

Dark matter in compact stars and galactic structure

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N3AS Seminar June 27



Dark matter in compact stars and galactic structure

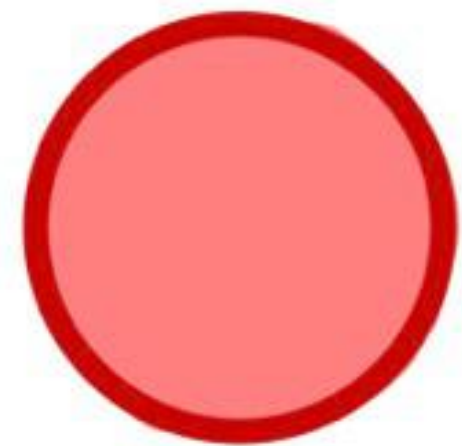
- *Dark matter heating neutron stars through either infall or annihilation*
- Asymmetric dark matter causing neutron stars to implode
 - Connections to missing pulsars (FRBs)
- White dwarf explosions from asymmetric dark matter / composites
 - Connections to type Ia supernovae, galactic structure

Neutron stars: nature's dark matter accelerators

- Neutron stars accelerate dark matter to beyond freezeout speeds

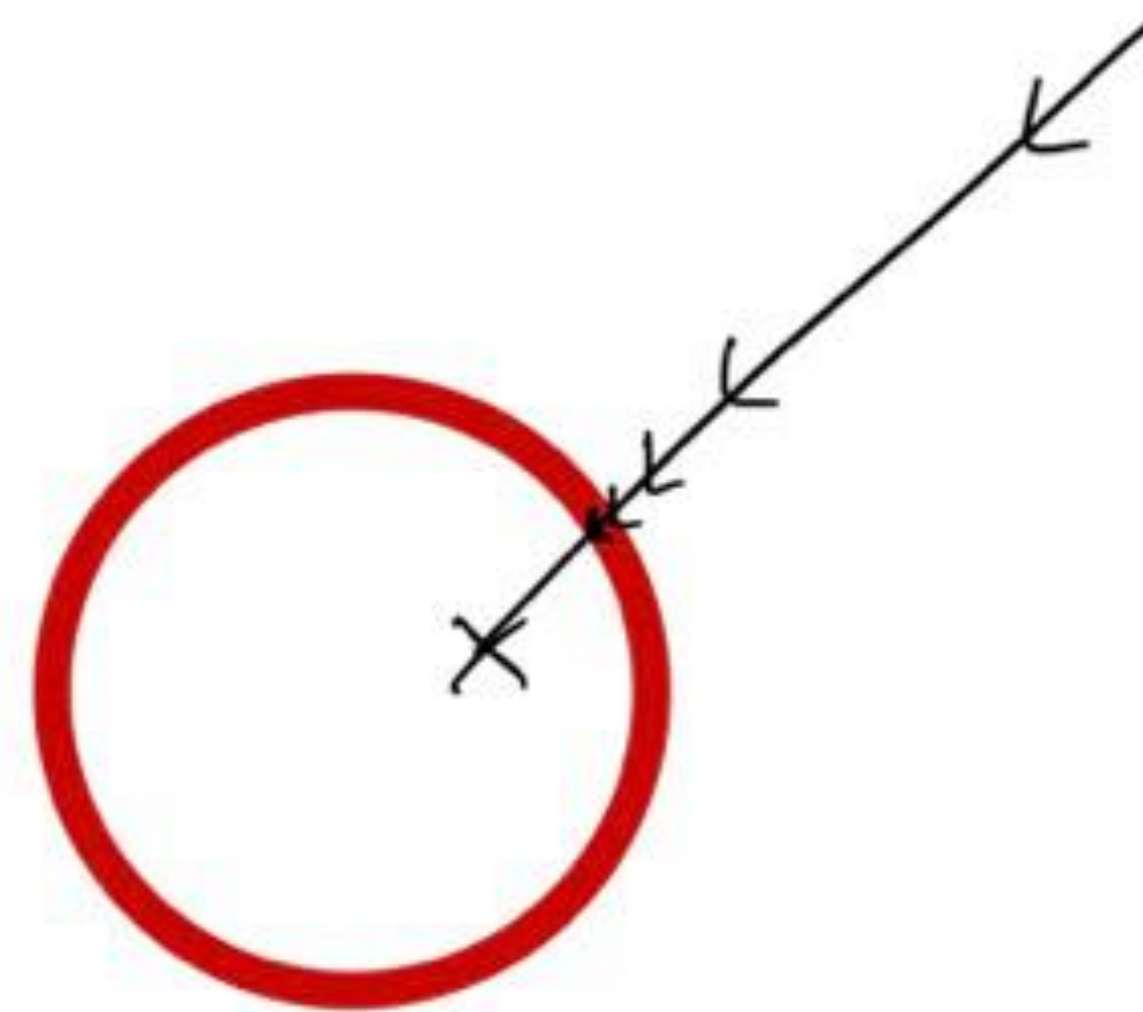
$$v_{esc} = \sqrt{\frac{2GM}{R}} \sim 0.7c$$

- Dense, accept a large DM flux



- fiducial mass of $\sim 10^{57}$ GeV
- neutrons:protons:electrons $\sim 10:1:1$
- flux of ~ 100 grams of DM/second

10 km

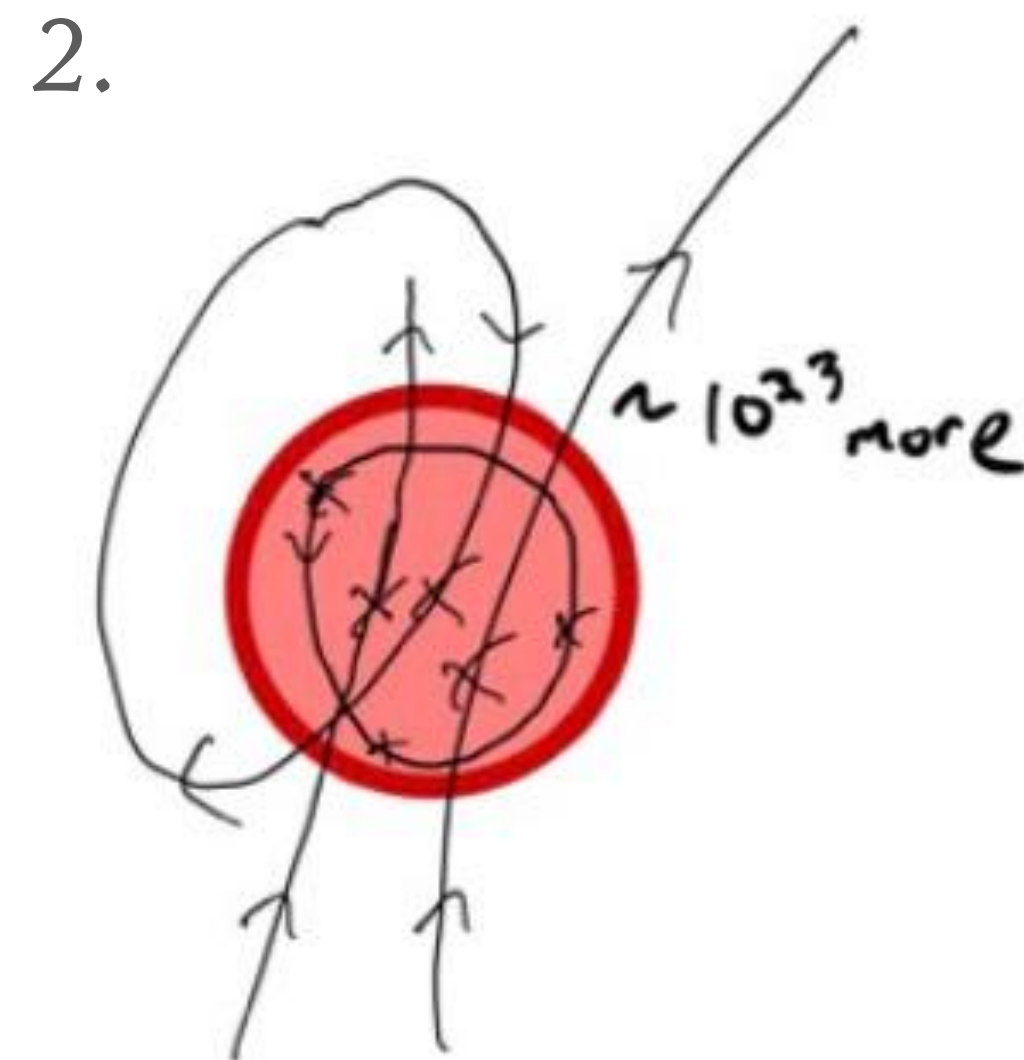
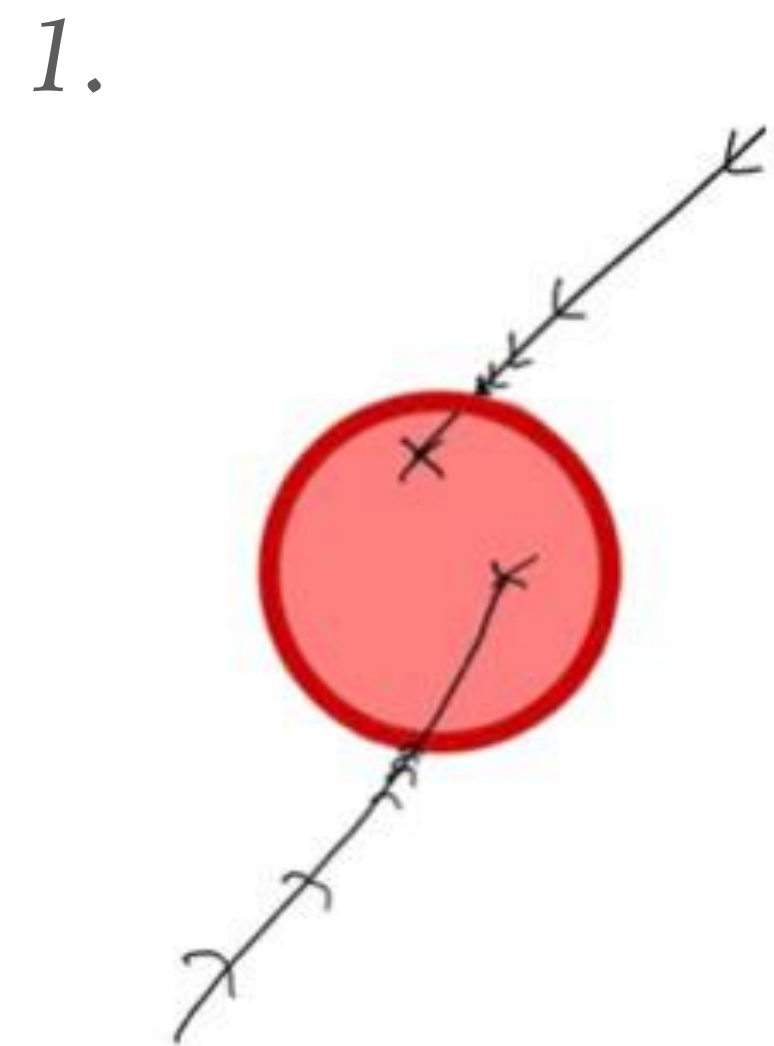


to scatter once with neutron per crossing

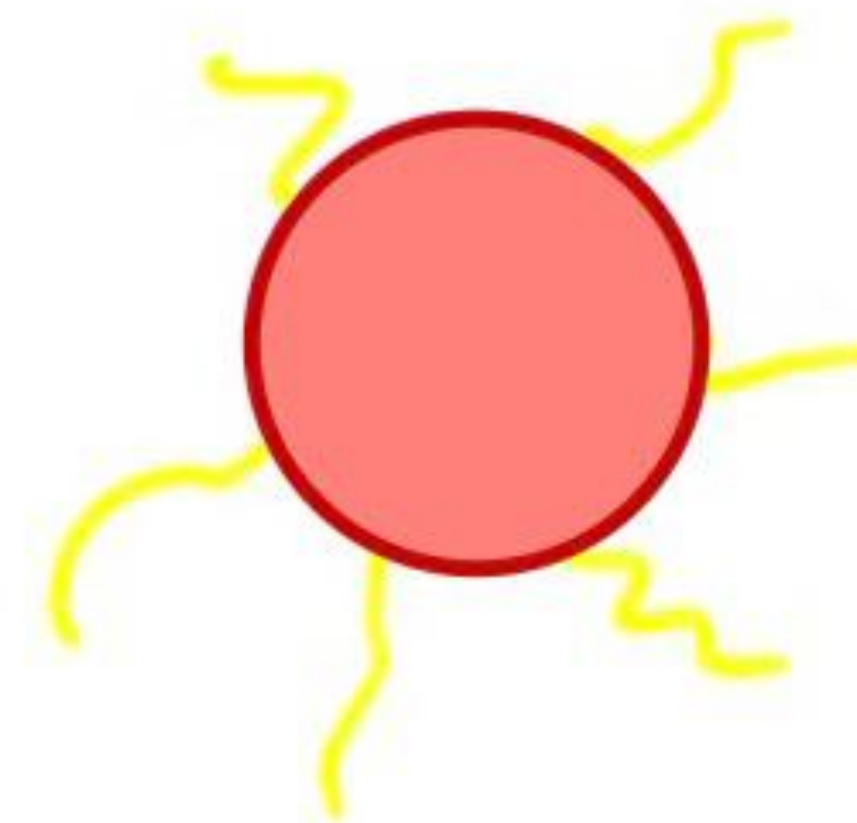
$$\sigma_n = \frac{R_{NS}^2}{N_n} \sim 10^{-45} \text{ cm}^2$$

Dark matter kinetic and annihilation heating of neutron stars

1. Dark matter accelerated to $\sim 0.7c$ by neutron star
2. DM deposits kinetic energy by scattering and re-scattering in the neutron star (may also annihilate in the NS)
3. Heats NS to 1750 K if all DM captured, 2500 K with annihilation ($0.4 \text{ GeV/cm}^3 \text{ DM}$)



3. $T \sim 1750 / 2500 \text{ K}$, for NS near Earth

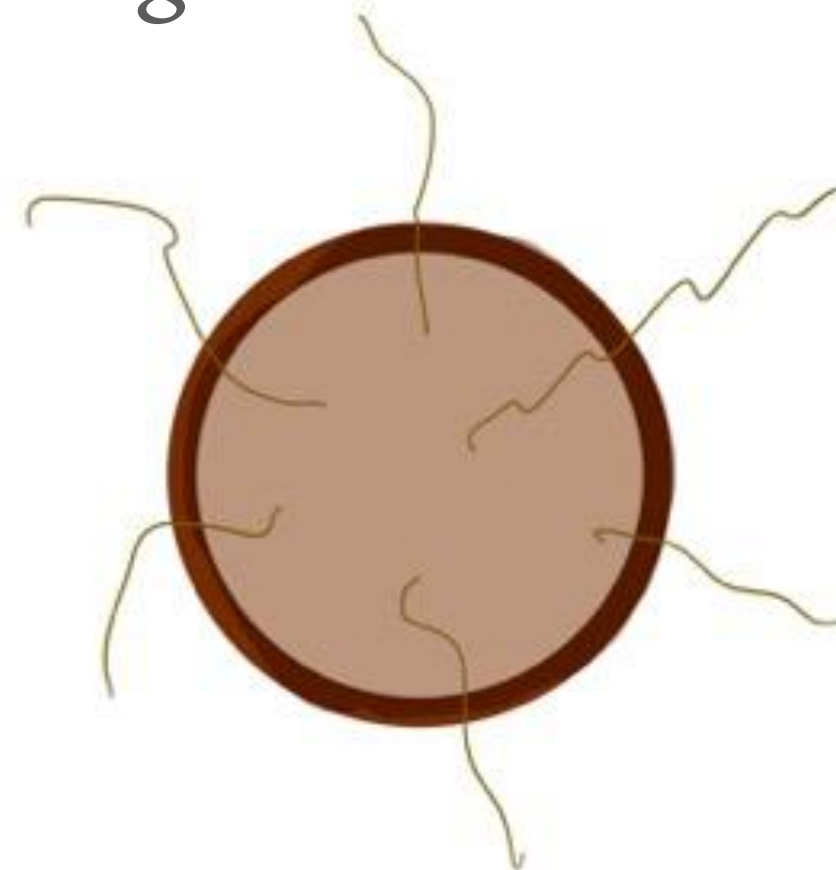


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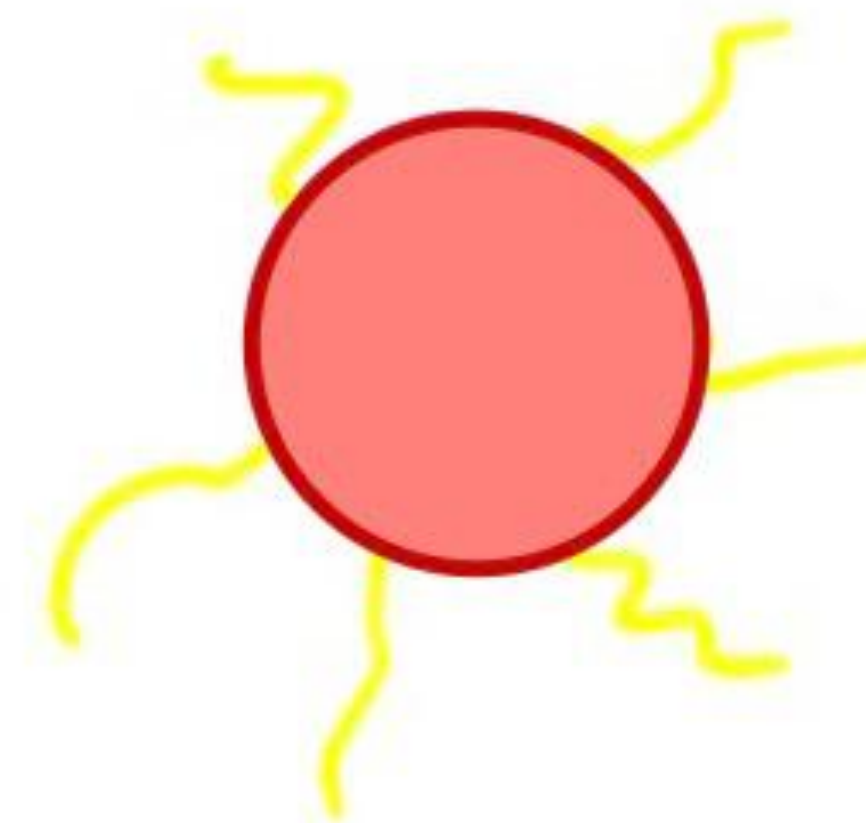
0. Compare to NS without DM heating

$$T_{eff}^{\infty} \sim 100 \text{ K} \left(\frac{\text{Gyr}}{t} \right)^{1/2}$$



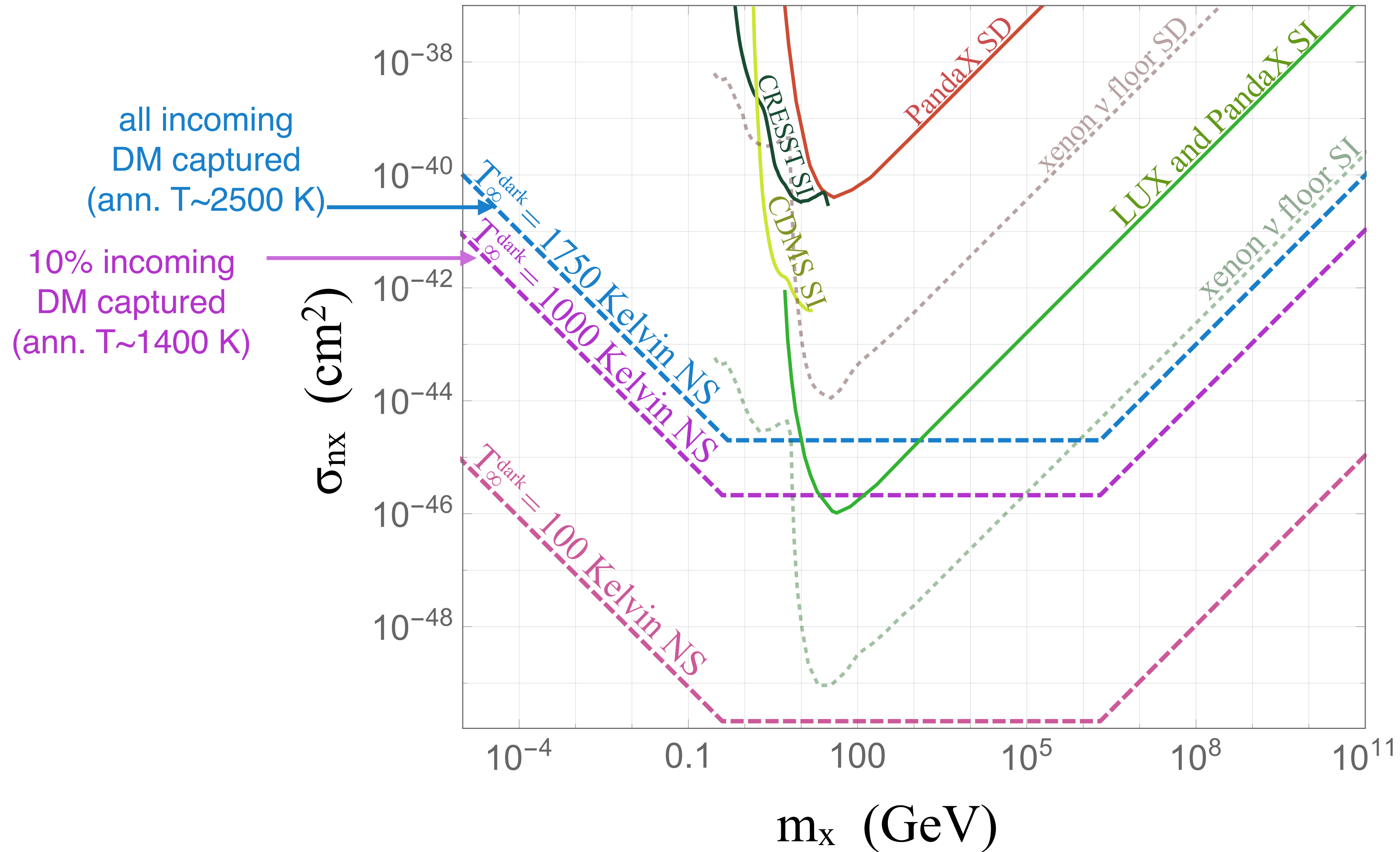
e.g. Yakovlev Pethick astro-ph/0402143
Page Lattimer et al. astro-ph/0403657

3. $T \sim 1750 / 2500 \text{ K}$, for NS near Earth



Neutron Star Dark Matter Heating Sensitivity

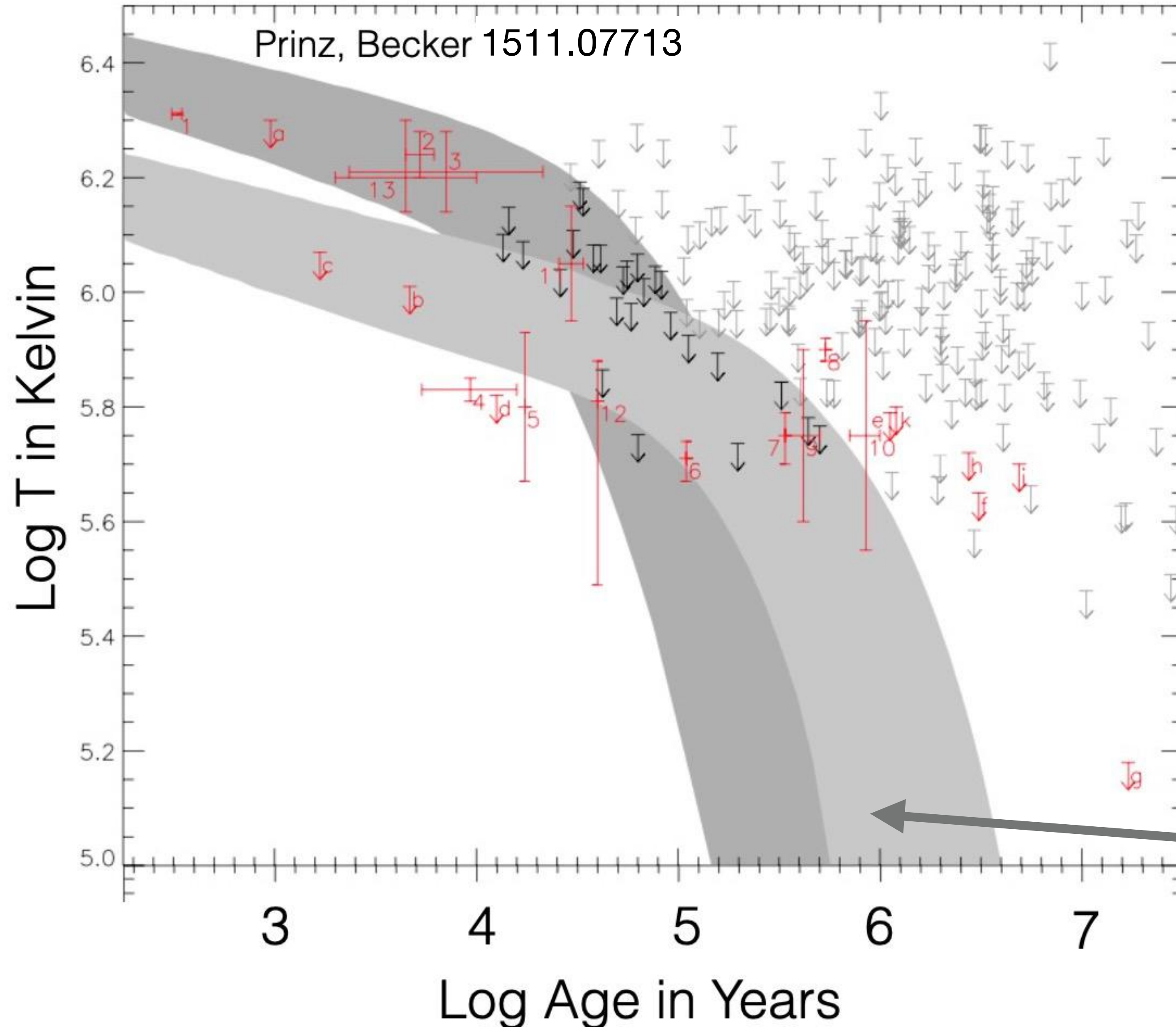
Baryakhtar, JB, Li, Linden, Raj 2017



Neutron stars, broad reach for particle dark matter

1. EFT, Spin-Dependent, Spin-Independent, Strongly Interacting, Electroweakino, Inelastic
2. Leptophilic dark matter
3. Self-interacting dark matter
4. Heavy DM, baryon and lepton annihilating DM, compressed WIMPs, co-annihilating DM
5. Winos, Higgsinos, Precision Capture, Pasta Capture
6. Muonphilic
7. Asymmetric - converts NSs into black holes

Kouvaris 2007 (more...)
Bertone, Fairbairn 2007
JB Delgado, Martin 2017
Baryakhtar, JB, Li, Linden, Raj 2017
Raj, Tanedo, Yu 2017
Acevedo, JB, Leane, Raj 2019
Bell, Busoni, Robles 2019
Joglekar, Raj, Tanedo, Yu 2019
Dasgupta, Gupta, Ray 2019
Kopp, Laha, Opferkuch, Shepherd 2018
Jin, Gao 2018
Hamaguchi, Nagata, Yanagi 2019
Garani, Genolini, Hambye 2018
Keung, Marfatia, Tseng 2020
Bai, Berger, Korwar, Orlofsky 2020
Camargo, Queiroz, Sturani 2019
Bell, Busoni, Robles 2020
Garani, Heeck 2019
Goldman, Nussinov 1989
Kouvaris, Tinyakov 2011
McDermott, Yu, Zurek 2011
JB, Fukushima, Kumar 2013
Bell, Melatos, Petraki 2013
Bertoni, Nelson, Reddy 2014
JB, Linden 2014
JB, Elahi 2015



-At early times, NS cooling depends on neutrino emission and crust opacity, along with magnetic heating effects

-At late times, NSs are isothermal, photon emission dominated

-Early temperature discrepancies disappear after ~ 10 Myr,

Yakovlev & Pethick astro-ph/0402143

Page, Lattimer et al. astro-ph/0403657

$$T_{eff}^{\infty} \sim 100 \text{ K} \left(\frac{\text{Gyr}}{t} \right)^{1/2}$$

Coldest Known Neutron Star

PSR J2144–3933

Coldest known neutron star, nearby, deep imaged by HST, observed no emission setting a bound (Guillot et al, 1901.07998)

$$T_s < 3 \times 10^4 \text{ K}$$

This isn't cold enough to set a bound on local non-subhalo DM, since 0.3 GeV/cm^3 maximally heats NS to $\sim 1750\text{-}2500 \text{ K}$



Looking for WIMPs with 30+ meter telescopes

ELT 2σ sensitivity estimates

annihilation of WIMPs, Higgsinos

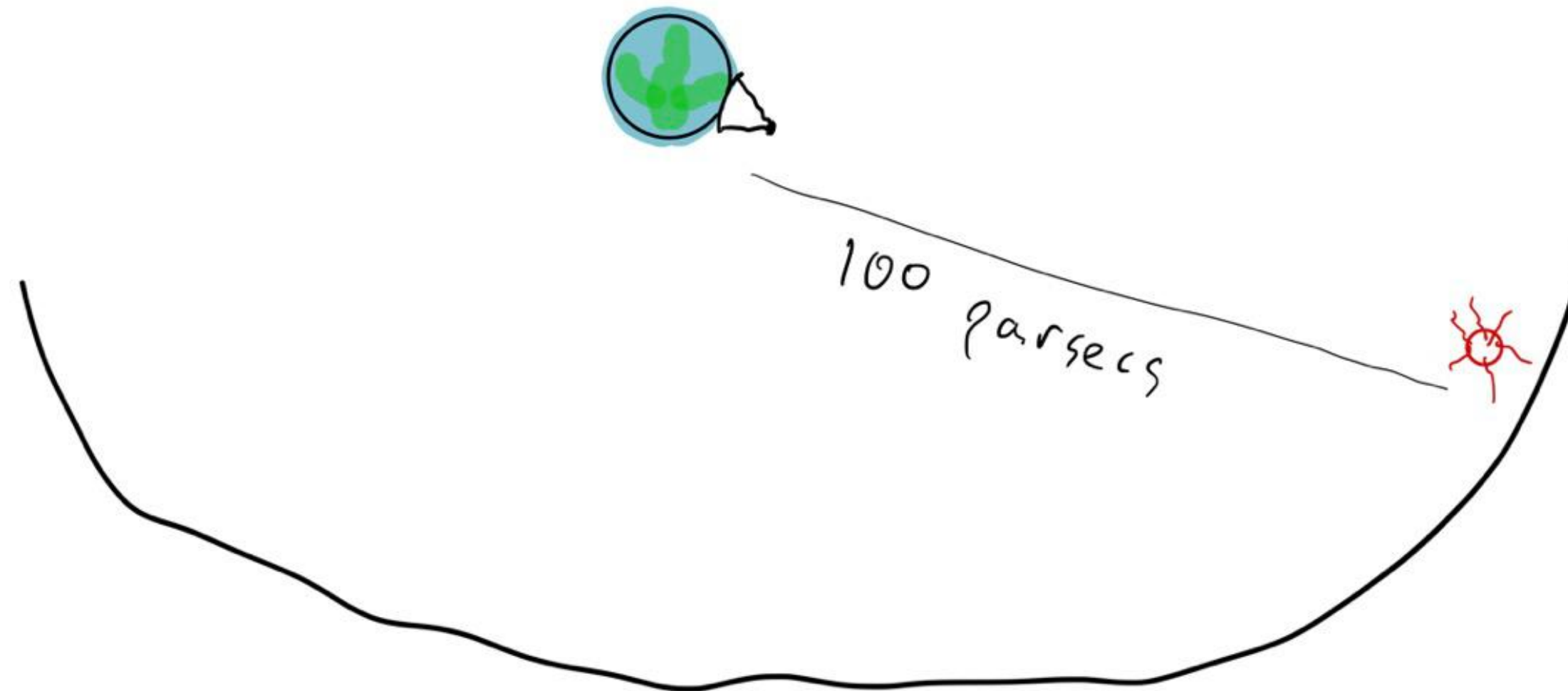
$$t \sim 3 \times 10^6 \text{ sec} \left(\frac{d}{100 \text{ pc}} \right)^4 \quad (\text{Y band})$$

kinetic only

$$t \sim 10^6 \text{ sec} \left(\frac{d}{30 \text{ pc}} \right)^4 \quad (\text{K band})$$

