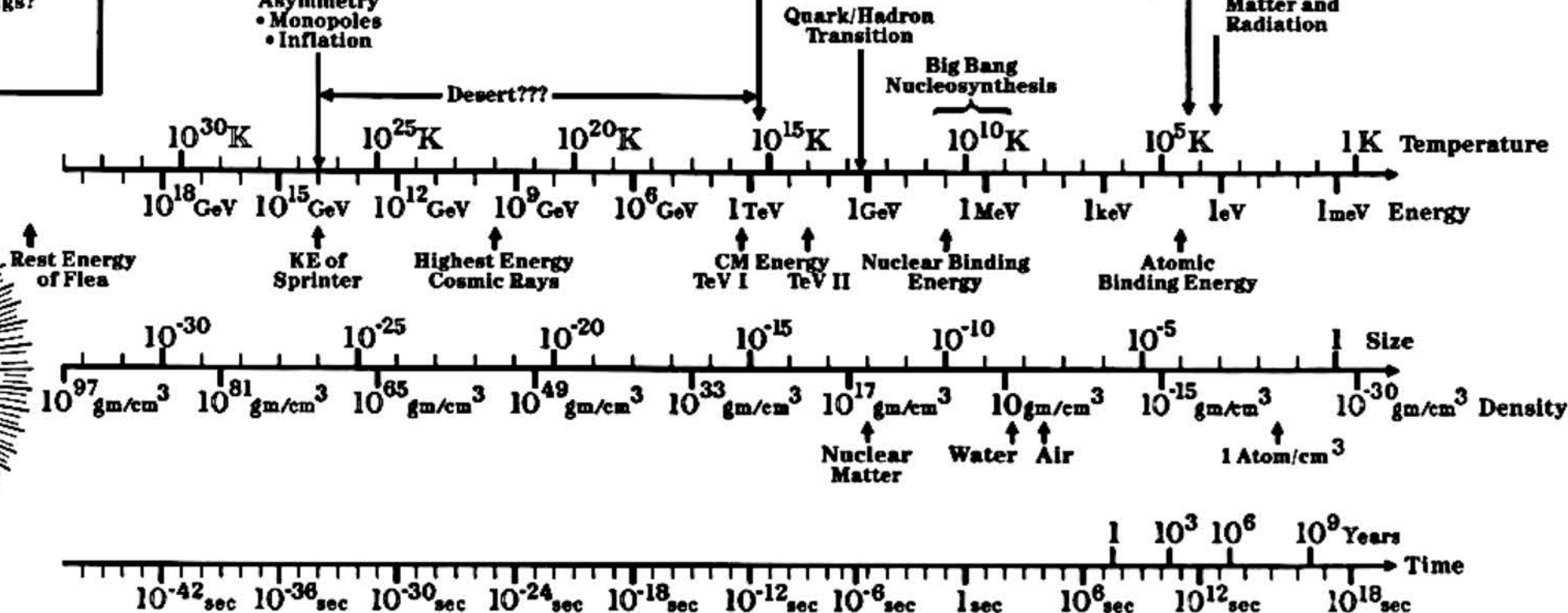
QUANTUM GRAVITY
 • Supergravity?
 • Extra Dimensions?
 • Supersymmetry?
 • Superstrings?

END OF GRAND UNIFICATION
 • Origin of Matter-Antimatter Asymmetry
 • Monopoles
 • Inflation

END OF ELECTROWEAK UNIFICATION
 • End of Supersymmetry?

MATTER DOMINATION
 • Formation of Structure Begins

• Formation of Atoms
 • Decoupling of Matter and Radiation



CONSTITUENTS

Leptons and Quarks
 $\left\{ \begin{array}{l} (\nu_e) (\nu_\mu) (\nu_\tau) \\ (e^-) (\mu^-) (\tau^-) \end{array} \right\} ???$
 $\left\{ \begin{array}{l} (u) (c) (t) \\ (d) (s) (b) \end{array} \right\} ???$

Gauge Bosons
 { GLUONS
 $W^\pm Z$
 $X, Y, ...??$

$\nu \bar{\nu}$
 $e^+ e^-$
 n, p
 $H^+, D^+, {}^3He^{++}, {}^4He^{++}, {}^7Li^{++}, e^-$
 $H, D, {}^3He, {}^4He, {}^7Li$

Ratio of Matter/Radiation = 5×10^{-10}

Photons γ

3K Microwave Background

ORDINARY MATTER: FROM QUARKS TO US

INFLATION
BARYOGENESIS



TRANSITION FROM
QUARKS → NEUTRONS, PROTONS



BIG-BANG
NUCLEOSYNTHESIS
Formation of H, D,
He, He-3, Li

$D/H = (3 \pm 0.2) \times 10^{-5}$
 $\Omega_B = 0.04 \pm 0.002$

FORMATION OF ATOMS
COSMIC MICROWAVE
BACKGROUND



CMB

RATIO OF FIRST-TO-
SECOND PEAKS: 2/1

$\Omega_B = 0.045 \pm 0.003$



1 BILLION YRS
FIRST QUASARS

QSO LIGHT

INTERGALACTIC GAS

ABSORPTION OF
QUASAR LIGHT
BY HYDROGEN

$\Omega_B \geq 0.04$

CMB

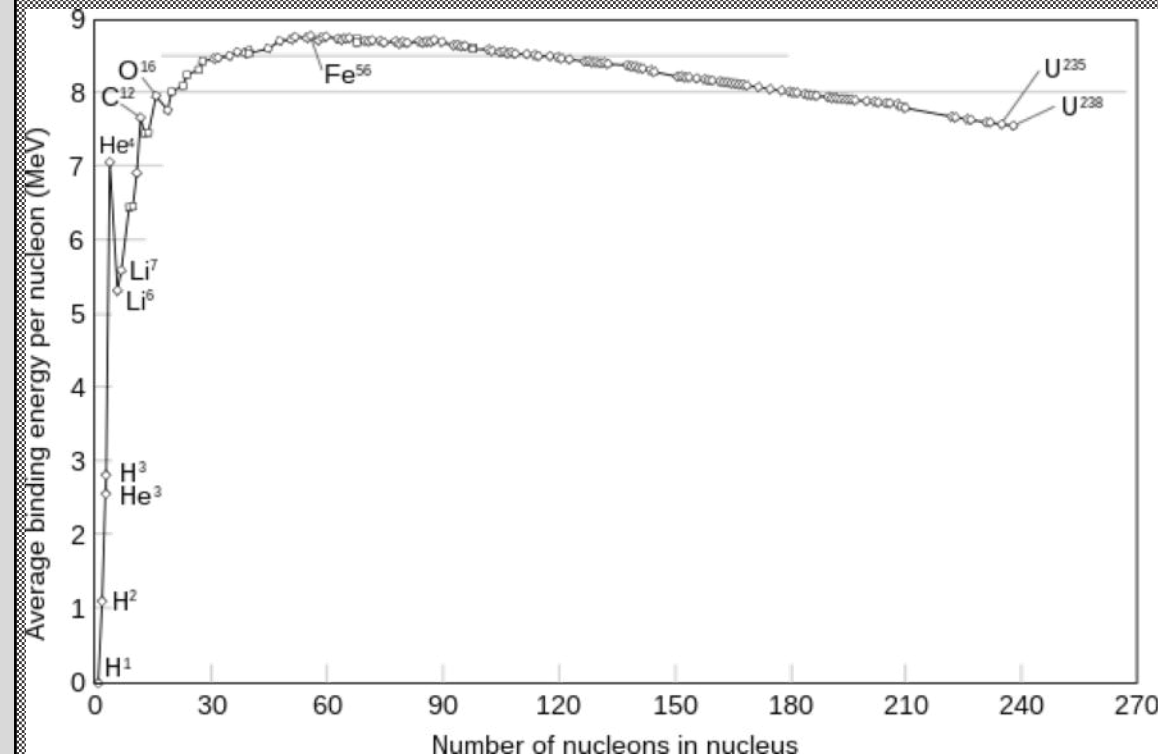
HERE & NOW

44 Billion YRS
stars, gas,
dust, ...

7. BBN and the high entropy of the Universe (small baryon number)

Most “accurate” description of BBN physics <https://arxiv.org/pdf/2111.14254.pdf>

- Boltzmann equations in the expanding Universe → the results of BBN depend upon the baryon to photon ratio $\eta = 6 \times 10^{-10}$ and nuclear data (cross sections, neutron lifetime) and $N_{\text{neutrinos}}$
- Large entropy per baryon (small η) plays a critical role in delaying BBN to a time when Coulomb barriers prevent nucleosynthesis beyond ${}^4\text{He}$ and NSE
- That is a good thing: a tremendous amount of nuclear free energy is left to power stars and life in the Universe 10^{52} ergs per solar mass!



- BBN begins with NSE (chemical equilibrium), followed by a series of departures from thermal equilibrium

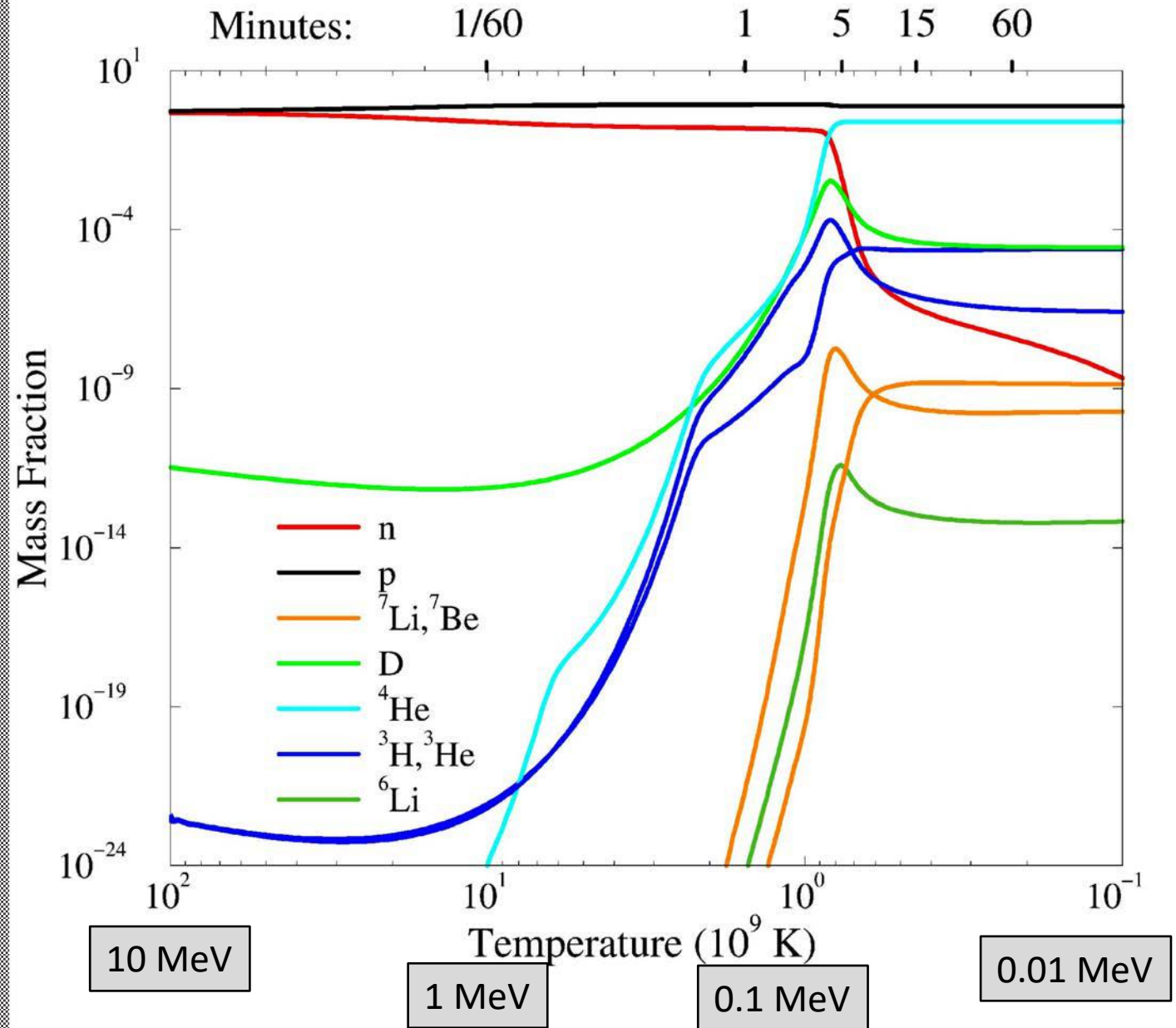
$$T_{\text{nuclei}} \simeq \frac{B_A/(A-1)}{\ln \eta^{-1} + 1.5 \ln(m/T)} \simeq 0.25 \text{ MeV}$$

- n/p ratio freezes in at a value of around 1/7 at $T \sim 1 \text{ MeV}$

$$(n/p)_{\text{eq}} = \exp[-\Delta m/T]$$

- Coulomb barriers ($T < 0.05 \text{ MeV}$) prevent NSE from being established and significant nucleosynthesis beyond ${}^4\text{He}$
- End result, lots of ${}^4\text{He}$ made from free neutrons, a little unburnt D and ${}^3\text{He}$ and a trace amount of ${}^7\text{Li}$

NB: without the non-thermal neutrons, first step of BBN would have to be $p + p \rightarrow D + \gamma$ (a weak interaction that would have time to take place!)



← NSE n/p freeze in | Nuclei favored → | Coulomb barriers →

Nuclear Statistical Equilibrium

- NSE abundance of Z_A :

$$n_A = g_A \left(\frac{m_A T}{2\pi} \right)^{3/2} \exp \left(\frac{\mu_A - m_A}{T} \right)$$

- Chemical Equilibrium

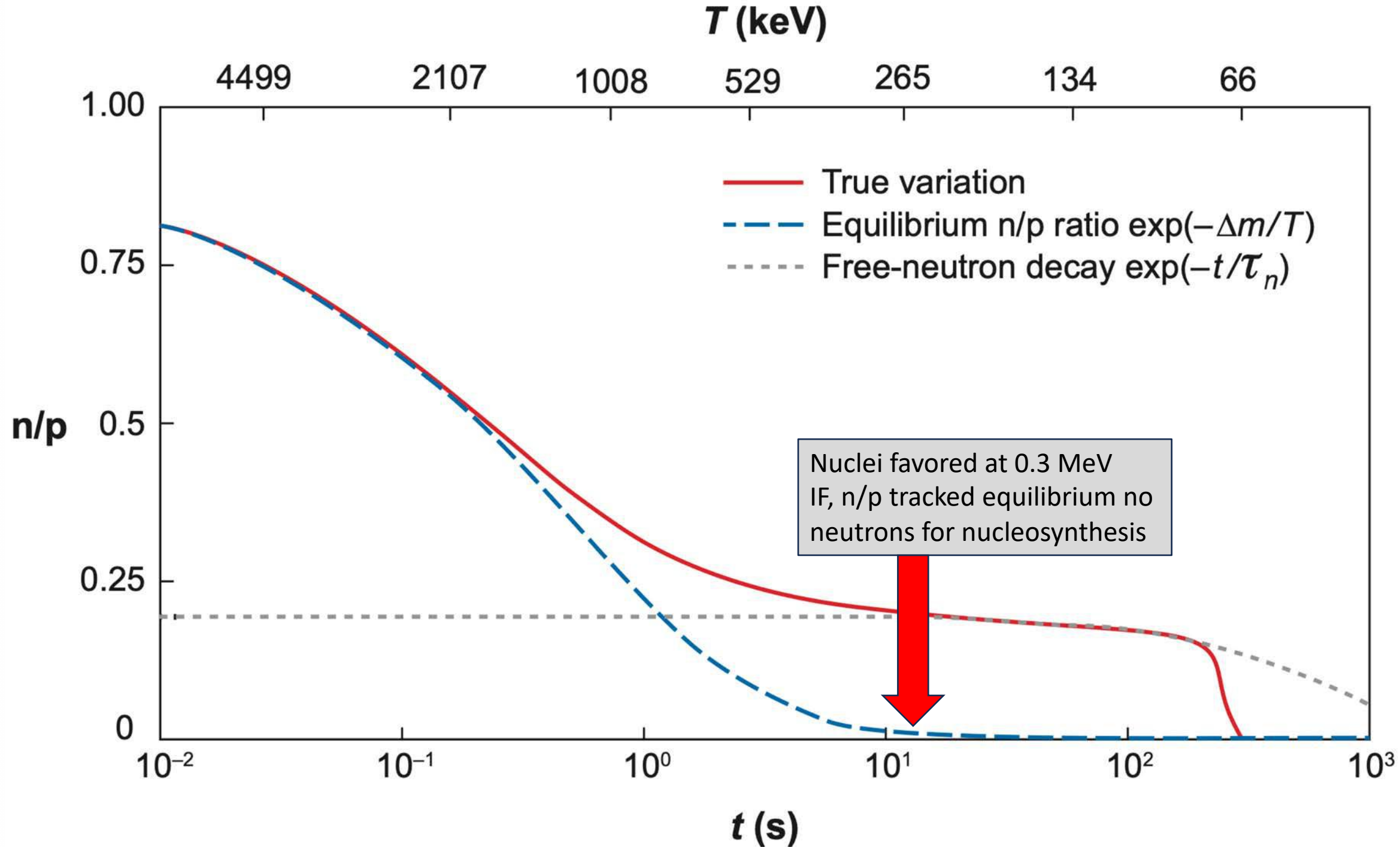
$$\mu_A = Z\mu_p + (A - Z)\mu_n.$$

- Binding energy

$$B_A \equiv Zm_p + (A - Z)m_n - m_A.$$

- Mass fraction of Z_A in NSE (after algebra)

$$X_A = g_A [\zeta(3)^{A-1} \pi^{(1-A)/2} 2^{(3A-5)/2}] A^{5/2} (T/m_N)^{3(A-1)/2} \\ \times \eta^{A-1} X_p^Z X_n^{A-Z} \exp(B_A/T),$$



Two big successes: D/H (vs. CMB) and ^4He and one big problem ^7Li

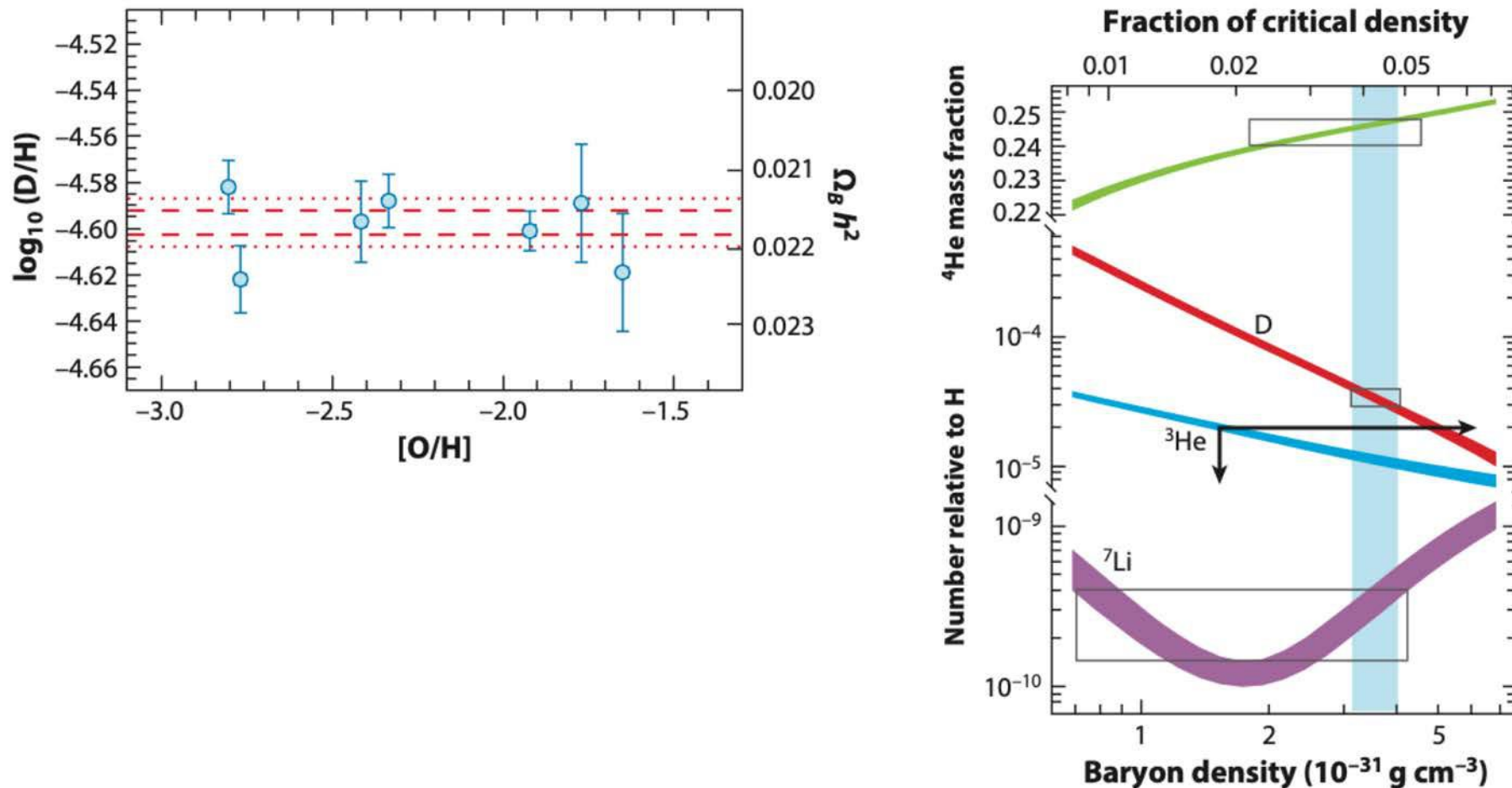


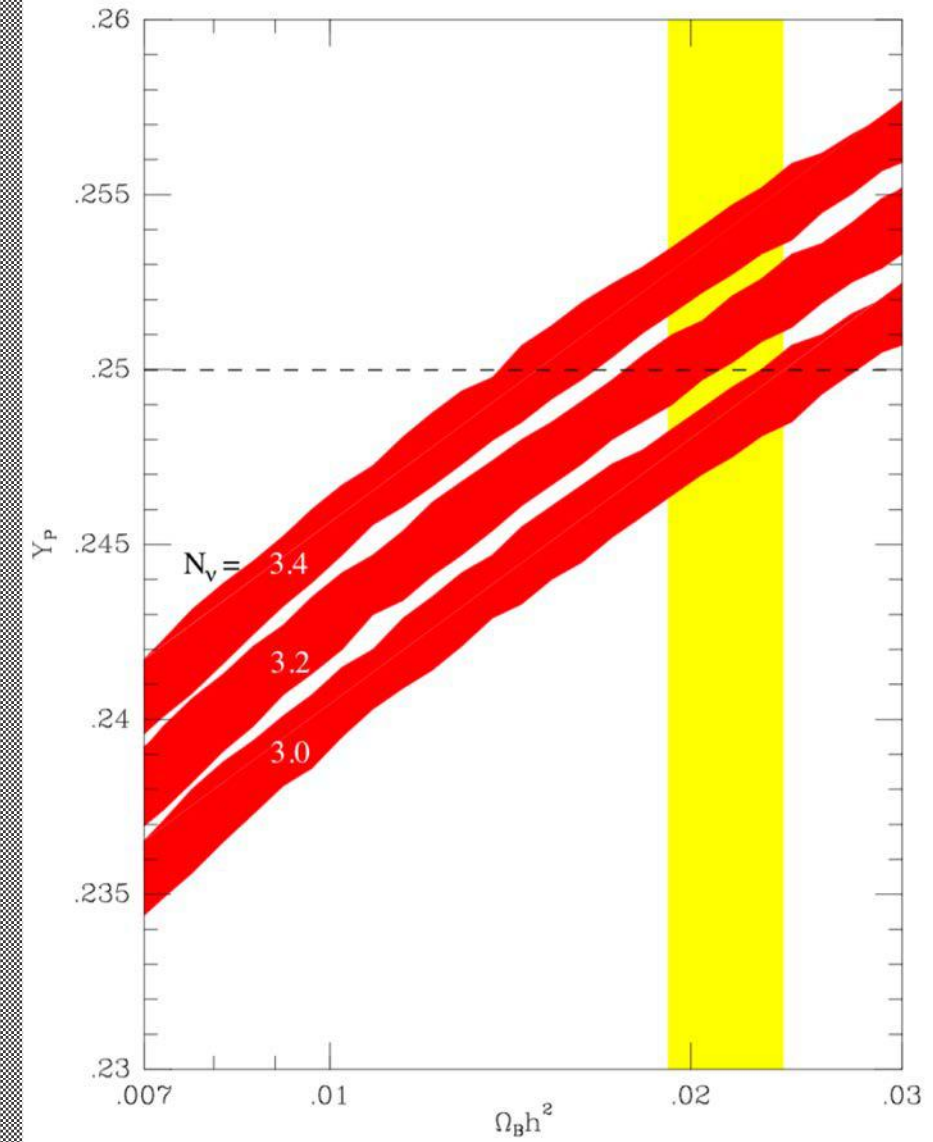
Figure 8

Big bang nucleosynthesis. (a) D/H determinations. Panel adapted with permission from Reference 178; copyright 2018 AAS. (b) The vertical band is the deuterium-determined baryon density, and the other bands are the 1σ predictions. The heights of the black boxes indicate the measured abundances with error estimates. The upper density scale assumes $H_0 = 65 \text{ km s}^{-1} \text{ Mpc}^{-1}$. Panel adapted from Reference 179.

Looking for more dark stuff

- N_{eff} : the number of relativistic species ($m \ll T$) expressed in as neutrino species*: portal to the TOE (dark sector, ...)
 - BBN: species less massive than 1 MeV & $\sigma_N = 0.05$
 - CMB: species less massive 1 eV & $\sigma_N = 0.03$
 - Don't have to agree

*Chicago school convention; NB: SM predicts 3.045 for CMB



8. Recombination and CMB last scattering

- Two separate events!
- Recombination (misnomer):

- Chemical equilibrium: rapid

$p + e \leftrightarrow H + \gamma$ ensures

$$\mu_p + \mu_e = \mu_H$$

- Ionization fraction in thermal equilibrium (X_e), $B = 13.6$ eV

$$\frac{1 - X_e}{X_e^2} \simeq \eta (T/m_e)^{3/2} \exp(B/T)$$

- Rec: $X_e \sim 0.5$ ($z = 1300$)

$$T_{\text{rec}} \simeq \frac{B}{-\ln \eta - 1.5 \ln(T_{\text{rec}}/m_e)} \simeq 0.3 \text{ eV}$$

- NB: freeze out of recombination leaves residual ionization of $X_e = 10^{-4}$ or so

$\ll 13.6 \text{ eV!!}$

Last-scattering (also called decoupling)

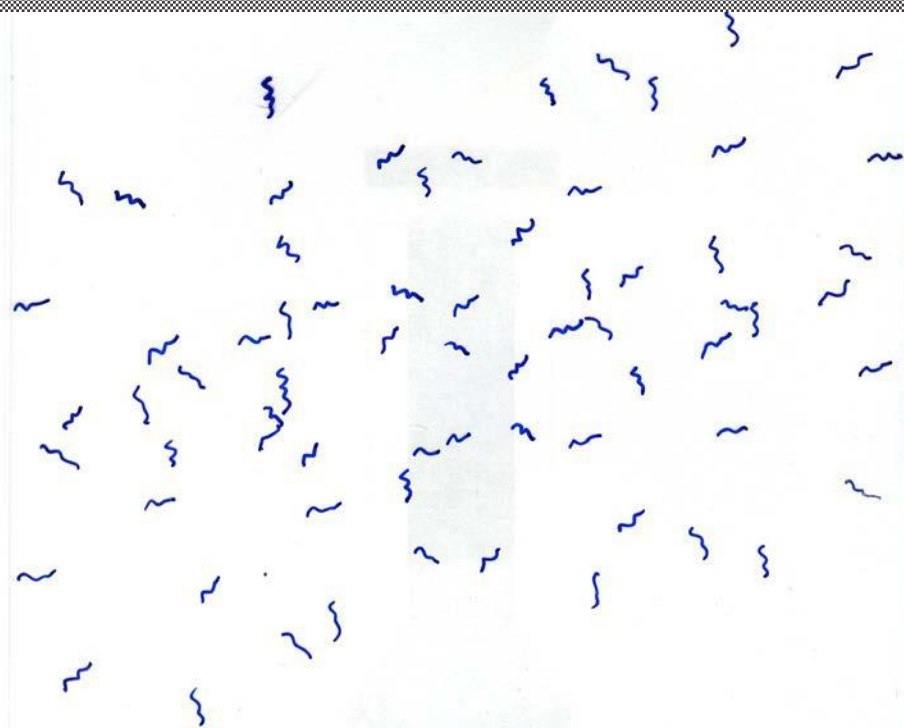
- Thomson scattering rate per photon:

$$\Gamma = n_e \sigma_T \simeq X_e \eta T^3 \alpha^2 / m_e^2$$

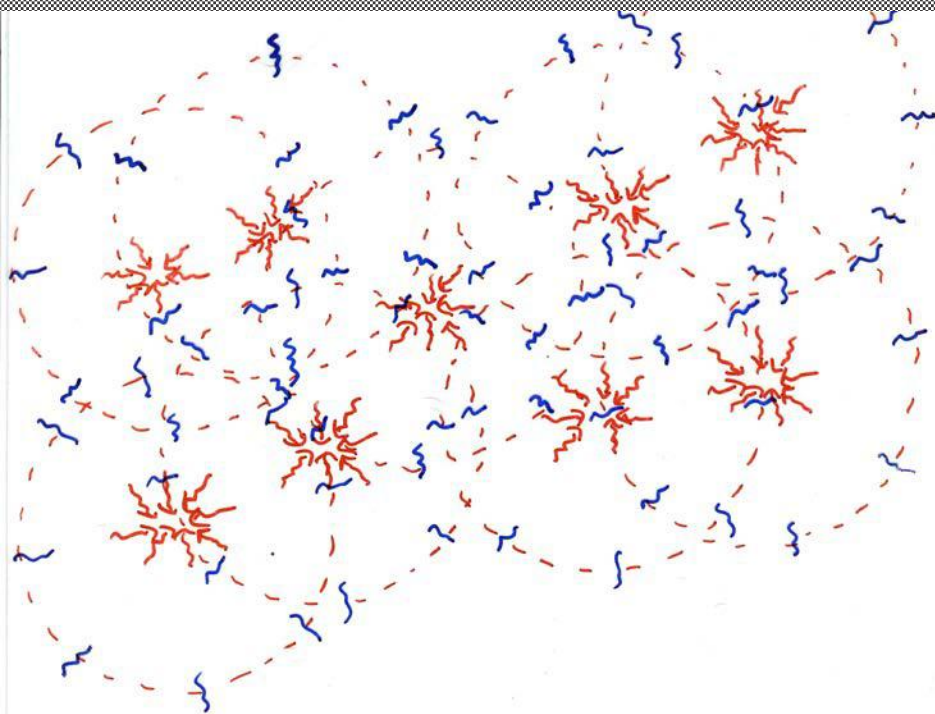
- Decoupling is driven by decreasing X_e

$$\Gamma/H \sim 300 X_e (T/\text{eV})$$

- Last scattering/decoupling occurs at $z = 1100$, shortly after – and driven by – recombination ($z = 1300$)



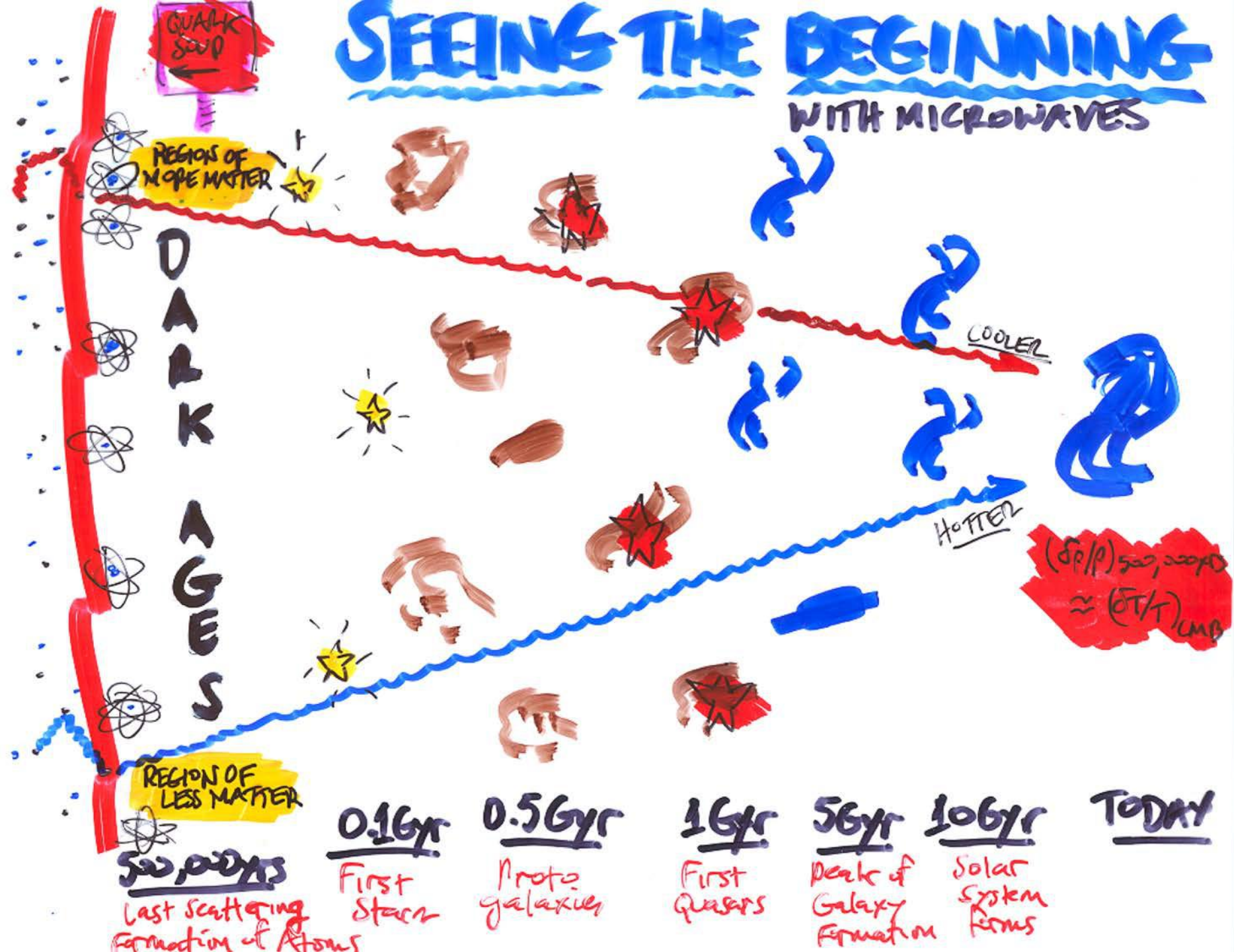
PHOTONS IN THE
400,000yr OLD UNIVERSE



PHOTONS IN THE
400,000yr OLD UNIVERSE
CMB \rightarrow SEEN BY
DIFF. OBSERVERS TODAY

SEEING THE BEGINNING

WITH MICROWAVES

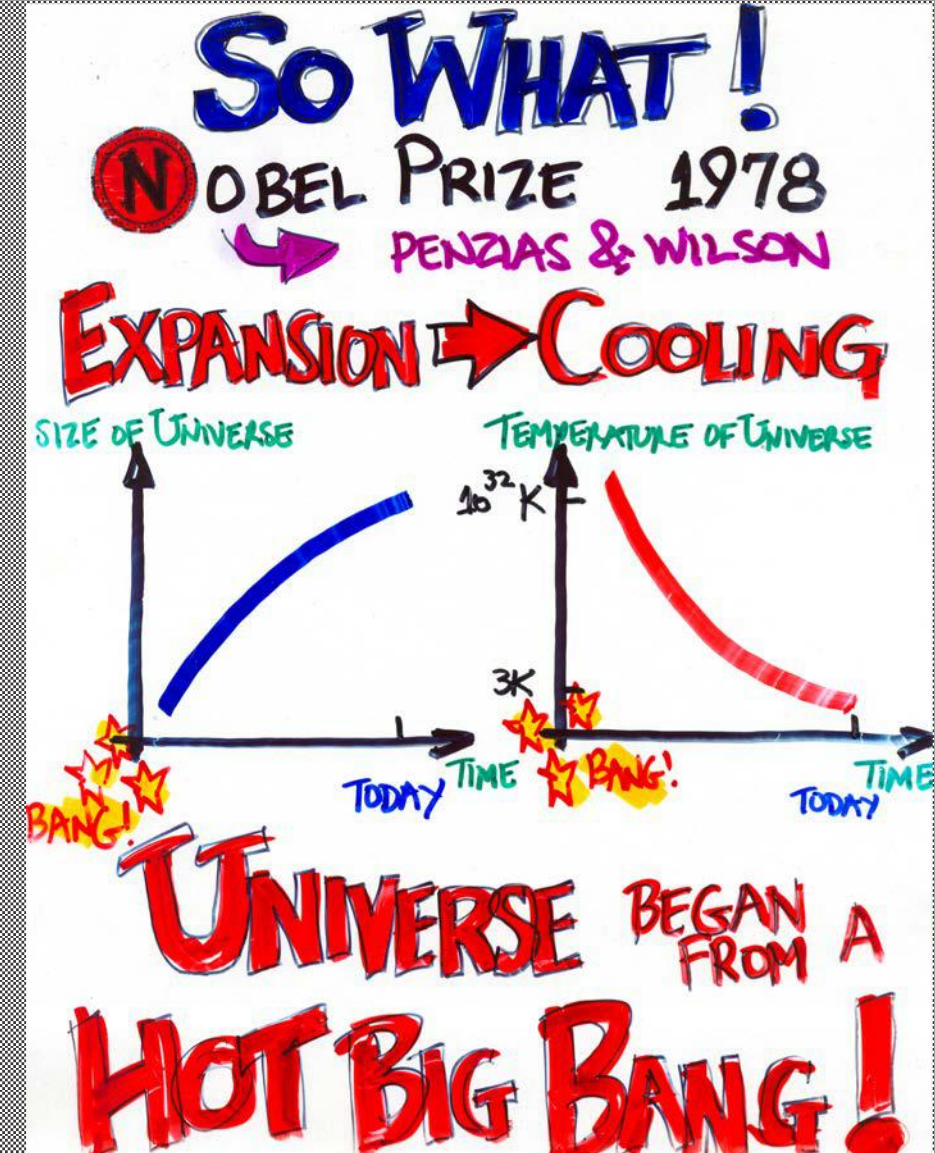


Hot Big Bang!

1964 Arno Penzias & Robert
accidentally discover the Cosmic
Microwave Background



© 2004 Thomson - Brooks/Cole



"Perfect" Blackbody

$$T = 2.7255 \pm 0.0006 \text{ K}$$

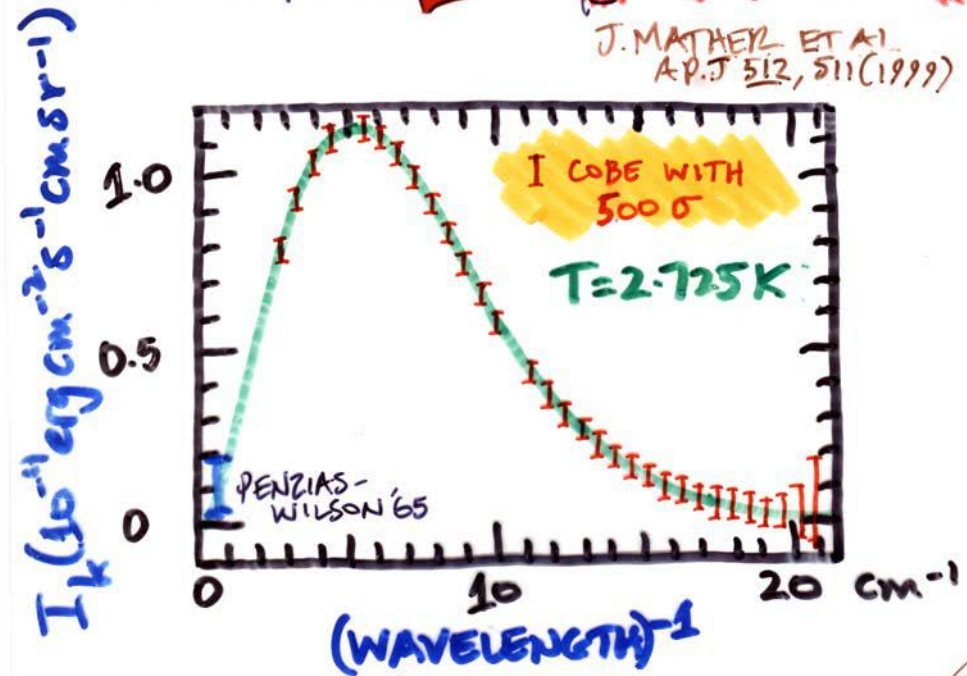
$$n_\gamma = 411 \text{ cm}^{-3}$$

$$\Omega_\gamma \simeq 5 \times 10^{-5}$$



COBE FIRAS

J. MATHER ET AL.
APJ 512, 511 (1999)



$$T = 2.725 \text{ K} \pm 0.00001 \text{ K} \pm 0.001 \text{ K}$$

$$\Delta I / I_{\text{max}} < 0.005 \%$$

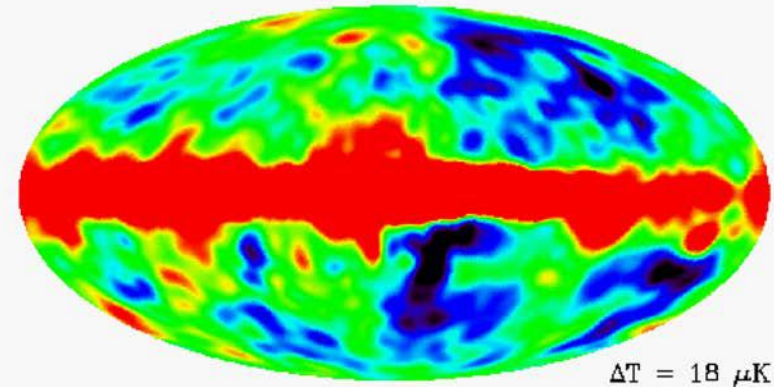
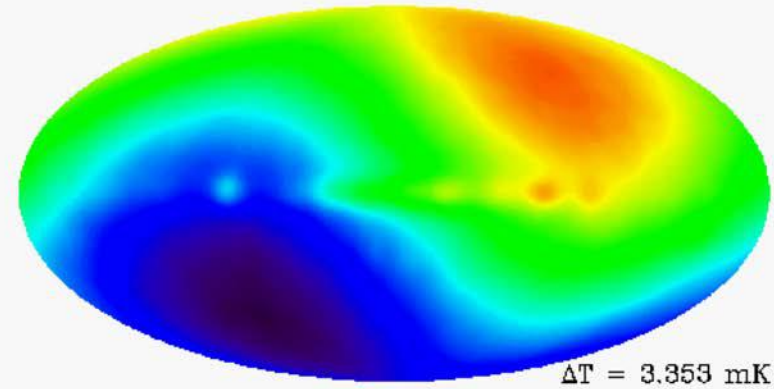
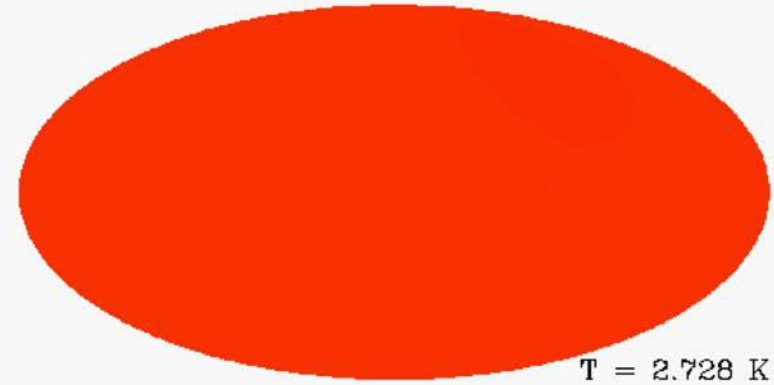
$$|u/kT| < 3.3 \times 10^{-4} \quad y < 2.5 \times 10^{-5} \quad (95\% \text{ c.l.})$$

"BEST BLACK BODY KNOWN"

← ACCURACY OF MEASUREMENT

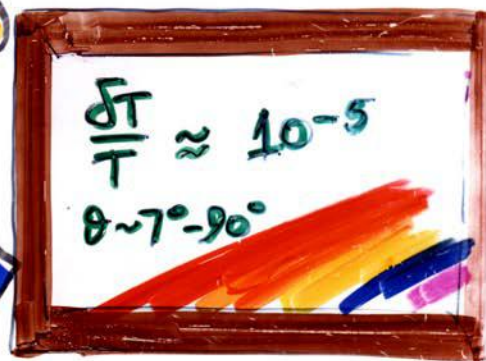
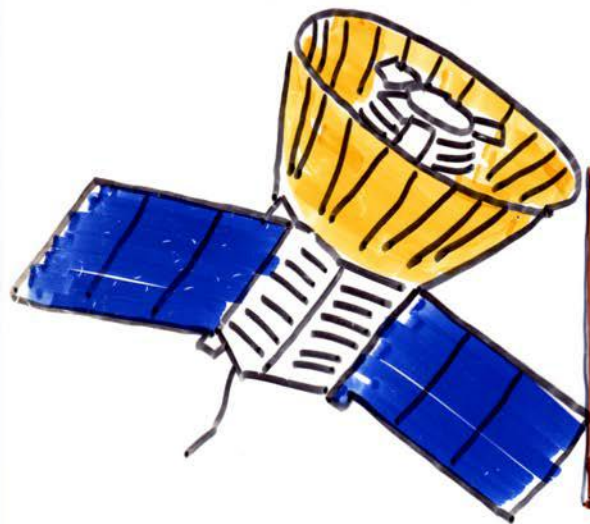
← TEMP. SCALE

1992: COBE DMR
discovers
CMB anisotropy



COBE

23 April 1992



Wow!