Baryakhtar, JB, Li, Linden, Raj 1704.01577

for recent JWST estimates see Garani et al. 2205.05048



Radio observations of nearby pulsars

-YMW16 dispersion measure distances

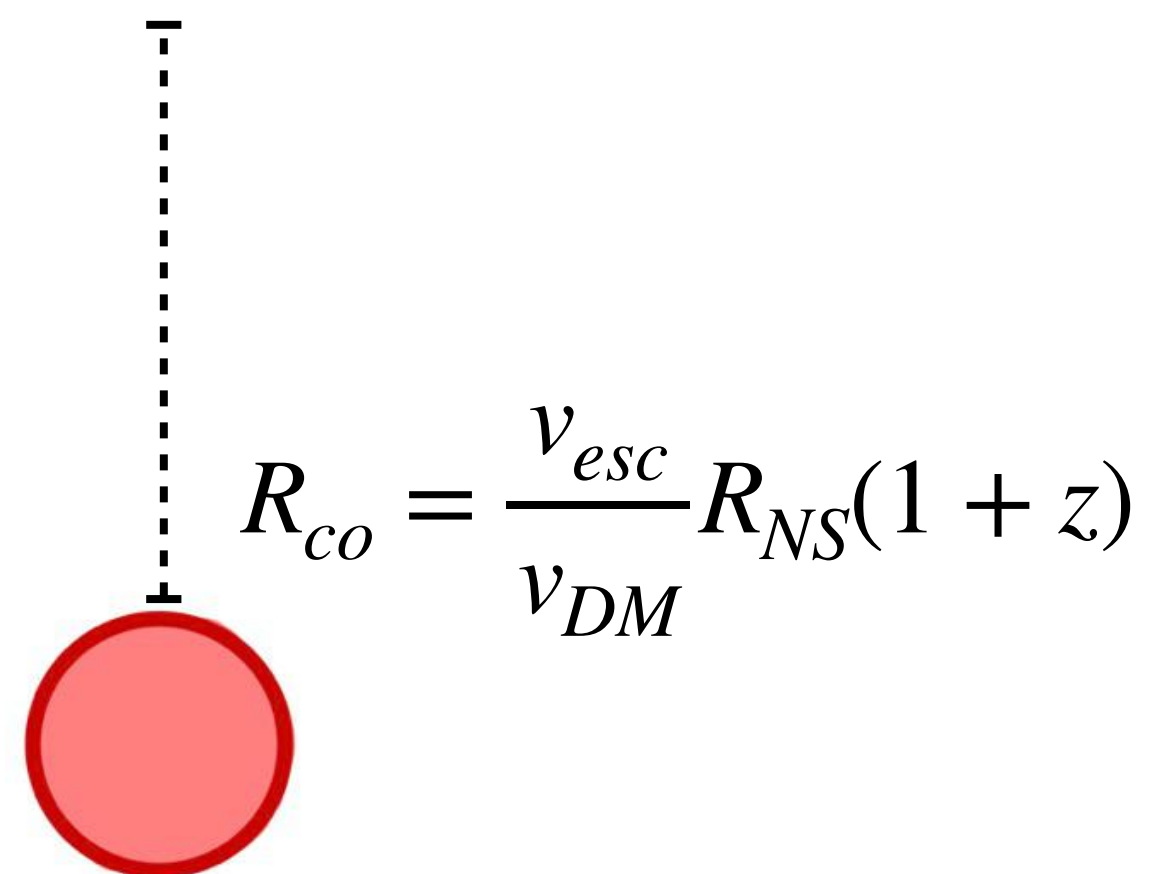
	<u>d (pc)</u>	<u>period (s)</u>
J1057-5226	90	0.19
J0736-6304	95	4.86
J0834-60	100	0.38
J0711-6830	110	0.005
J0749-68	110	0.91
J0924-5814	110	0.71

Accretion onto NSs, generalized

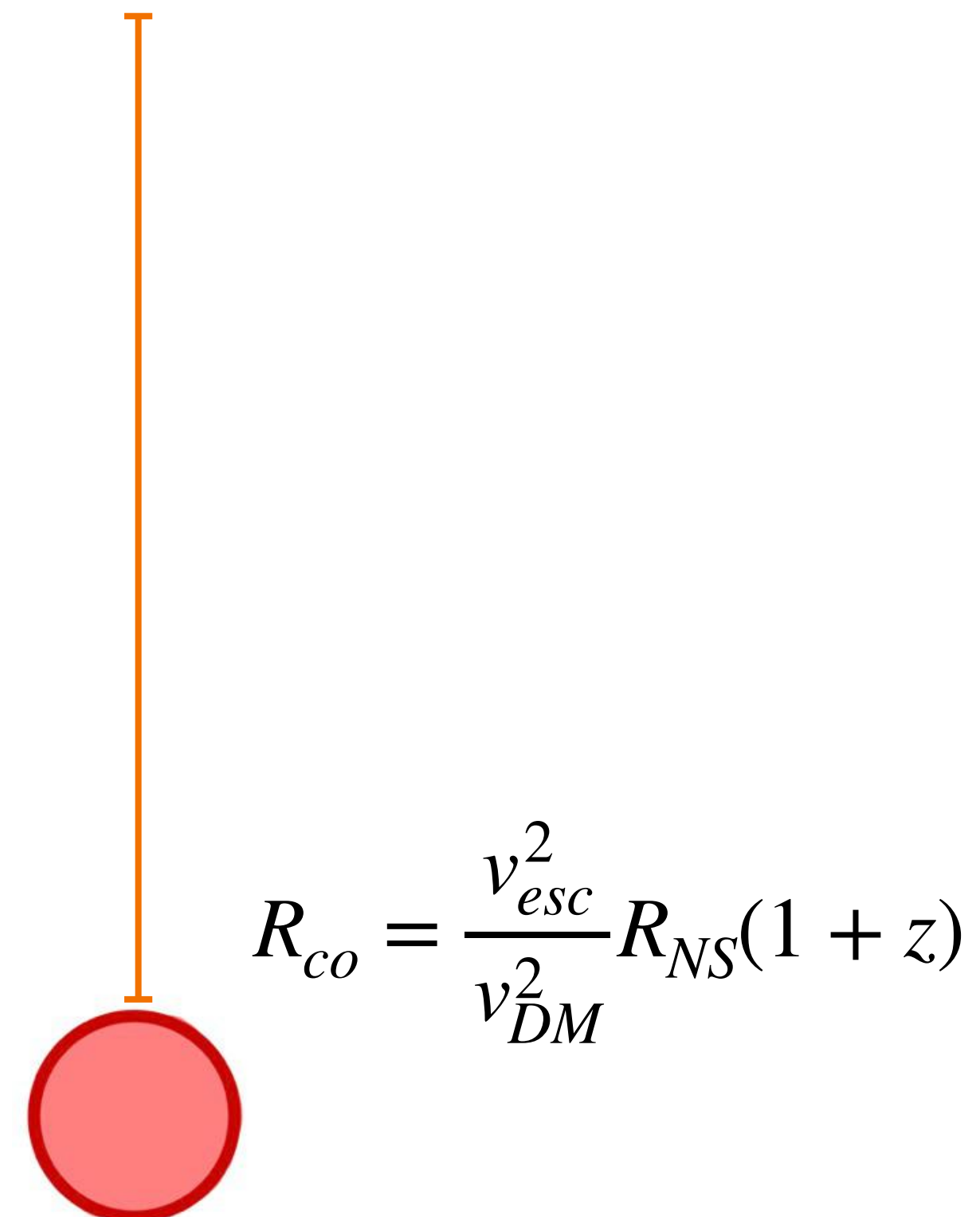
JB, Kavanagh, Raj 2109.04582

- Radius for DM accretion onto NS depends on size of DM subhalo, and whether accretion is collisionless or collisional

○ Collisionless

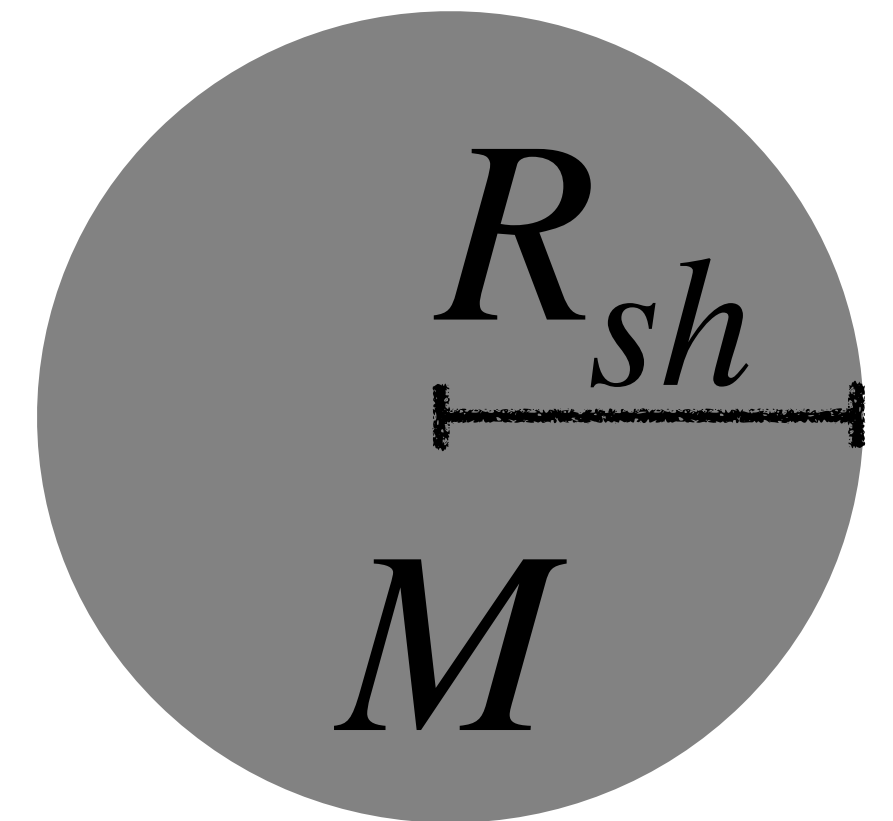


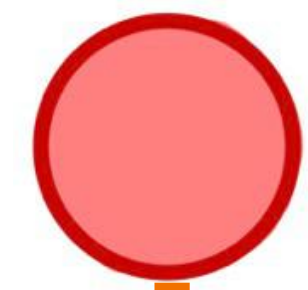
○ Collisional/Bondi



Subhalo DM

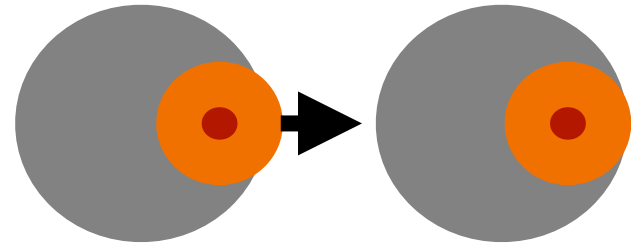
- Spherical subhalos, constant density, single mass/radius
- Subhalo DM composes most DM
- Consider both collisionless and collisional
- Future exploration:
 - Nontrivial profiles (e.g. NFW, Einasto, boson star)
 - Spectrum of subhalo masses



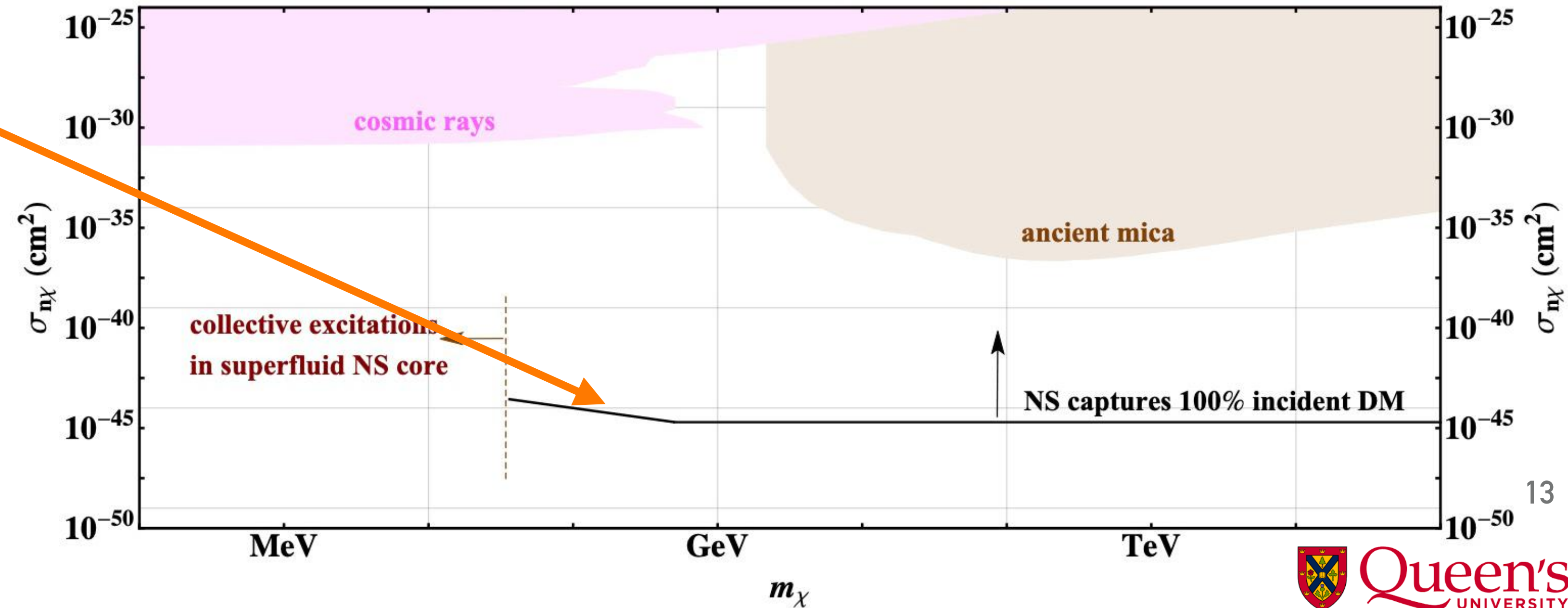
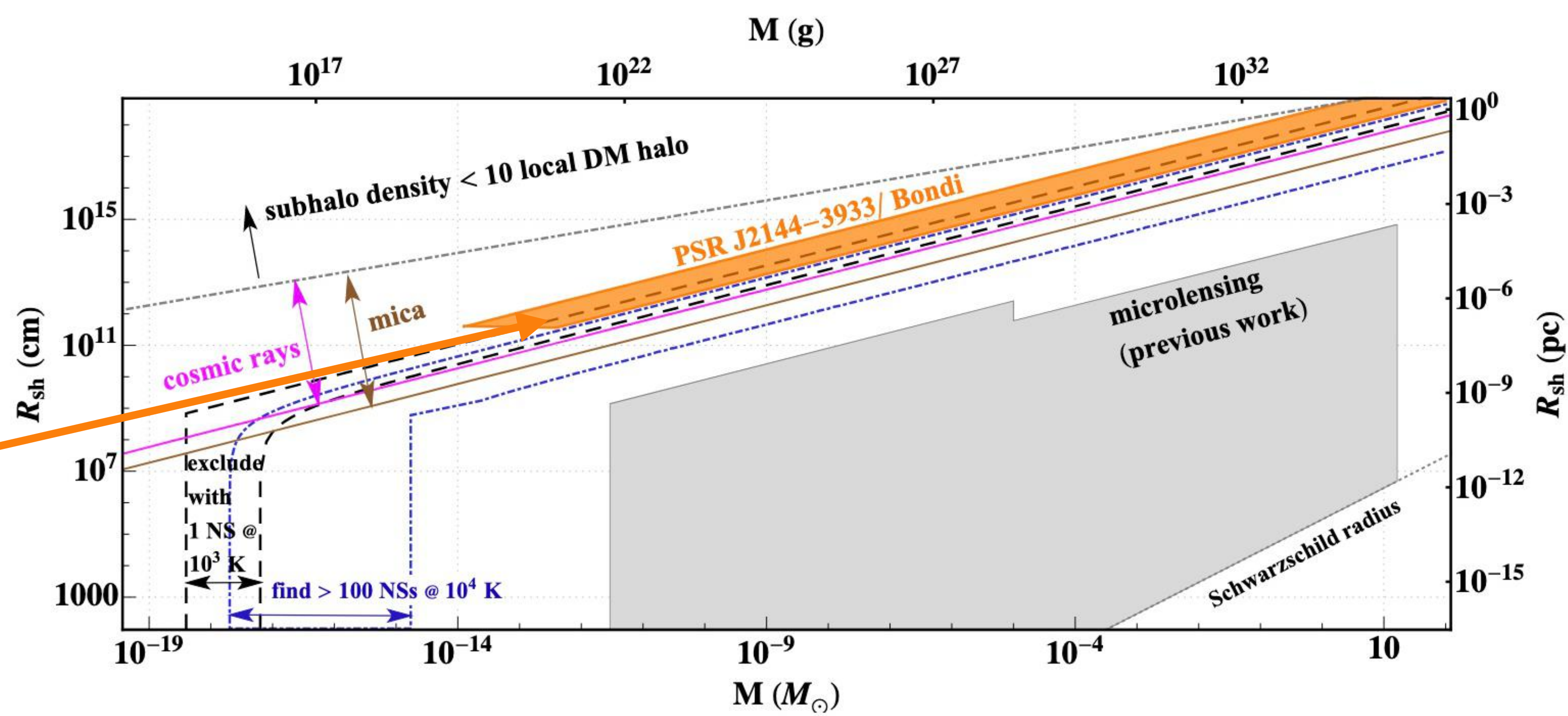


$$R_{co} = \frac{v_{esc}^2}{v_{DM}^2} R_{NS}(1+z)$$

PSR J2144 $T < 3 \times 10^4$ K
excludes this region
for Bondi accretion



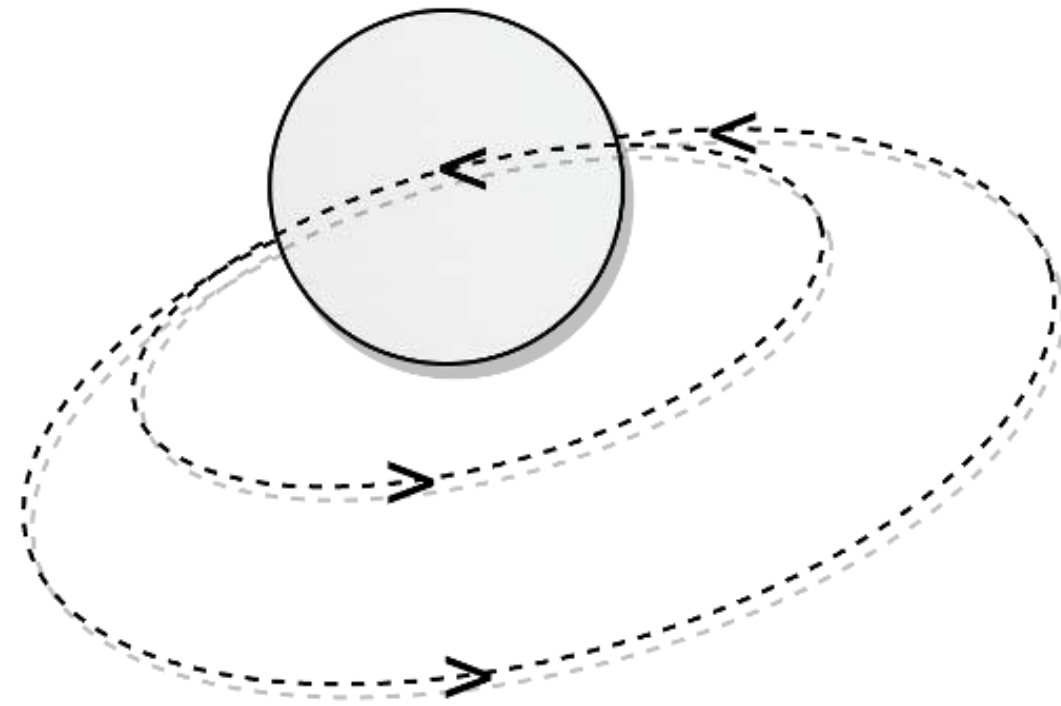
for this DM-nucleon cross-section



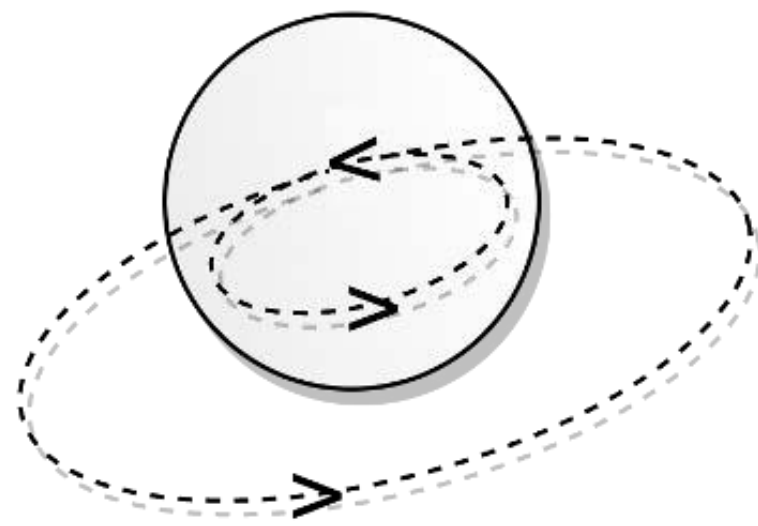
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DM thermalizing in compact stars



Dark matter is captured, undergoes successive scatters decreasing its orbital radius



Dark matter thermalizes to a radius $R < R_{NS}$, called t_1 thermalization



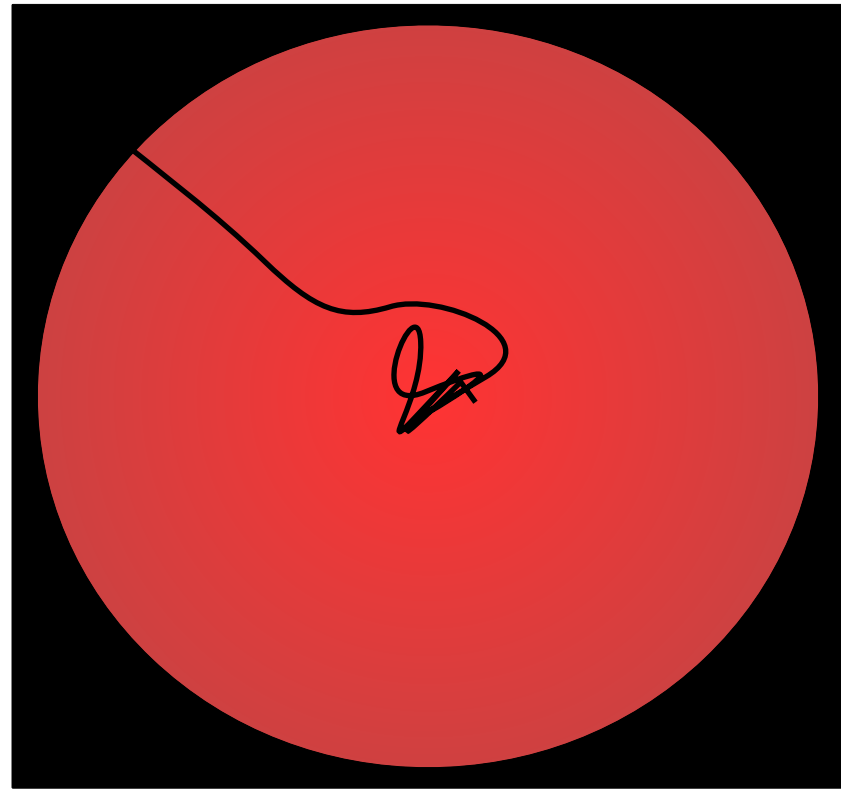
Dark matter continues to thermalize until $T_x = T_{NS}$ after a time called t_2 thermalization



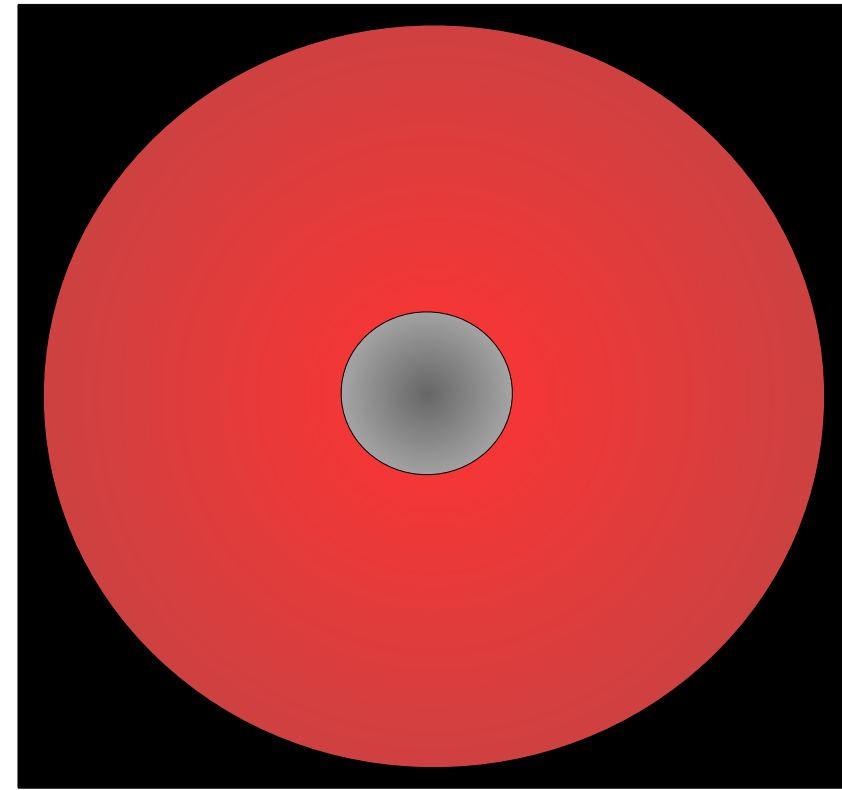
Narayani Tyagi

Asymmetric dark matter in compact stars

1. DM captured



2. DM thermalizes

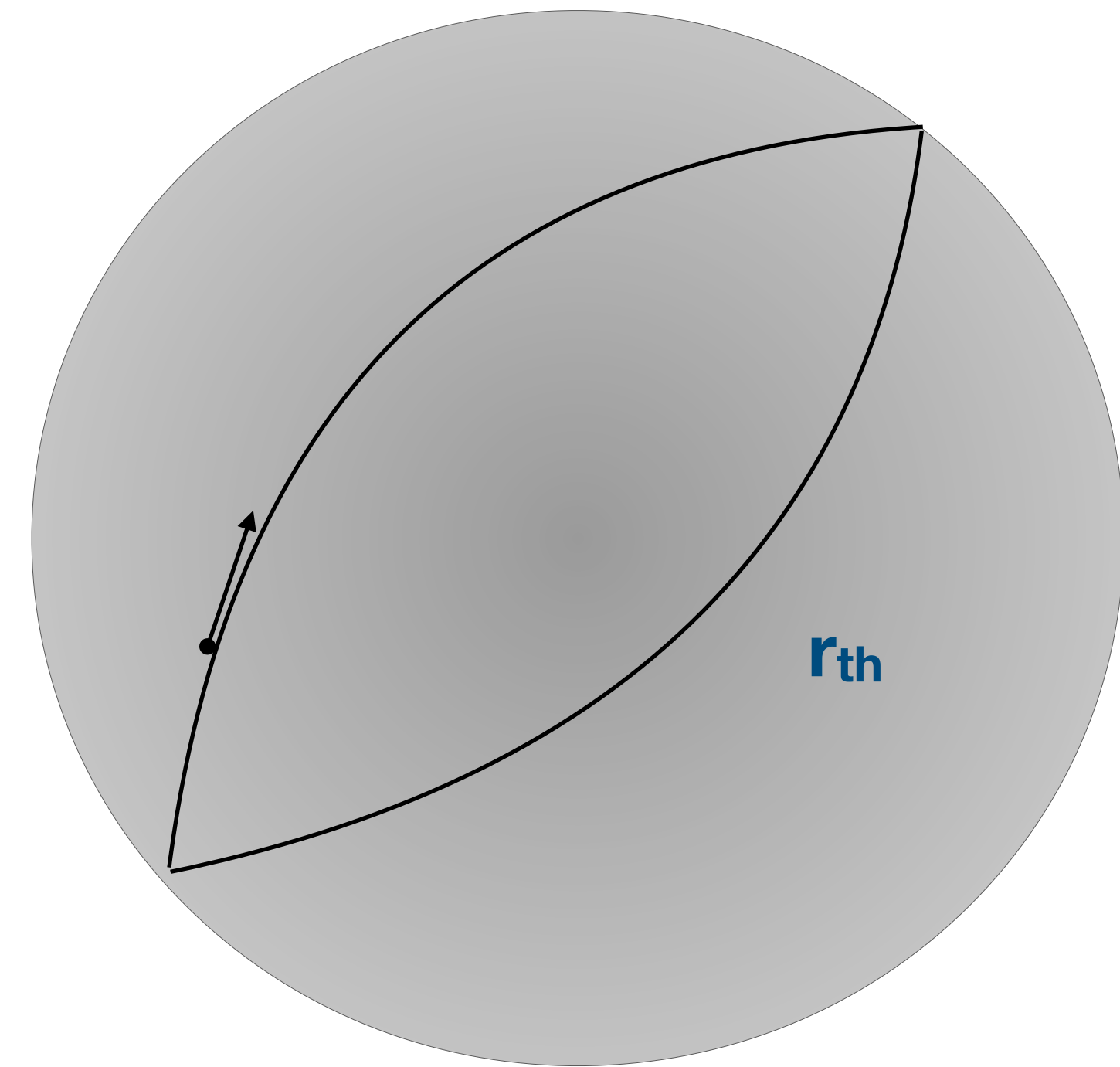


Harmonic oscillator potential

$$k_B T \sim G \rho_{wd} m_x r_{th}^2$$

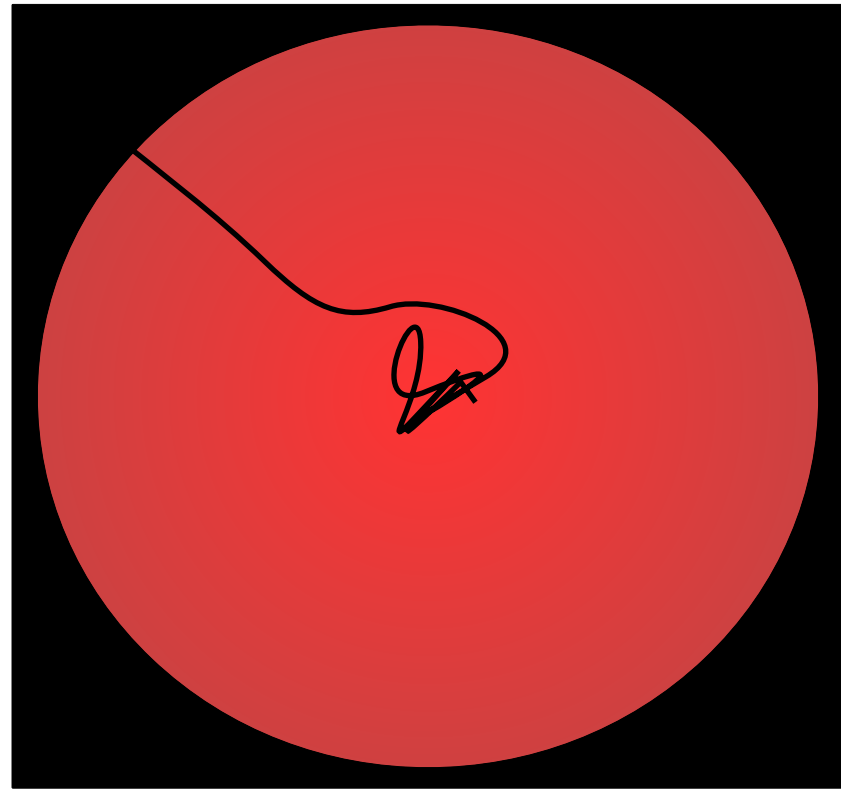
Thermalization radius

$$r_{th} \sim 1 \text{ millimeter} \sqrt{\frac{\text{PeV}}{m_x}} \sqrt{\frac{T_{ns}}{10^5 \text{ K}}}$$