"RIPPLES" IN THE MICROWAVE ECHO

TEMPERATURE VARIATION $\approx 30 \mu\text{K}$



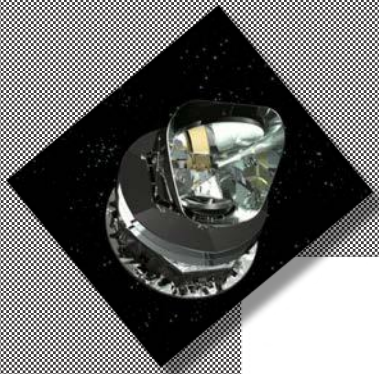
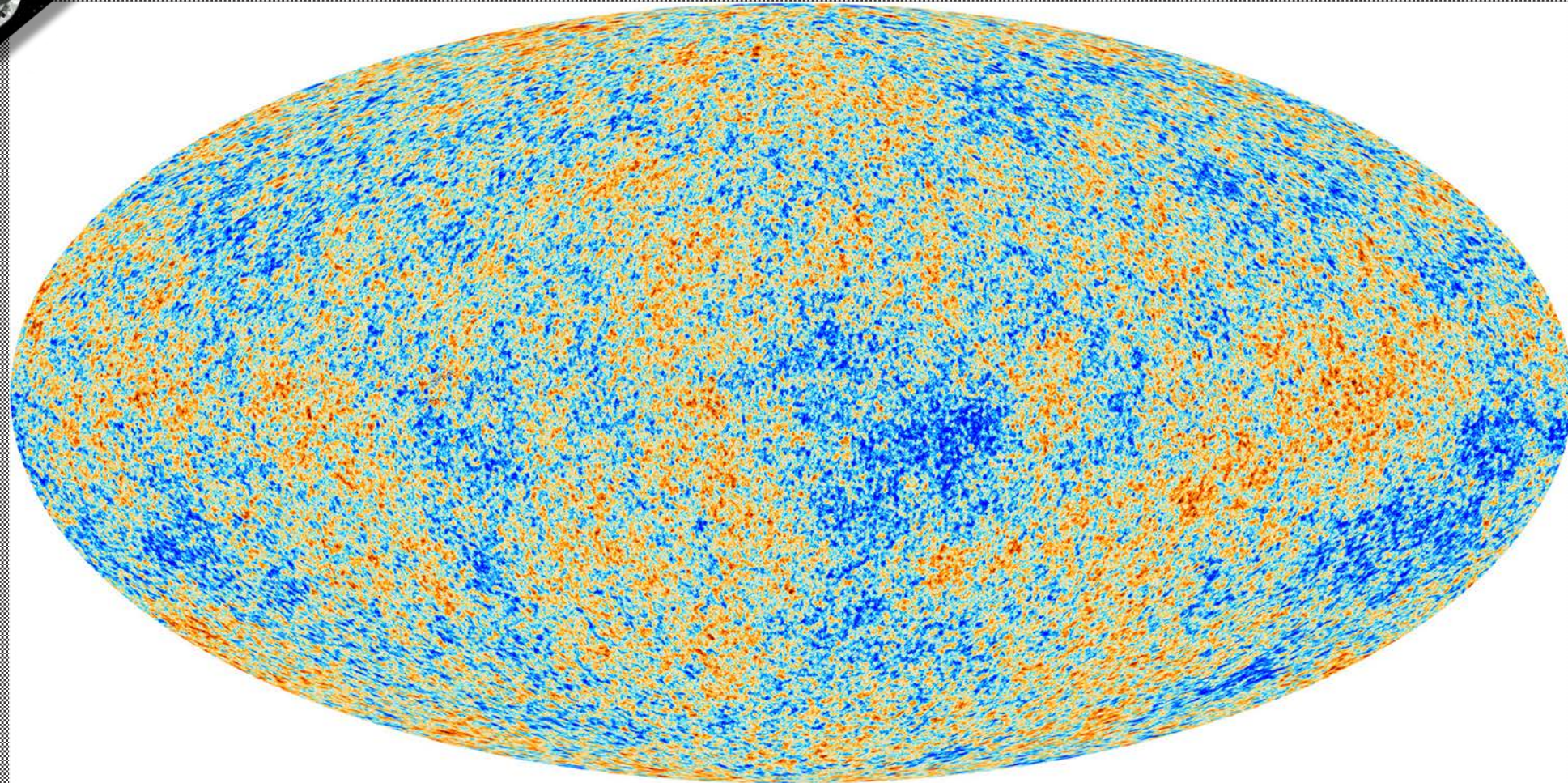
EVIDENCE FOR "PRIMEVAL
LUMPINESS" THAT SEEDS
STRUCTURE (STARS,
GALAXIES, CLUSTERS OF
GALAXIES, SUPERCLUSTERS
VOIDS, WALLS, ...)

S. HAWKING: "GREATEST DISCOVERY
OF ALL TIME"

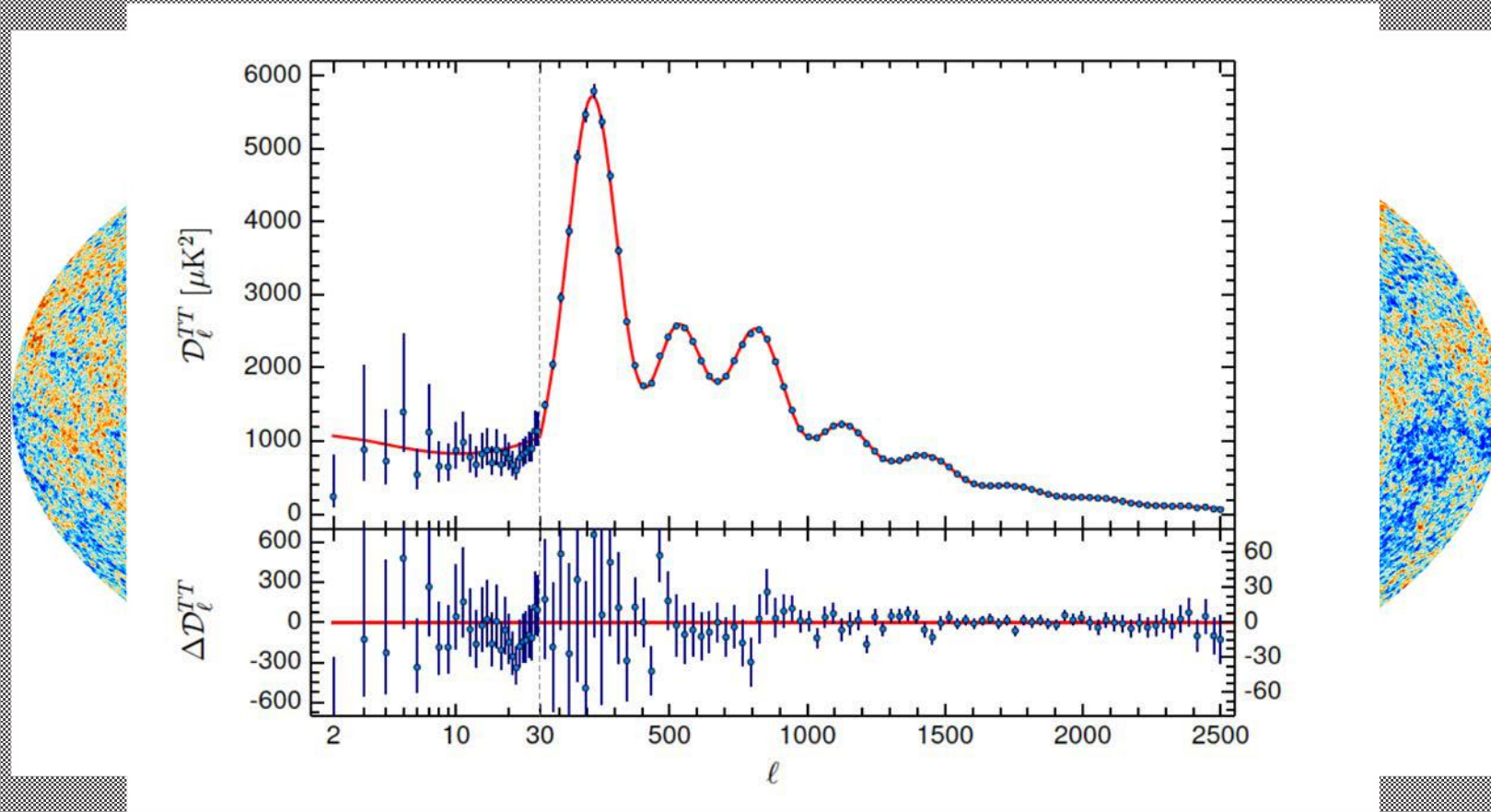


ONLY A SLIGHT OVERSTATEMENT

The Universe at 380,000 years



Best evidence for Λ CDM



6 numbers describe the Universe from
the big bang until today

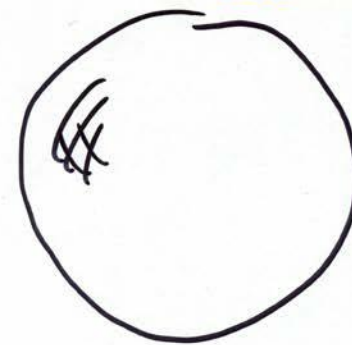
- | | |
|-----------------------------------|------------------|
| 1. Baryon mass density | 4. Tilt |
| 2. CDM mass density | 5. Sound horizon |
| 3. Density perturbation amplitude | 6. Optical depth |

$$R_{\text{curv}} \equiv \frac{H_0^{-1}}{\sqrt{\Omega_0 - 1}}$$

$$\rho_{\text{critical}} = \frac{3H_0^2}{8\pi G} \simeq 10^{-29} \text{ g/cm}^3$$

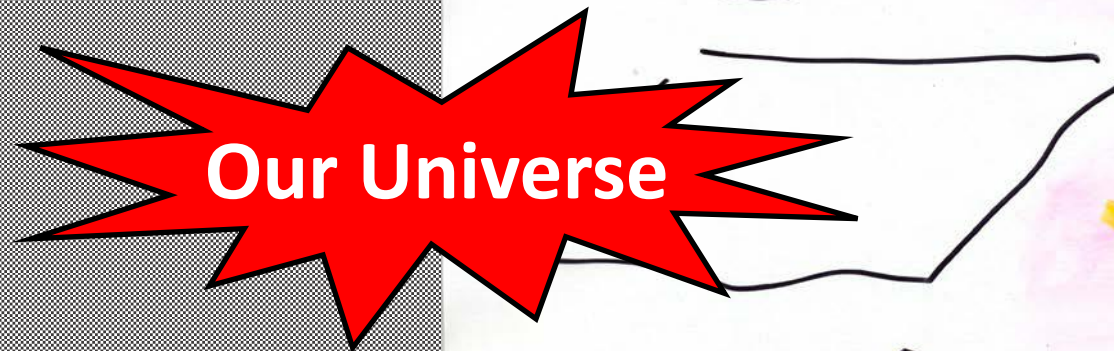
3 KINDS OF BIG BANG UNIVERSES

TWO DIMENSIONAL ANALOGUES



CURVES BACK ON ITSELF (like the surface of a balloon)

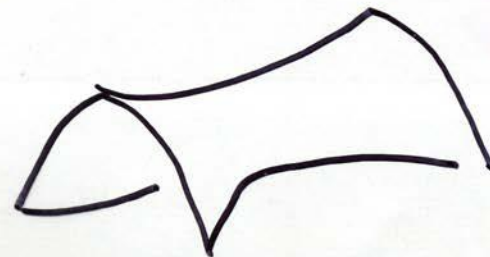
HIGH DENSITY OF MATTER/ENERGY



Our Universe

UNCURVED

"CRITICAL" DENSITY



CURVED LIKE SADDLE

LOW DENSITY

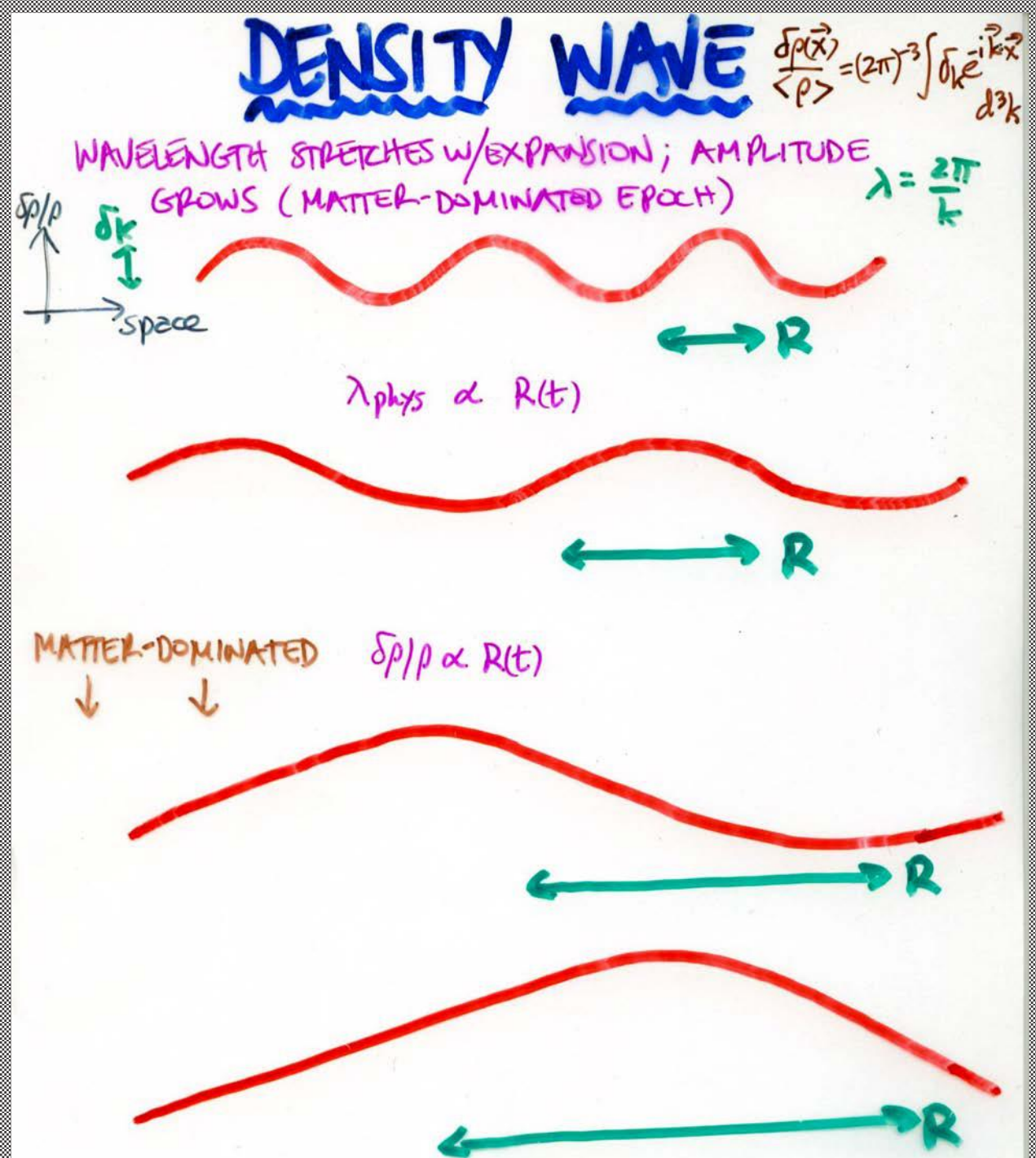
9. The perturbed Universe: beyond homogeneity and isotropy

$$\frac{\delta\rho(\vec{x})}{\bar{\rho}} = \frac{1}{(2\pi)^3} \int \delta_k \exp(-i\vec{k} \cdot \vec{x}) d^3k$$

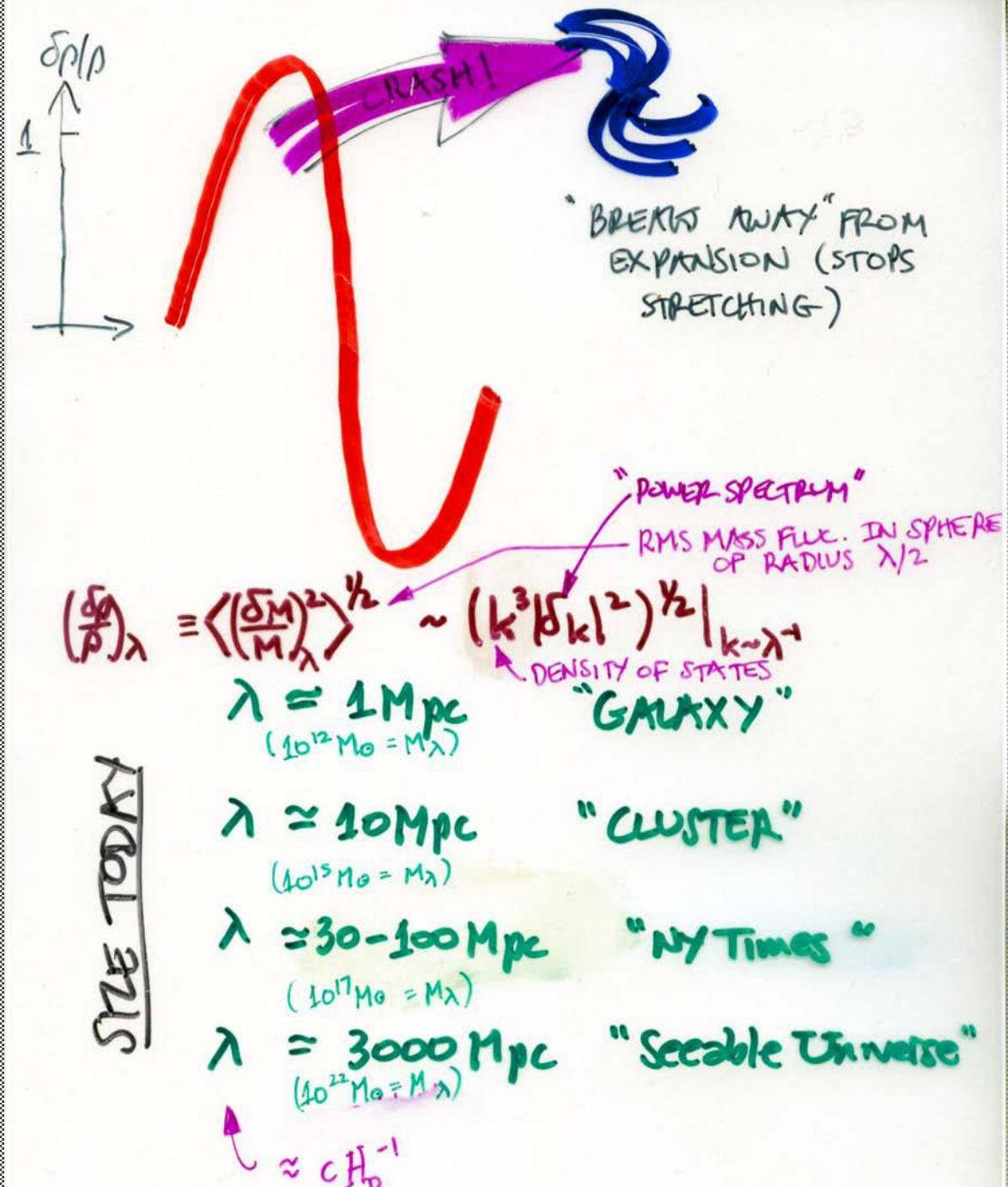
$$\ddot{\delta}_k + 2H\dot{\delta}_k - 4\pi G\rho_M\delta_k = 0$$

1. $H = 0 \Rightarrow \delta_k \propto \exp[(4\pi G\rho_M)^{1/2}t]$ Classic Jeans' instability
2. $H = 2/3t$ (MD) $\delta_k \propto t^{2/3} \propto a(t)$ GROWTH!
3. $H = 1/2t$ (RD) $\delta_k \propto \ln(a)$ No growth
4. $H = \text{const}$ (Λ) $\delta_k \propto \text{const}$ No growth

Expand density field in comoving Fourier components (which contain fixed amount of matter) but whose physical wavelength grows with time



During matter-dominated era, wave amplitudes grow with time (as the scale factor), reach unity and bound structures form and cease expanding



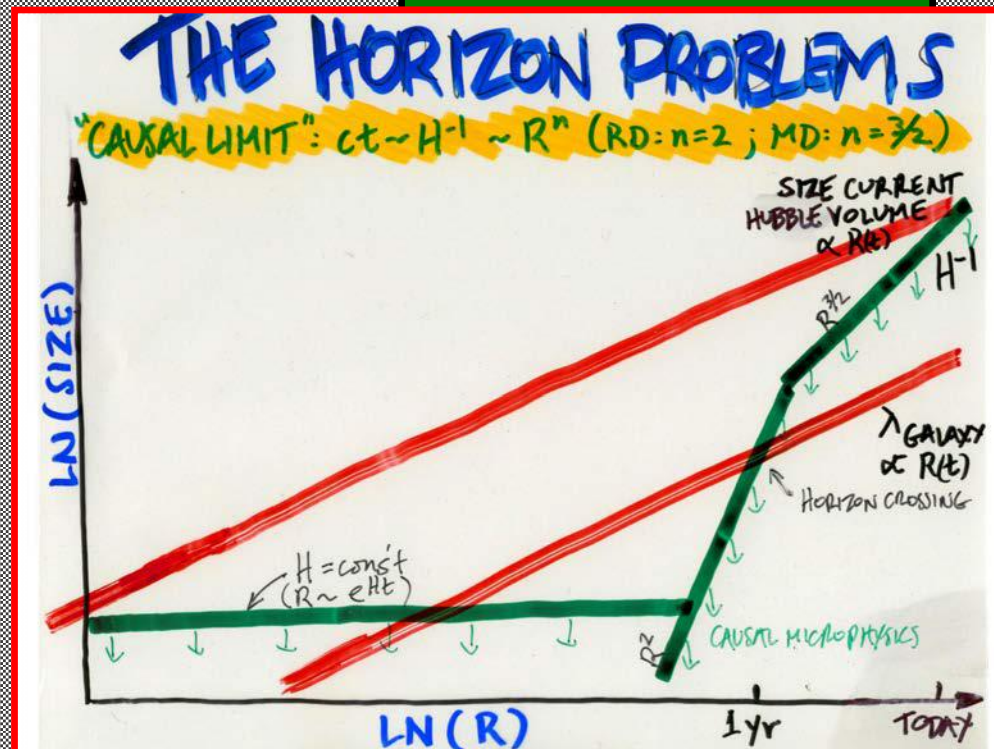
The Physics of Structure Formation (simplified)

1. Creation of
Density
Inhomogeneities

2. Outside
Horizon
 $\lambda > H^{-1}$
Kinematic

3. Inside Horizon
 $\lambda < H^{-1}$
Dynamic
Evolution

Inflation



$$\ddot{\delta}_k + 2H\dot{\delta}_k + (v_s^2 k^2 / R^2)\delta_k - 4\pi G\rho_M\delta_k = 0$$

10. Inflation! The most important idea since the big bang and a pillar of Λ CDM

- “Original intent” (Guth, 1981): First-order phase transition can
 - Solve the horizon, flatness and monopole problems
 - Add on: create the seed inhomogeneities for structure formation (1983)
- Remarkable paper: also proved his idea didn’t work (and then proven it in more detail and rigor with Erick Weinberg, 1983!)
- The slow-roll work around (Linde and Albrecht & Steinhardt, 1982)

PHYSICAL REVIEW D

VOLUME 23, NUMBER 2

15 JANUARY 1981

Inflationary universe: A possible solution to the horizon and flatness problems

Alan H. Guth*

Stanford Linear Accelerator Center, Stanford University, Stanford, California 94305

(Received 11 August 1980)

The standard model of hot big-bang cosmology requires initial conditions which are problematic in two ways: (1) The early universe is assumed to be highly homogeneous, in spite of the fact that separated regions were causally disconnected (horizon problem); and (2) the initial value of the Hubble constant must be fine tuned to extraordinary accuracy to produce a universe as flat (i.e., near critical mass density) as the one we see today (flatness problem). These problems would disappear if, in its early history, the universe supercooled to temperatures 28 or more orders of magnitude below the critical temperature for some phase transition. A huge expansion factor would then result from a period of exponential growth, and the entropy of the universe would be multiplied by a huge factor when the latent heat is released. Such a scenario is completely natural in the context of grand unified models of elementary-particle interactions. In such models, the supercooling is also relevant to the problem of monopole suppression. Unfortunately, the scenario seems to lead to some unacceptable consequences, so modifications must be sought.



TINY ($\ll 1\text{cm}$) BIT
OF UNIVERSE IS FLAT
& SMOOTH (but too
small to contain all we
see today)



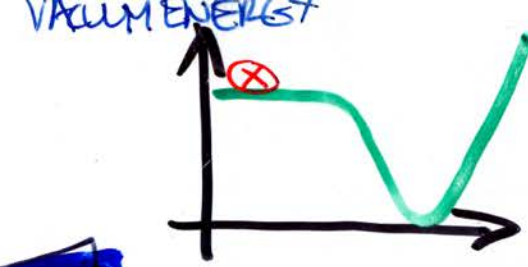
ALL THAT
WE CAN
SEE TODAY
(STILL SMOOTH
& FLAT)

Solving the Flatness,
Horizon Problems

COSMIC

Guth; Linde;
Albrecht-Steinhardt

INFLATION



SHORT PERIOD OF RAPID EXPANSION DRIVEN
BY "FALSE-VACUUM" ENERGY

★ MORE EXPANSION IN 10^{-32} sec THAN NEXT 15 BYR

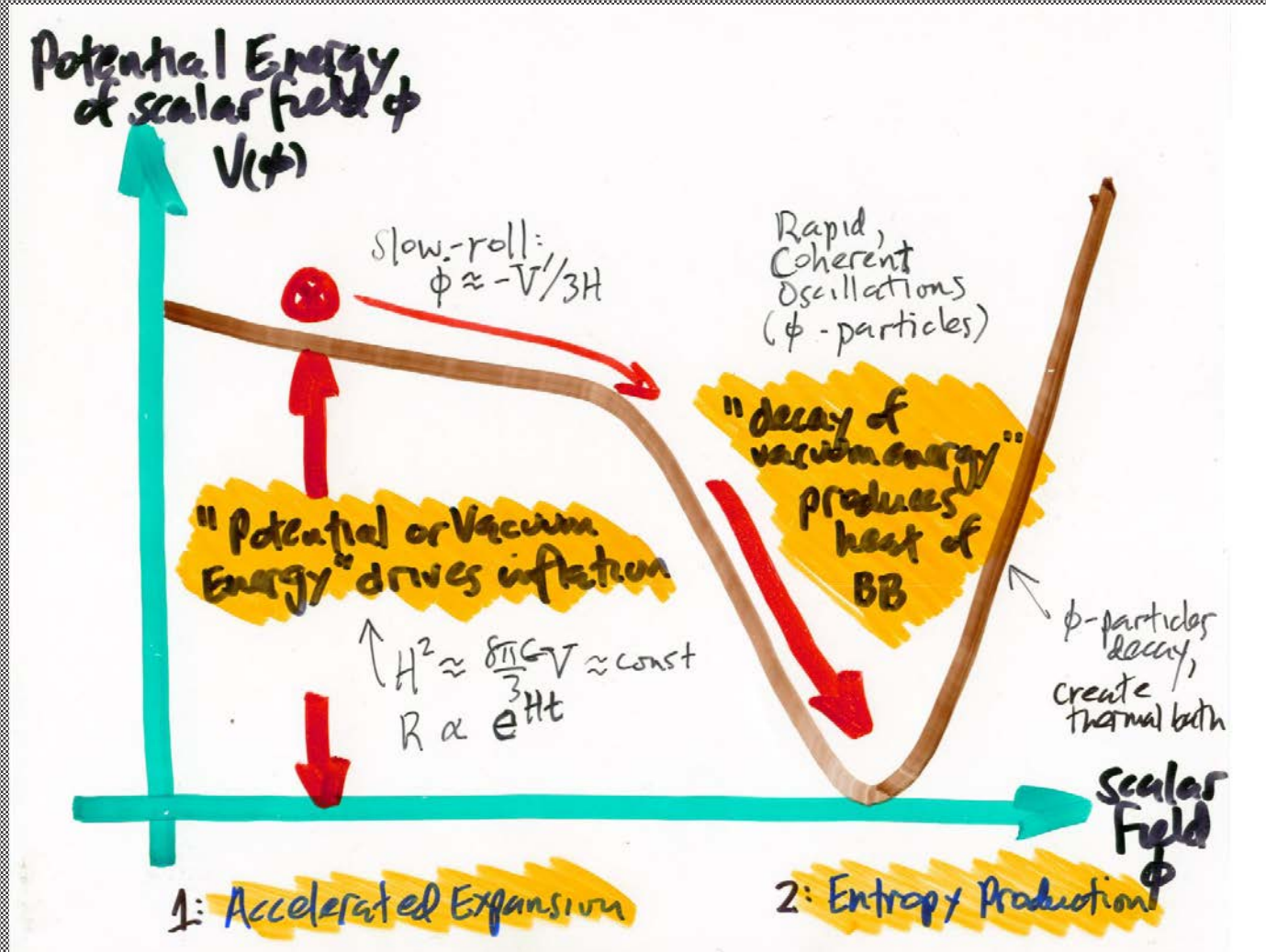
★ OBSERVED UNIVERSE BEGAN FROM INCREDIBLY
SMALL PATCH → FLAT & SMOOTH



OBSERVED UNIVERSE

★ MICROSCOPIC PHENOMENA (QUANTUM FLUCTUATIONS)
CAN INFLUENCE MACROSCOPIC SCALES
→ QUANTUM ORIGIN OF "LUMPINESS"

Slow-roll inflation: scalar-field dynamics



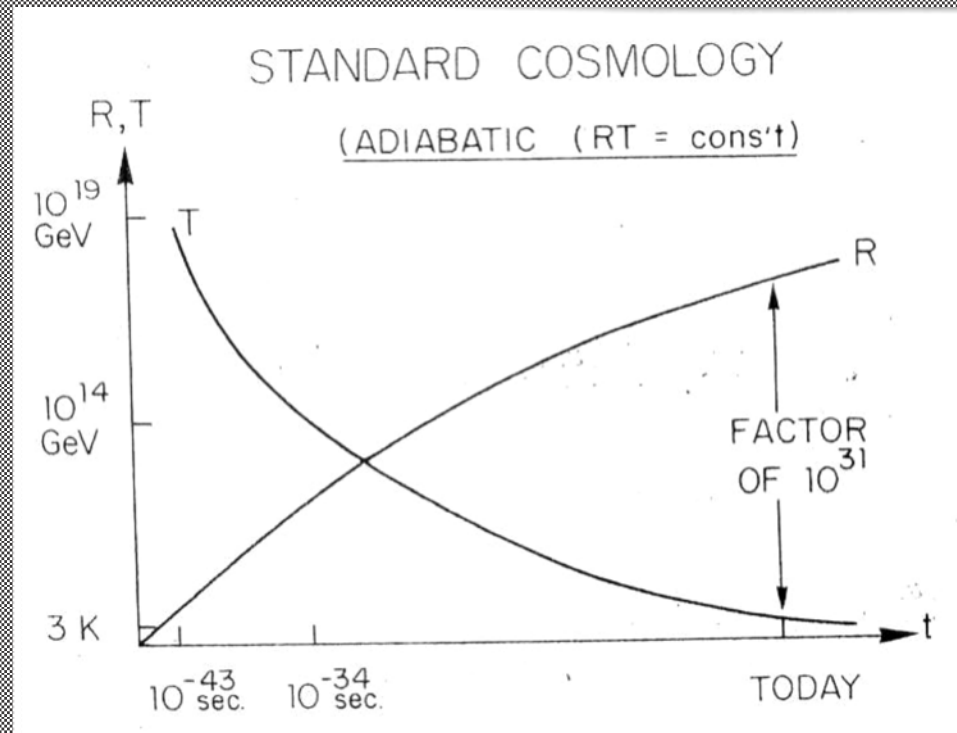
$$\rho = \frac{1}{2}\dot{\phi}^2 + V(\phi) \quad p = \frac{1}{2}\dot{\phi}^2 - V(\phi)$$

$$w = \frac{\frac{1}{2}\dot{\phi}^2 - V(\phi)}{\frac{1}{2}\dot{\phi}^2 + V(\phi)}$$

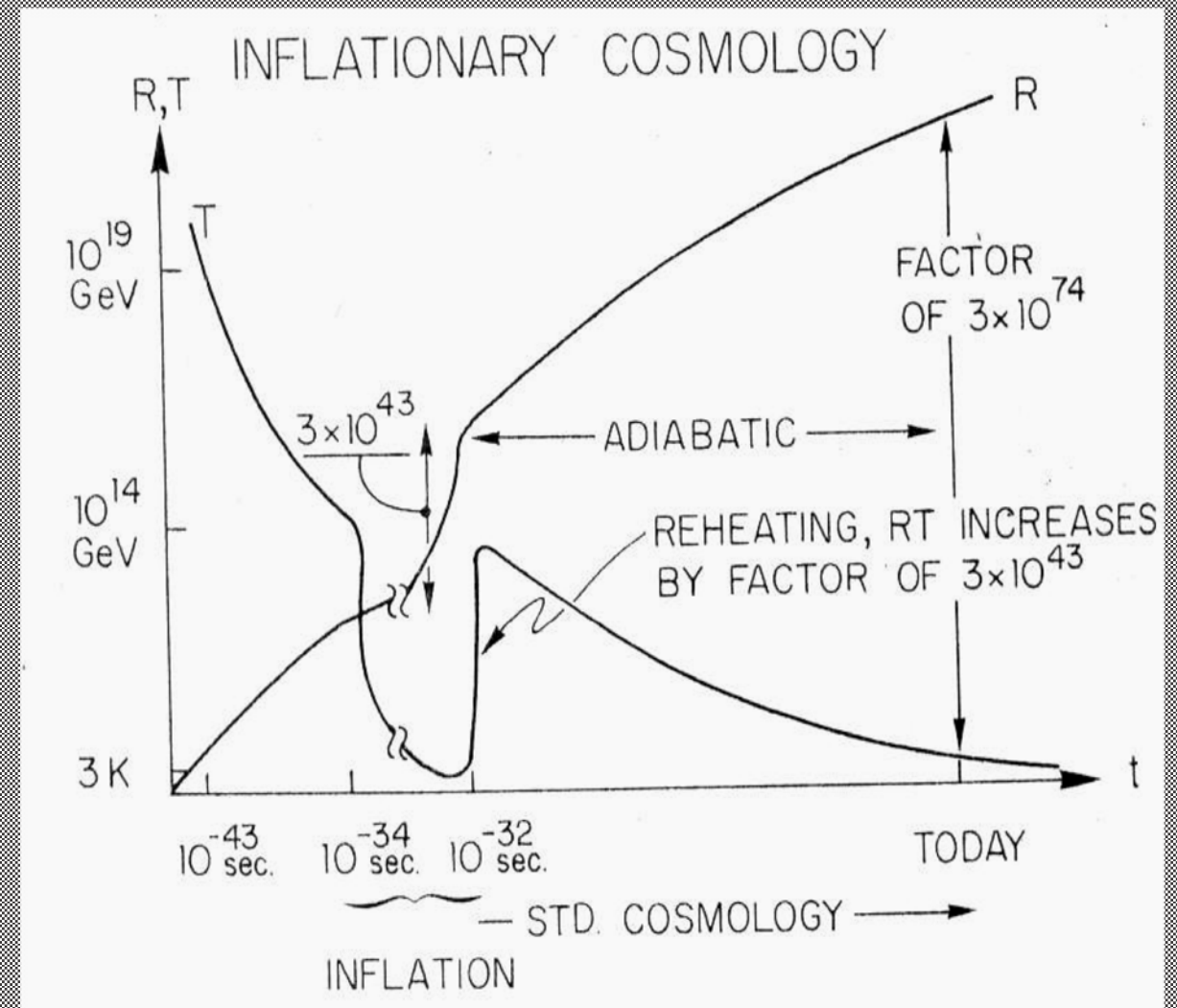
$$\ddot{\phi} + 3H\dot{\phi} + V'(\phi) [+ \Gamma\dot{\phi}] = 0$$

Entropy Production/Reheating

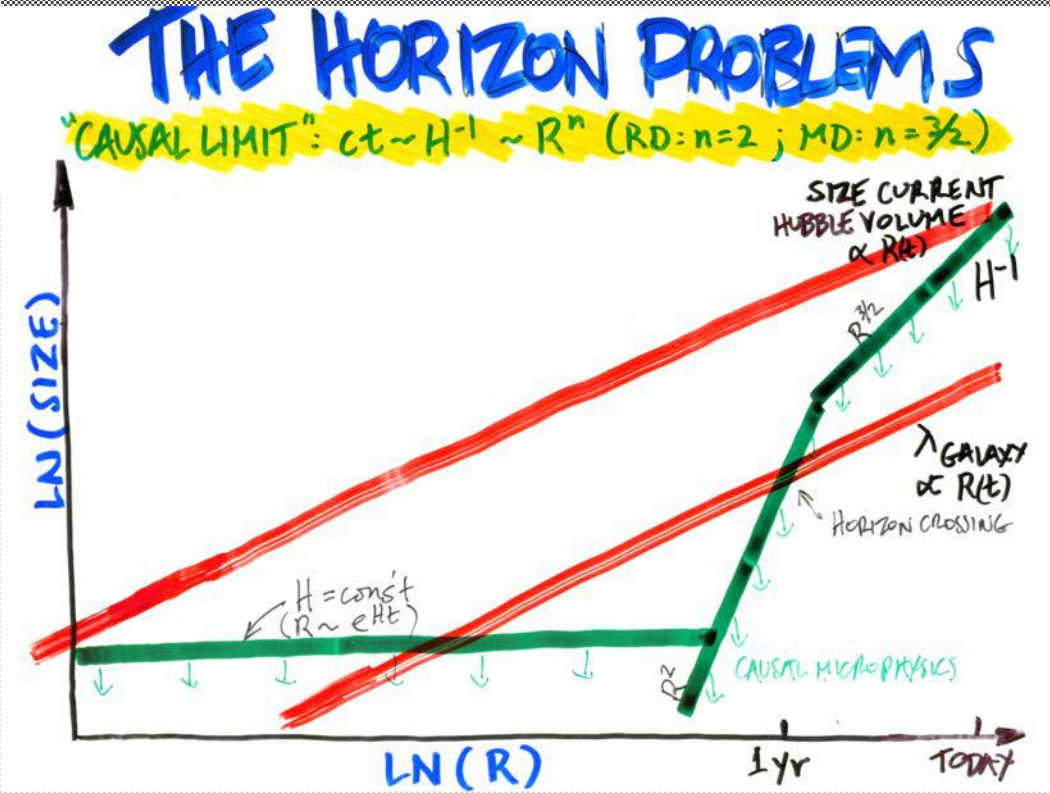
Recall: entropy per comoving volume $\sim R^3 T^3$



Adiabatic: Constant Number of Photons per co-moving Volume, i.e., $RT = \text{const}$



Quantum Fluctuations Seed Density Perturbations



QM FLUCTUATIONS

in the "INFLATON" ϕ

$\Delta\phi \approx H/2\pi$

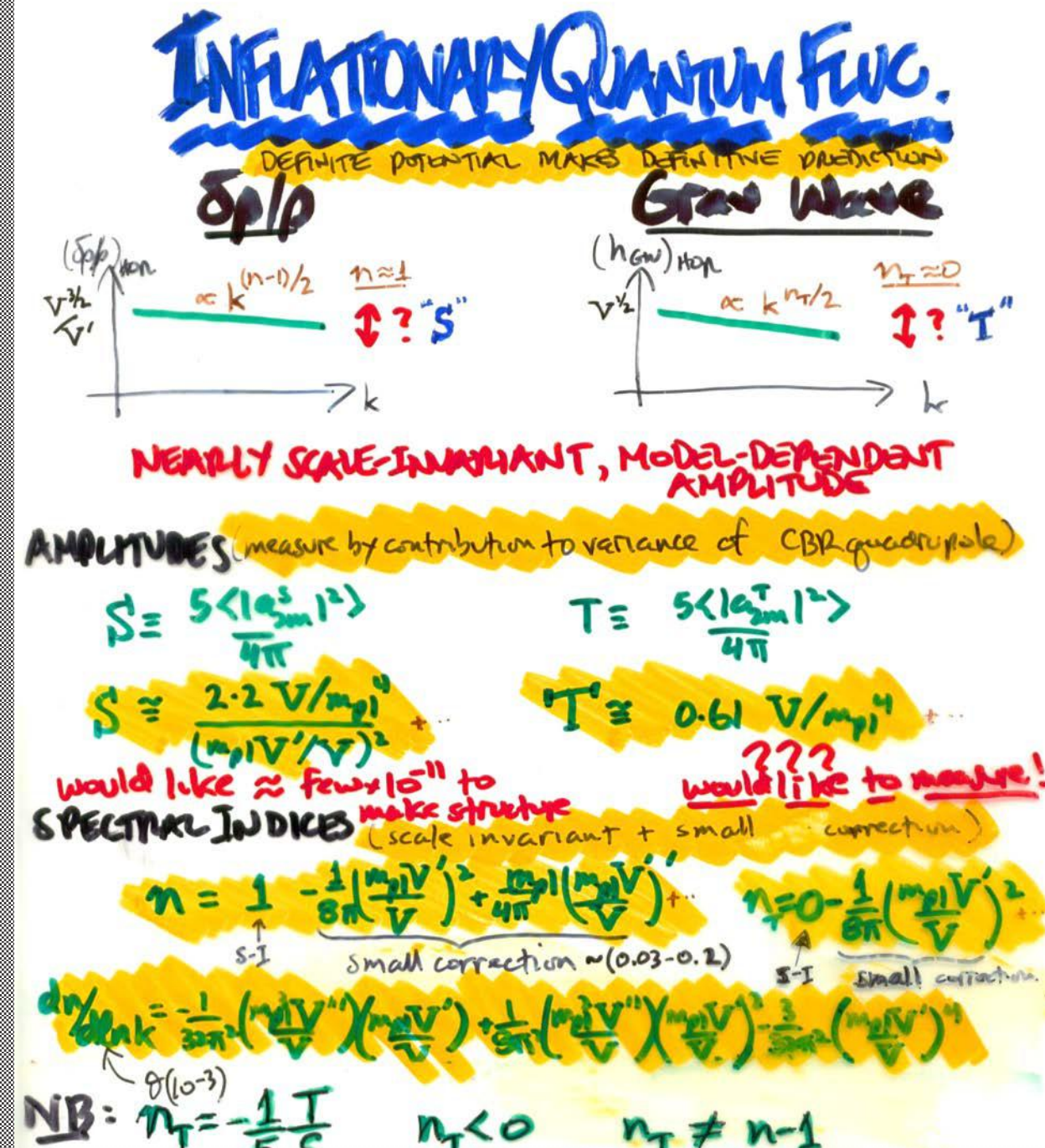
ENERGY DENSITY PERTURBATIONS

$\delta\rho_\phi = V'\Delta\phi$

DENSITY PERTURBATIONS AFTER REHEATING

Good news, bad news:

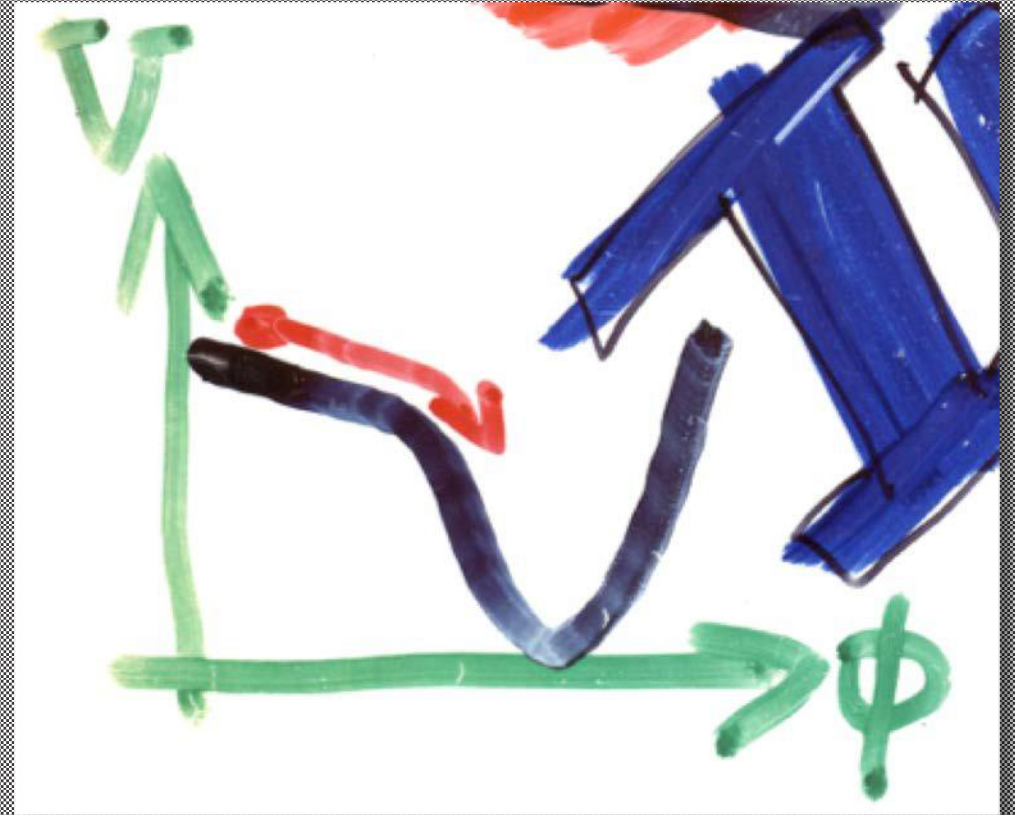
Given a scalar potential $V(\phi)$,
can compute all
observables in terms of
 V , V' and V''