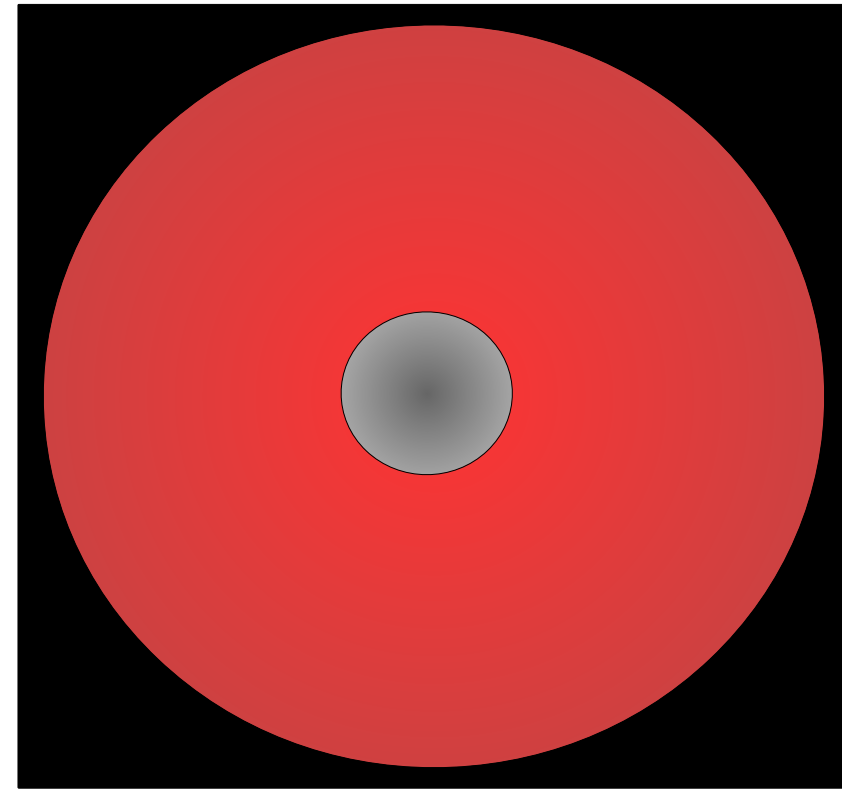


Asymmetric dark matter in compact stars

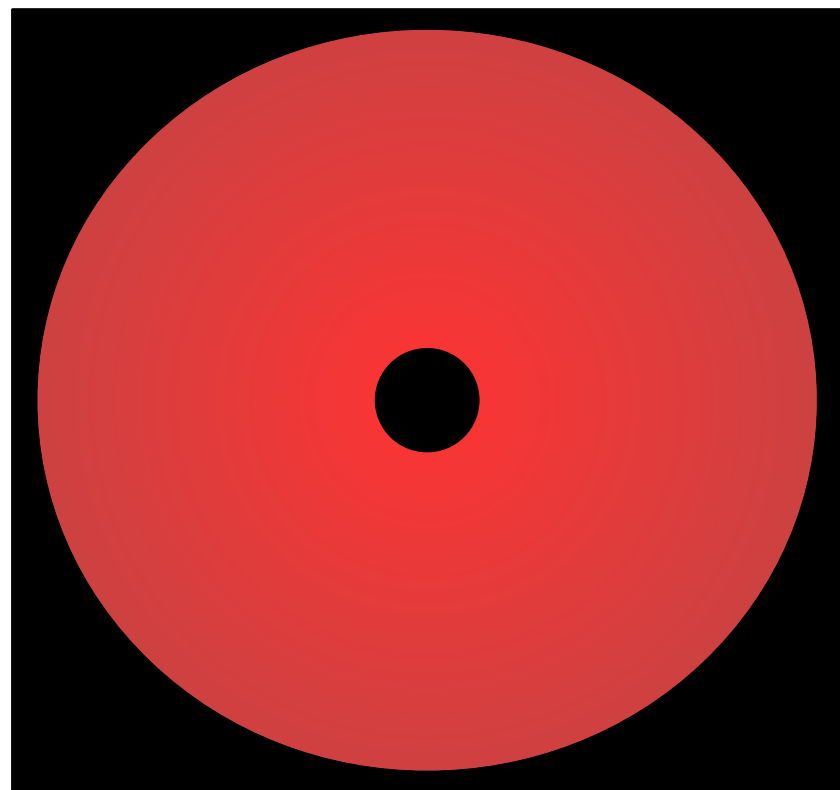
1. DM captured



2. DM thermalizes



3. DM collapses



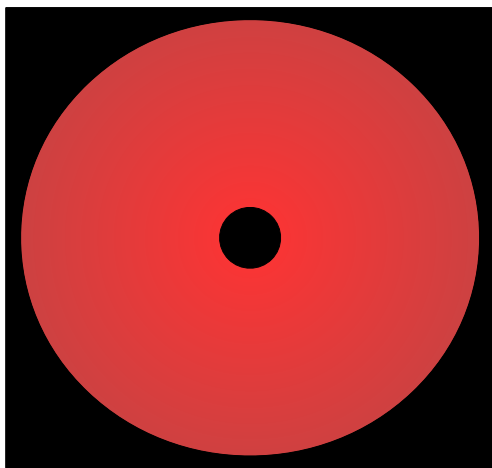
DM will collapse to a black hole if it

1. Self-gravitates $\rho_{DM} > \rho_{ns}$

2. Exceeds its own degeneracy pressure

$$M_{crit}^{ferm} \simeq M_{pl}^3 / m_X^2$$

($\sim 10^{-12}$ solar masses for PeV mass DM)



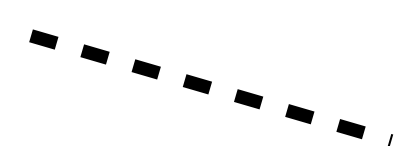
Dark matter that implodes pulsars

X ~GeV mass, asymmetric dark fermions — degeneracy pressure stabilizes up to a solar mass of dark matter.

m_X

KeV-PeV

✓ Bosonic dark matter without repulsive self interactions — requires very small effective quartic ($\lambda < 10^{-15}$).



$(\lambda \times)$

Bosonic DM Constraints
JB, Kumar, et al. 2013

PeV-EeV

✓ **Heavy** dark matter, fermionic or bosonic — fewer particles required for collapse.

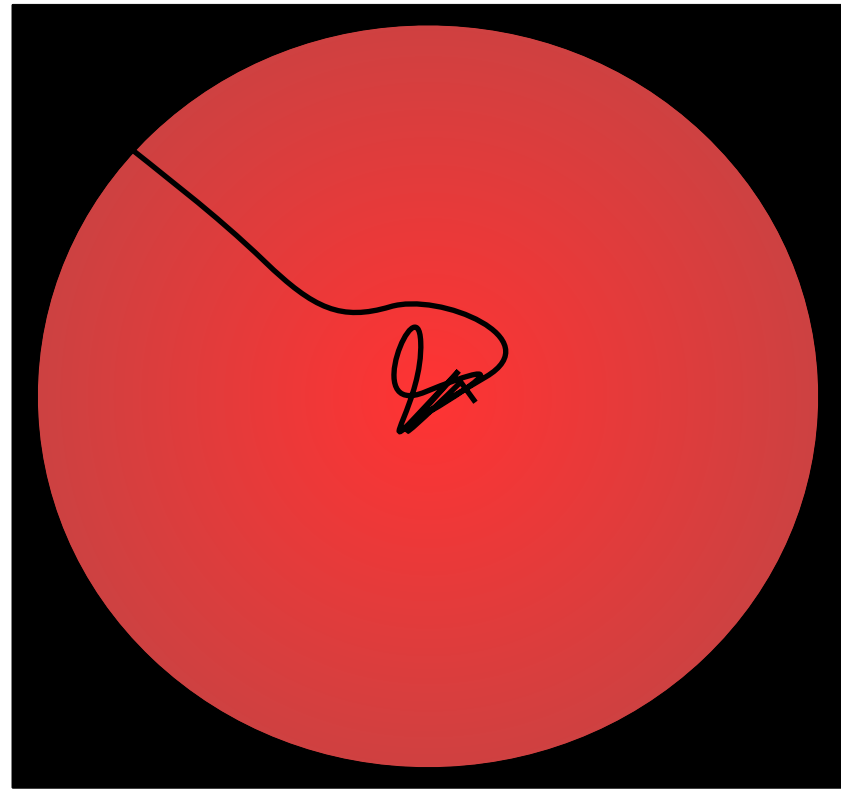
$$M_{crit}^{ferm} \simeq M_{pl}^3 / m_X^2$$
$$M_{crit}^{bos} \simeq \sqrt{\lambda} M_{pl}^3 / m_X^2$$

$\sim 10^{-12}$ solar masses

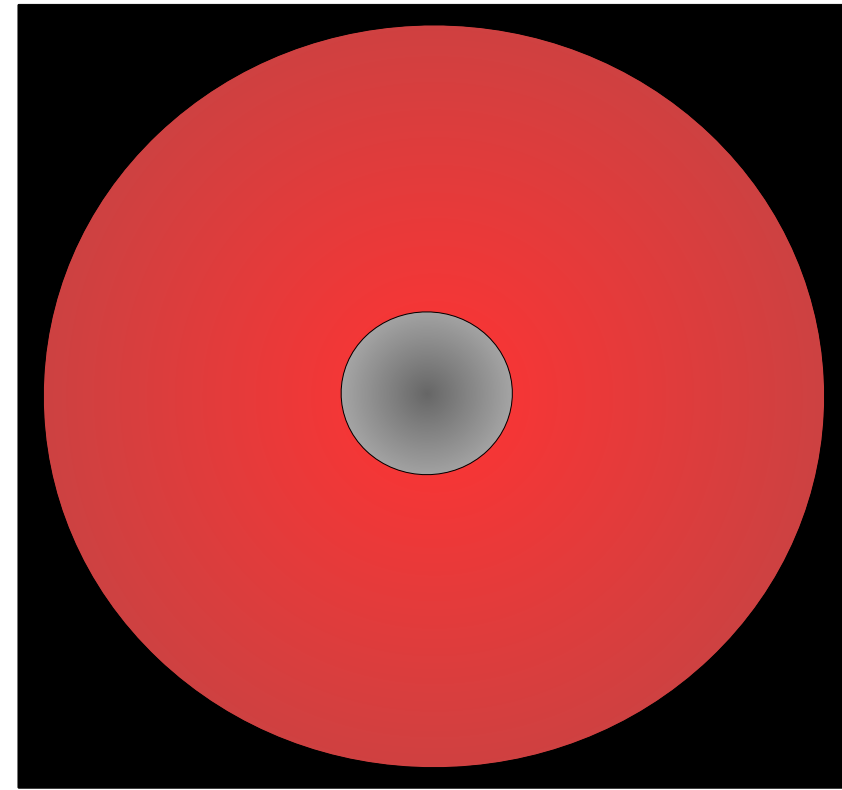


Asymmetric dark matter in compact stars

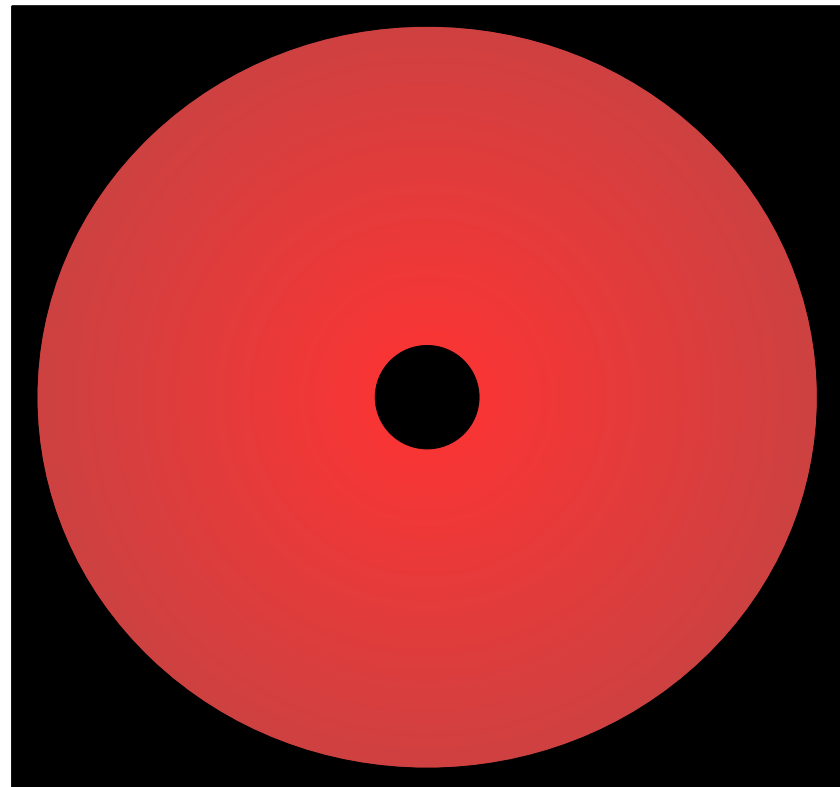
1. DM captured



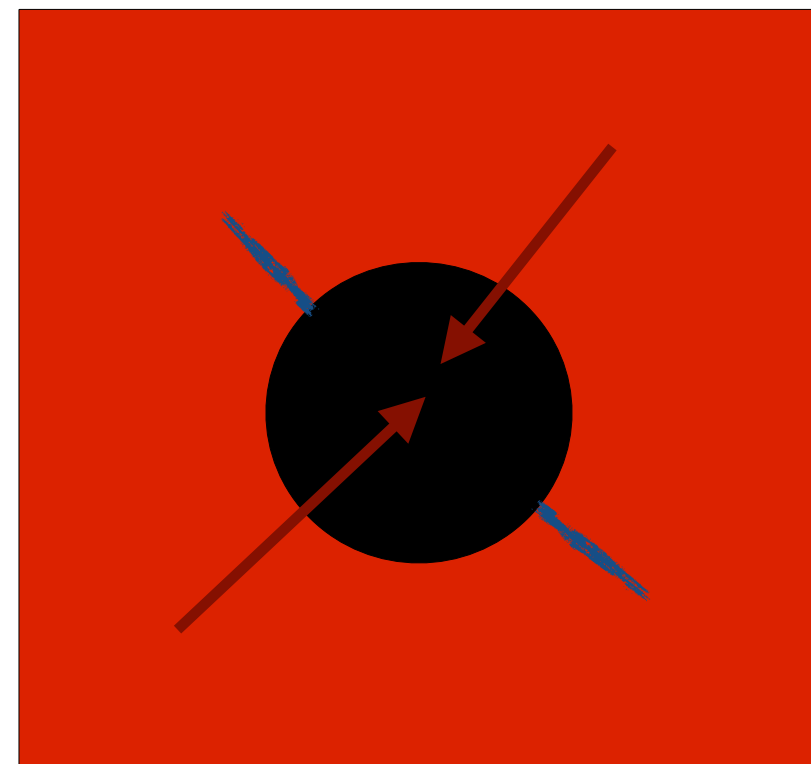
2. DM thermalizes



3. DM collapses



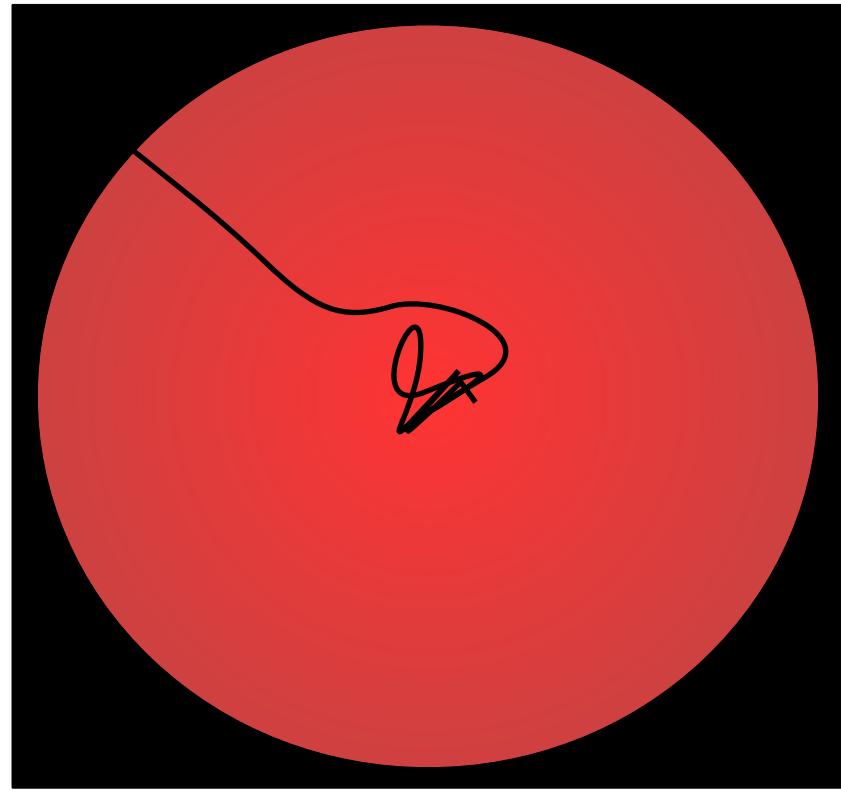
4. BH consumes neutron star



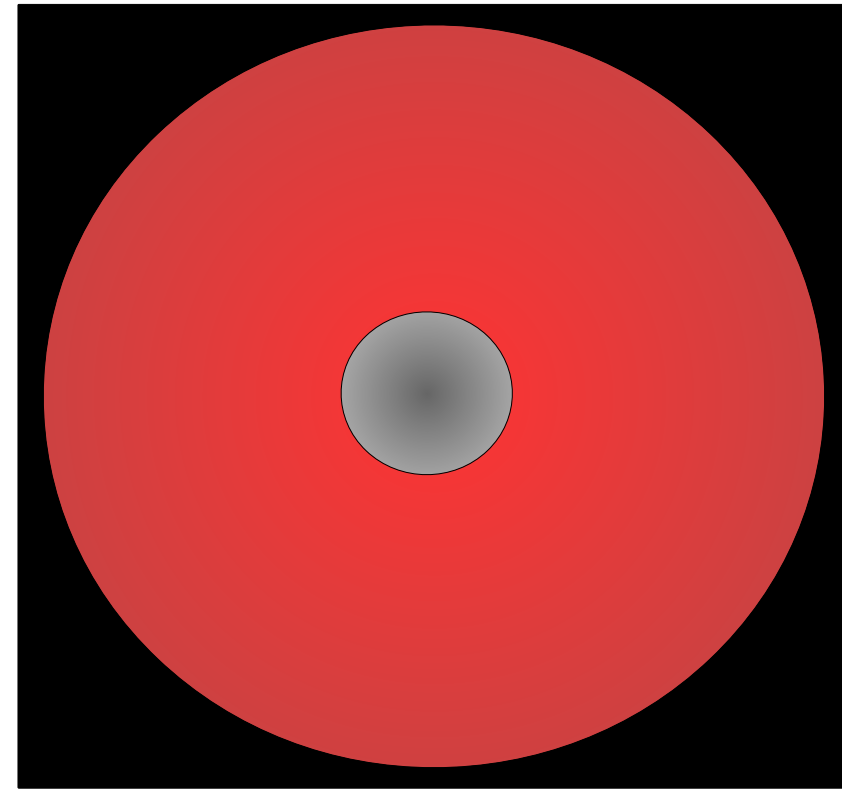
$$\frac{dM_{bh}}{dt} \approx \frac{4\pi\rho_{ns}(GM_{bh})^2}{v_s^3} - \frac{1}{15360\pi(GM_{bh})^2}$$

Asymmetric dark matter in compact stars

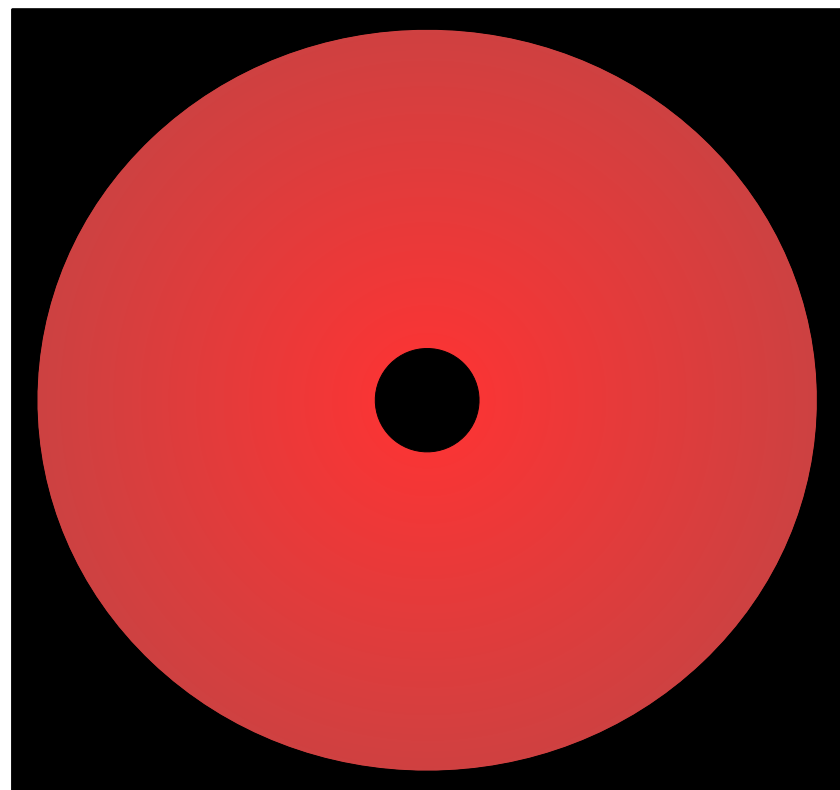
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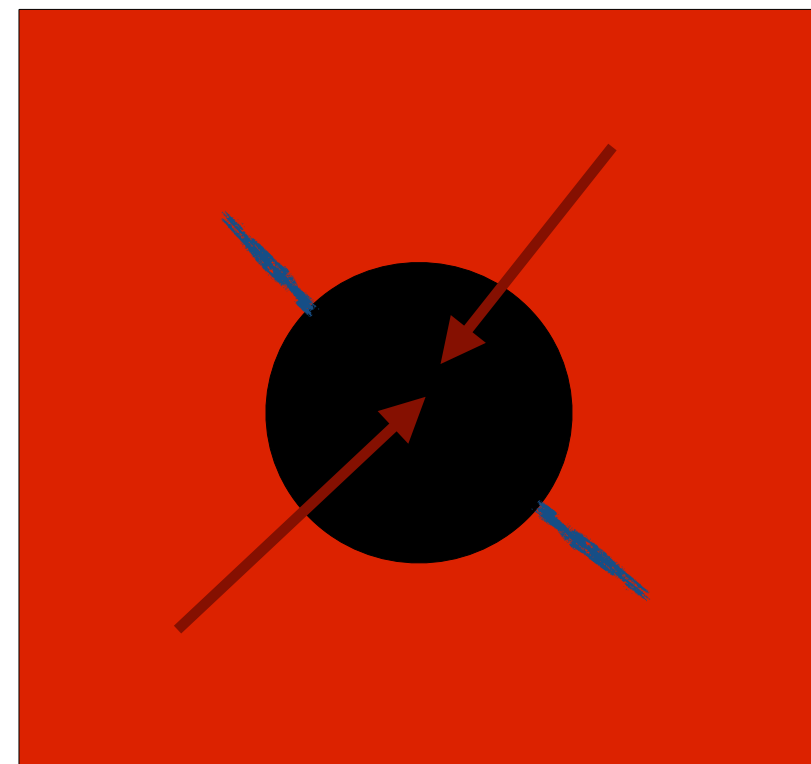
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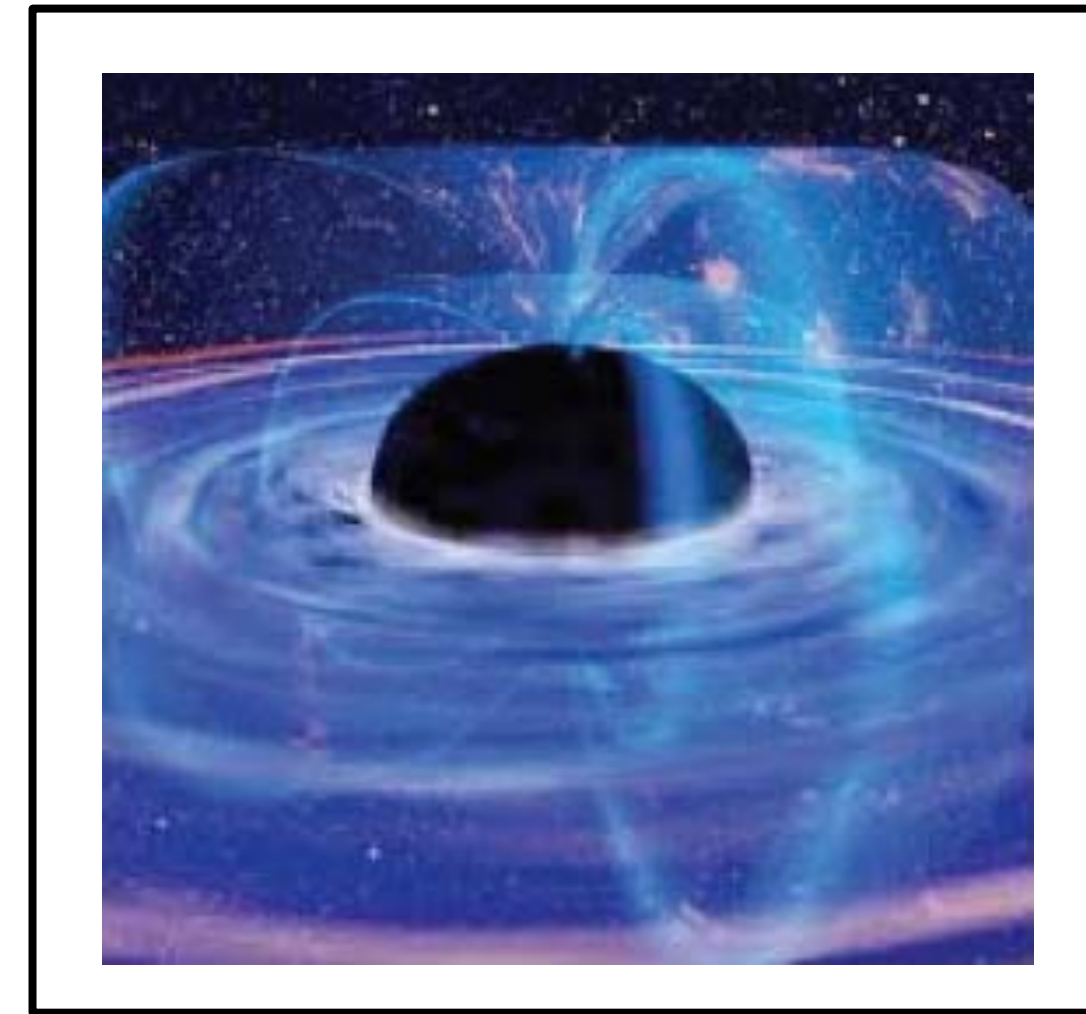
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4. BH consumes neutron star

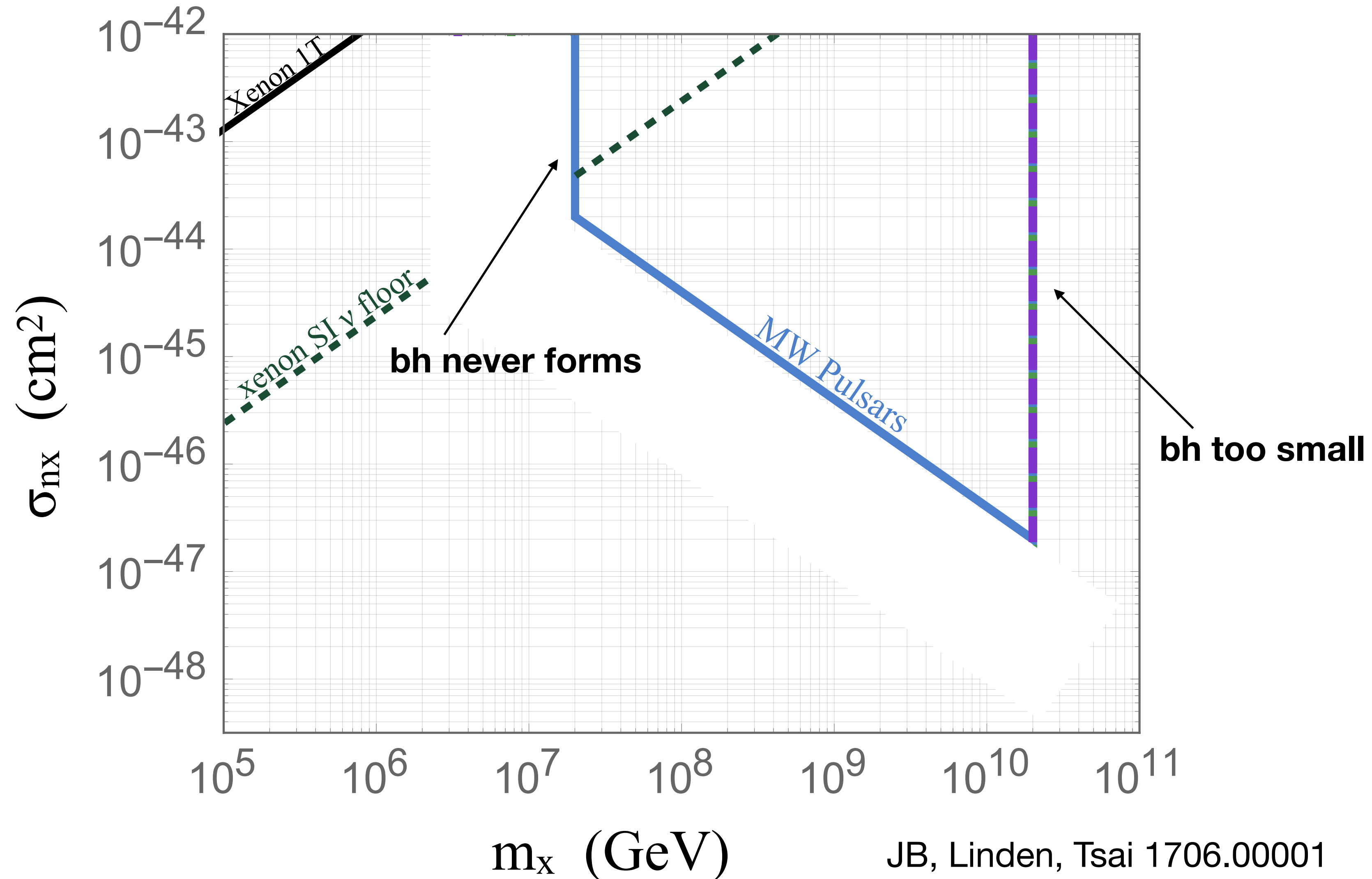


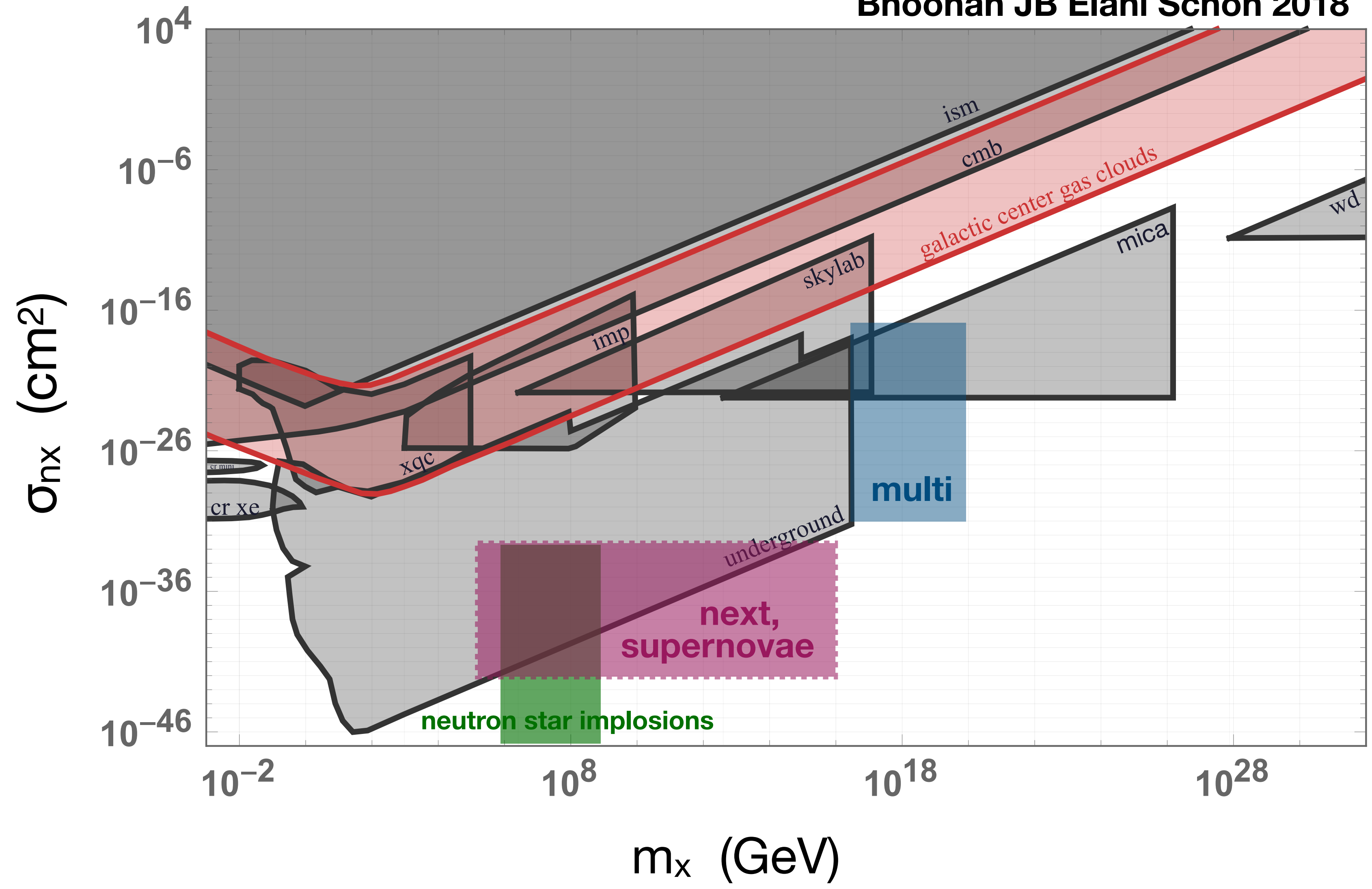
5. Form solar mass BH



$$\frac{dM_{bh}}{dt} \approx \frac{4\pi\rho_{ns}(GM_{bh})^2}{v_s^3} - \frac{1}{15360\pi(GM_{bh})^2}$$

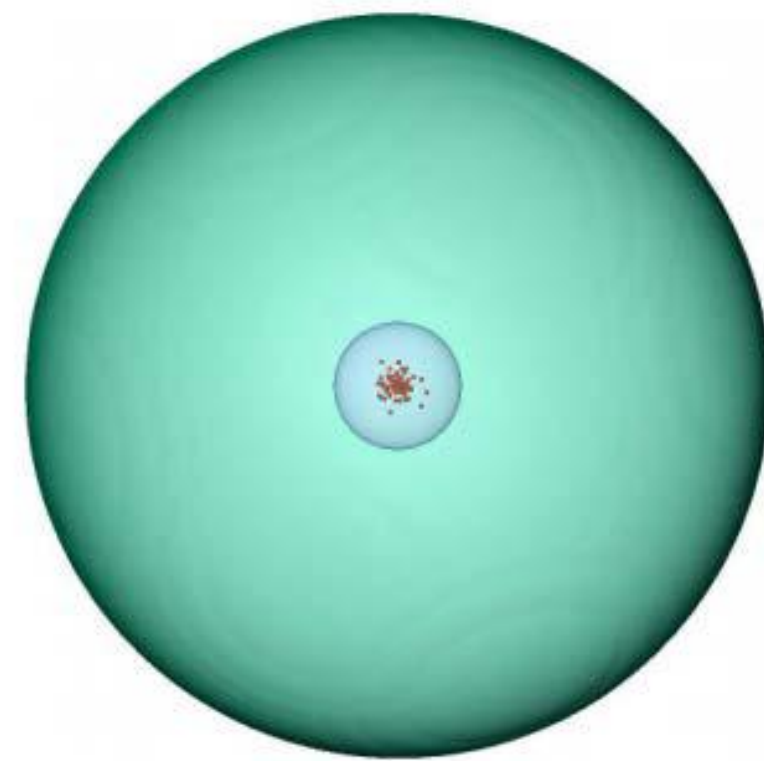
Using an old pulsar in the Milky Way, the best bound so far...





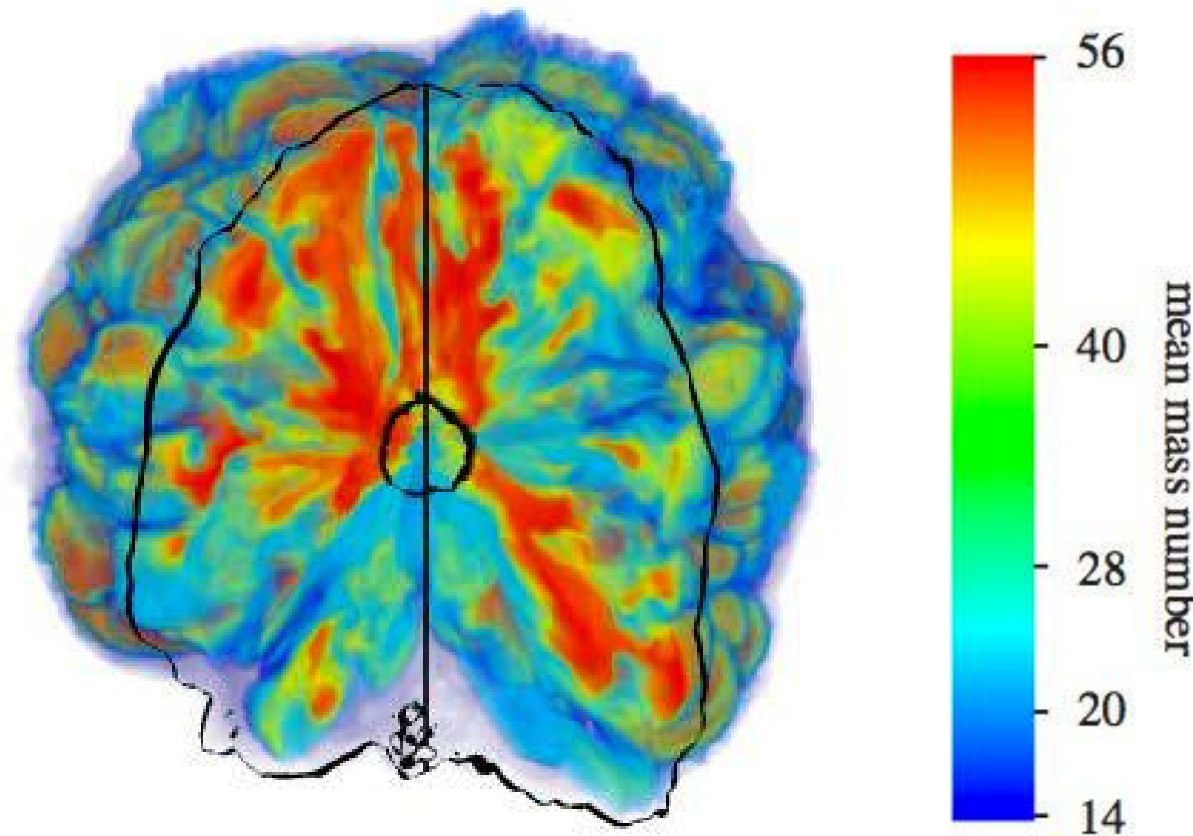
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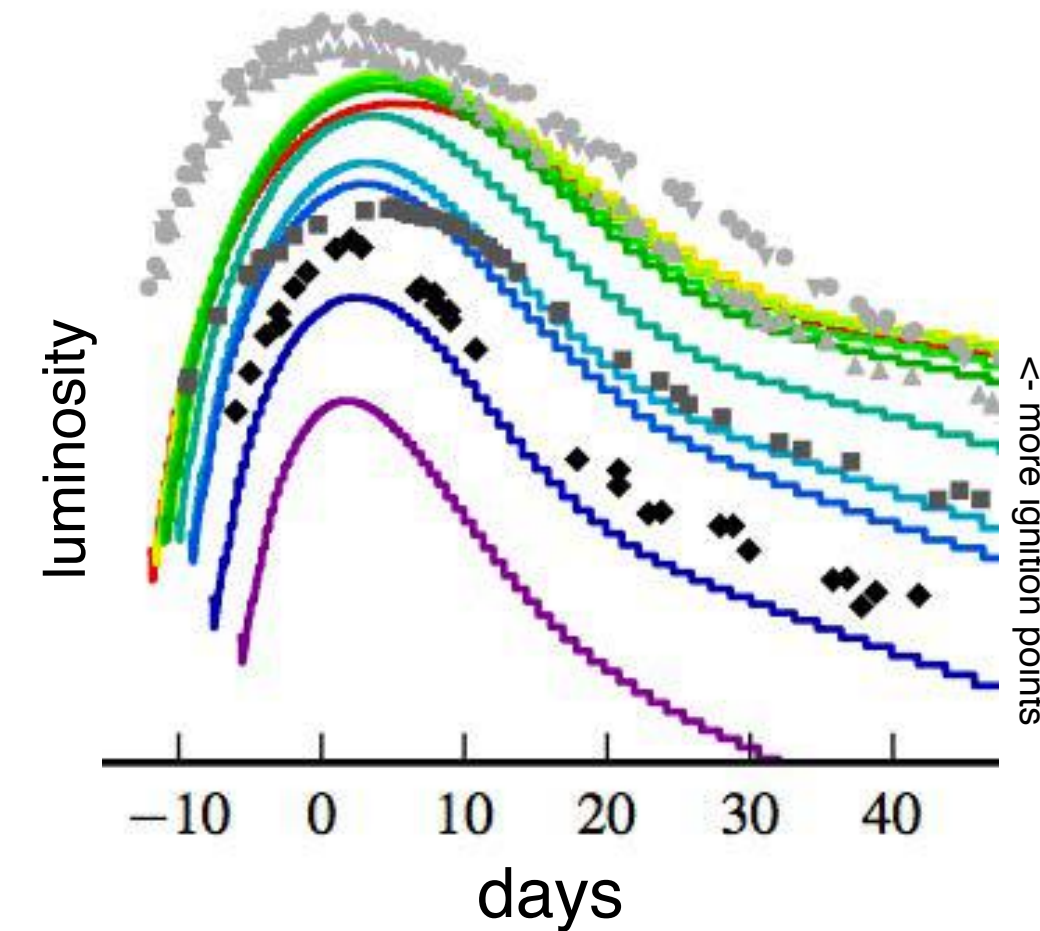


Ignition region inside WD

ignition



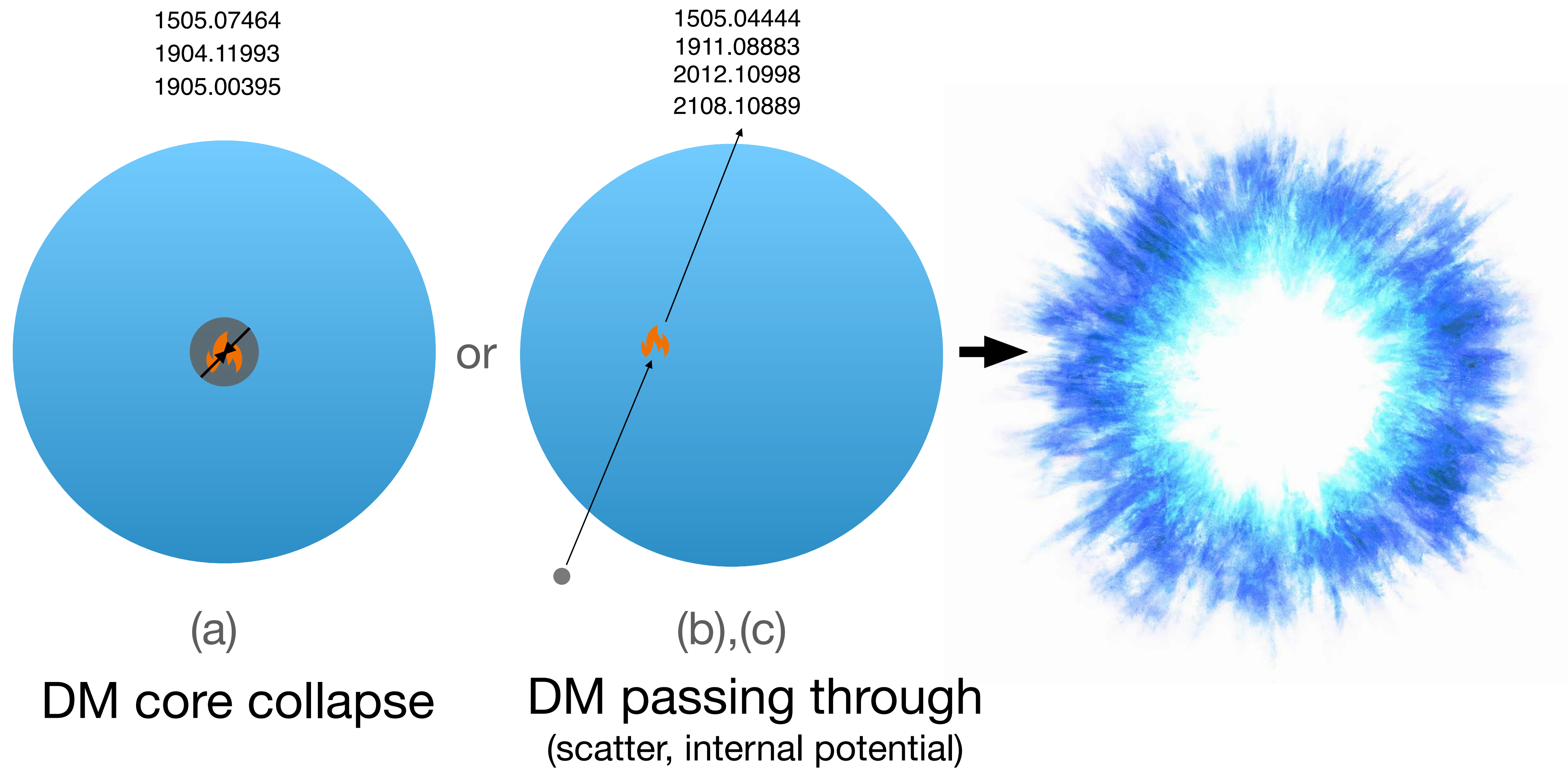
nuclei fuse (^{56}Ni)



nuclei decay

(figures from 1308.3257)

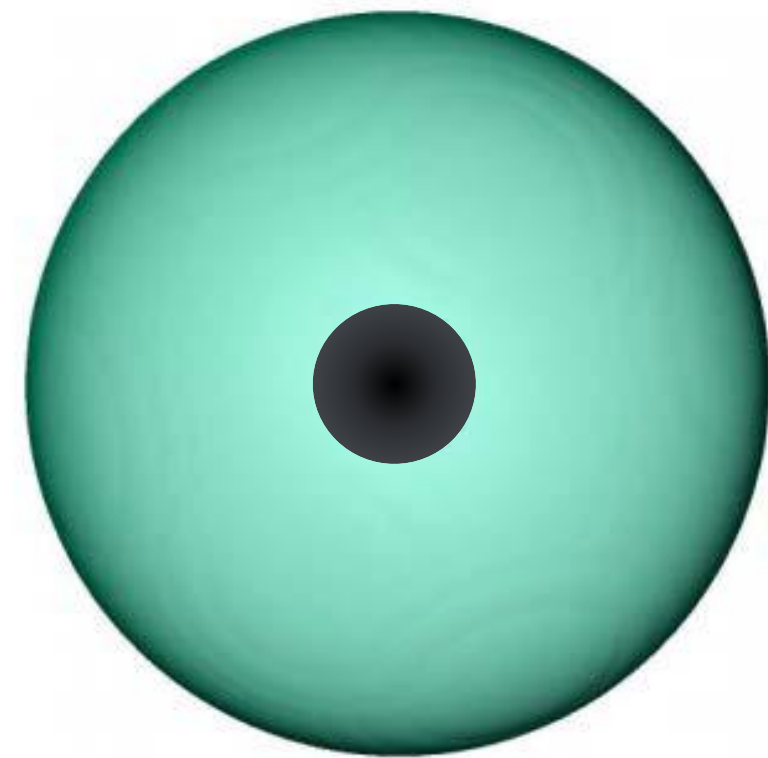
Type Ia supernovae erupt when a portion of white dwarf becomes heated, igniting a thermonuclear flame-front that sweeps through the star, followed by an explosion.



Two primary methods to explode WDs with dark matter.

Heavy Dark Matter Ignition of Type Ia Supernovae

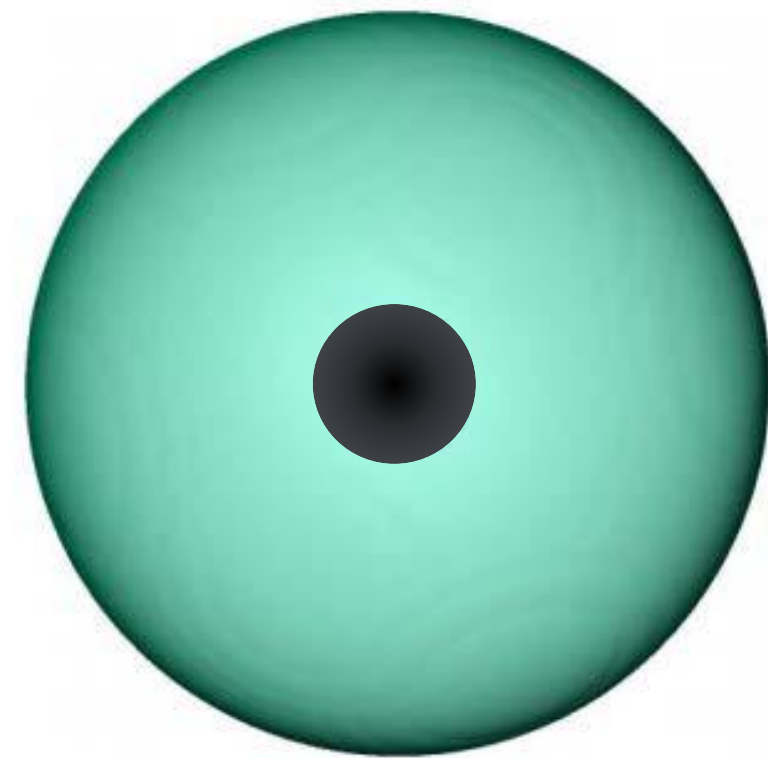
In order to ignite a carbon-oxygen white dwarf, the dark matter must be **heavy** so that it thermalizes inside a small volume within the white dwarf, and collects to the point of collapse within $\sim 10^{10}$ years.



DM collects to the point of self-gravitation.

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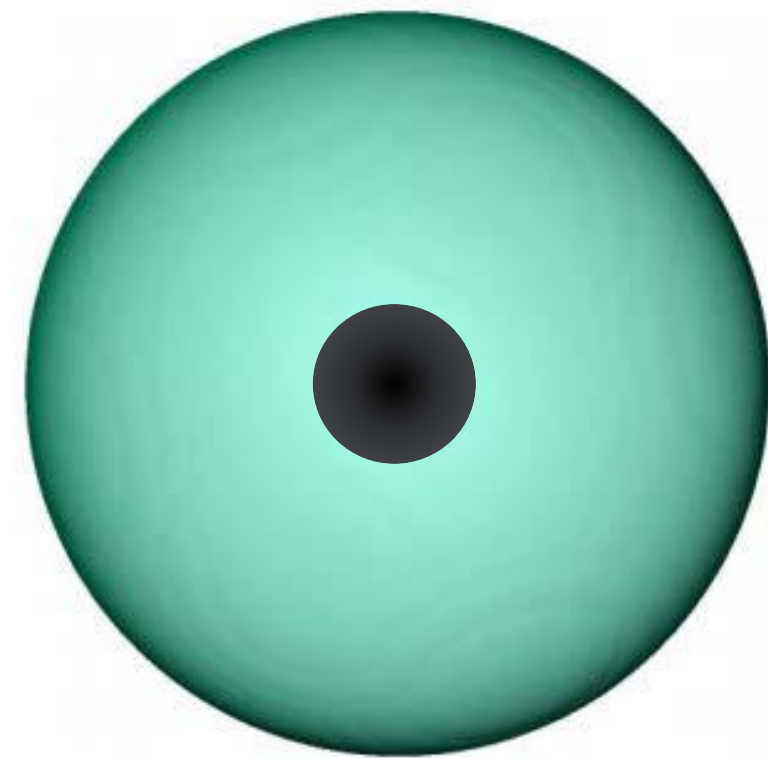
DM collects to the point of self-gravitation.

Harmonic Oscillator potential

$$k_B T \sim G \rho_{wd} m_x r_{th}^2$$

Heavy Dark Matter Ignition of Type Ia Supernovae

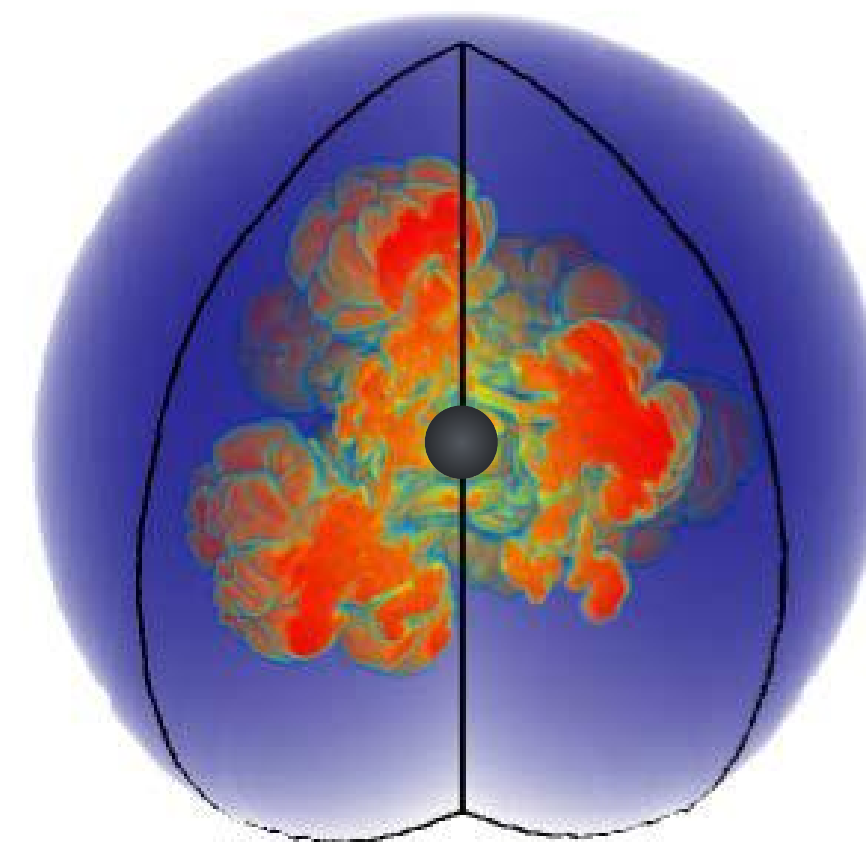
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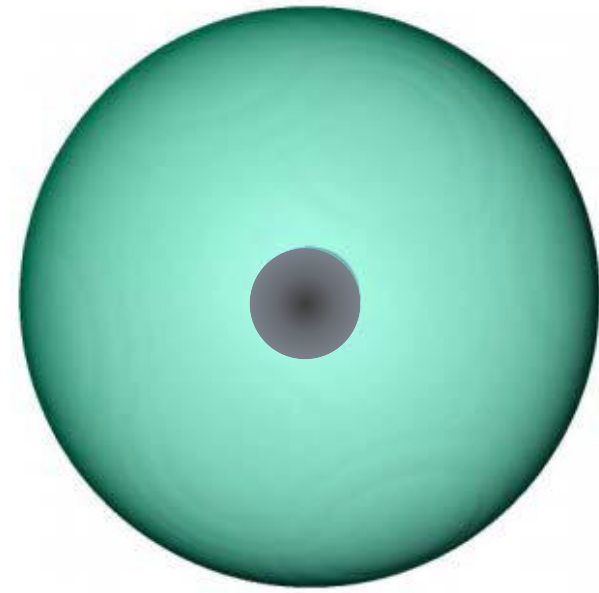
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Harmonic Oscillator potential

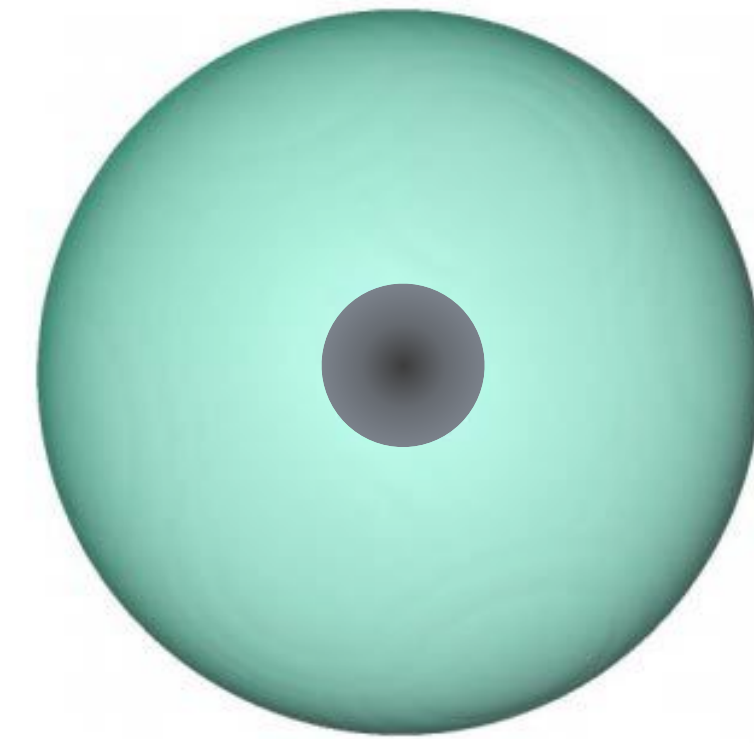
$$k_B T \sim G \rho_{wd} m_x r_{th}^2$$



DM collapses, shedding gravitational potential energy through **scattering**, igniting a SNIa.



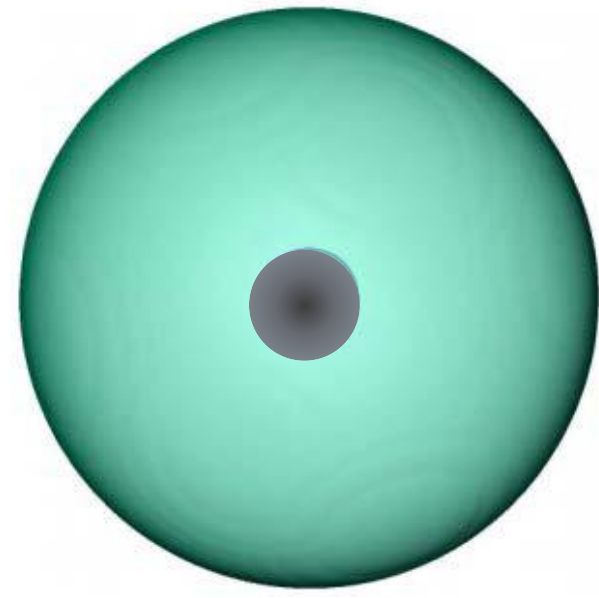
more massive WD,
DM collapses sooner



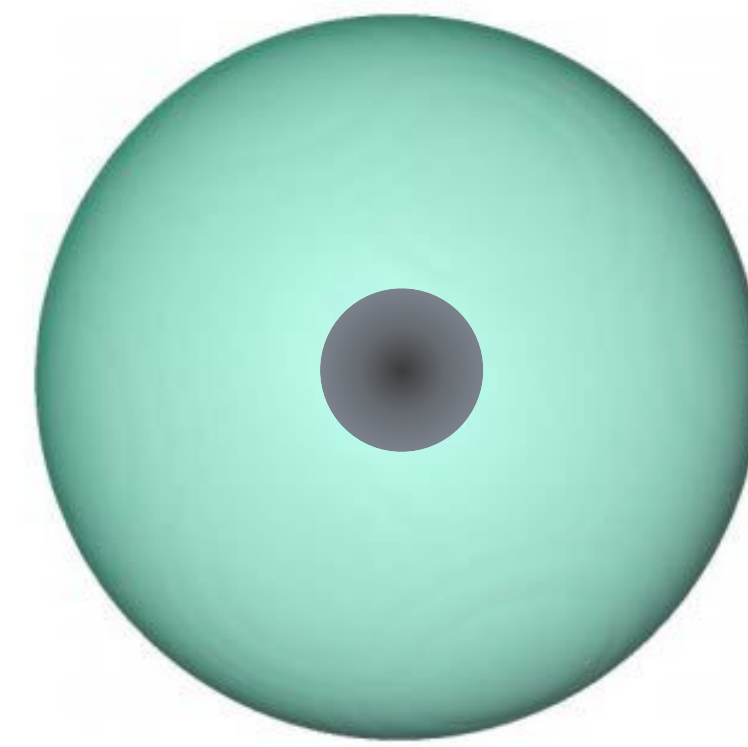
less massive WD,
DM collapses later

A more massive white dwarf is denser, so dark matter collects into a smaller ball, and collapses sooner.

Altogether, this shortens the time for dark matter collapse in more massive white dwarfs.