Inflation's predictions

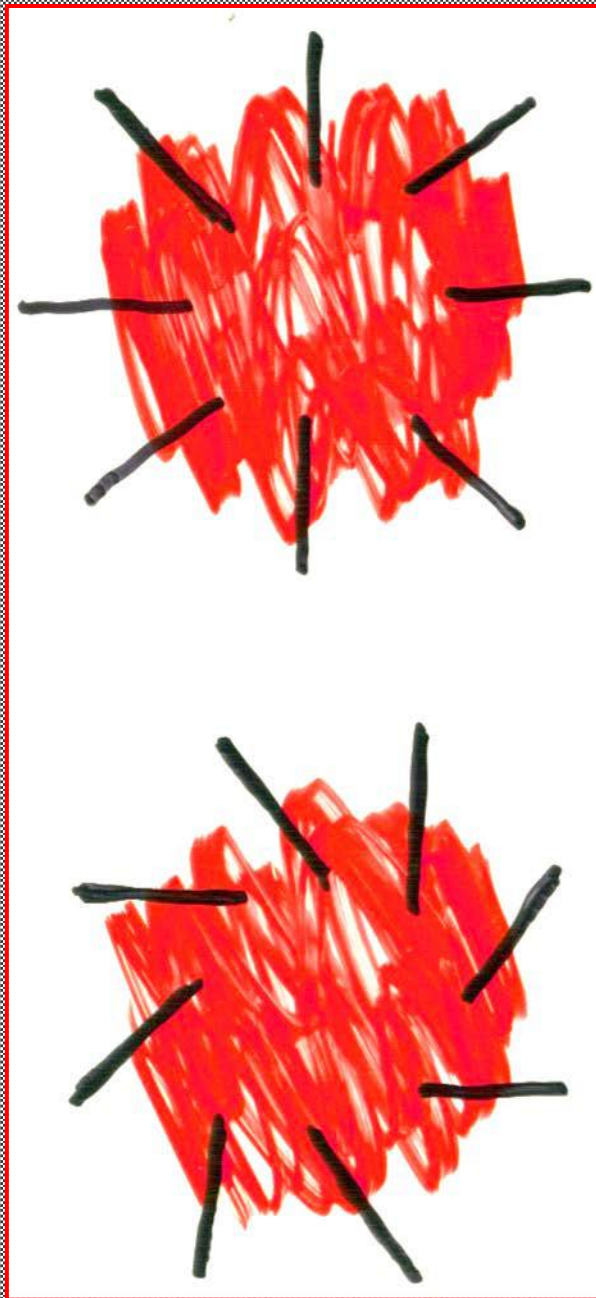
$$\begin{aligned}
 P(k) &= \frac{1024\pi^3}{75} \frac{k}{H_0^4} \frac{V_*^3}{m_{\text{Pl}}^6 V_*'^2} \left(\frac{k}{k_*}\right)^{n-1} T^2(k) \\
 n-1 &= -\frac{1}{8\pi} \left(\frac{m_{\text{Pl}} V_*'}{V_*}\right)^2 + \frac{m_{\text{Pl}}}{4\pi} \left(\frac{m_{\text{Pl}} V_*'}{V_*}\right)' \\
 \frac{dn}{d \ln k} &= -\frac{1}{32\pi^2} \left(\frac{m_{\text{Pl}}^3 V_*'''}{V_*}\right) \left(\frac{m_{\text{Pl}} V_*'}{V_*}\right) \\
 &\quad + \frac{1}{8\pi^2} \left(\frac{m_{\text{Pl}}^2 V_*''}{V_*}\right) \left(\frac{m_{\text{Pl}} V_*'}{V_*}\right)^2 - \frac{3}{32\pi^2} \left(m_{\text{Pl}} \frac{V_*'}{V_*}\right)^4 \\
 T(q) &= \frac{\ln(1 + 2.34q)/2.34q}{[1 + 3.89q + (16.1q)^2 + (5.46q)^3 + (6.71q)^4]^{1/4}},
 \end{aligned}$$

$$\begin{aligned}
 P_T(k) &\equiv \langle |h_k|^2 \rangle = \frac{8}{3\pi} \frac{V_*}{m_{\text{Pl}}^4} \left(\frac{k}{k_*}\right)^{n_T-3} T_T^2(k) \\
 n_T &= -\frac{1}{8\pi} \left(\frac{m_{\text{Pl}} V_*'}{V_*}\right)^2 \\
 \frac{dn_T}{d \ln k} &= \frac{1}{32\pi^2} \left(\frac{m_{\text{Pl}}^2 V_*''}{V_*}\right) \left(\frac{m_{\text{Pl}} V_*'}{V_*}\right)^2 - \frac{1}{32\pi^2} \left(\frac{m_{\text{Pl}} V_*'}{V_*}\right)^4 = -n_T[(n-1) - n_T] \\
 T_T(k) &\simeq \left[1 + \frac{4}{3} \frac{k}{k_{\text{EQ}}} + \frac{5}{2} \left(\frac{k}{k_{\text{EQ}}}\right)^2\right]^{1/2},
 \end{aligned}$$

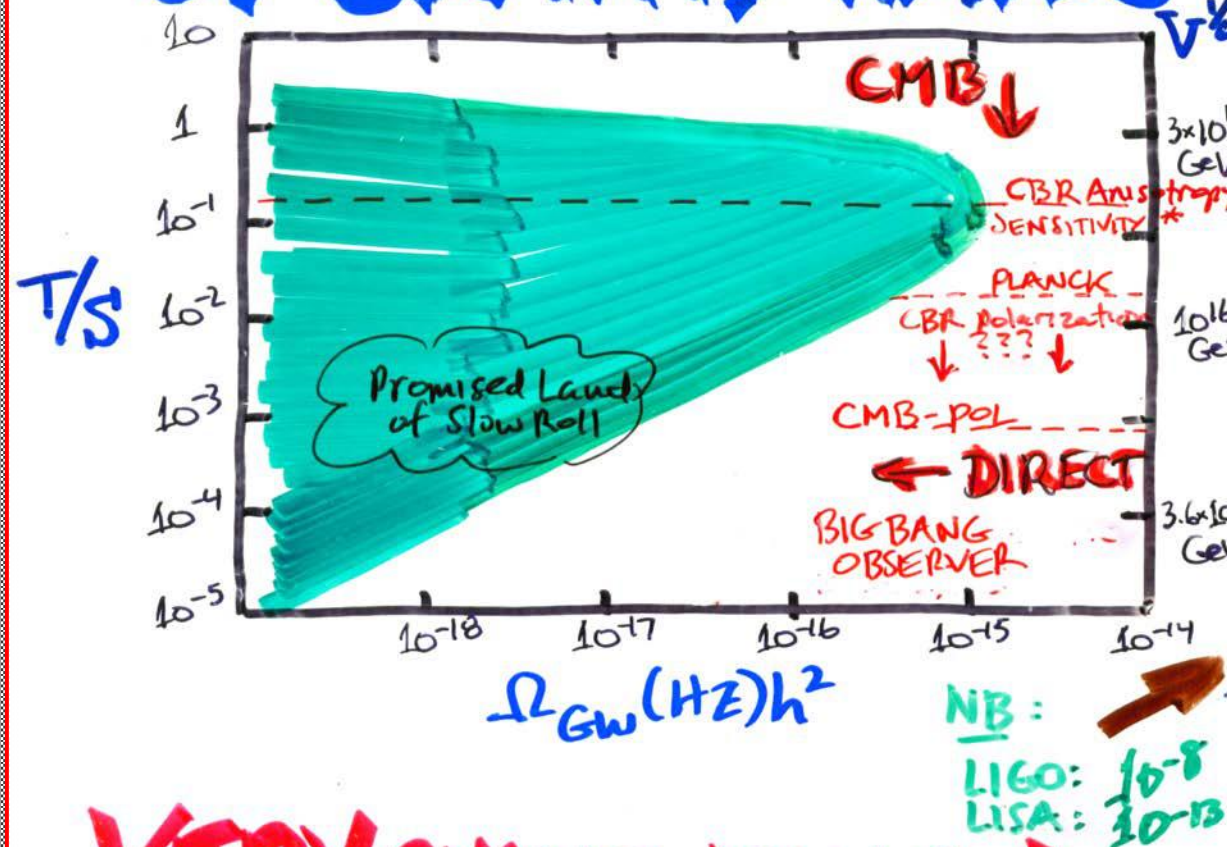


Key predictions of inflation

- Flat Universe (at a time when the data said: $\Omega_0 = 0.1!$)
- Almost scale-invariant, Gaussian curvature perturbations (not precisely a power-law and n not precisely 1.0)
- Almost scale-invariant spectrum of GWs: in the B-mode polarization
- Consistency relationship: $T/S = 7n_T$
- ... and it explains the isotropy/homogeneity, quark soup and absence of superheavy magnetic monopoles
- BUT, it is an incomplete theory: no std model, temporary fix, what about the BB singularity, and on and on



DETECTION OF GRAVITY WAVES



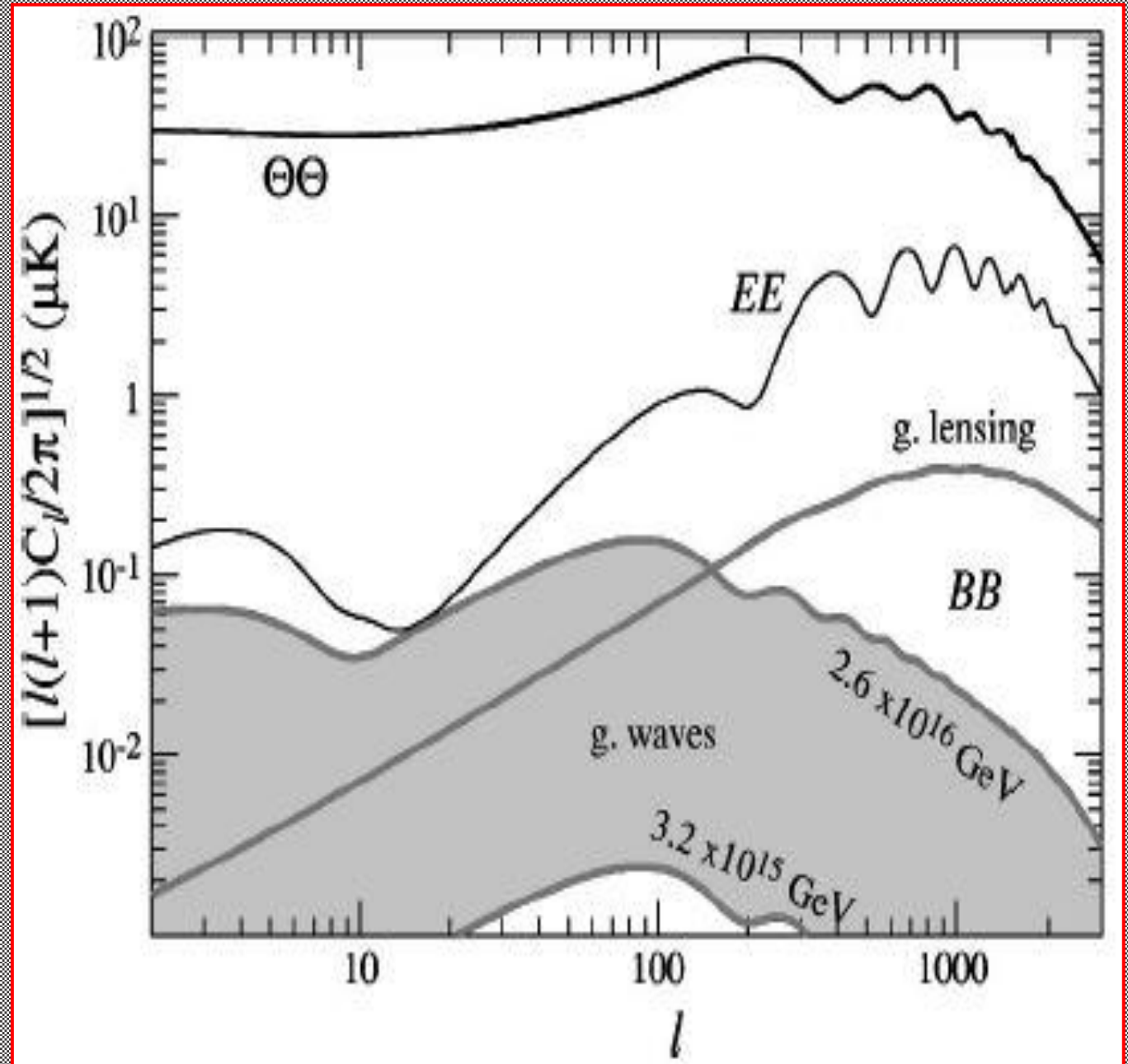
VERY CHALLENGING!

"DOUBLE DETECTION" $\Rightarrow n_t$ to ± 0.03

NanoGrav: 10^{-8} at 10^{-9} Hz

CMB Anisotropy from Gravity Waves

- $\Theta\Theta$ = GW temp
- EE = E mode (scalar)
- g lensing: grav lensing of EE
- BB/g waves = GW B-mode
- nanoKelvin cosmology!



GWs and B-mode CMB polarization

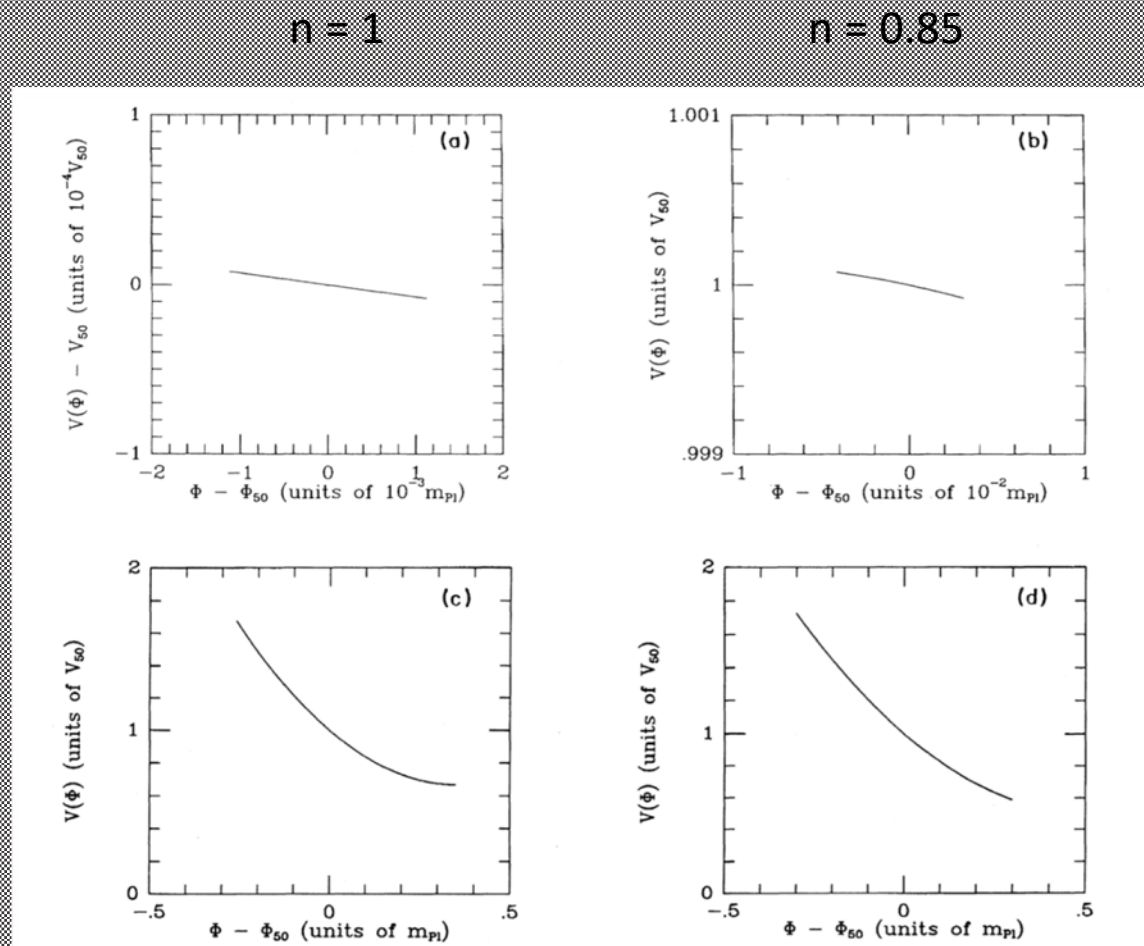
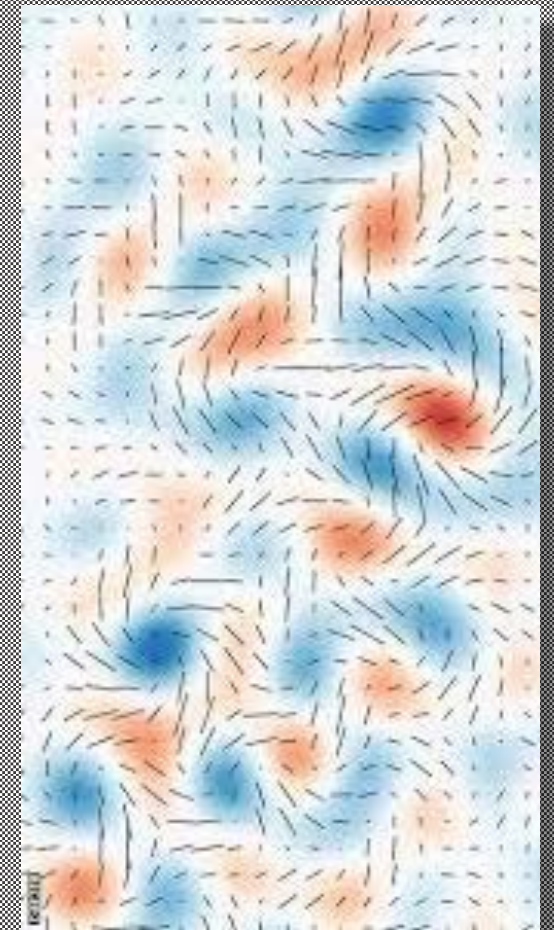


FIG. 5. The four generic inflationary potentials: (a) $n=1$ and $T/S=1.4 \times 10^{-5}$, with the COBE DMR normalization $V_{50}^{1/4}=2.0 \times 10^{15}$ GeV; (b) $n=0.85$ and $T/S=1.4 \times 10^{-4}$, $V_{50}^{1/4}=3.6 \times 10^{15}$ GeV; (c) $n=1$ and $T/S=1$, $V_{50}^{1/4}=2.9 \times 10^{16}$ GeV; and (d) $n=0.85$ and $T/S=1$, $V_{50}^{1/4}=2.9 \times 10^{16}$ GeV. (a)–(d) correspond to cases (1)–(4) in the text.

$r \ll 1$

$r = 1$



PHYSICAL REVIEW D

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Recovering the inflationary potential

Michael S. Turner

Departments of Physics and of Astronomy and Astrophysics, Enrico Fermi Institute, The University of Chicago, Chicago, Illinois 60637-1433

and NASA/Fermilab Astrophysics Center, Fermi National Accelerator Laboratory, Batavia, Illinois 60510-0500
(Received 27 July 1993)

A procedure is developed for the recovery of the inflationary potential over the interval that affects astrophysical scales (≈ 1 Mpc to 10^8 Mpc). The amplitudes of the scalar and tensor metric perturbations and their power-spectrum indices, which in principle can be inferred from large-angle CBR anisotropy and other cosmological data, determine the value of the inflationary potential and its first two derivatives. From these, the inflationary potential can be reconstructed in a Taylor series and the consistency of the inflationary hypothesis tested. Examples are presented, and the effect of observational uncertainties is discussed.

BICEP/Keck leading the way

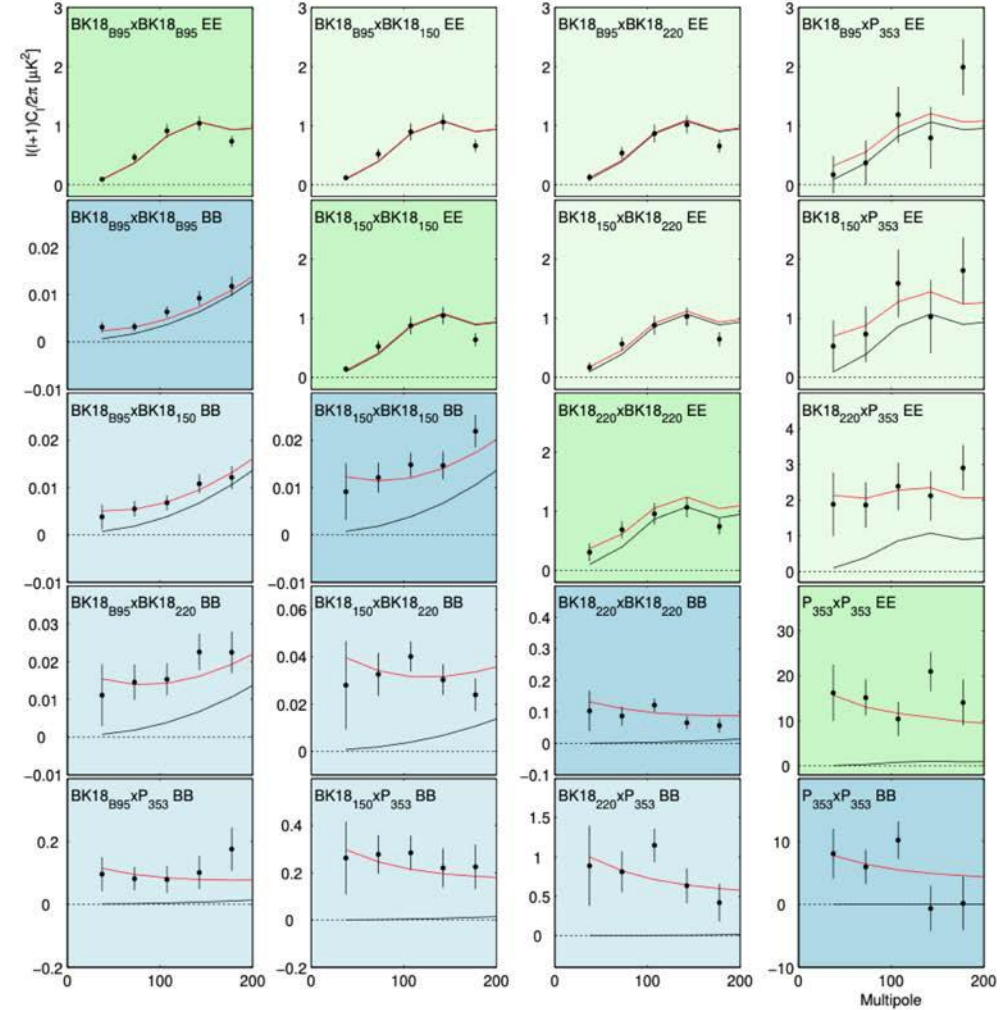
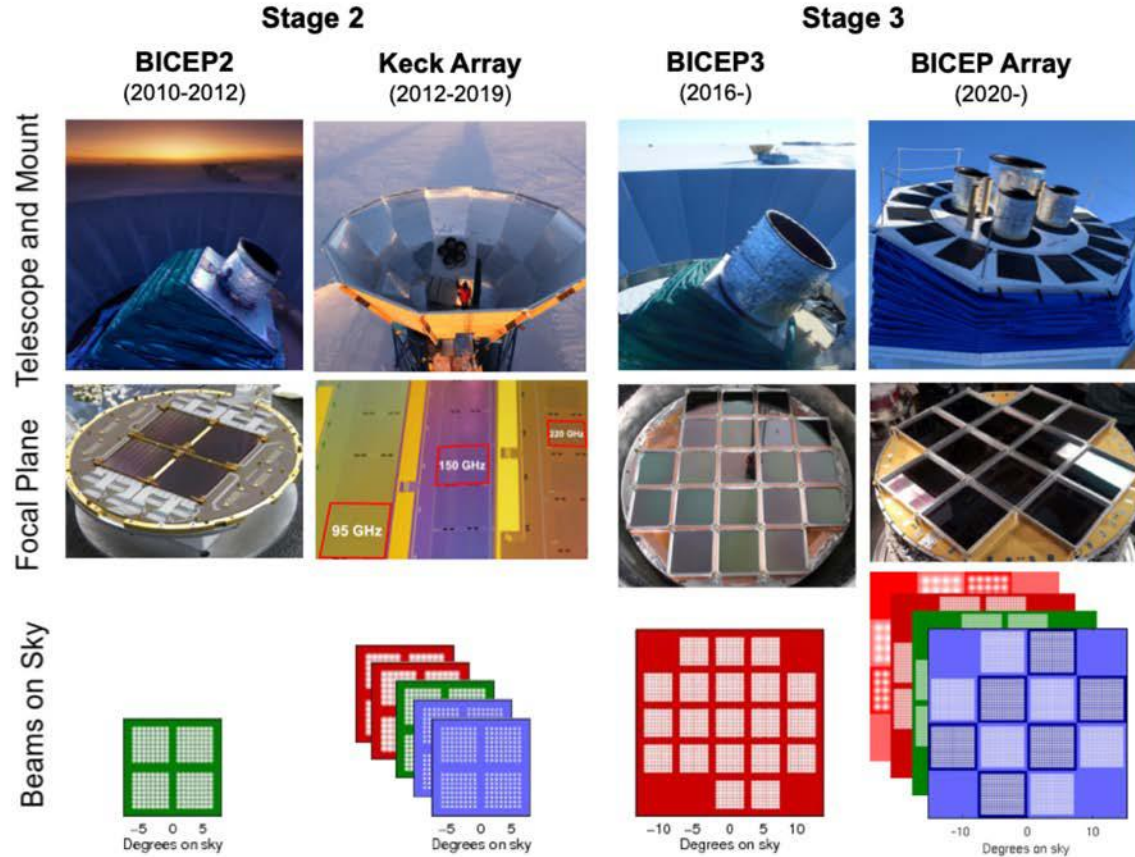


Figure 4. *BB* (blue) and *EE* (green) auto- and cross-spectra from BICEP3 95 GHz, BICEP2/Keck 150 GHz, Keck 220 GHz and Planck 353 GHz maps. The black lines are the Λ CDM model expectation values, while the red lines are the Λ CDM+foreground expectation values from the foreground best fit of our previous BK15 analysis. The *EE* spectra are computed as a demonstration under the assumption $EE/BB = 2$ for dust. *EE* spectra are not included in our likelihood analysis.

The Latest Constraints on Inflationary B-modes from the BICEP/Keck Telescopes

The BICEP/Keck Collaboration: P. A. R. Ade^a, Z. Ahmed^b, M. Amiri^c, D. Barkats^d, R. Basu Thakur^e, C. A. Bischoff^f, D. Beck^{b,g}, J. J. Bock^{e,h}, H. Boenish^d, E. Bullockⁱ, V. Buza^j, J. R. Cheshire IVⁱ, J. Connors^d, J. Cornelison^d, M. Crumrine^k, A. Cukierman^{g,b}, E. V. Denison^l, M. Dierickx^d, L. Duband^m, M. Eiben^d, S. Fatigoni^c, J. P. Filippini^{n,o}, S. Fliescher^k, C. Giannakopoulos^f, N. Goeckner-Wald^g, D. C. Goldfinger^d, J. Grayson^g, P. Grimes^d, G. Hall^k, G. Halal^g, M. Halpern^c, E. Hand^f, S. Harrison^d, S. Henderson^b, S. R. Hildebrandt^{e,h}, G. C. Hilton^l, J. Hubmayr^l, H. Hui^e, K. D. Irwin^{g,b,l}, J. Kang^{g,e}, K. S. Karkare^{d,j}, E. Karpel^g, S. Kefeli^e, S. A. Kernasovskiy^g, J. M. Kovac^{d,p}, C. L. Kuo^{g,b}, K. Lau^{k,*}, E. M. Leitch^l, A. Lennoxⁿ, K. G. Megerian^h, L. Minutolo^e, L. Moncelsi^e, Y. Nakato^g, T. Namikawa^q, H. T. Nguyen^h, R. O'Brien^{e,h}, R. W. Ogburn IV^{g,b}, S. Palladino^f, M. Petroff^d, T. Prouve^m, C. Pryke^{k,i}, B. Racine^{d,r}, C. D. Reintsema^l, S. Richter^d, A. Schillaci^e, R. Schwarz^k, B. L. Schmitt^d, C. D. Sheehy^a, B. Singari^l, A. Soliman^e, T. St. Germaine^{d,p}, B. Steinbach^e, R. V. Sudiwala^a, G. P. Teply^e, K. L. Thompson^{g,b}, J. E. Tolan^g, C. Tucker^a, A. D. Turner^h, C. Umiltà^{f,n}, C. Vergès^d, A. G. Viereggs^j, A. Wandui^e, A. C. Weber^h, D. V. Wiebe^e, J. Willmert^k, C. L. Wong^{d,p}, W. L. K. Wu^b, H. Yang^g, K. W. Yoon^{g,b}, E. Young^{g,b}, C. Yu^g, L. Zeng^d, C. Zhang^e, and S. Zhang^e

^aSchool of Physics and Astronomy, Cardiff University, Cardiff, CF24 3AA, United Kingdom

^bKavli Institute for Particle Astrophysics and Cosmology, SLAC National Accelerator Laboratory, 2575 Sand Hill Rd, Menlo Park, CA 94025, USA

^cDepartment of Physics and Astronomy, University of British Columbia, Vancouver, British Columbia V6T 1Z1, Canada

^dCenter for Astrophysics, Harvard & Smithsonian, Cambridge, MA 02138, USA

^eDepartment of Physics, California Institute of Technology, Pasadena, CA 91125, USA

^fDepartment of Physics, University of Cincinnati, Cincinnati, OH 45221, USA

^gDepartment of Physics, Stanford University, Stanford, CA 94305, USA

^hJet Propulsion Laboratory, Pasadena, CA 91109, USA

ⁱMinnesota Institute for Astrophysics, University of Minnesota, Minneapolis, MN 55455, USA

^jKavli Institute for Cosmological Physics, University of Chicago, Chicago, IL 60637, USA

^kSchool of Physics and Astronomy, University of Minnesota, Minneapolis, MN 55455, USA

^lNational Institute of Standards and Technology, Boulder, CO 80305, USA

^mService des Basses Températures, Commissariat à l'Energie Atomique, 38054 Grenoble, France

ⁿDepartment of Physics, University of Illinois at Urbana-Champaign, Urbana, IL 61801, USA

^oDepartment of Astronomy, University of Illinois at Urbana-Champaign, Urbana, IL 61801, USA

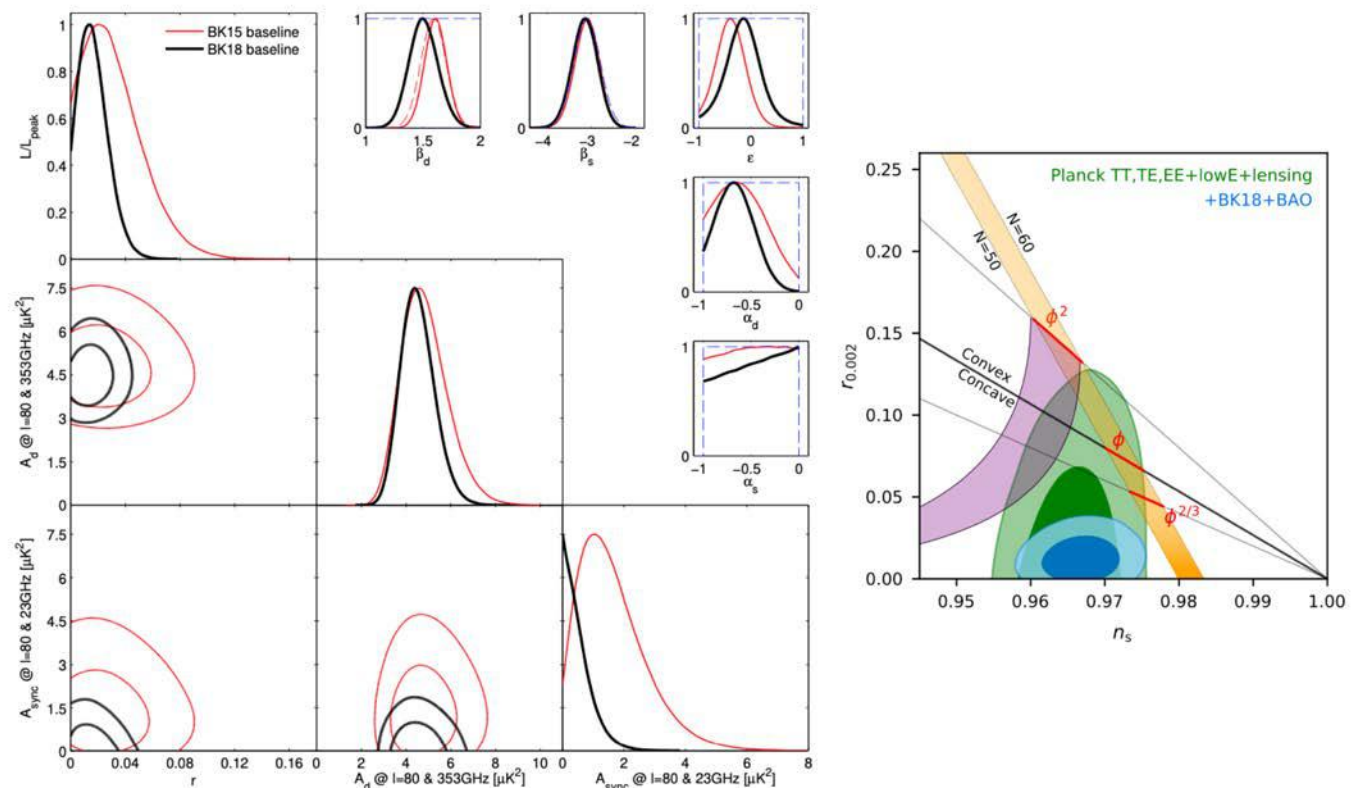
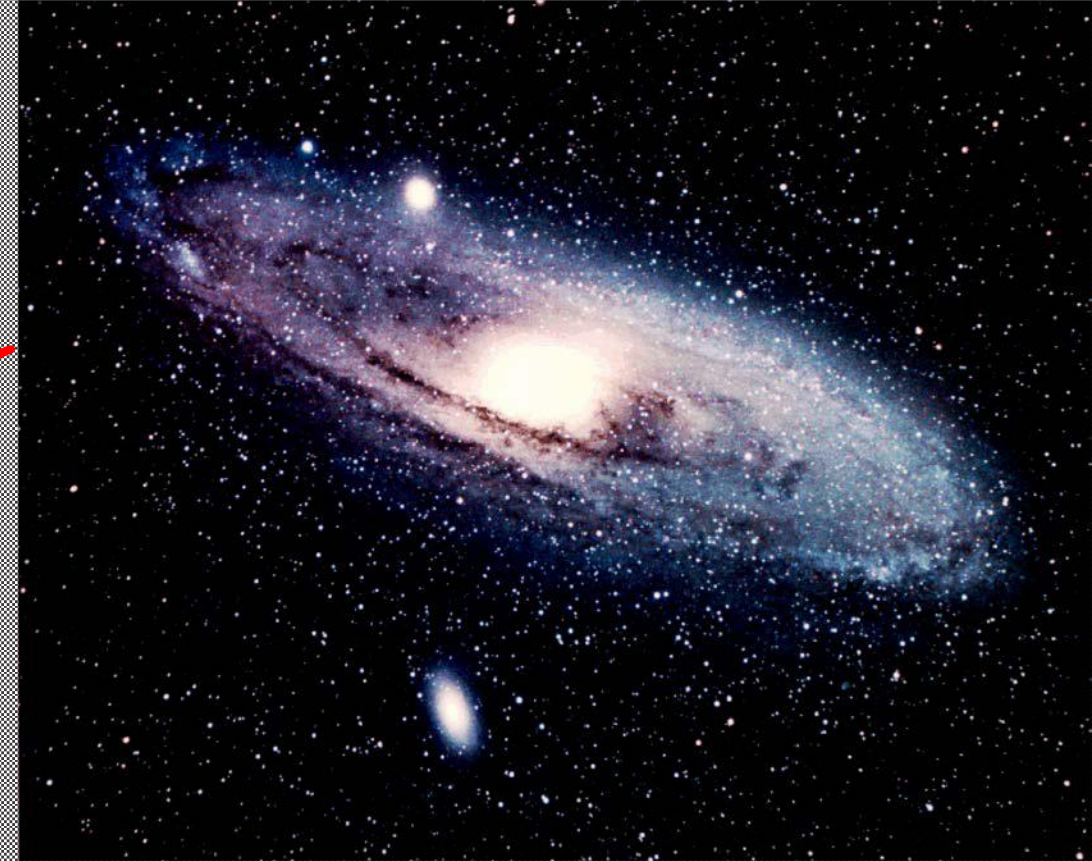
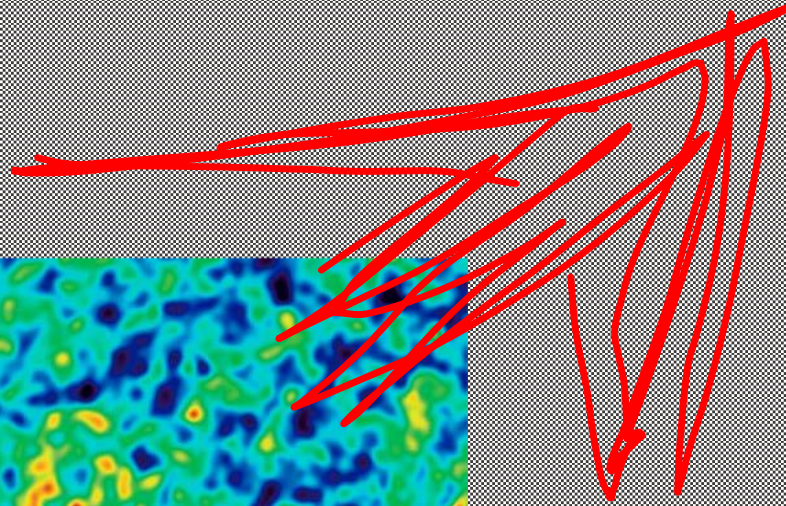
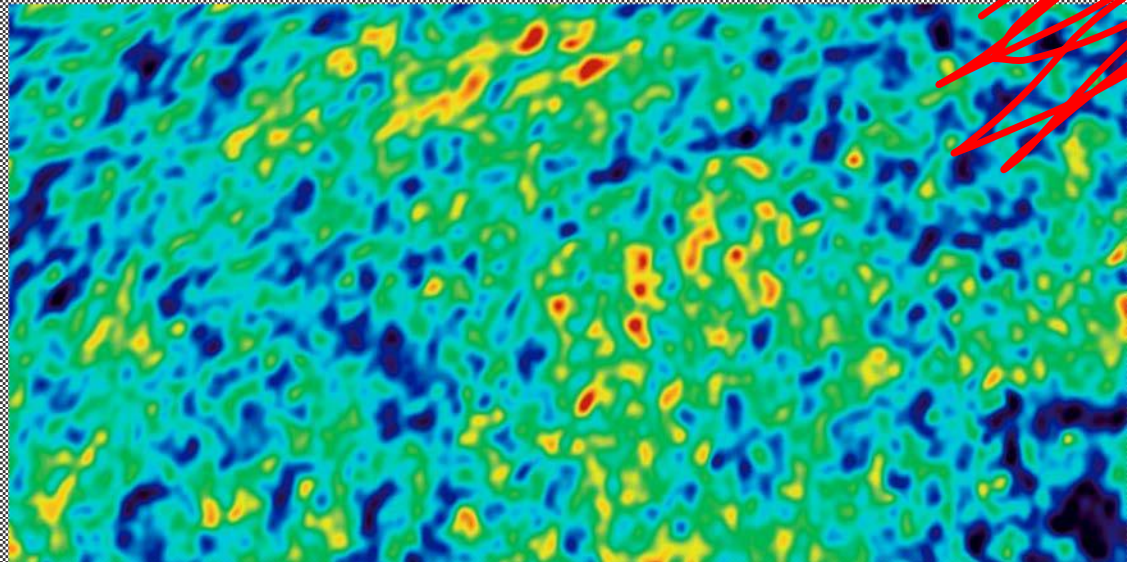


Figure 5. *Left:* CosmoMC likelihood results for the BICEP/Keck baseline model. Selected 1D and 2D marginalized posteriors are shown. The red faint curves are the results from BK15 while the black solid curves are the results of BK18. The dashed blue and red lines show priors on foreground parameters. The analysis method is the same as in BK15, except the β_d prior based on *Planck* data from other regions of the sky is removed this time due to the improved sensitivity of BK18. *Right:* Constraints in the r vs. n_s plane. The purple and orange bands are natural inflation and monomial inflation respectively. The blue contour shows the updated constraint after adding BK18 and BAO data to the *Planck* baseline analysis. The r posterior is tightened from $r_{0.05} < 0.11$ to $r_{0.05} < 0.035$ at 95% confidence.

Triumphs and tests

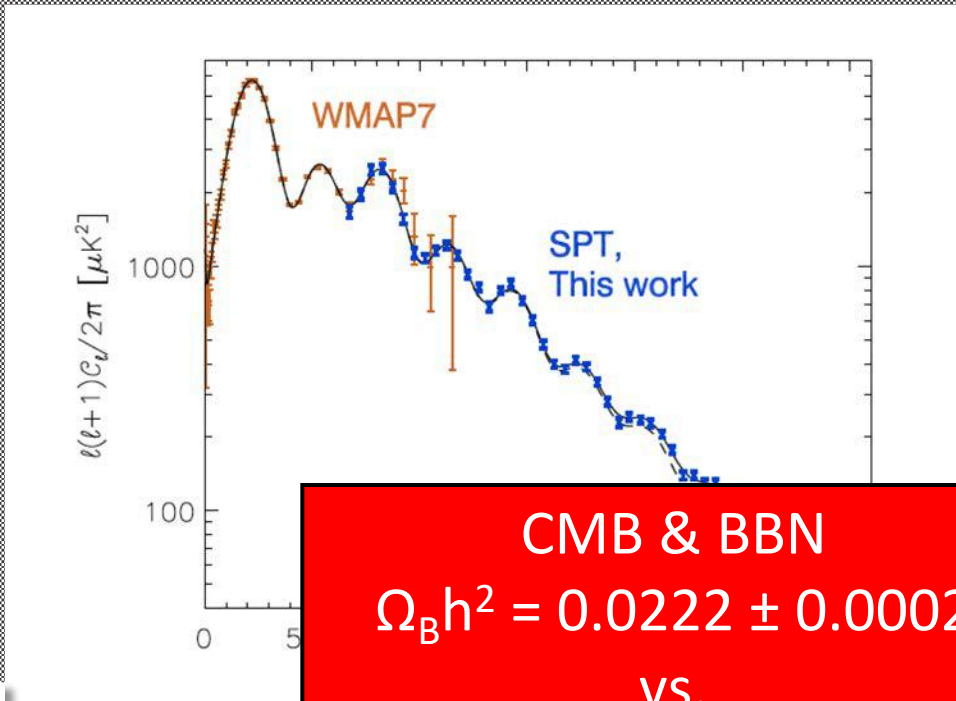
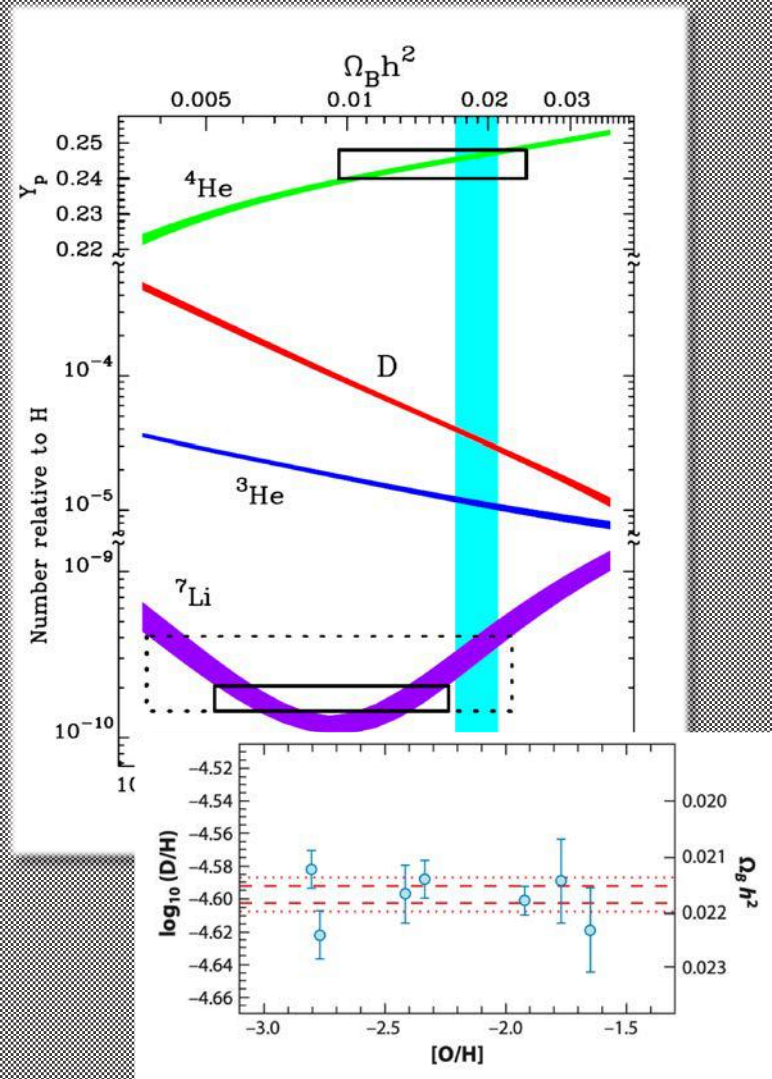
- Big bang nucleosynthesis (no need to posit large primordial ^4He abundance) and predicted ^4He ($Y_p = 0.2469 \pm 0.0002$) vs observed ($Y_p = 0.245 \pm 0.00034$) vs CMB inferred ($Y_p = 0.242 \pm 0.024$)
- BBN baryon density ($\Omega_B h^2 = 0.02166 \pm 0.00015$) and CMB baryon density ($\Omega_B h^2 = 0.02237 \pm 0.00015$)
- Structure formation: perturbations measured in CMB + gravity (numerical simulation) = the Universe we see today
- Crosschecks: H_0 , σ_8 and others (tensions and opportunities??)
- Basic predictions of inflation verified: almost scale-invariant, Gaussian curvature fluctuations, flat Universe, and coming soon GWs
- Precision set of cosmological parameters