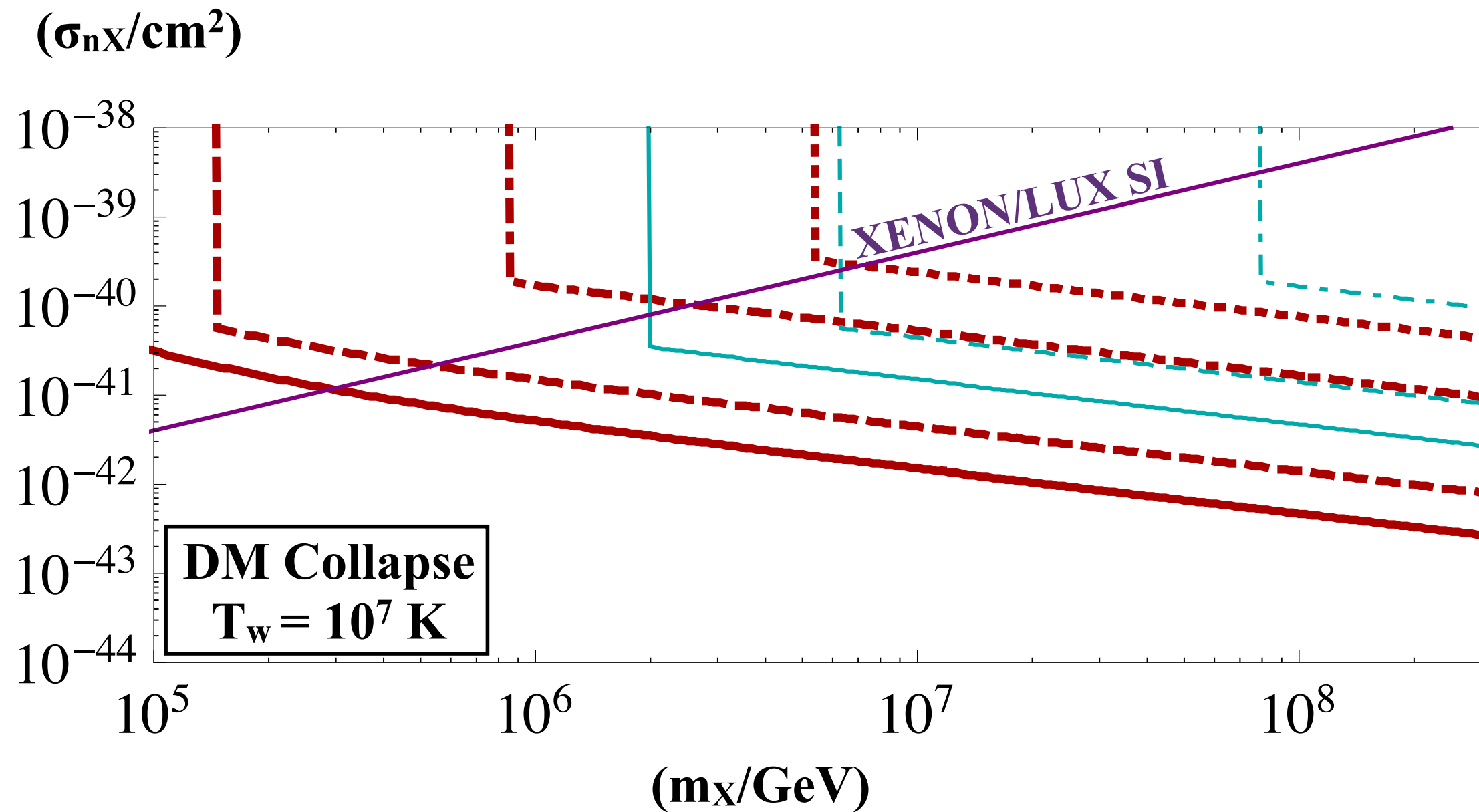


DM collapses sooner

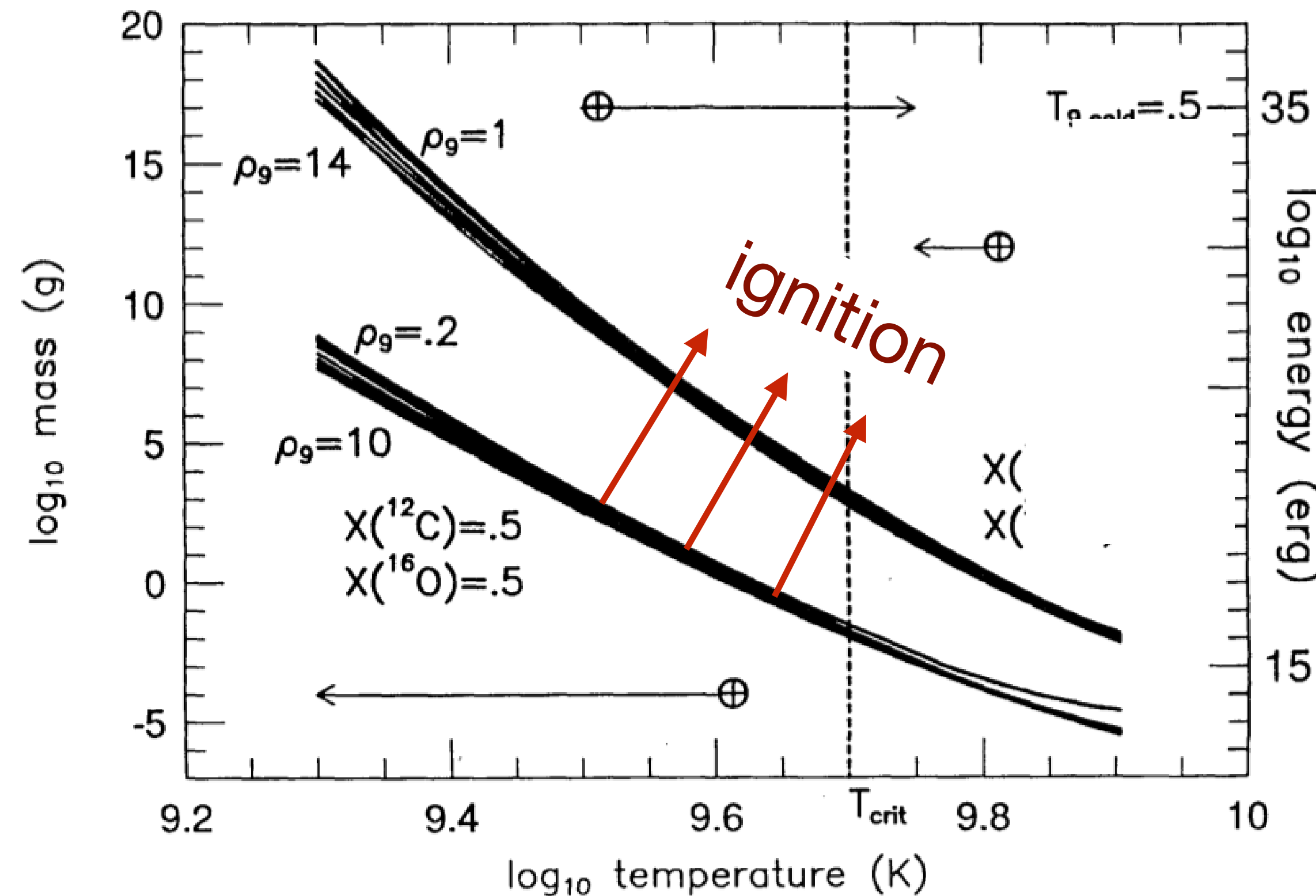


DM collapses later

Line	M_w/M_\odot	$R_w/(10^3 \text{ km})$	$\rho_w/(10^7 \text{ g/cm}^3)$	Line	t_w/Gyr
—	1.4	2.5	100	—	5
- · - · -	1.3	3	40	- · - · -	0.5
- · - -	1.1	5	6	- · - -	
- · · · -	0.9	6	2	- · · · -	











As it collapses to its minimum energy state, dark matter will shed gravitational potential energy.

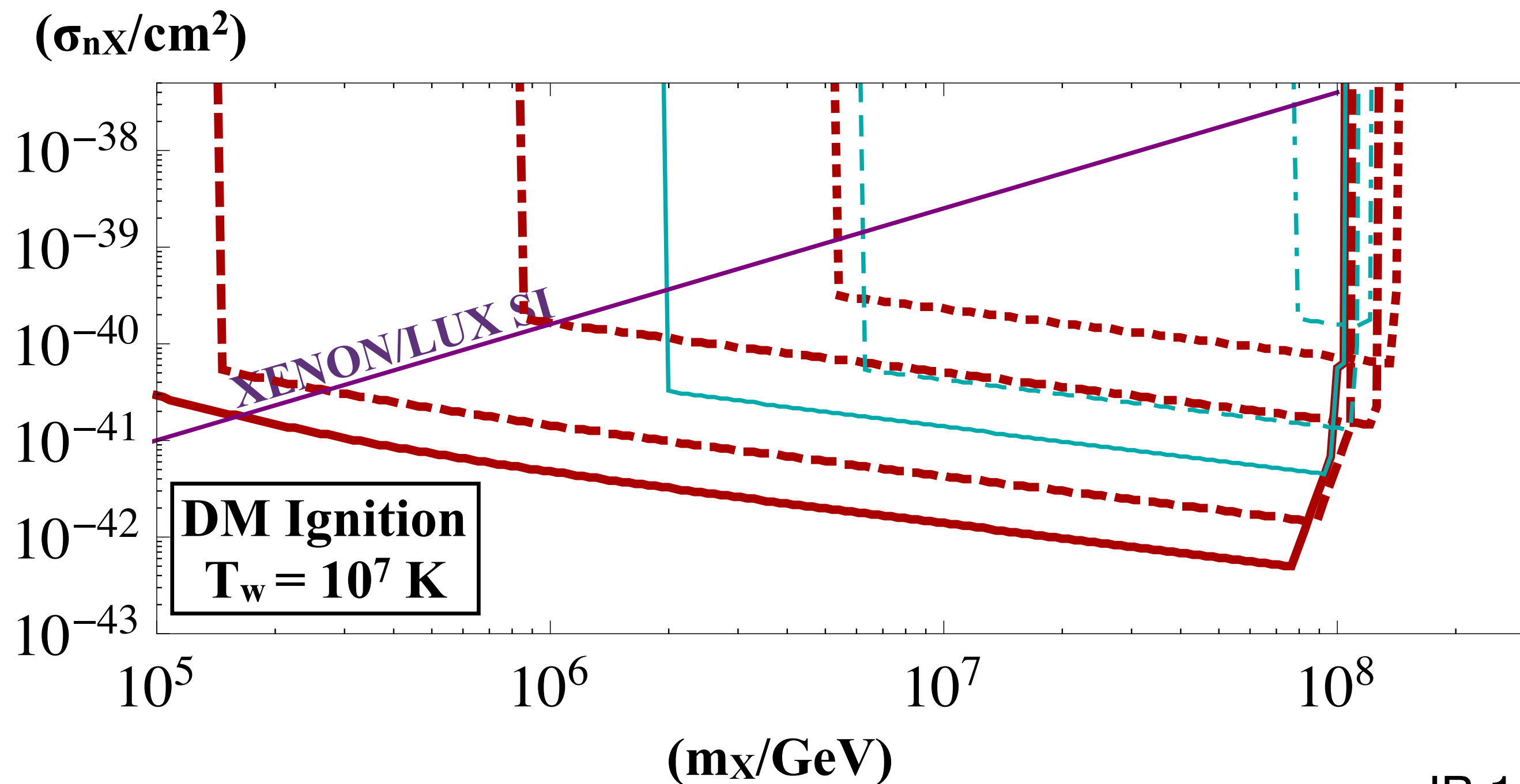


Timmes & Woosley 1992

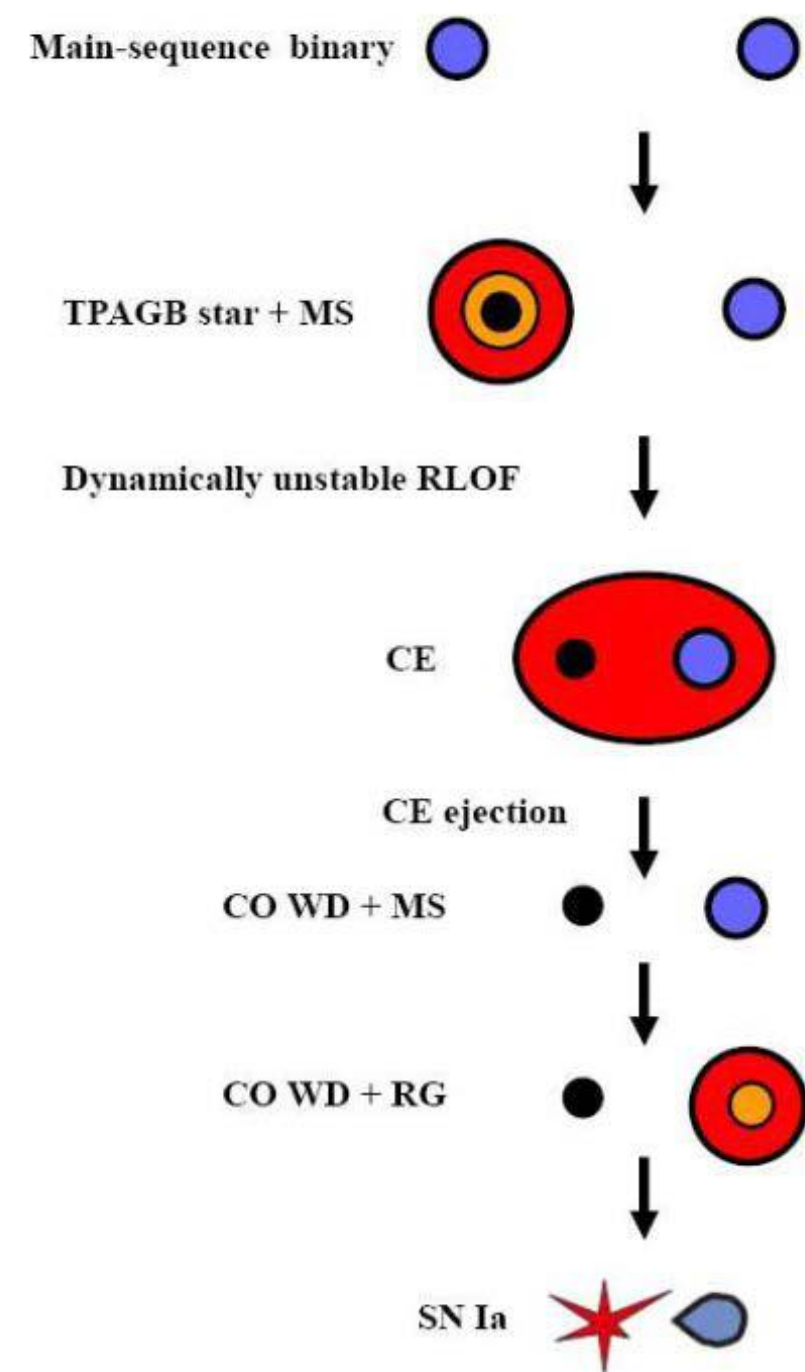
For ignition to occur, at a given temperature, the speed of nuclear burning across a fixed region must exceed the white dwarf's electron conduction diffusion rate.

The limiting factor for heavy DM-induced-ignition of white dwarfs is the mass of the DM particle. Heavier DM collapses after fewer particles have collected, and fewer collapsing particles transfer less heat.

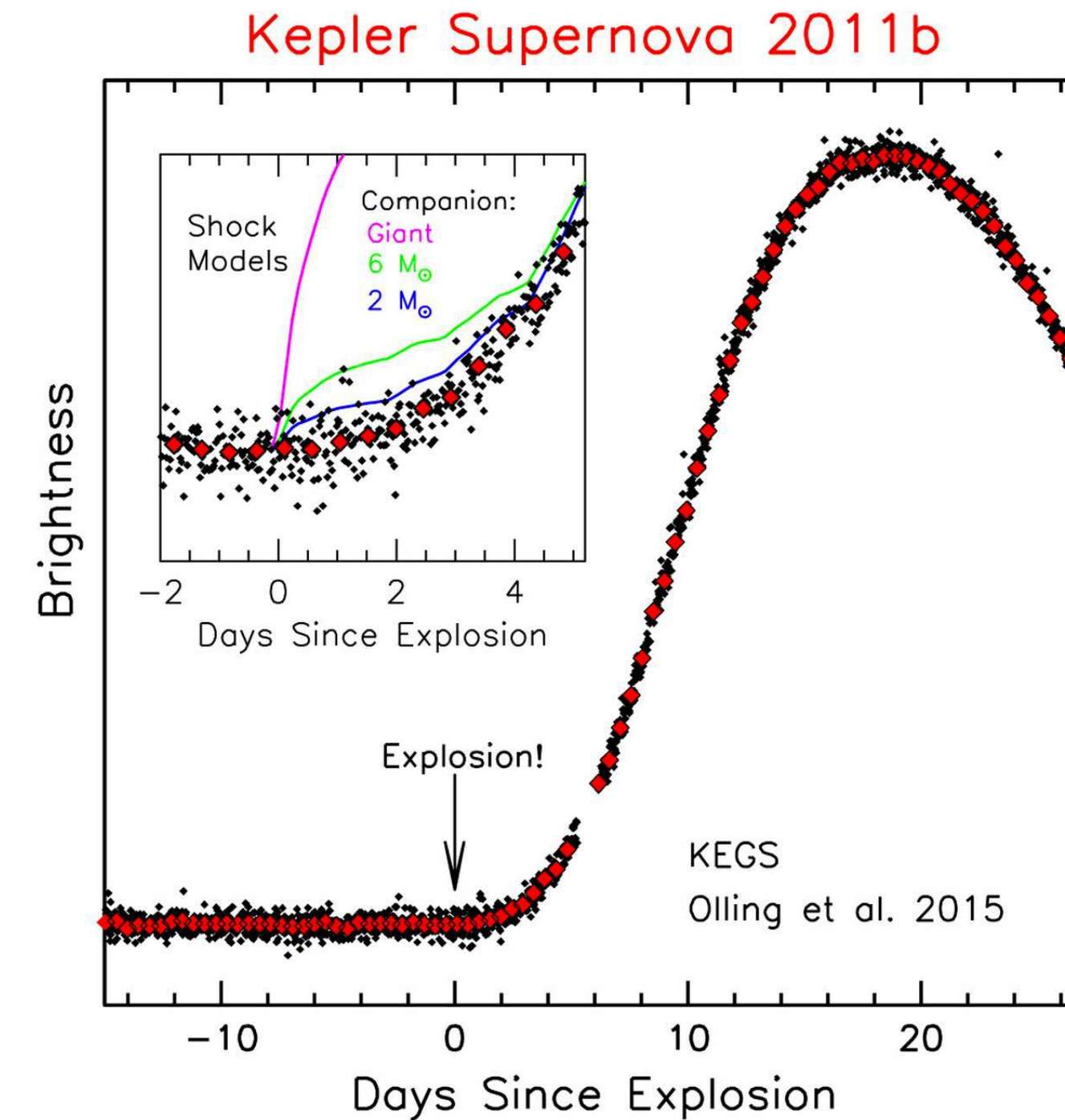
<u>Line</u>	<u>M_w/M_\odot</u>	<u>$R_w/(10^3 \text{ km})$</u>	<u>$\rho_w/(10^7 \text{ g/cm}^3)$</u>	<u>Line</u>	<u>t_w/Gyr</u>
	1.4	2.5	100		5
	1.3	3	40		5
	1.1	5	6		0.5
	0.9	6	2		0.5



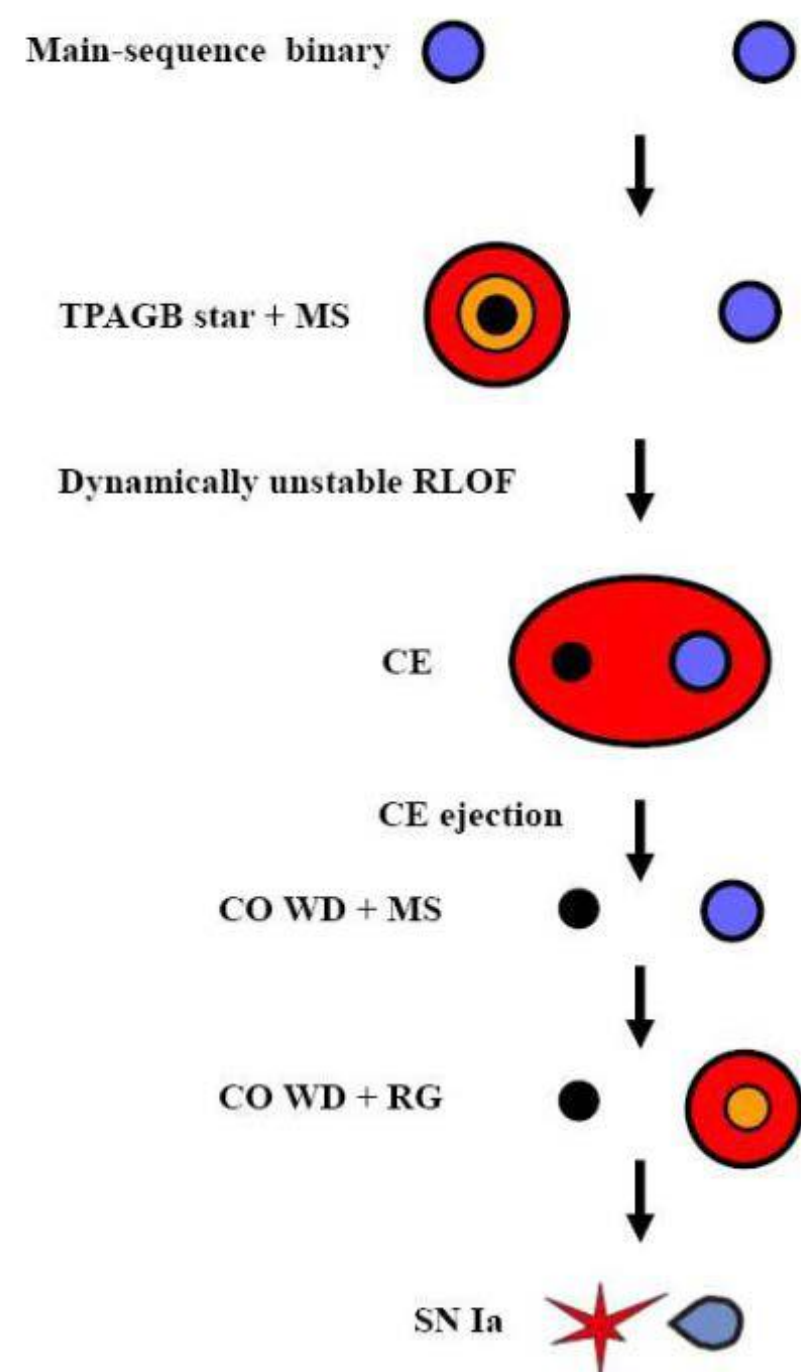
It is not clear what ignites type Ia supernovae. Candidates include binary accretion to criticality (a.k.a. Chandrasekhar mass), and white dwarfs merging.



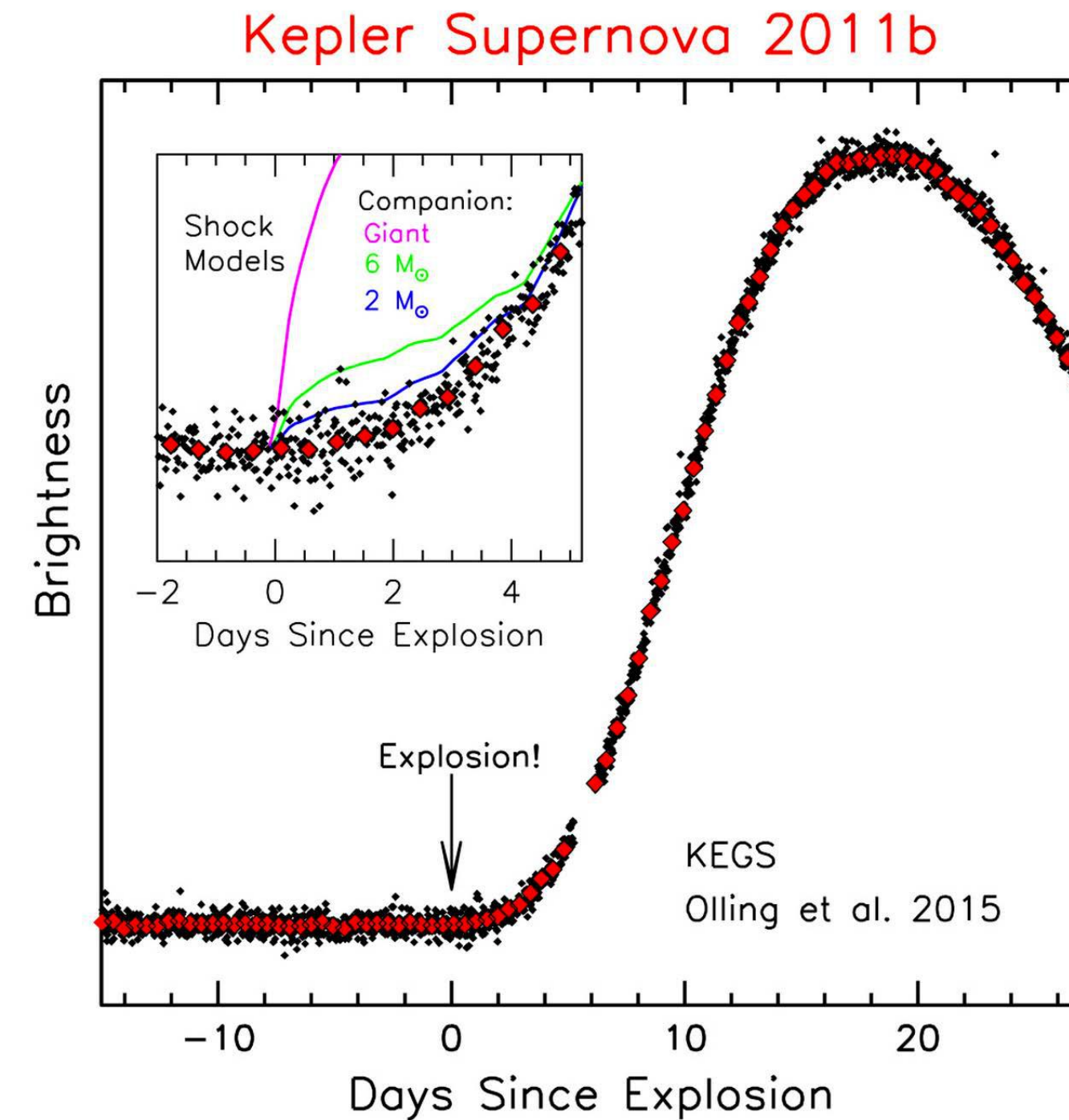
implied shocks
not seen



It is not clear what ignites type Ia supernovae. Candidates include binary accretion to criticality (a.k.a. Chandrasekhar mass), and white dwarfs merging.



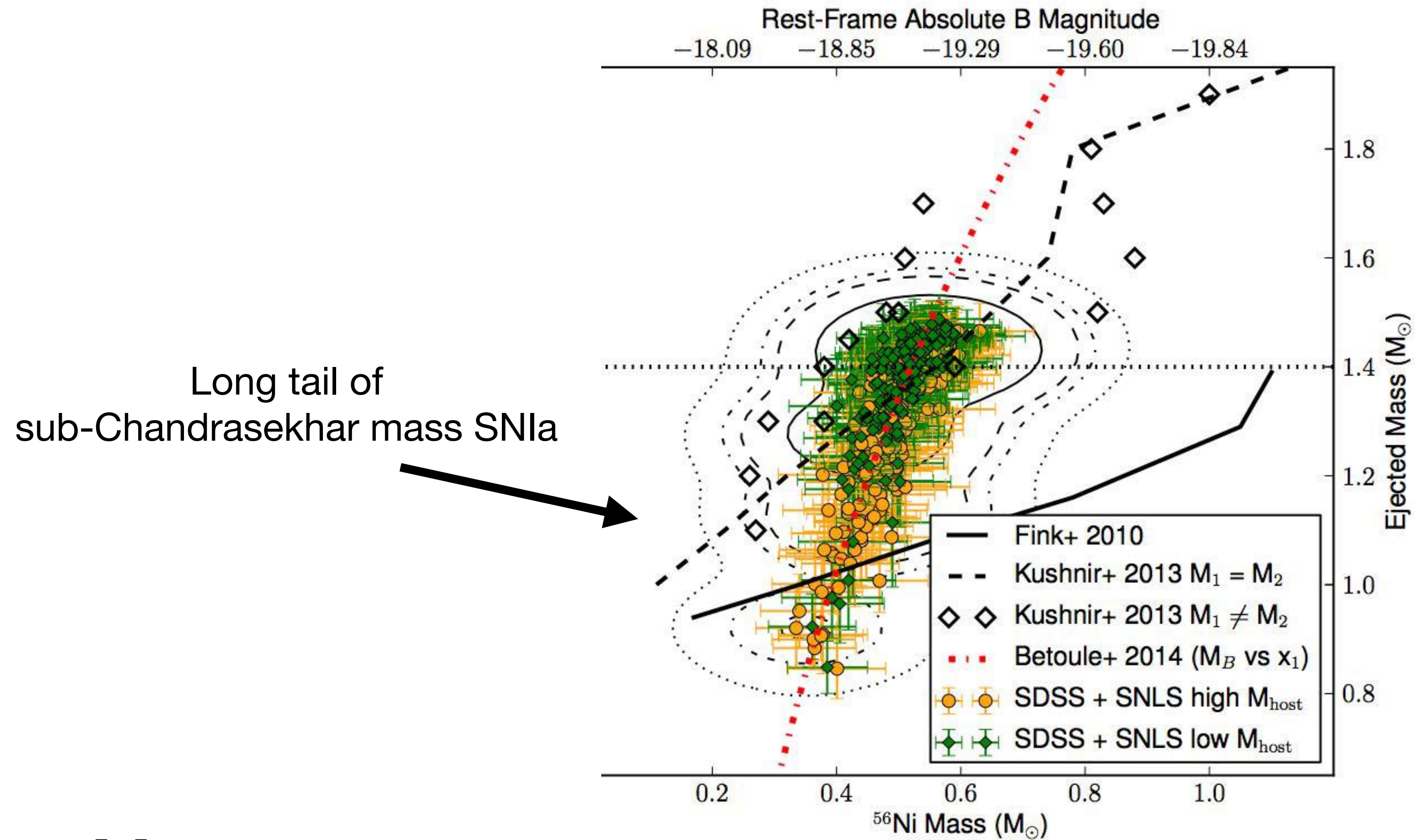
implied shocks
not seen



Binary ignition is now disfavored, by a lack of companion star "shocks" in SNIa light curves, along with the non-observation of H or He lines in any type Ia spectra. White dwarf mergers may not match the high rate of type Ia supernovae.

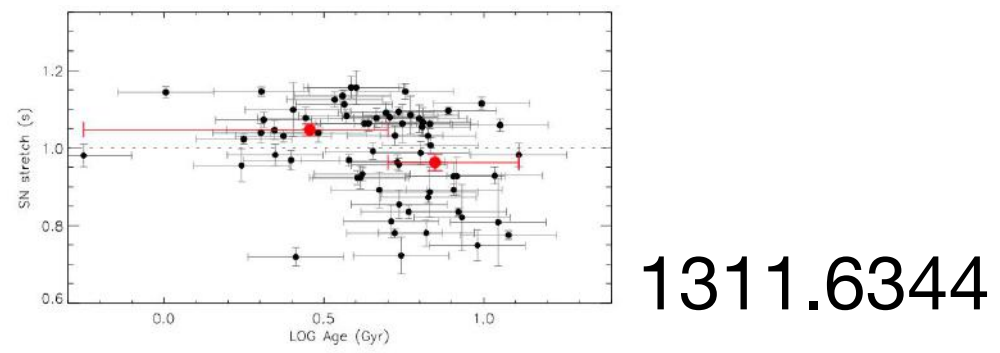
Maoz et al. 1312.0628 (Review), Olling et al. Nature 2015

The existence of sub-Chandrasekhar supernovae presents a quandary that binary accretion has trouble accounting for.

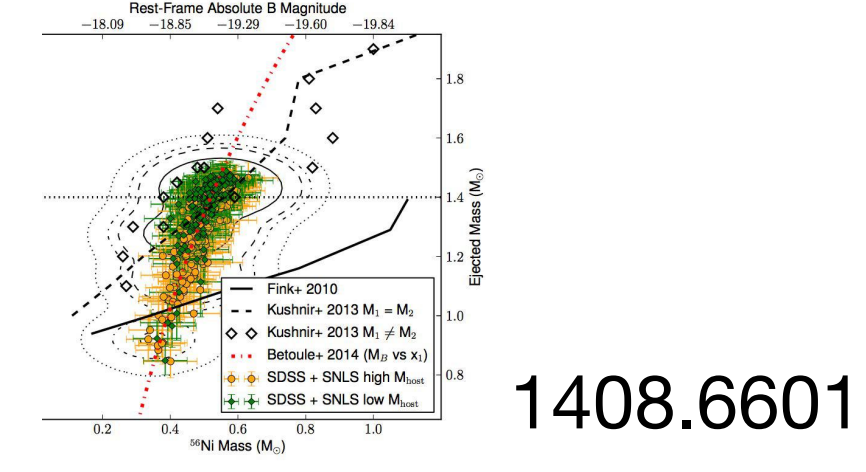


1408.6601

$$\frac{M_{\text{ej}}}{M_{\odot}} = (1.322 \pm 0.022) + (0.185 \pm 0.018)x_1, \quad x_1 = [-3, 2]$$



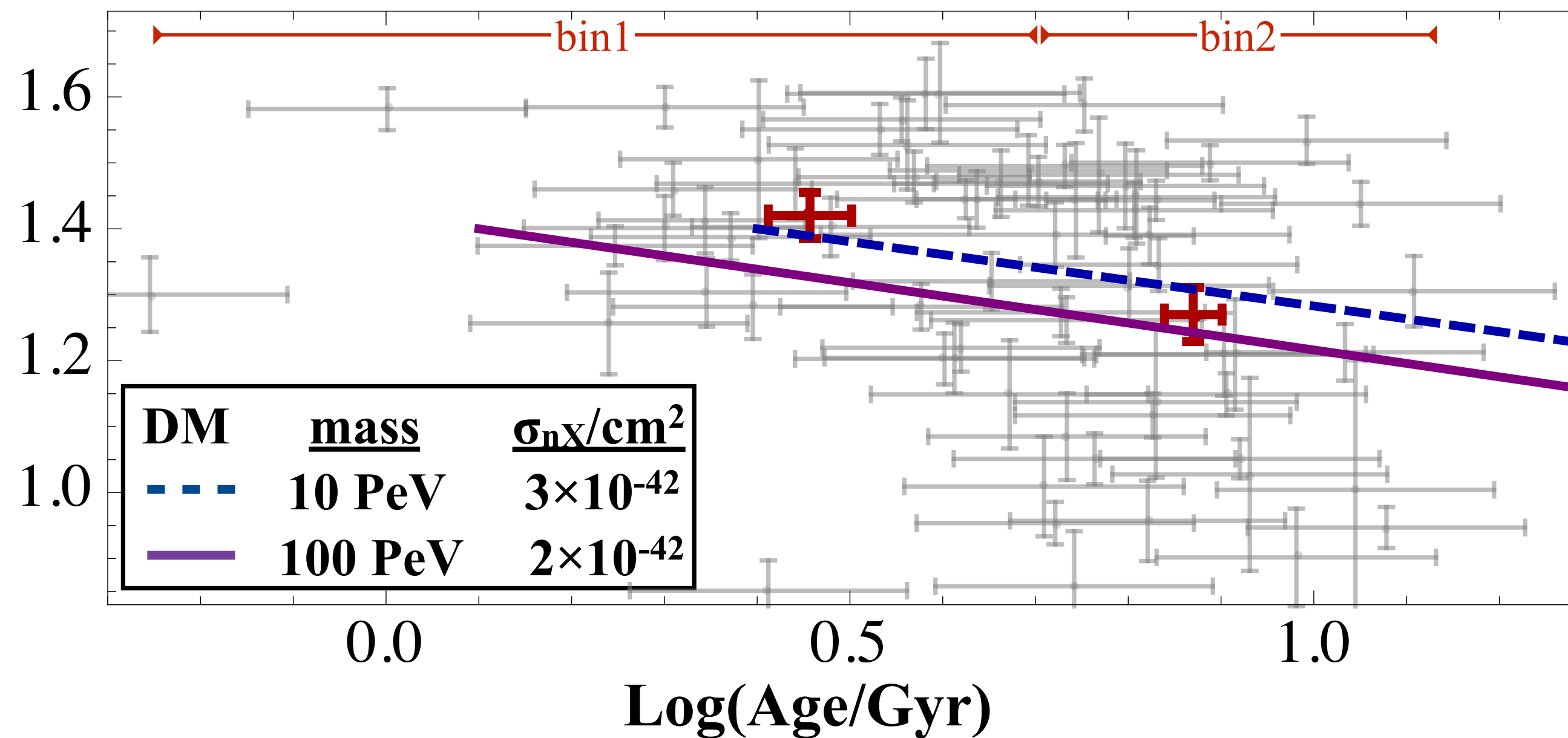
Data on the ages of stars adjacent to type Ia supernovae



Data on the mass of type Ias inferred from luminosity

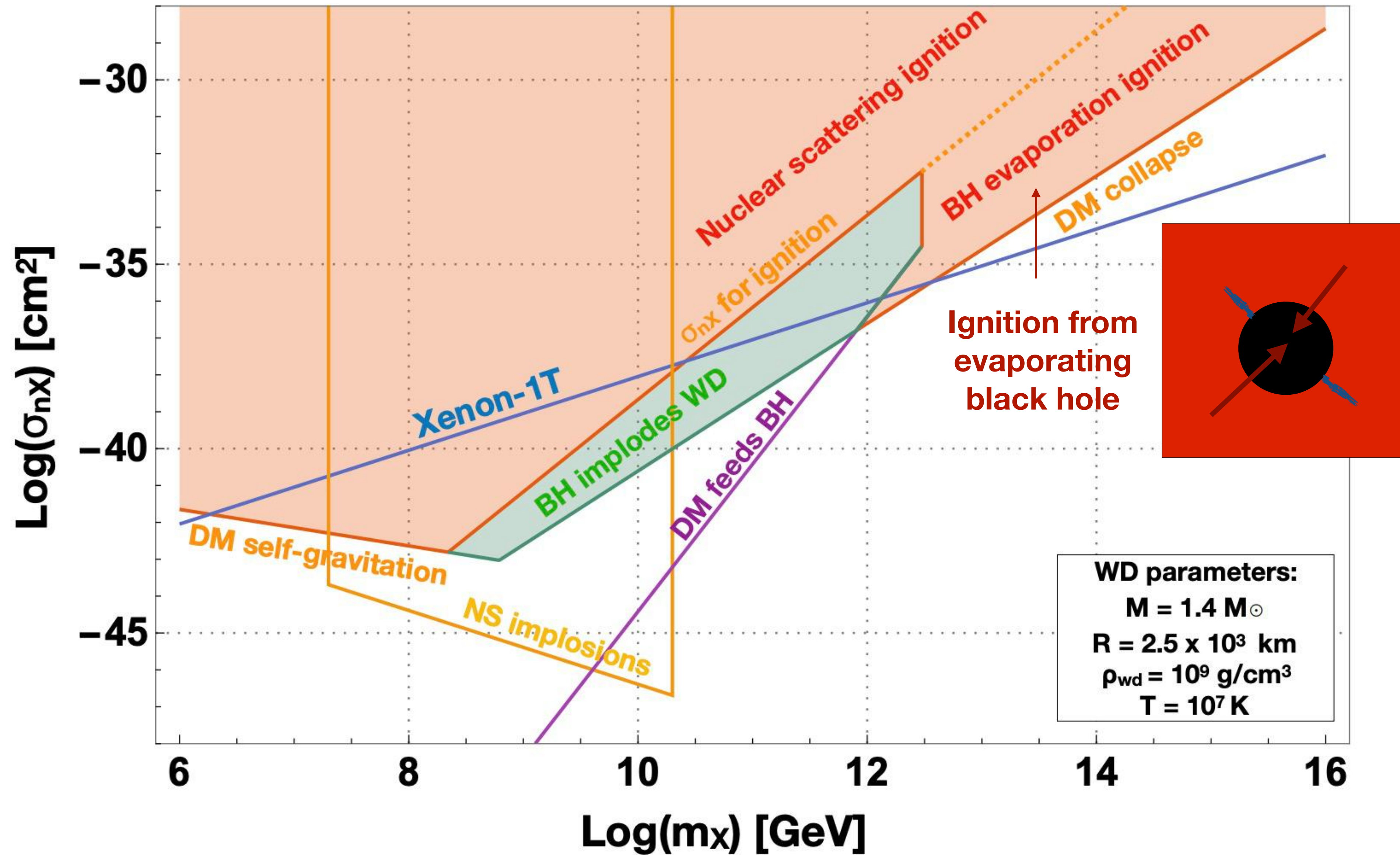
SNIa progenitor age versus mass

$(M_{\text{SN1a}} / M_{\odot})$



Interesting trend — more massive WDs explode sooner.

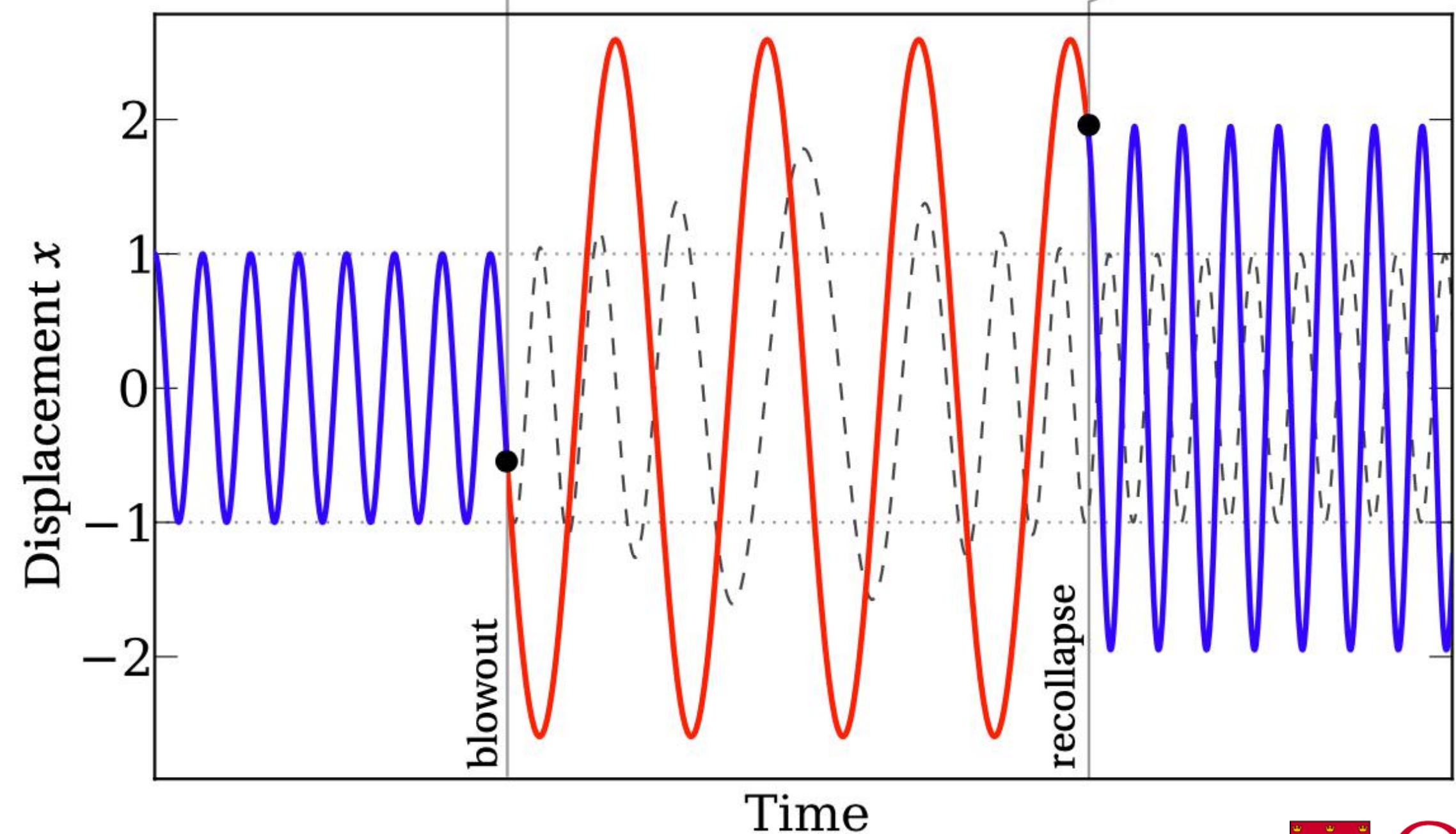
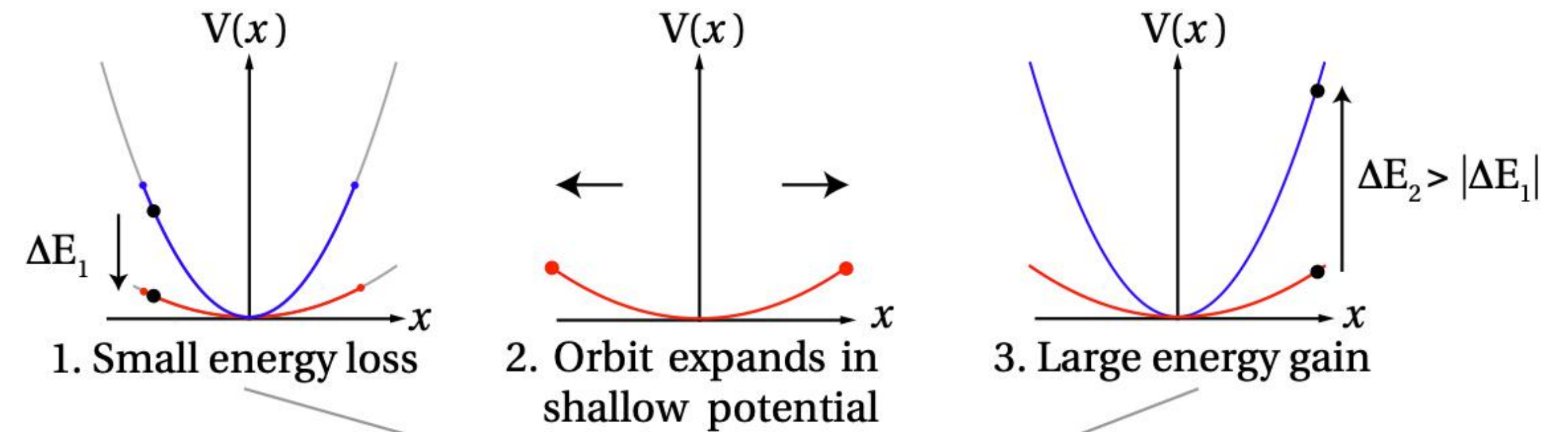
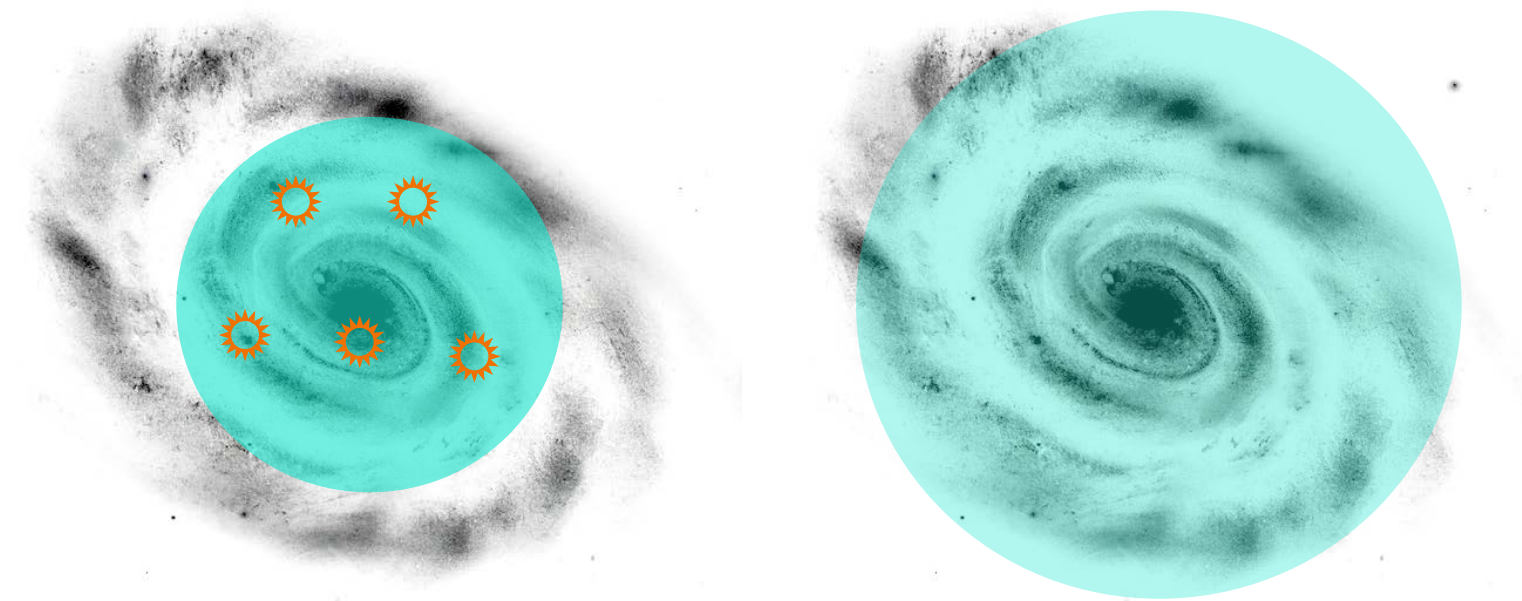
Bounds on dark matter from an old GAIA White Dwarf



Javier Acevedo

BARYONIC FEEDBACK CREATING HALO PROFILE CORES IN GALAXIES

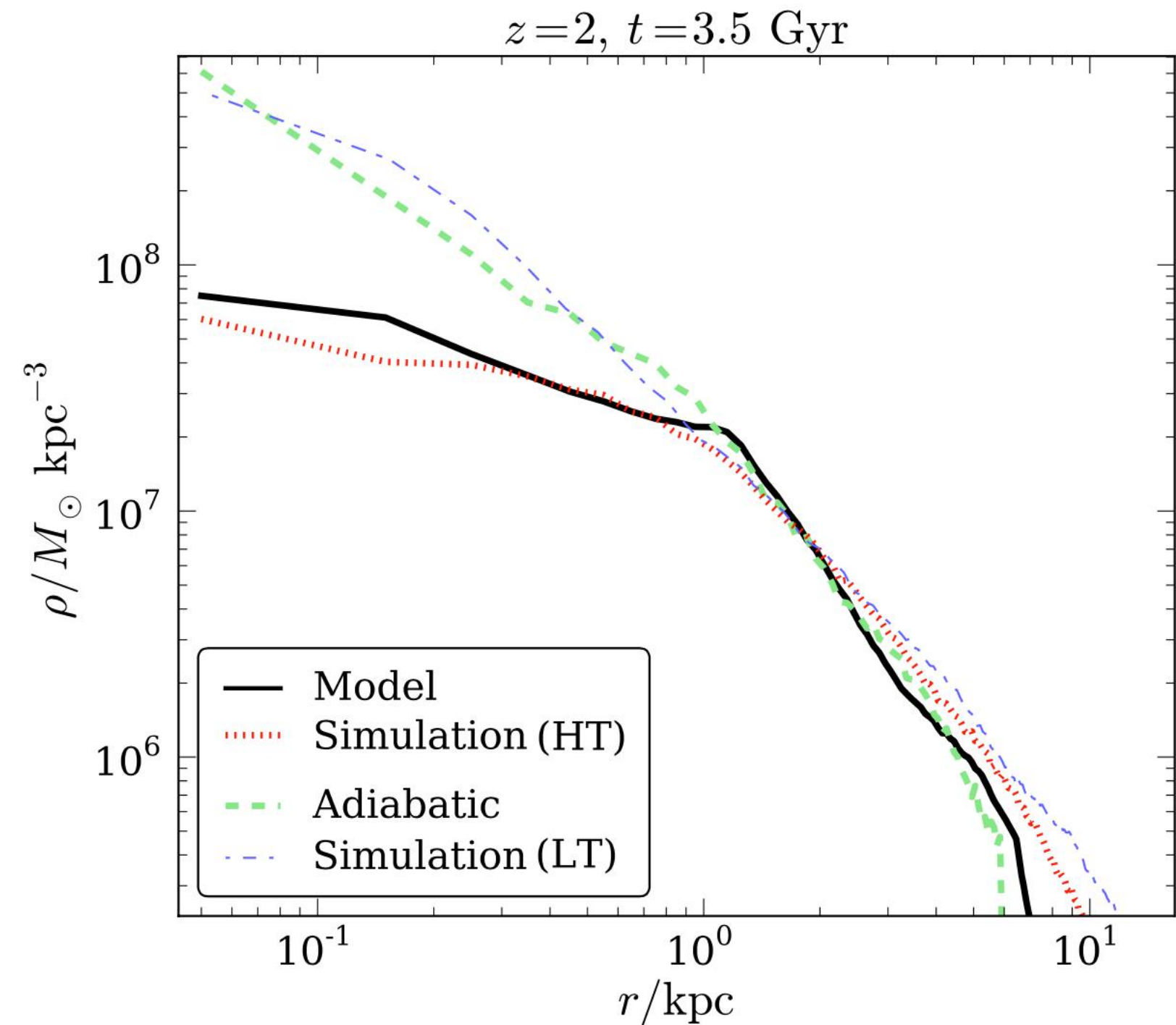
Pontzen and Governato 1106.0499



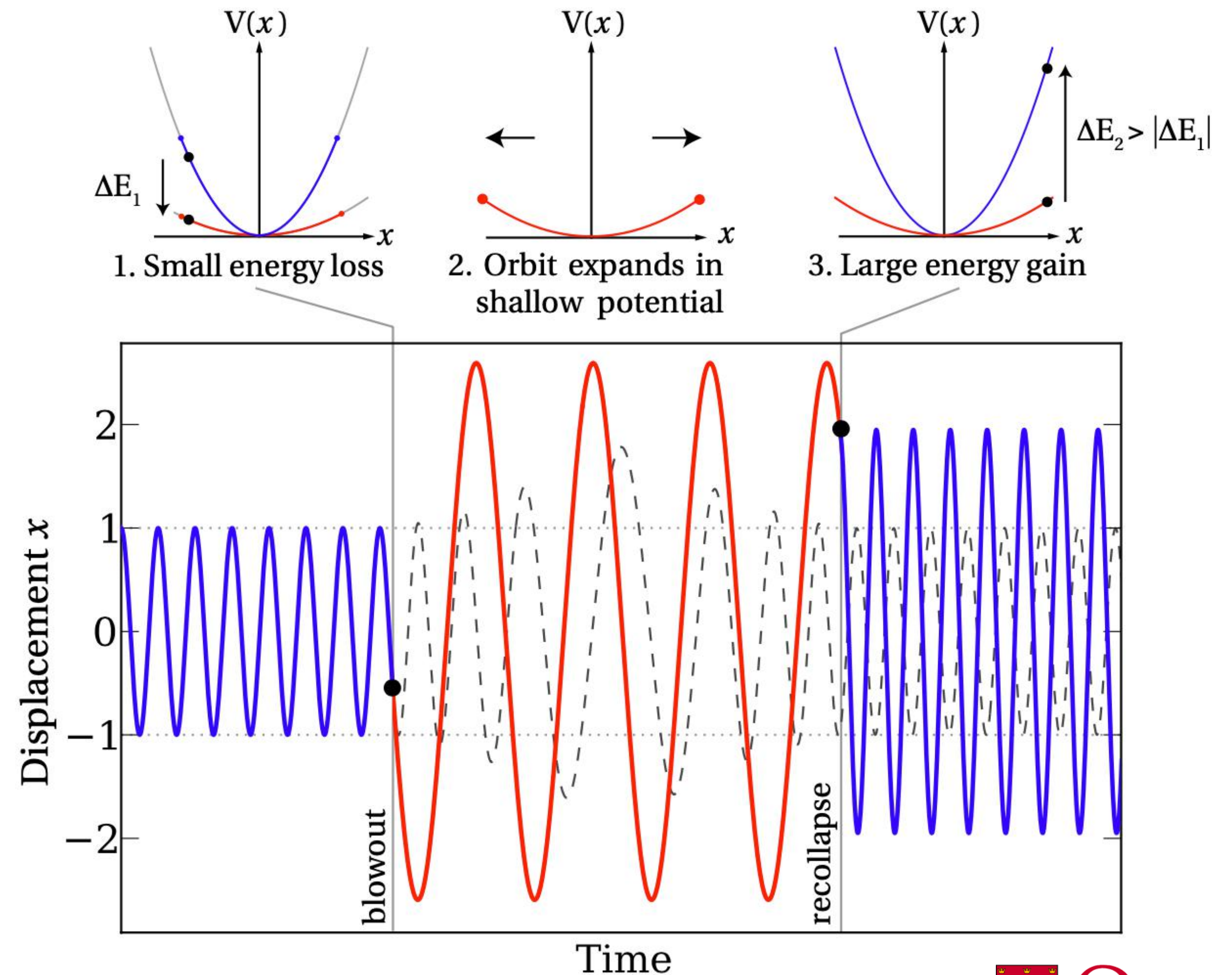
- Supernovae in galaxies blowout gas, non-adiabatic shift in gas potential
- Particles at the end of orbital phase end up on wider orbits, leading to a cored dark matter halo profile

BARYONIC FEEDBACK CREATING HALO PROFILE CORES IN GALAXIES

Pontzen and Governato 1106.0499



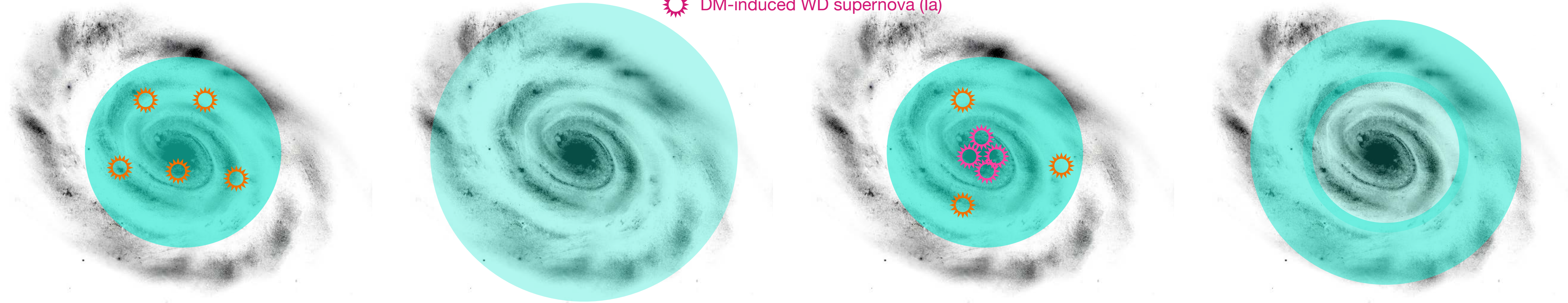
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DARK MATTER INDUCED BARYONIC FEEDBACK

An, Acevedo, Boukhtouchen, JB, Richardson, Sansom

 core collapse supernova
 DM-induced WD supernova (Ia)