The grand connection between big and small



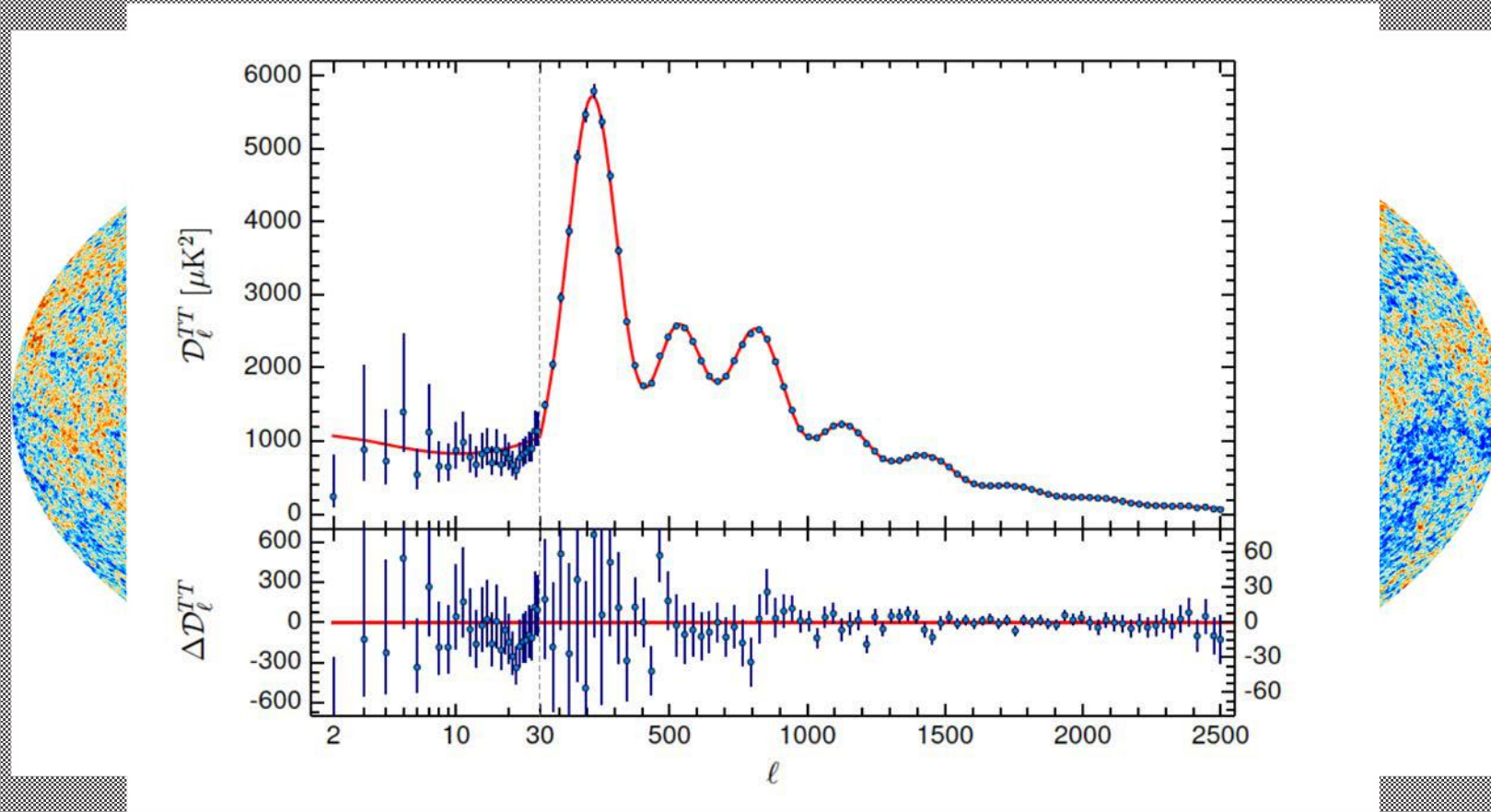
Quantum fluctuations on unimaginably small scales lead to structure on cosmic scales

The airtight evidence for nonbaryonic DM



CMB & BBN
 $\Omega_B h^2 = 0.0222 \pm 0.0002$
vs.
CMB/SDSS
 $\Omega_M h^2 = 0.143 \pm 0.001$
 $> 50\sigma$ discrepancy

Best evidence for Λ CDM



6 numbers describe the Universe from
the big bang until today

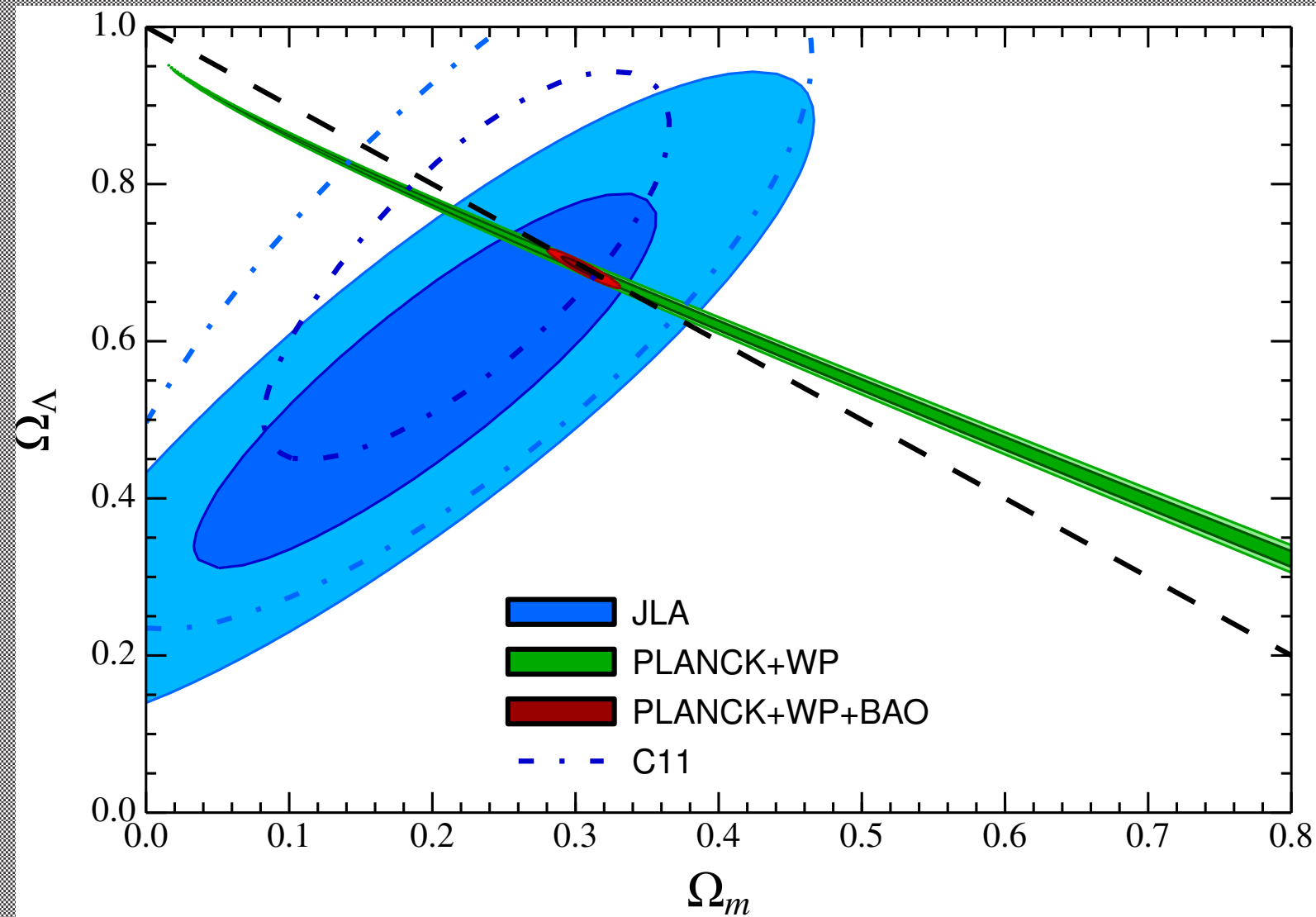
- | | |
|-----------------------------------|------------------|
| 1. Baryon mass density | 4. Tilt |
| 2. CDM mass density | 5. Sound horizon |
| 3. Density perturbation amplitude | 6. Optical depth |

Era of Precision Cosmology (plenty of well measured numbers)

$$\begin{aligned}T_0 &= 2.7255 \pm 0.00057 \text{ K} \\t_0 &= 13.8 \pm 0.02 \text{ Gyr} \\\Omega_0 &= 1.00 \pm 0.002 \\H_0 &= 67.4 \pm 0.5 \text{ km/s/Mpc} \\H_0 &= 73.5 \pm 2 \text{ km/s/Mpc} \\N_\nu &= 2.99 \pm 0.17 \\n_s &= 0.965 \pm 0.004 \\r &< 0.07 \pm 0.03 \\w &= -1.03 \pm 0.04 \\w_a &= -0.22 \pm 0.41 \\\Omega_B h^2 &= 0.0222 \pm 0.0002 \\\Omega_M h^2 &= 0.142 \pm 0.0013 \\\sigma_8 &= 0.811 \pm 0.006 \\\theta_{MC} &= 1.04092 \pm 0.0003 \times 10^{-2} \\\tau &= 0.0544 \pm 0.0073 \\A_S &= 2.10 \pm 0.03 \times 10^{-9} \\z_{rec} &= 1090 \pm 0.2 \\z_{eq} &= 3387 \pm 27 \\\dots &= \dots\end{aligned}$$

and more to come – more tests
and hopefully some surprises

Λ fits perfectly!



We have come a long way since 1970, when Allan Sandage said:



COSMOLOGY: A SEARCH FOR TWO NUMBERS

Precision measurements of the rate of expansion and the deceleration of the universe may soon provide a major test of cosmological models

ALLAN R. SANDAGE



Allan Sandage has been a staff member at the Mount Wilson and Palomar Observatories since he received his PhD from Cal Tech in 1953. His main interests are stellar evolution, observational cosmology, form of the redshift laws, quasars and distance scales. In 1960 Sandage and Thomas Matthews were the first to isolate the quasars.

Hubble troubles, part 2



Direct: 73 ± 1.1 vs. Indirect: 67.5 ± 0.4 km/s/Mpc

- Direct measurement:

$$H_0 = \frac{\dot{R}}{R} = \frac{\text{galaxy velocity}}{\text{galaxy distance}} \quad R = \text{size of Universe}$$

- NB: “v easy, d hard”
- Distance ladder: standard candles – Cepheids, TRB, SNe1a
- Time delay (jump ladder)
- Both agree

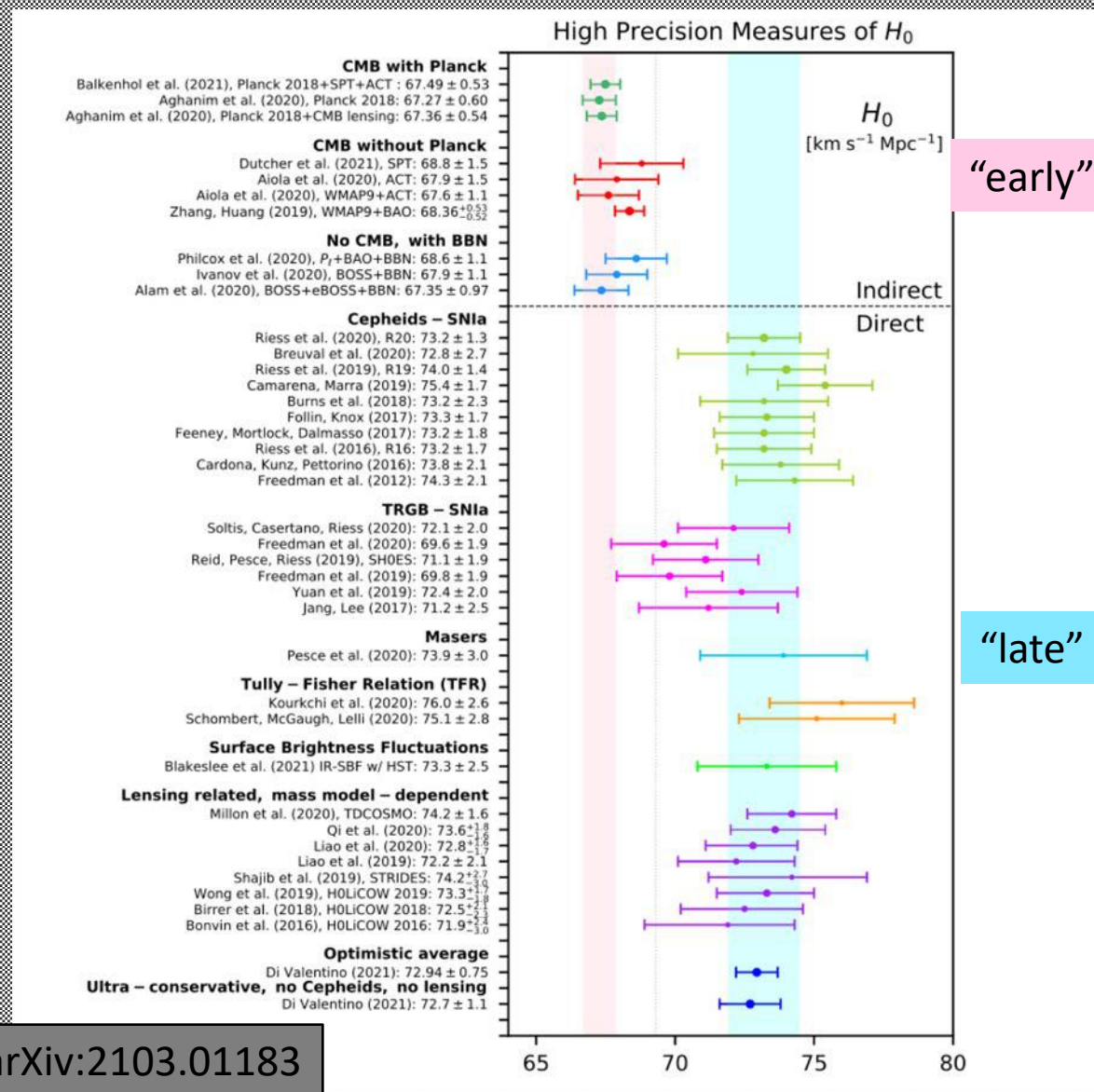
$$\text{distance} \sim \frac{\text{time delay}}{\theta^2}$$

- Indirect (CMB)

$$d_{\text{CMB}} = \frac{v_{\text{sound}} t_{\text{CMB}}}{\theta_{\text{sound}}} = H_0^{-1} \int \frac{dz}{[\Omega_M(1+z)^3 + \Omega_\Lambda]^{1/2}}$$

- Direct and indirect could both be correct and paradigm wrong! Or, one or both measurements could be wrong and Λ CDM correct

Hubble troubles or opportunities!



- Indirect (pink): 67.5 ± 0.5 km/s/Mpc



George Efstathiou (Planck)

- Direct (cyan): 73.2 ± 1.3 km/s/Mpc



Adam Riess (SH0ES)

- 5-sigma difference!

A Comprehensive Measurement of the Local Value of the Hubble Constant with $1 \text{ km s}^{-1} \text{ Mpc}^{-1}$ Uncertainty from the *Hubble Space Telescope* and the SHOES Team

ADAM G. RIESS,^{1,2} WENLONG YUAN,² LUCAS M. MACRI,³ DAN SCOLNIC,⁴ DILLON BROUT,⁵ STEFANO CASERTANO,¹ DAVID O. JONES,⁶ YUKEI MURAKAMI,² GAGANDEEP S. ANAND,¹ LOUISE BREUVAL,^{2,7} THOMAS G. BRINK,⁸ ALEXEI V. FILIPPENKO,^{8,9} SAMANTHA HOFFMANN,¹ SAURABH W. JHA,¹⁰ W. D'ARCY KENWORTHY,² JOHN MACKENTY,¹ BENJAMIN E. STAHL,⁸ AND WEIKANG ZHENG⁸

¹Space Telescope Science Institute, 3700 San Martin Drive, Baltimore, MD 21218, USA

²Department of Physics and Astronomy, Johns Hopkins University, Baltimore, MD 21218, USA

³George P. and Cynthia W. Mitchell Institute for Fundamental Physics and Astronomy,

Department of Physics & Astronomy, Texas A&M University, College Station, TX 77843, USA

⁴Department of Physics, Duke University, Durham, NC 27708, USA

⁵Center for Astrophysics, Harvard & Smithsonian, 60 Garden St., Cambridge, MA 02138, USA

⁶Einstein Fellow, Department of Astronomy & Astrophysics, University of California, Santa Cruz, CA 95064, USA

⁷LESIA, Observatoire de Paris, Université PSL, CNRS, Sorbonne Université, Université de Paris, 5 place Jules Janssen, F-92195 Meudon France

⁸Department of Astronomy, University of California, Berkeley, CA 94720-3411, USA

⁹Miller Institute for Basic Research in Science, University of California, Berkeley, CA 94720, USA

¹⁰Department of Physics and Astronomy, Rutgers, the State University of New Jersey, Piscataway, NJ 08854, USA

ABSTRACT

We report observations from the *Hubble Space Telescope* (*HST*) of Cepheid variables in the host galaxies of 42 Type Ia supernovae (SNe Ia) used to calibrate the Hubble constant (H_0). These include the complete sample of all suitable SNe Ia discovered in the last four decades at redshift $z \leq 0.01$, collected and calibrated from ≥ 1000 *HST* orbits, more than doubling the sample whose size limits the precision of the direct determination of H_0 . The Cepheids are calibrated geometrically from *Gaia* EDR3 parallaxes, masers in NGC 4258 (here tripling that sample of Cepheids), and detached eclipsing binaries in the Large Magellanic Cloud. All Cepheids in these anchors and SN Ia hosts were measured with the same instrument (WFC3) and filters (*F555W*, *F814W*, *F160W*) to negate zeropoint errors.

We present multiple verifications of Cepheid photometry and six tests of background determinations that show Cepheid measurements are accurate in the presence of crowded backgrounds. The SNe Ia in these hosts calibrate the magnitude–redshift relation from the revised Pantheon+ compilation, accounting here for covariance between all SN data and with host properties and SN surveys matched throughout to negate systematics. We decrease the uncertainty in the local determination of H_0 to $1 \text{ km s}^{-1} \text{ Mpc}^{-1}$ including systematics. We present results for a comprehensive set of nearly 70 analysis variants to explore the sensitivity of H_0 to selections of anchors, SN surveys, redshift ranges, the treatment of Cepheid dust, metallicity, form of the period–luminosity relation, SN color, peculiar-velocity corrections, sample bifurcations, and simultaneous measurement of the expansion history.

Our baseline result from the Cepheid–SN Ia sample is $H_0 = 73.04 \pm 1.04 \text{ km s}^{-1} \text{ Mpc}^{-1}$, which includes systematic uncertainties and lies near the median of all analysis variants. We demonstrate consistency with measures from *HST* of the TRGB between SN Ia hosts and NGC 4258, and include them *simultaneously* to yield $72.53 \pm 0.99 \text{ km s}^{-1} \text{ Mpc}^{-1}$. The inclusion of high-redshift SNe Ia yields $H_0 = 73.30 \pm 1.04 \text{ km s}^{-1} \text{ Mpc}^{-1}$ and $q_0 = -0.51 \pm 0.024$. We find a 5σ difference with the prediction of H_0 from *Planck* CMB observations under Λ CDM, with no indication that the discrepancy arises from measurement uncertainties or analysis variations considered to date. The source of this now long-standing discrepancy between direct and cosmological routes to determining H_0 remains unknown.