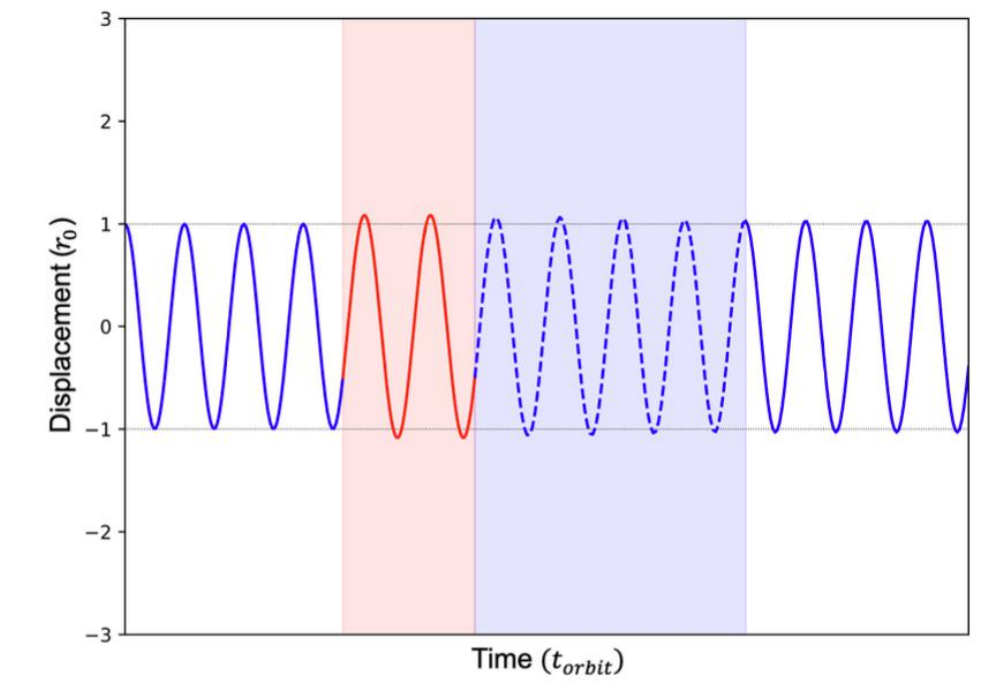
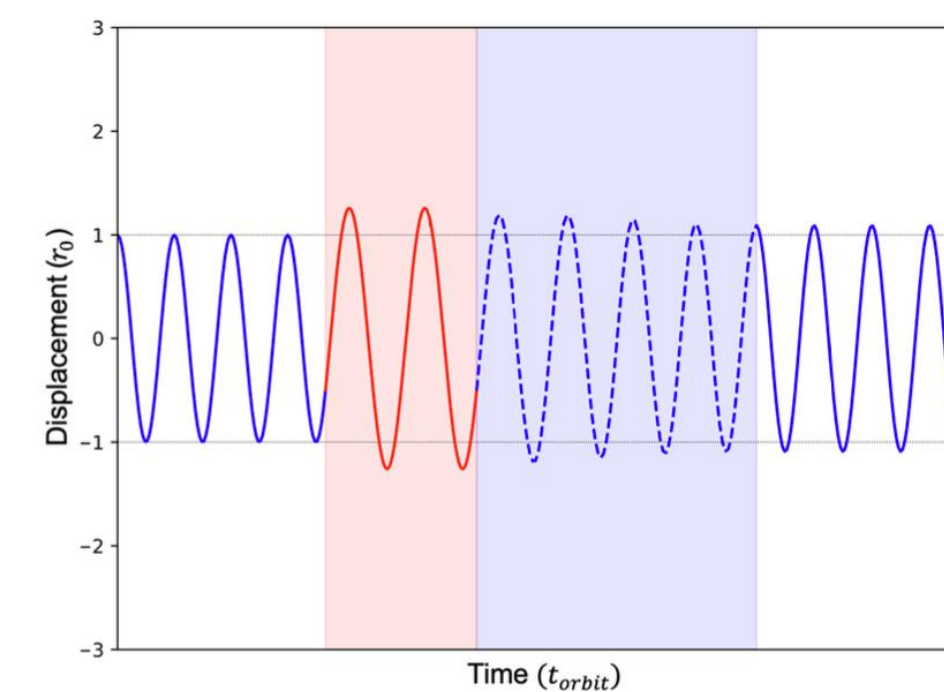
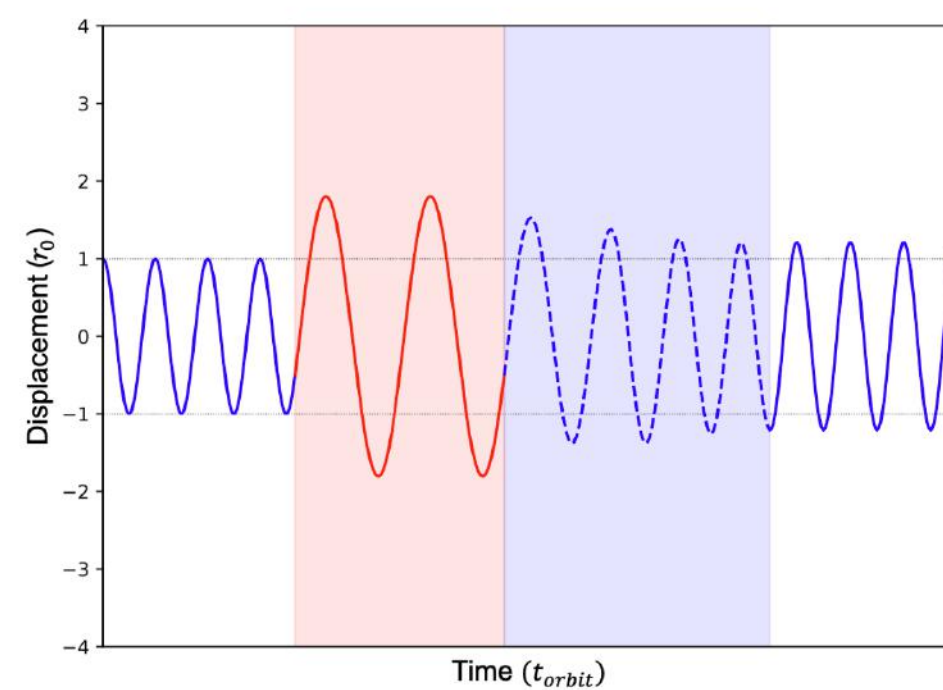
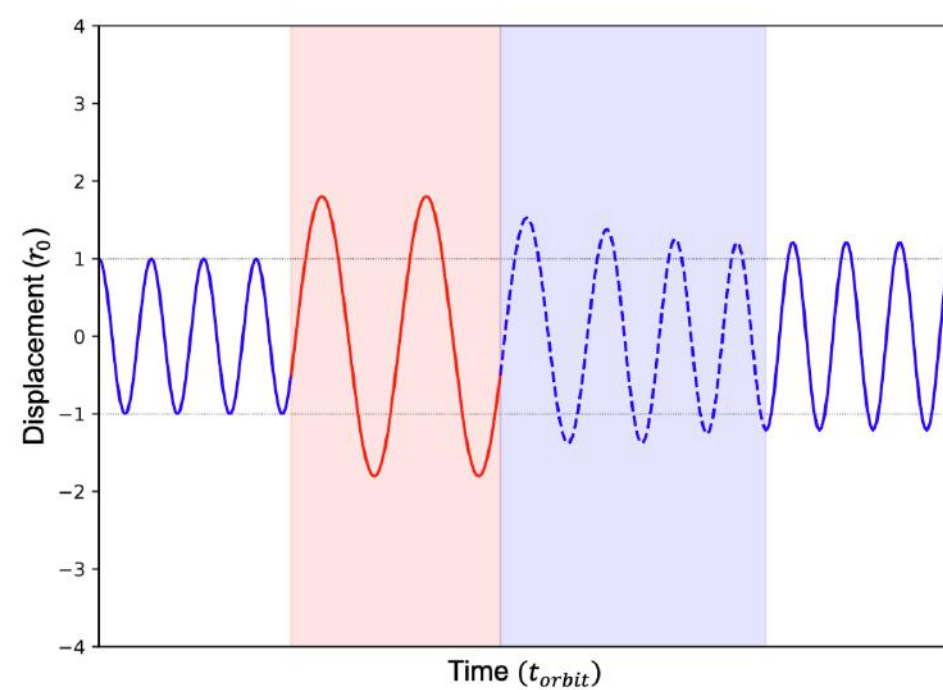
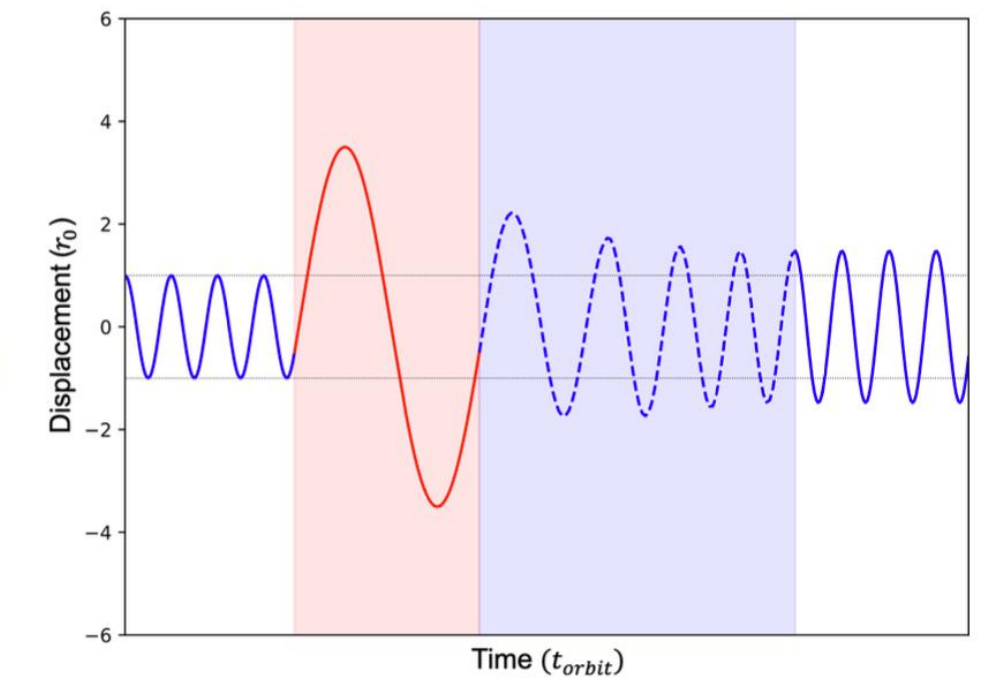
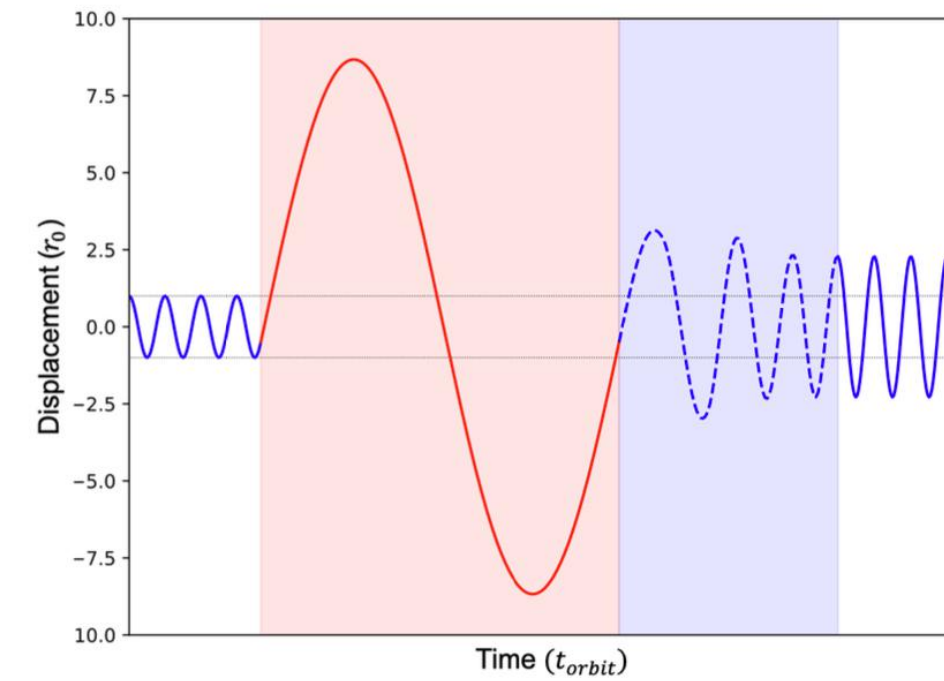
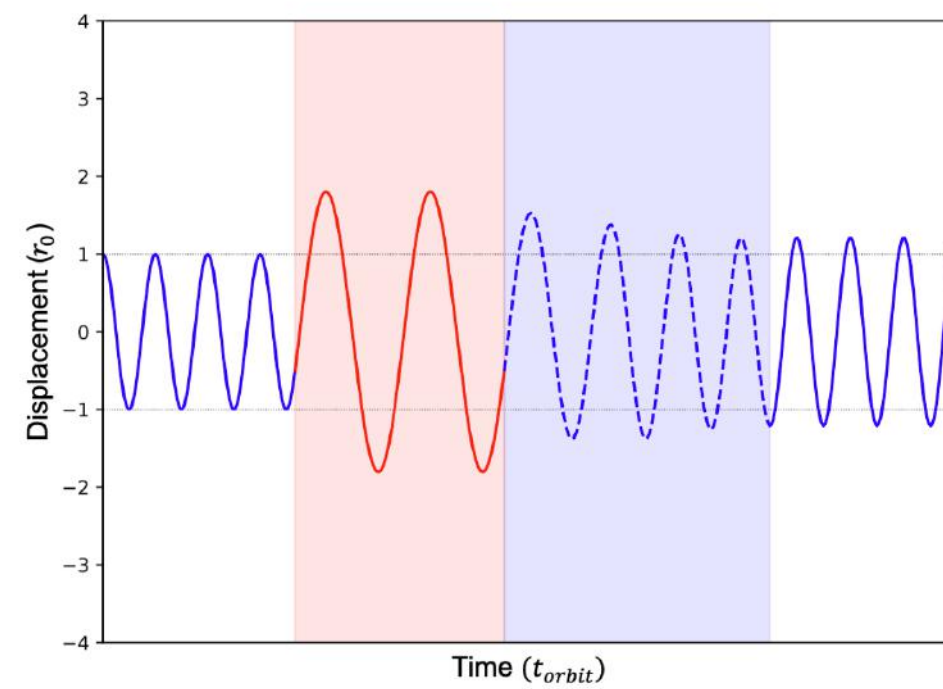
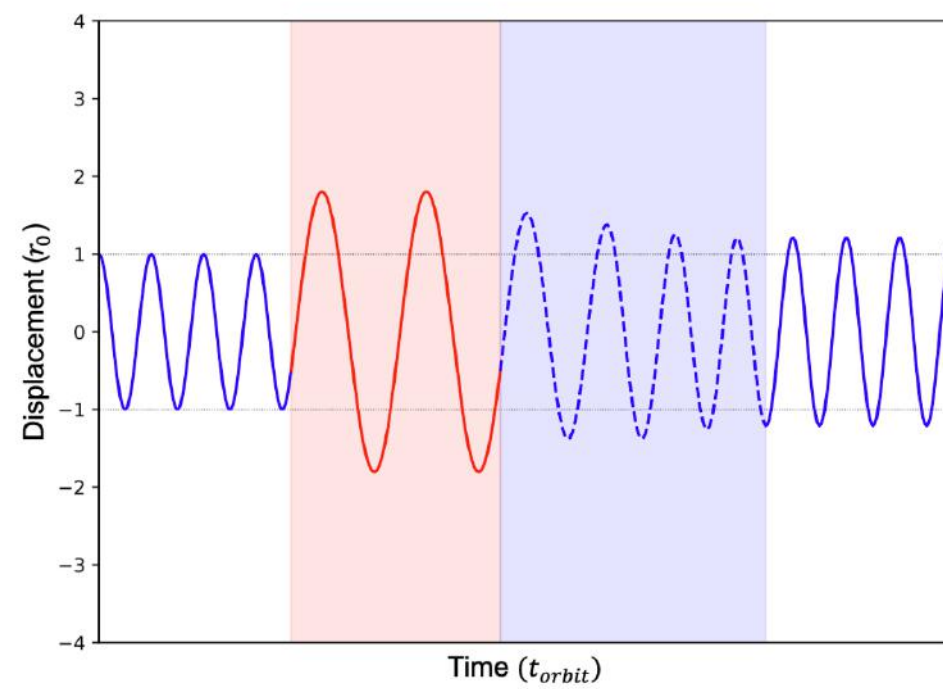


Yilda Boukhtouchen

DARK MATTER INDUCED BARYONIC FEEDBACK

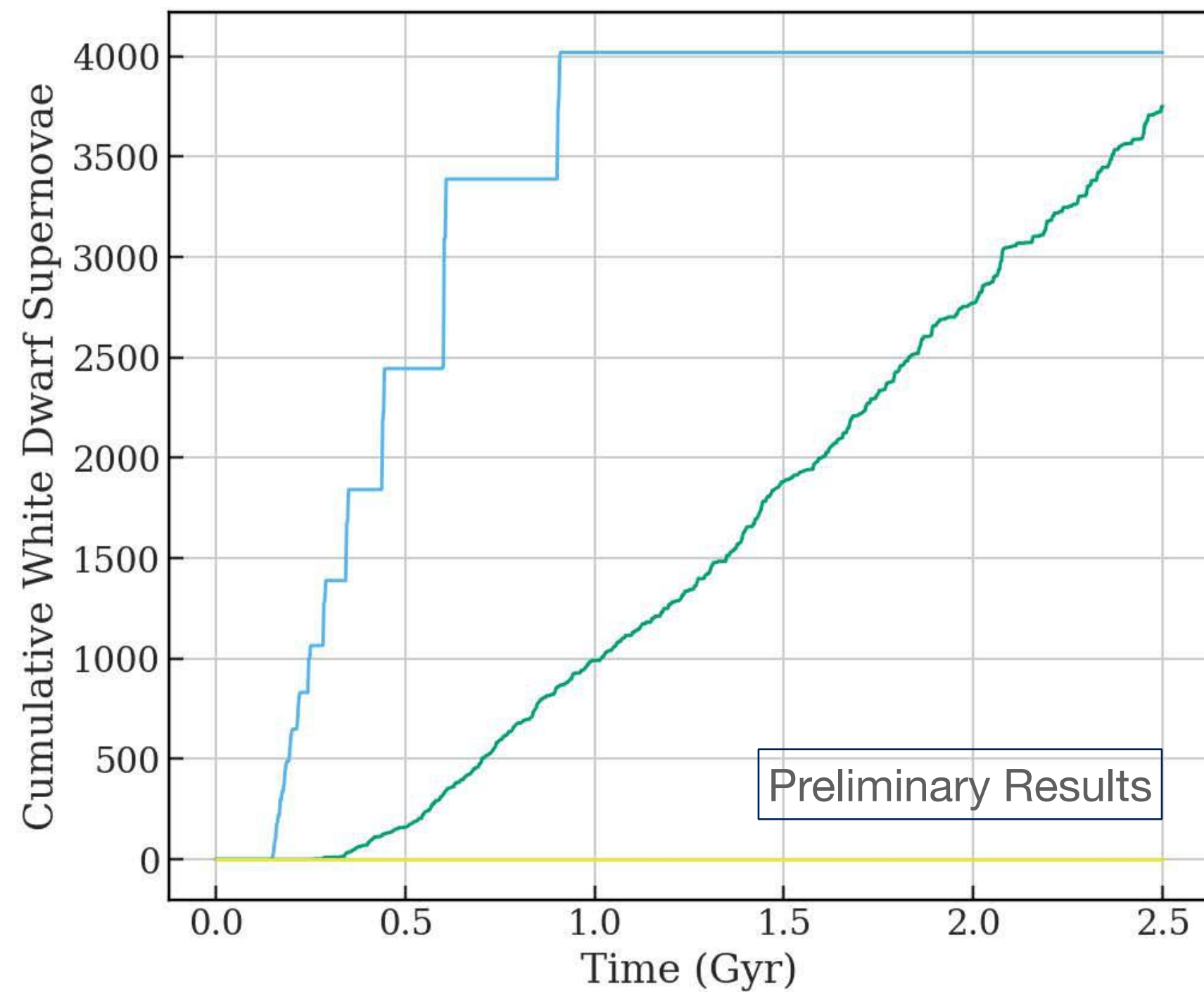
An, Acevedo, Boukhtouchen, JB, Richardson, Sansom

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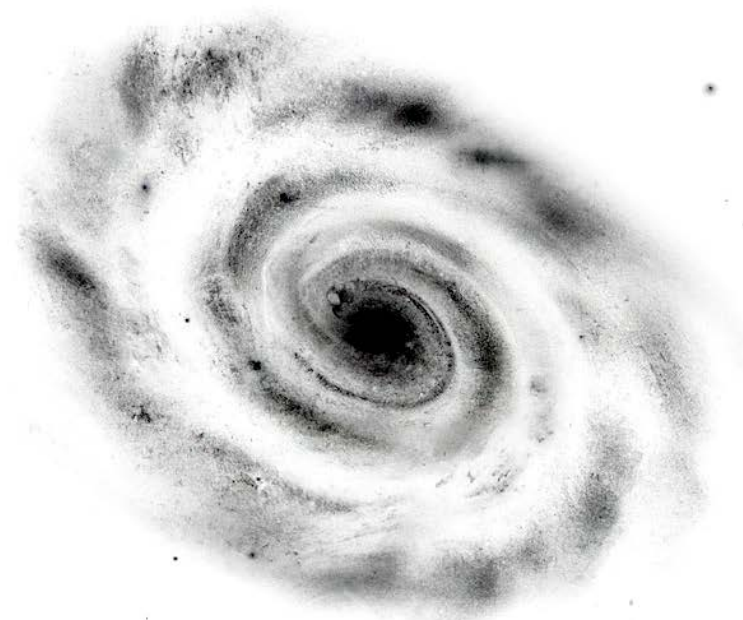
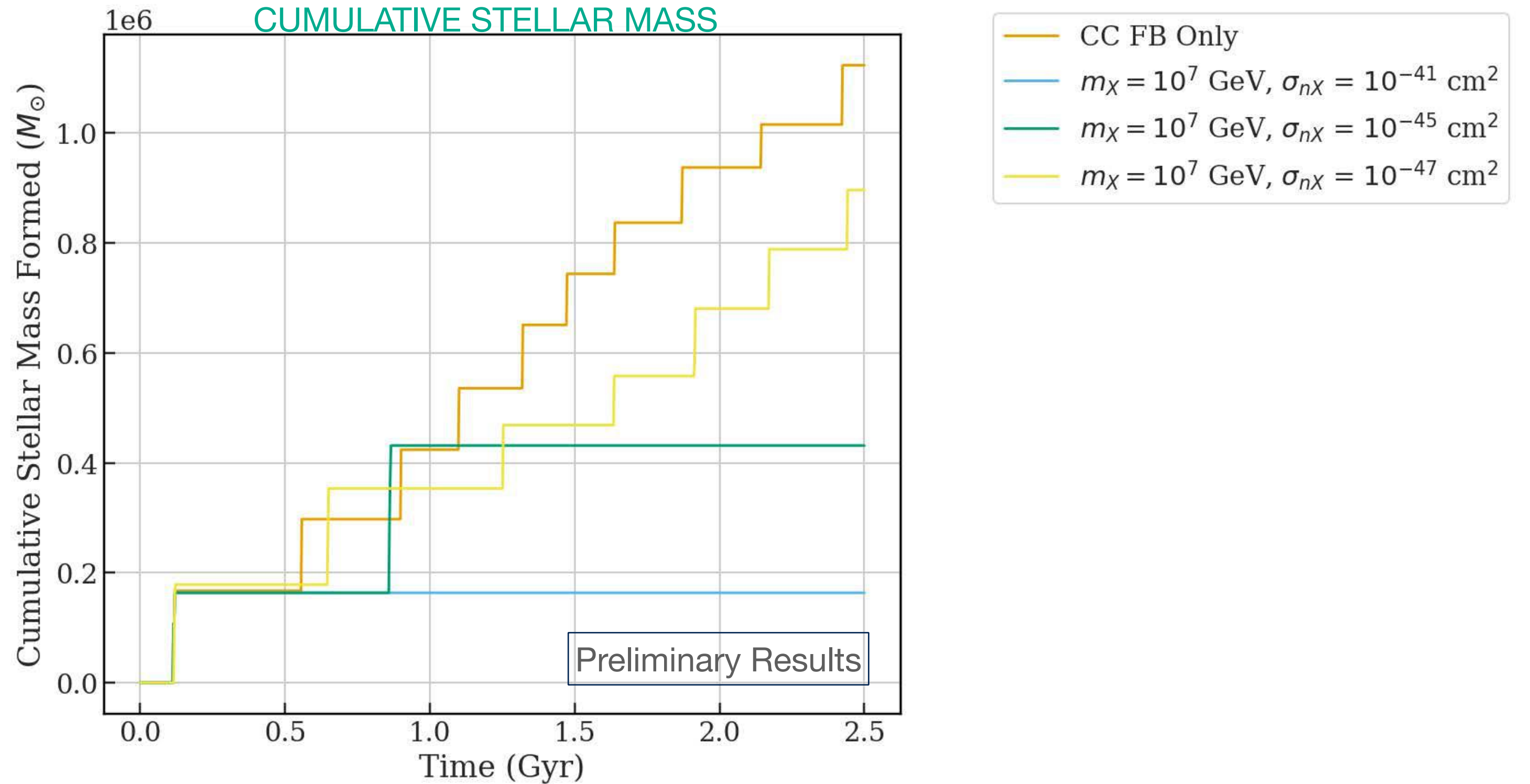


DARK MATTER INDUCED BARYONIC FEEDBACK – SMALL GALAXY RESULTS

WHITE DWARF EXPLOSIONS (Type Ia)



CUMULATIVE STELLAR MASS



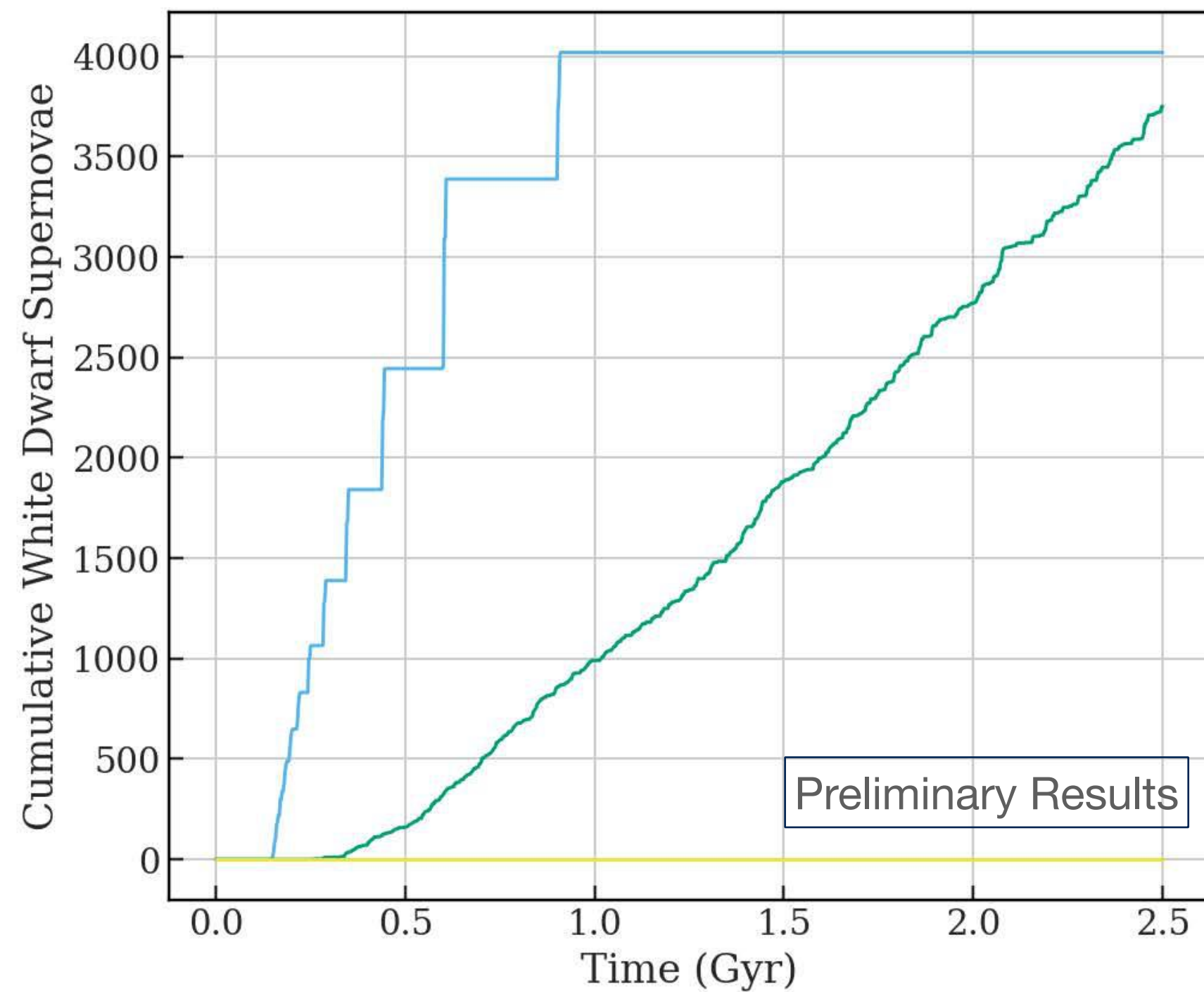
Total Galaxy Mass: $1.86 \cdot 10^9 M_{\text{sun}}$
 DM mass fraction: 0.9615
 Gas mass fraction: 0.0175
 Stellar disk mass fraction: 0.0175
 Stellar bulge mass fraction: 0.0035

all particles have mass of 3721 M_{sun}
 modified GIZMO codebase

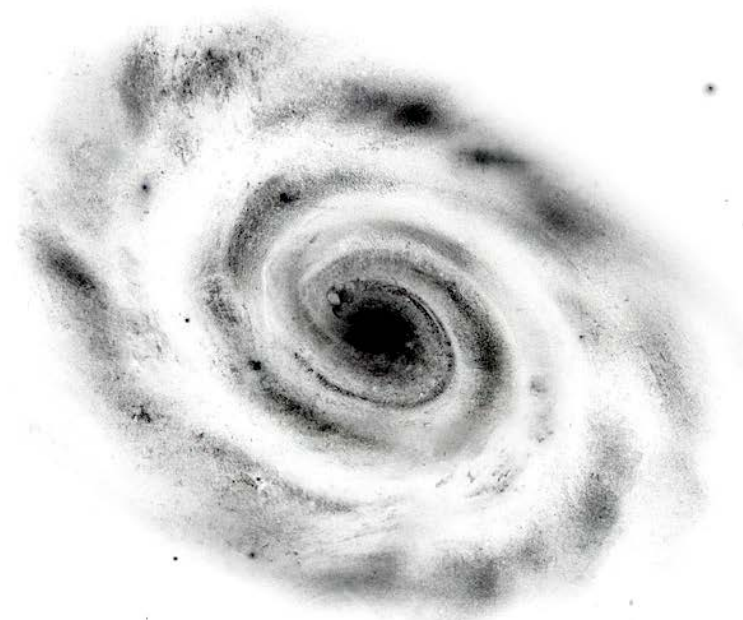
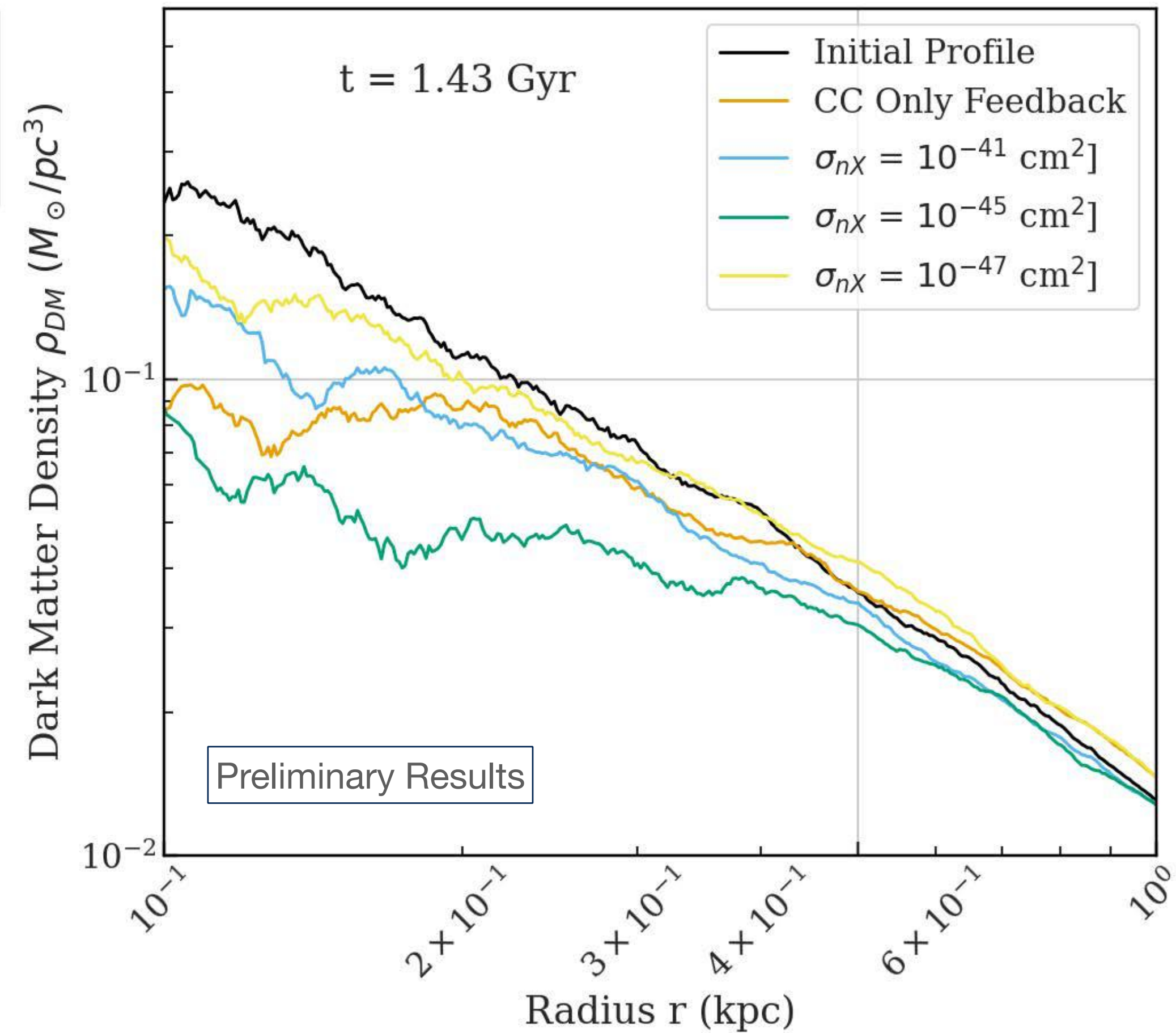
An, Acevedo, Boukhtouchen, JB, Richardson, Sansom

DARK MATTER INDUCED BARYONIC FEEDBACK – SMALL GALAXY RESULTS

WHITE DWARF EXPLOSIONS (Type Ia)



- CC FB Only
- $m_\chi = 10^7$ GeV, $\sigma_{n\chi} = 10^{-41}$ cm²
- $m_\chi = 10^7$ GeV, $\sigma_{n\chi} = 10^{-45}$ cm²
- $m_\chi = 10^7$ GeV, $\sigma_{n\chi} = 10^{-47}$ cm²



Total Galaxy Mass: $1.86 \cdot 10^9$ Msun
 DM mass fraction: 0.9615
 Gas mass fraction: 0.0175
 Stellar disk mass fraction: 0.0175
 Stellar bulge mass fraction: 0.0035

all particles have mass of 3721 M_{sun}
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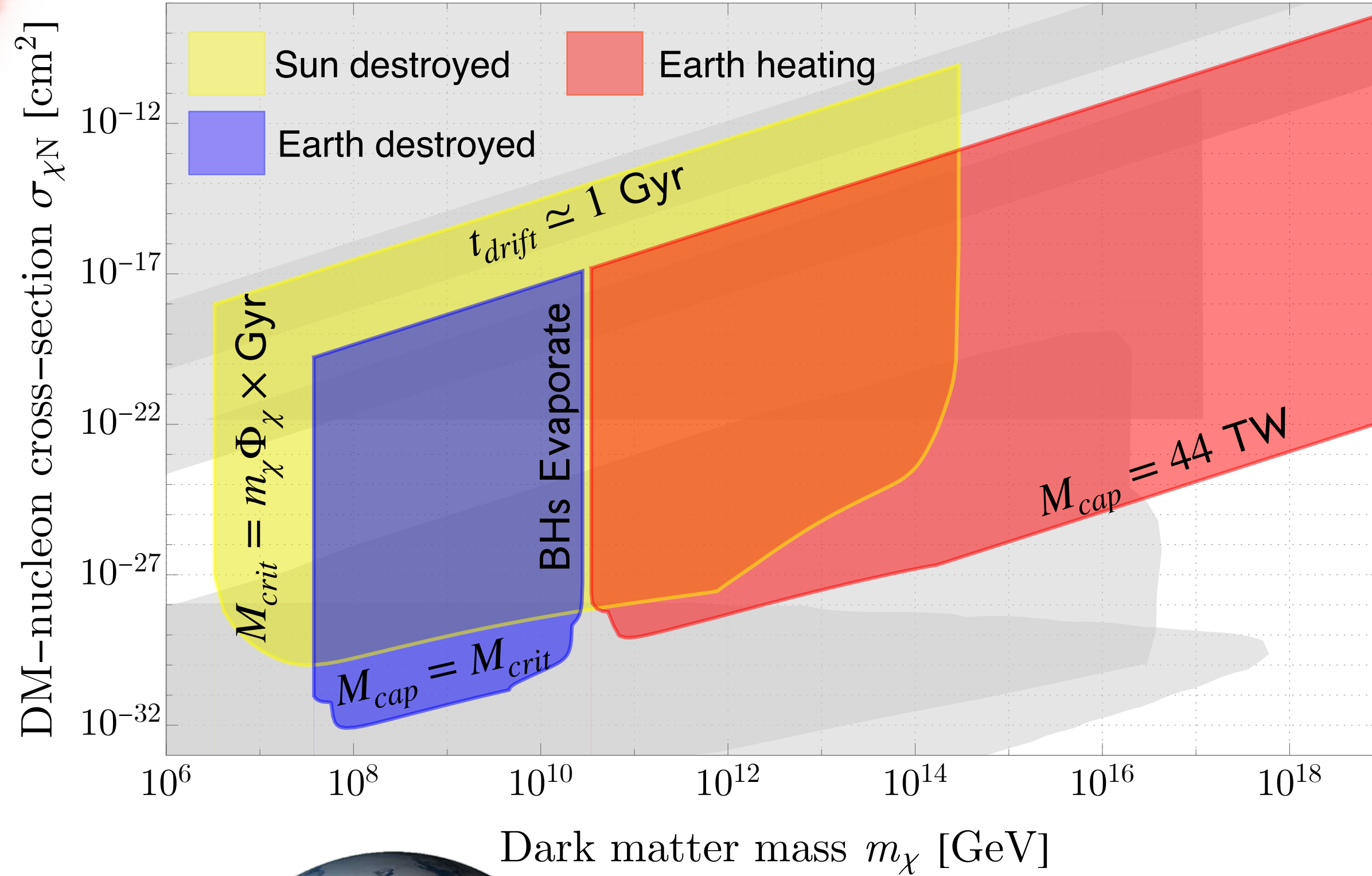
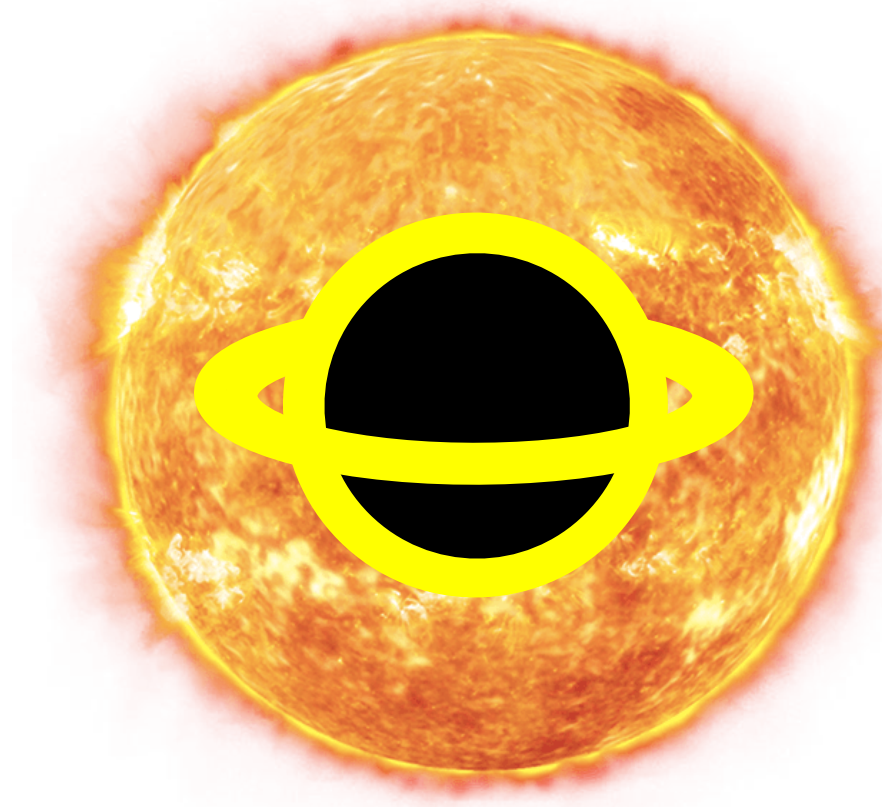
An, Acevedo, Boukhtouchen, JB, Richardson, Sansom

Dark matter in compact stars and galactic structure

- Dark matter heating neutron stars through either infall or annihilation
- Asymmetric dark matter causing neutron stars to implode
 - Connections to missing pulsars (FRBs)
- White dwarf explosions from asymmetric dark matter / composites
 - Connections to type Ia supernovae, galactic structure

Thanks!