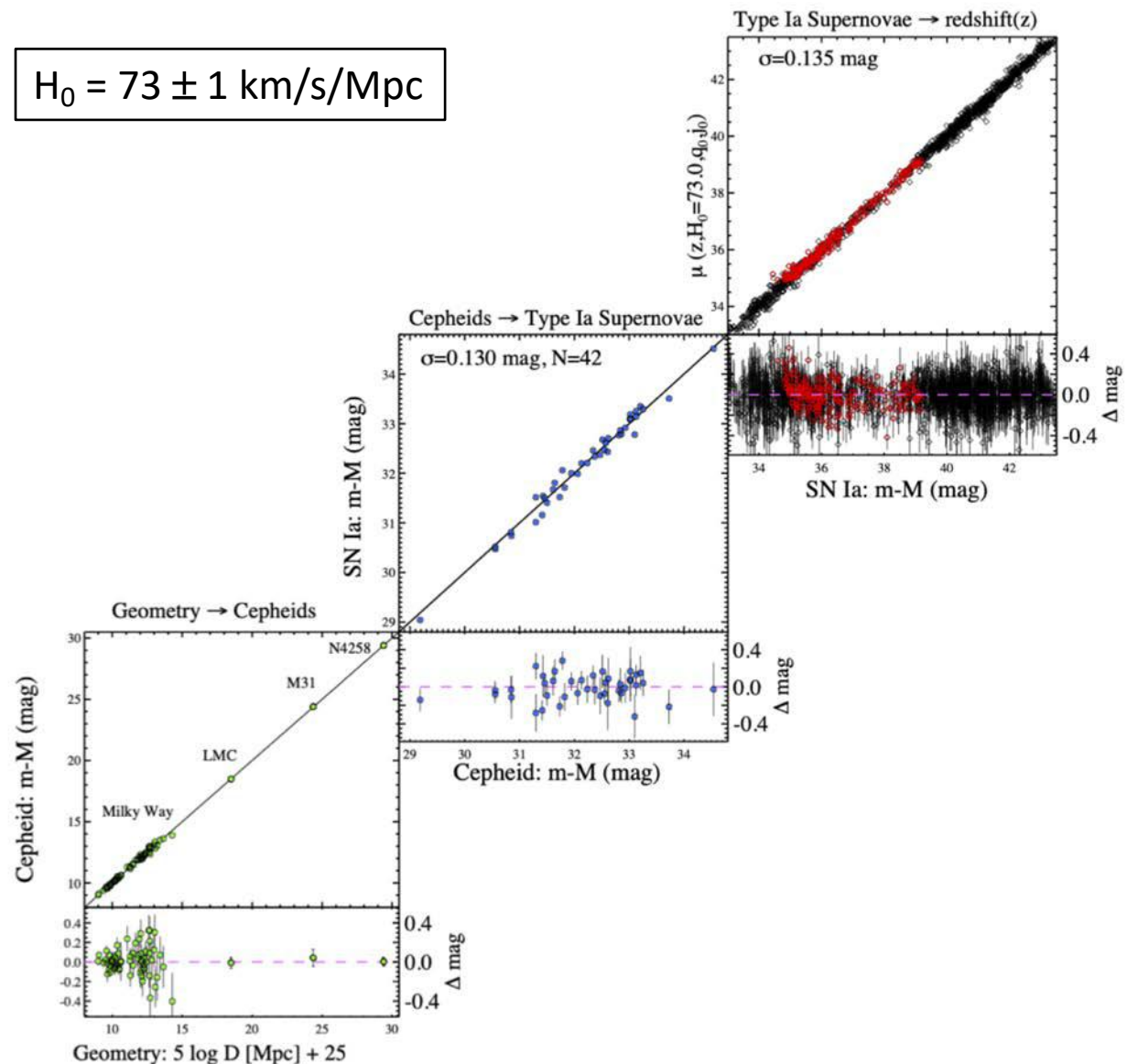
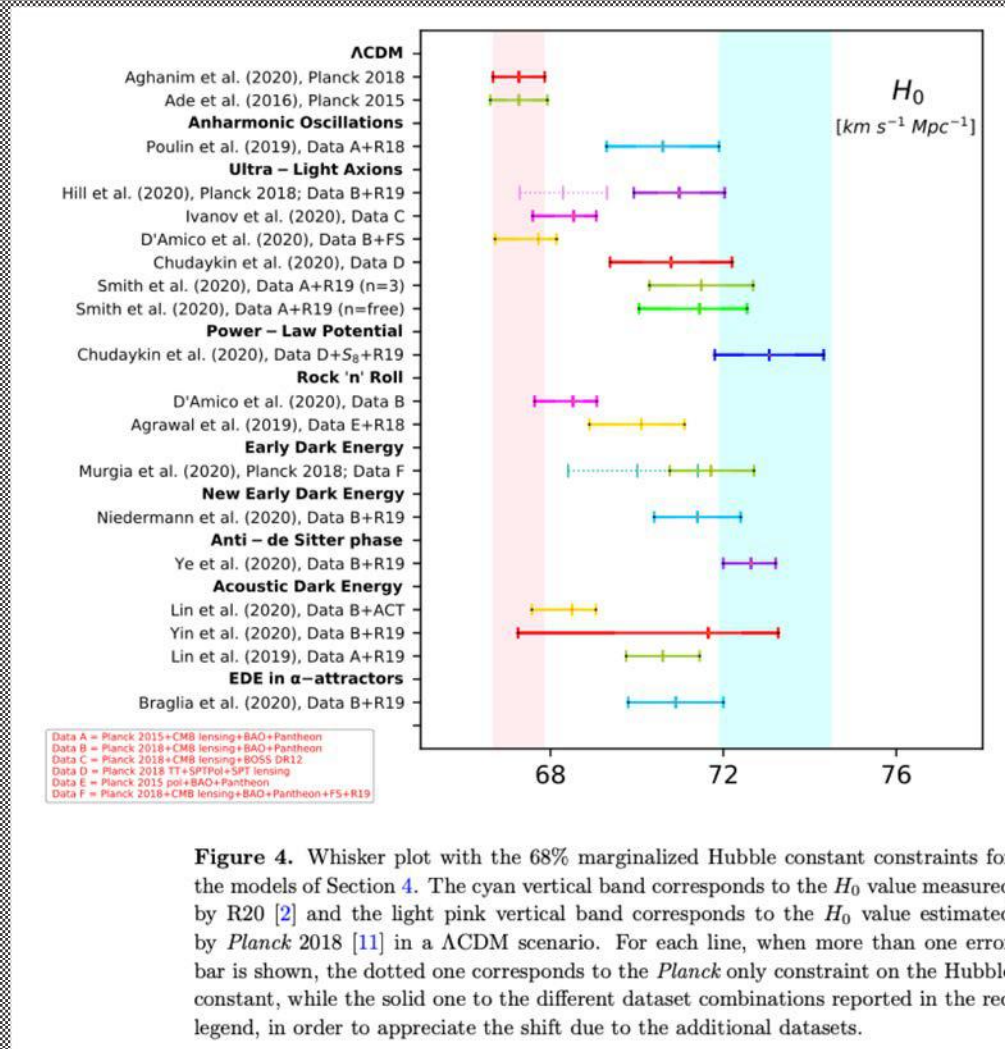


Figure 12. Complete distance ladder. The simultaneous agreement of distance pairs: geometric and Cepheid-based (lower left), Cepheid- and SN-based (middle), and SN- and redshift-based (top right) provides the measurement of H_0 . For each step, geometric or calibrated distances on the abscissa serve to calibrate a relative distance indicator on the ordinate through the determination of M_B or H_0 . Results shown are an approximation to the global fit as discussed in the text. Red SN points are at $0.0233 < z < 0.15$, with the lower-redshift bound producing the *appearance* of asymmetric residuals when plotted against distance.

“New physics”

- The two discrepant measurements could both could be right if Λ CDM is wrong!
- New ingredient(s) to Λ CDM
 - Early dark energy
 - Extra radiation
 - None compelling yet



Or one or both measurements could be wrong or NEW PHYSICS! Big mystery; stay tuned!

Λ CDM paradigm shift: adding ONE (odious) thing, solved FIVE problems with Inflation + CDM. H_0 fixes not as compelling – yet!

General Relativity and Gravitation, Vol. 27, No. 11, 1995

The Cosmological Constant Is Back[†]

Lawrence M. Krauss¹ and Michael S. Turner^{2,3}

A diverse set of observations now compellingly suggest that the universe possesses a nonzero cosmological constant. In the context of quantum-field theory a cosmological constant corresponds to the energy density of the vacuum, and the favored value for the cosmological constant corresponds to a very tiny vacuum energy density. We discuss future observational tests for a cosmological constant as well as the fundamental theoretical challenges — and opportunities — that this poses for particle physics and for extending our understanding of the evolution of the universe back to the earliest moments.

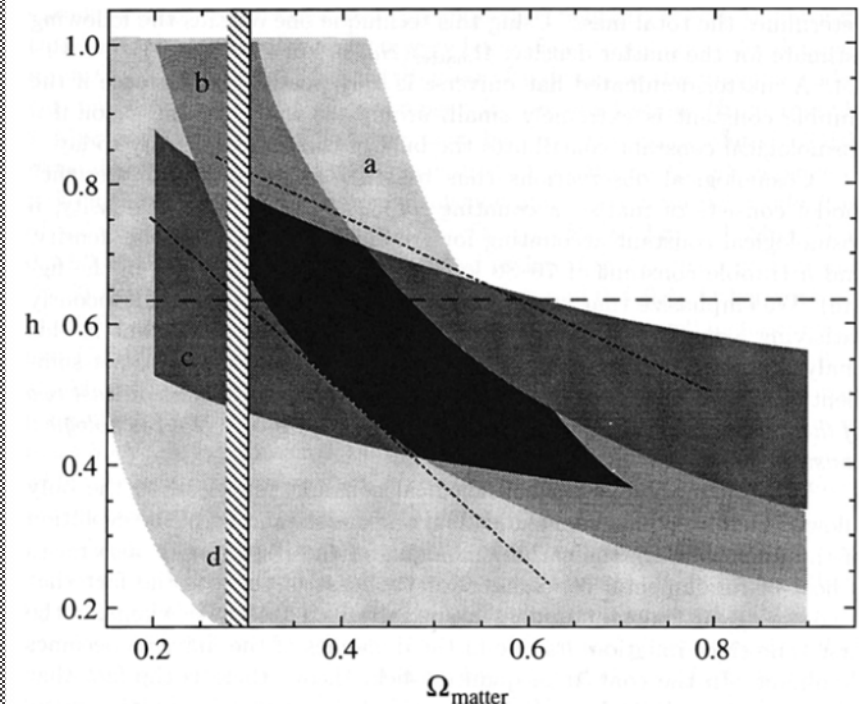


Figure 1. Constraints on the matter density in a flat universe as a function of the Hubble constant $H_0 = 100 h \text{ km s}^{-1} \text{ Mpc}^{-1}$. Shaded regions indicate allowed regions of parameter space. Region (a) comes from combining big-bang nucleosynthesis limits with x-ray observations of clusters, (b) arises from considerations of clustering on large scales, (c) is based on age determinations of globular clusters, (d) is a lower limit based on virial estimates of the density of clustered matter on large scales. The horizontal dashed line is a one sigma lower limit on the Hubble constant from recent Hubble Space Telescope measurements. The diagonal dashed lines represent the allowed limits of phase space based on combining COBE normalization of cold dark matter models with estimates of matter density fluctuations on galactic and cluster scales. The dark shaded region indicates the region allowed by all constraints.

Big questions and grand
aspirations

... the pillars of the Λ CDM paradigm!

- Don't understand Dark Energy ($69 \pm 0.6\%$) and why Λ (quantum vacuum energy) is so small
- The physics of Inflation or when it took place
- What Dark Matter ($31\% \pm 0.6\%$ less baryons) is comprised of
- How Baryons ($4.8 \pm 0.06\%$) survived annihilation (baryogenesis)



?!

Both!

Cosmic acceleration

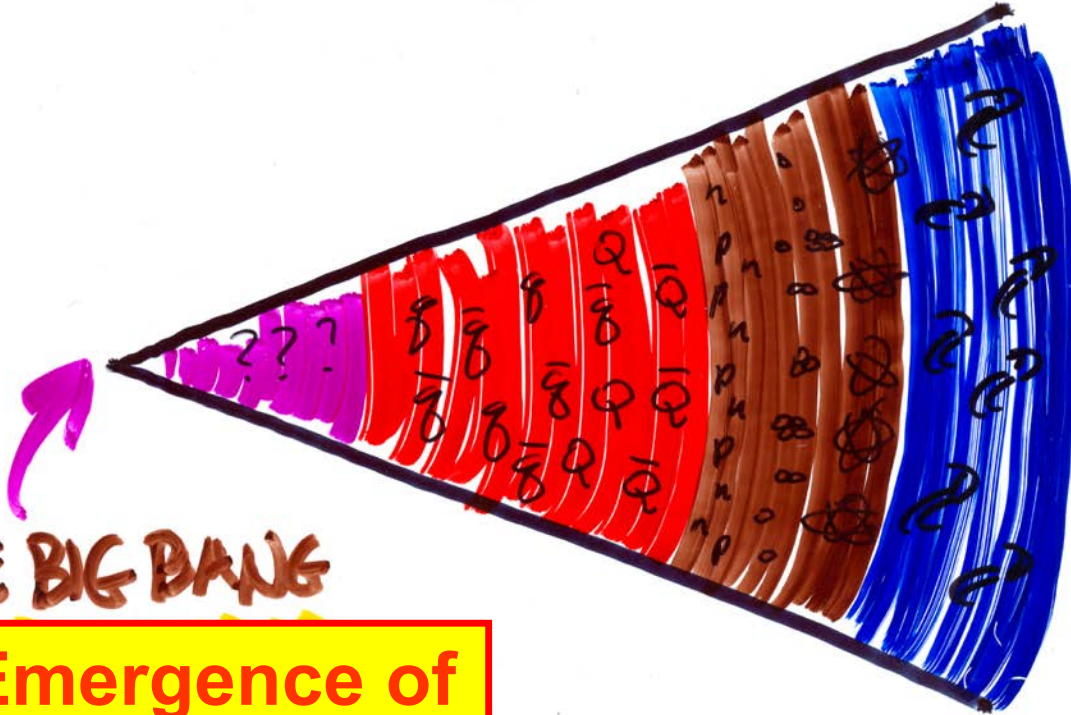
- How often?
- When?
- Why?
- Something beyond GR needed?
- Opportunity to unify inflation and dark energy?

Grand aspirations

1. Origin of the space, time and the Universe
2. Before the big bang (related to #1?)
3. Destiny of the Universe
4. Self-booting Universe (given the TOE, everything else follows automatically)
5. Making sense of the multiverse or getting rid of it
6. Why something rather nothing and, where did the laws of physics come from?

Einstein got the
right answer
for the wrong
reason?

EINSTEIN'S BIG BANG



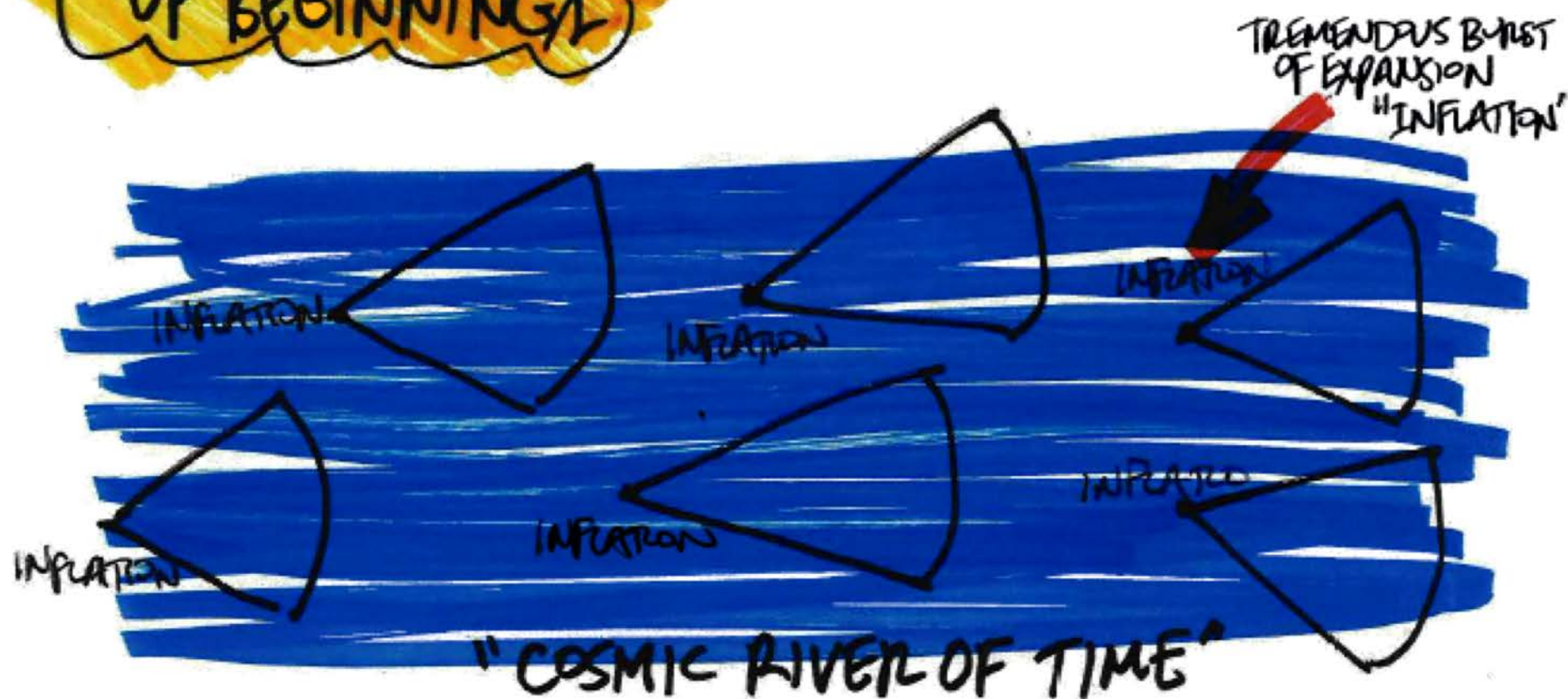
THE BIG BANG

= Emergence of
space and time

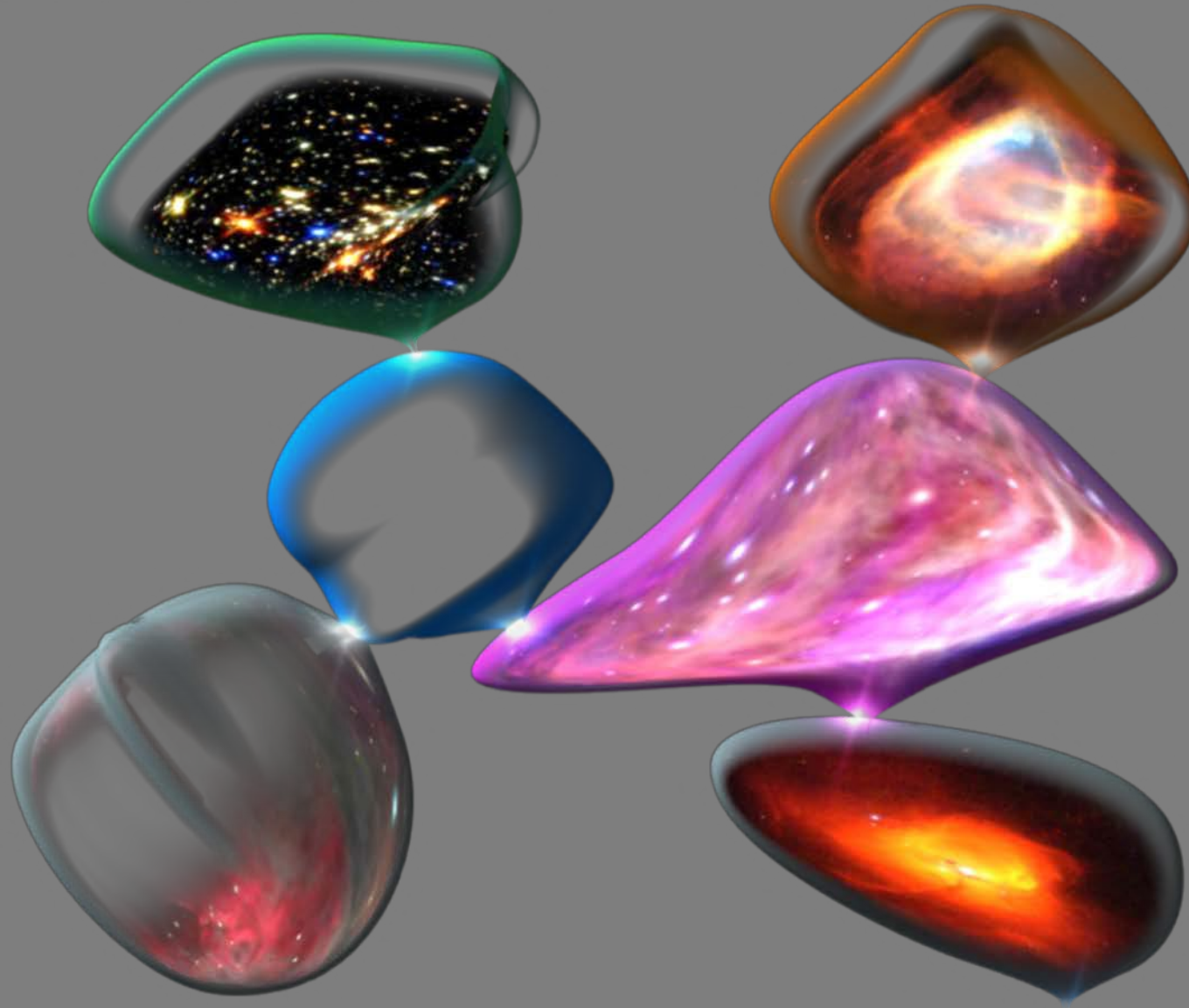
• NO BEFORE THE BIG BANG

INFLATIONARY MULTIVERSE

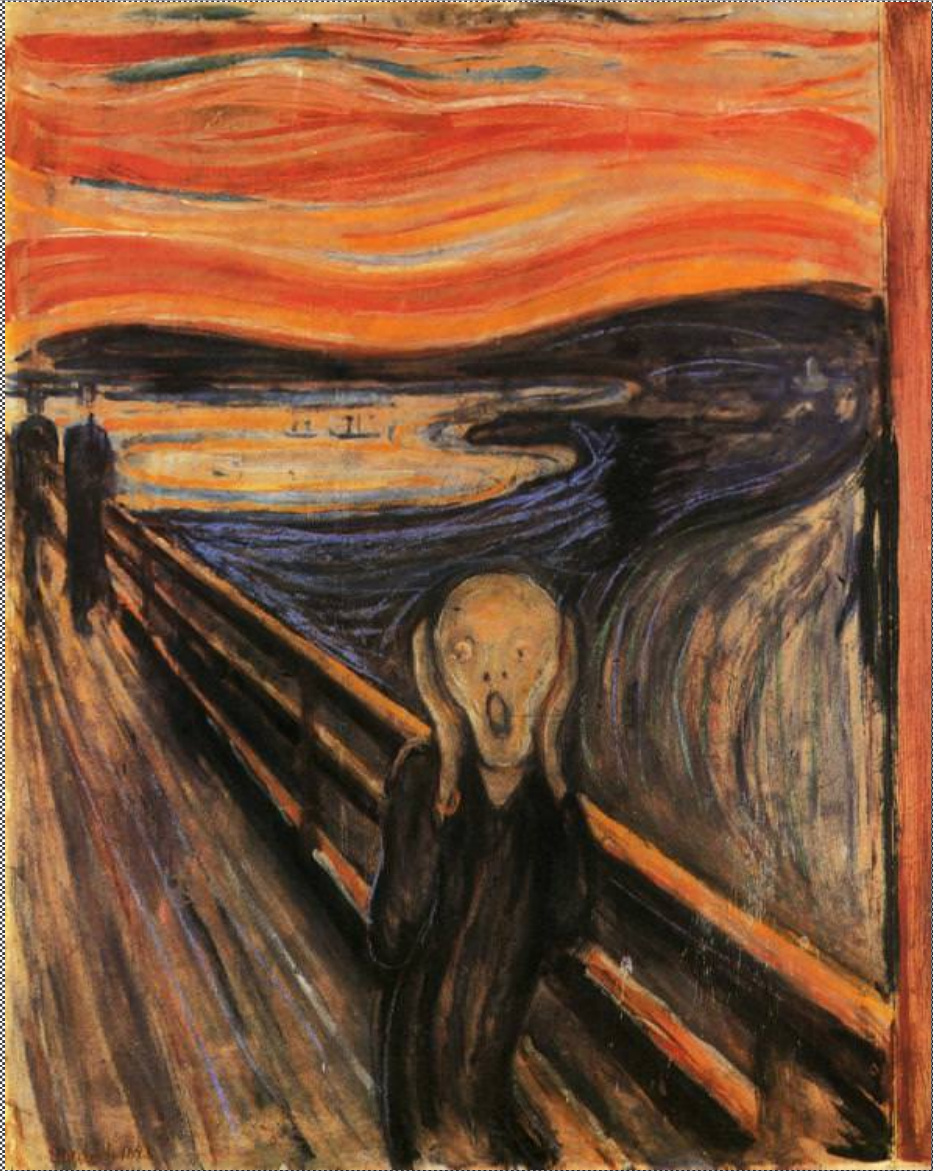
INFINITE NUMBER
OF BEGINNINGS



The multiverse



What to do about the multiverse



- Most important “discovery” since Copernicus?
- But is it science? (not testable – yet)