Also BHs from DM in the sun, earth

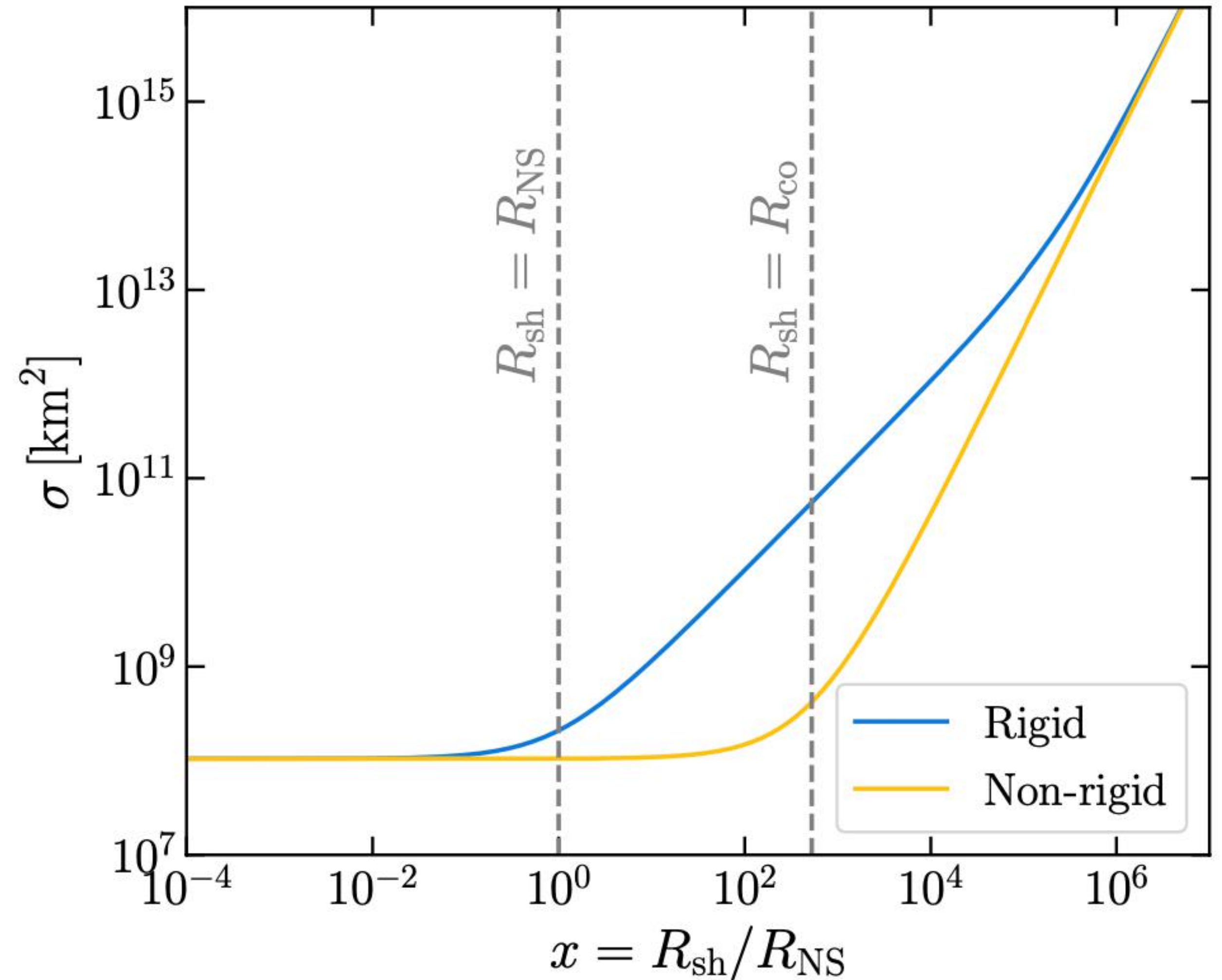
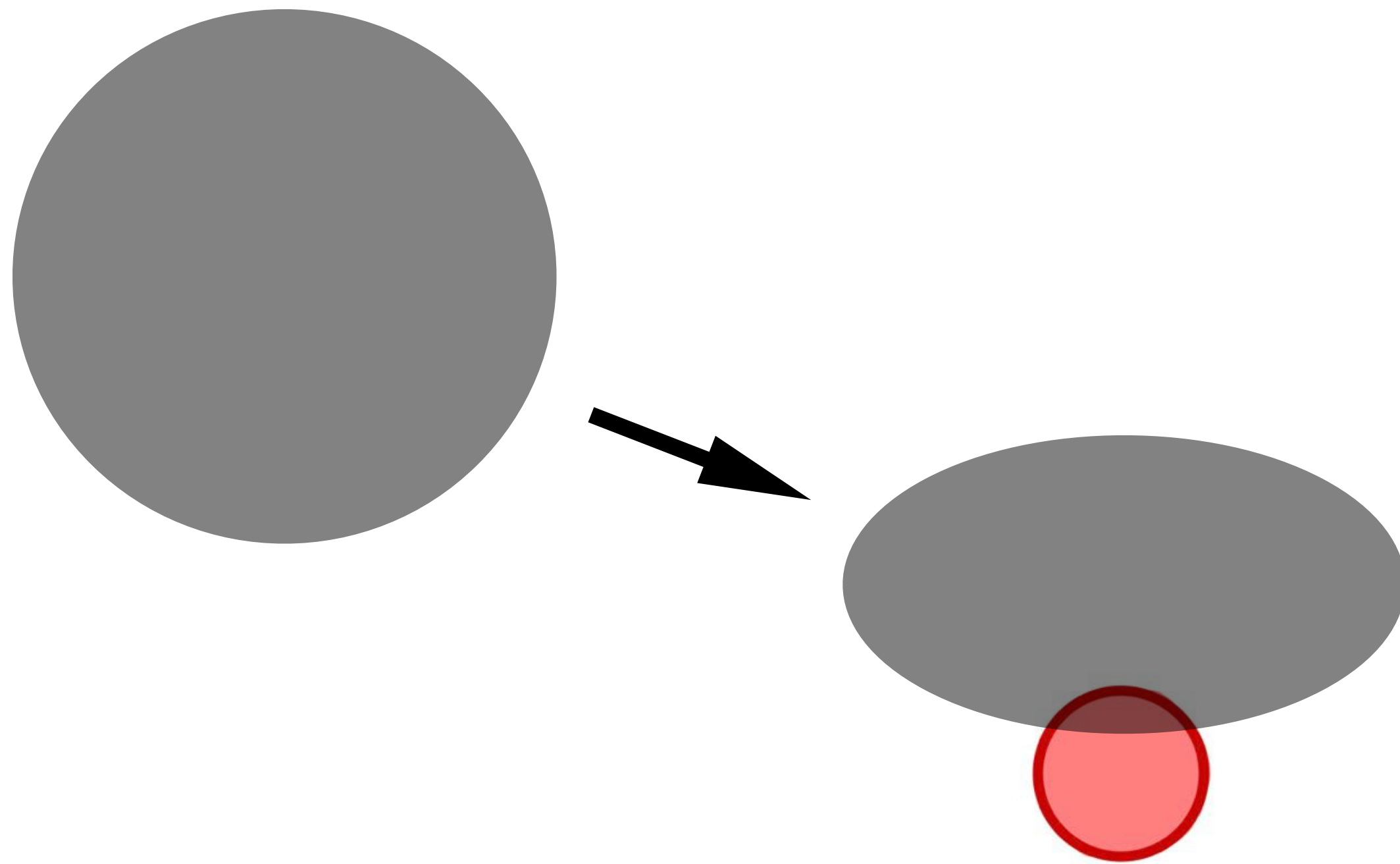


Acevedo, JB, Goodman, Kopp, Opferkuch 2012.09176

Accretion: rigidity

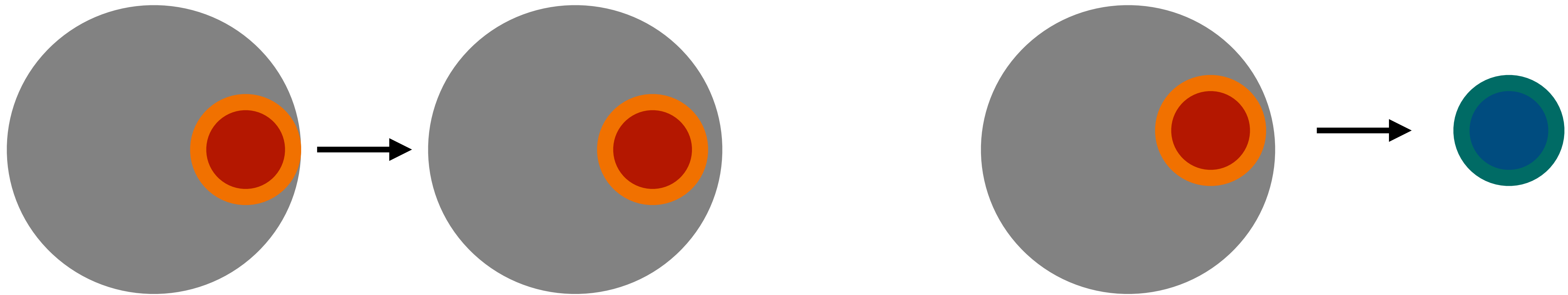
JB, Kavanagh, Raj 2109.04582

- In what follows we assume deformable subhalos



Neutron stars either sporadically flash heated or multi-subhalo steady-state heated

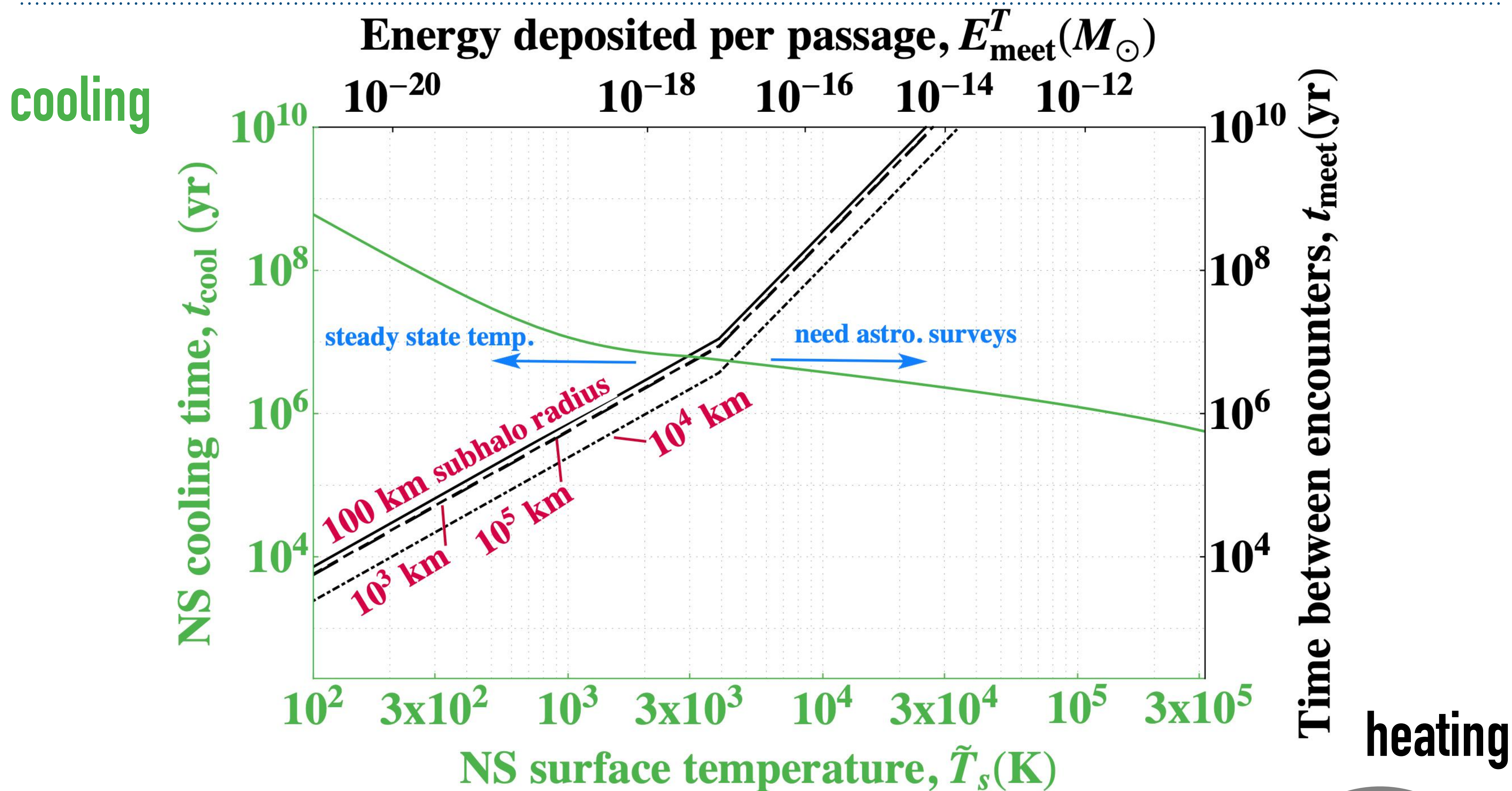
- Neutron stars heated to higher temperatures as they pass by/through dense DM subhalos



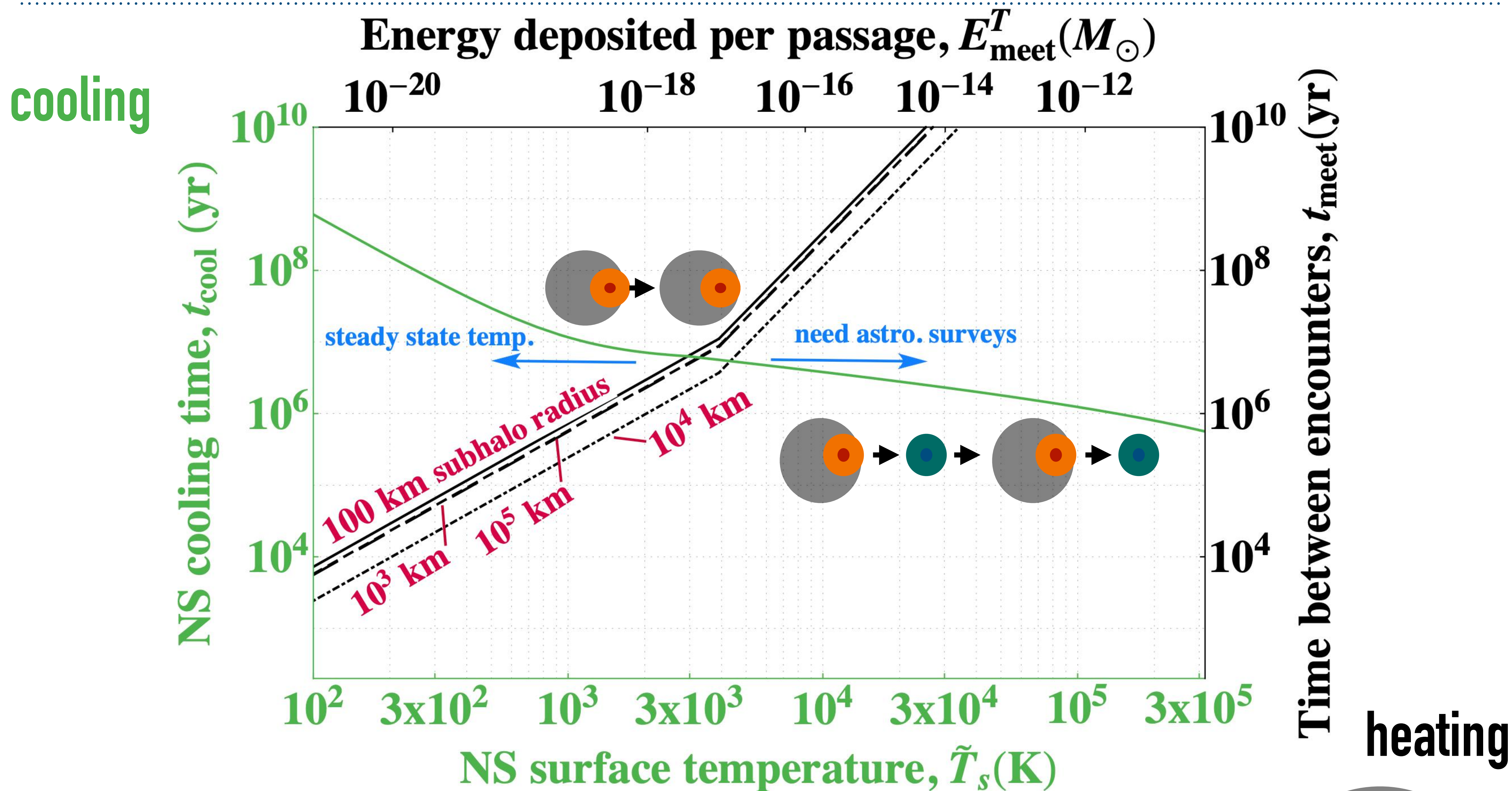
$$\left(\frac{\tilde{T}_{\text{hot}}}{10^4 \text{ K}}\right)^2 = \left(\frac{\tilde{T}_{\text{cold}}}{10^4 \text{ K}}\right)^2 + \frac{E_{\text{meet}}}{6.2 \times 10^{-18} M_{\odot}}$$

$$E_{\text{meet}} = zM \min \left[1, \left(\frac{R_{\text{co}}}{R_{\text{sh}}}\right)^2 \right]$$

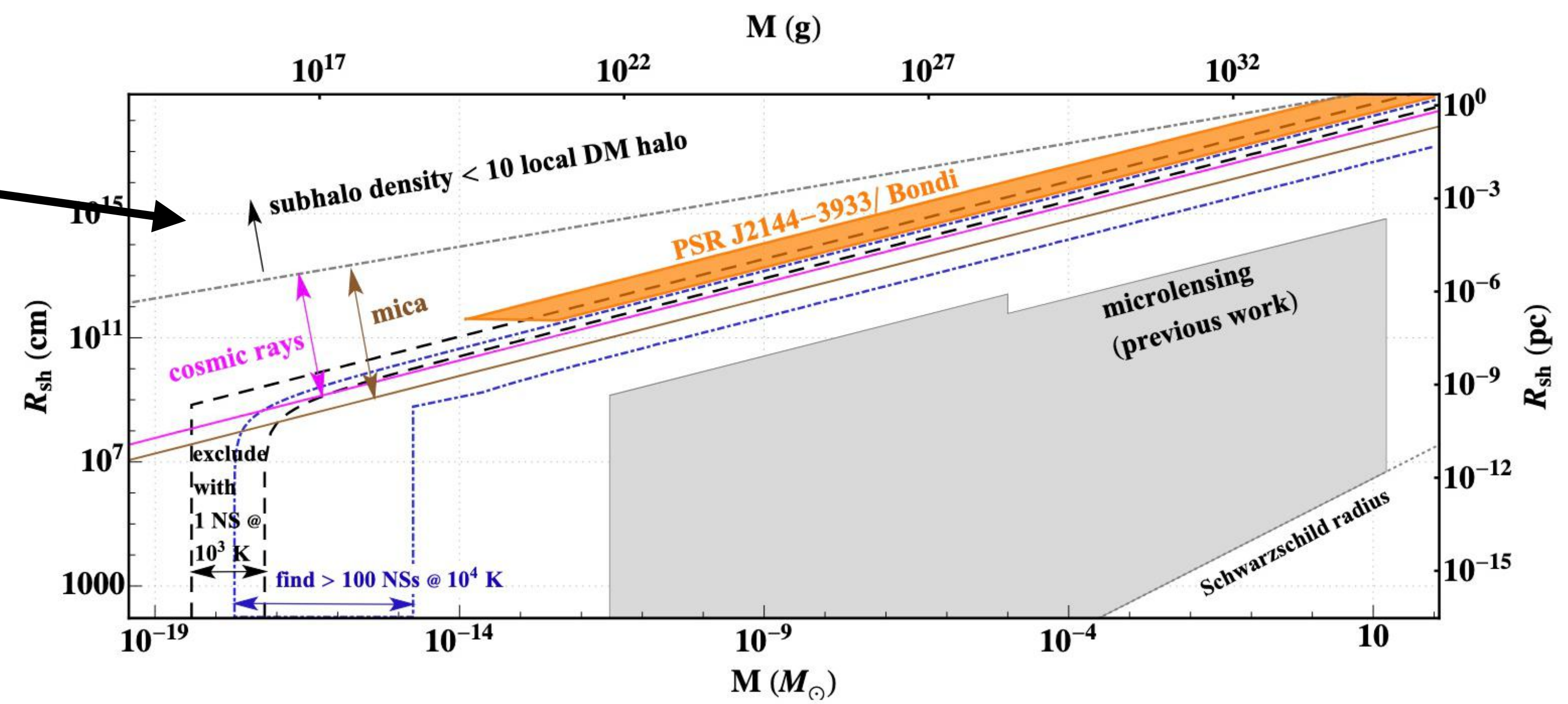
Neutron stars cooling and heating



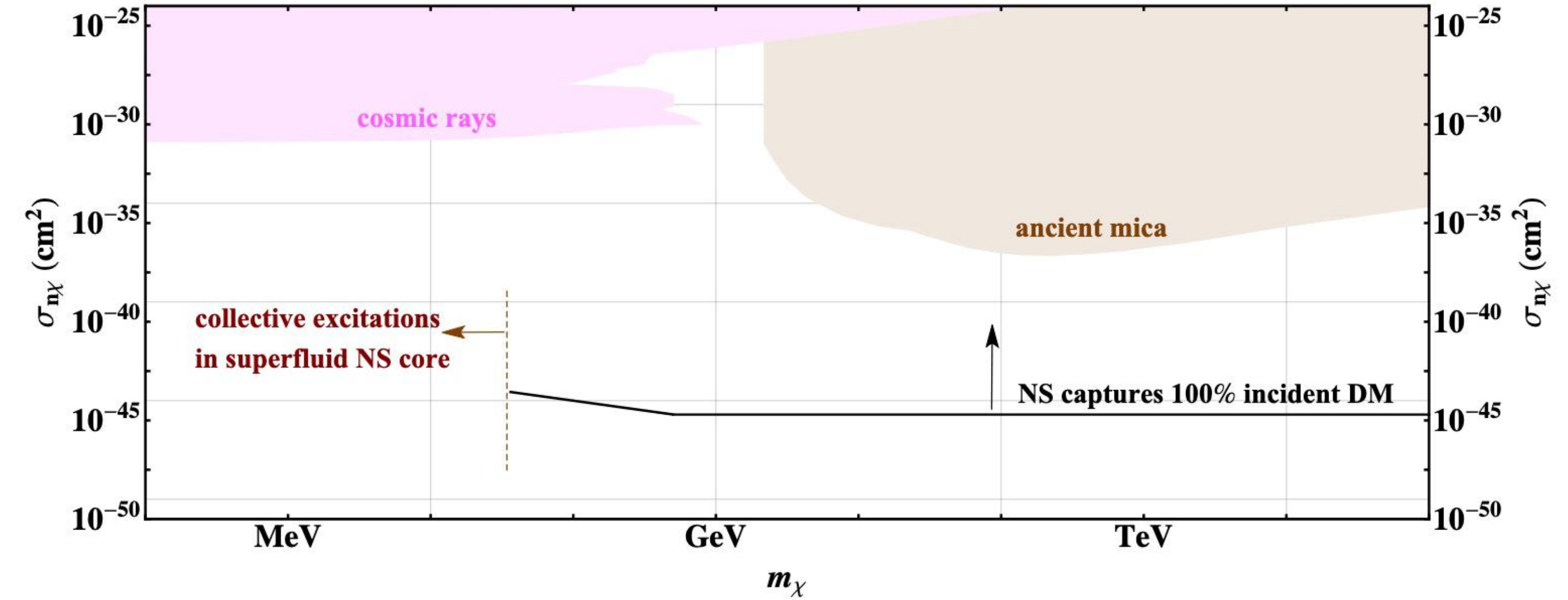
Neutron stars cooling and heating



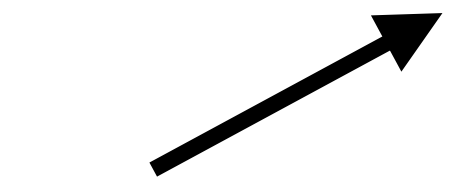
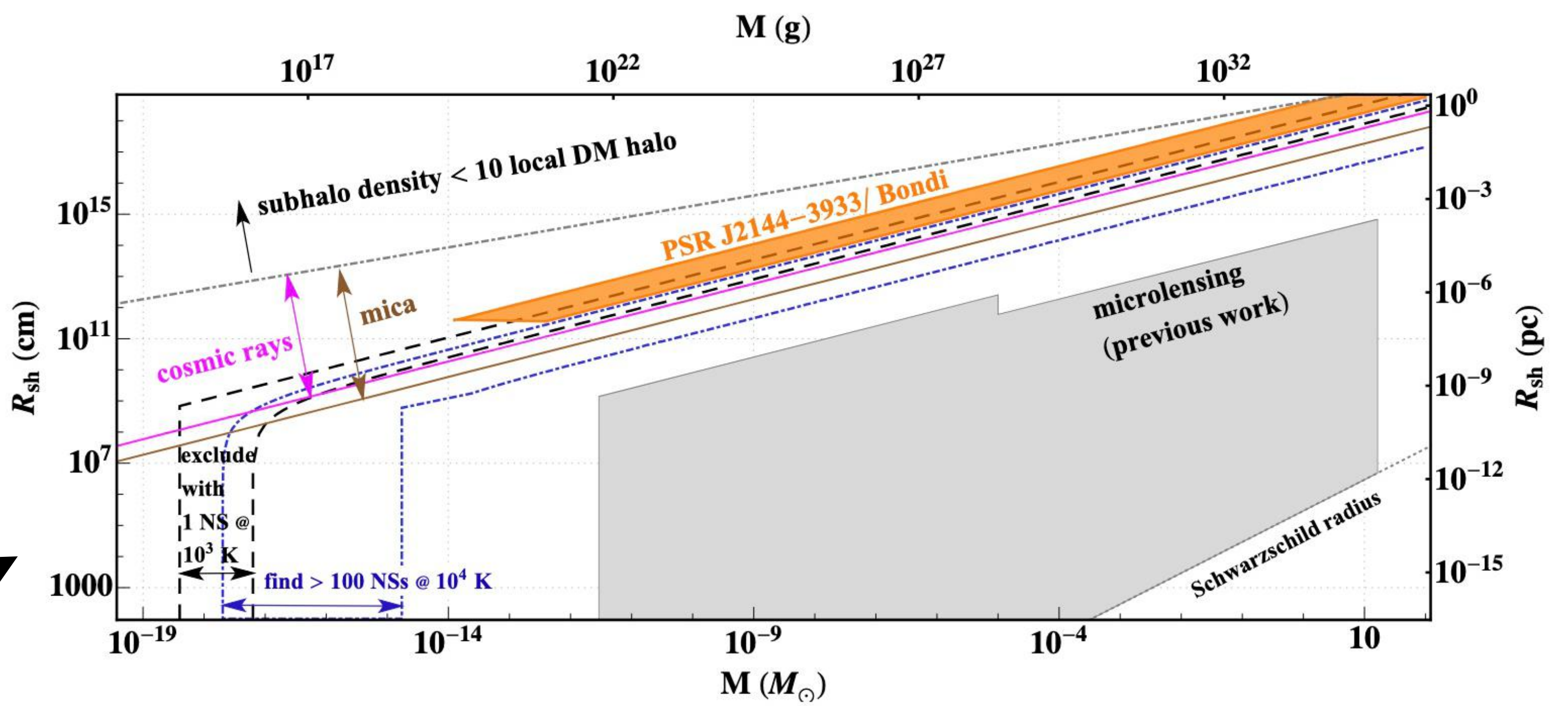
subhalo density < halo density



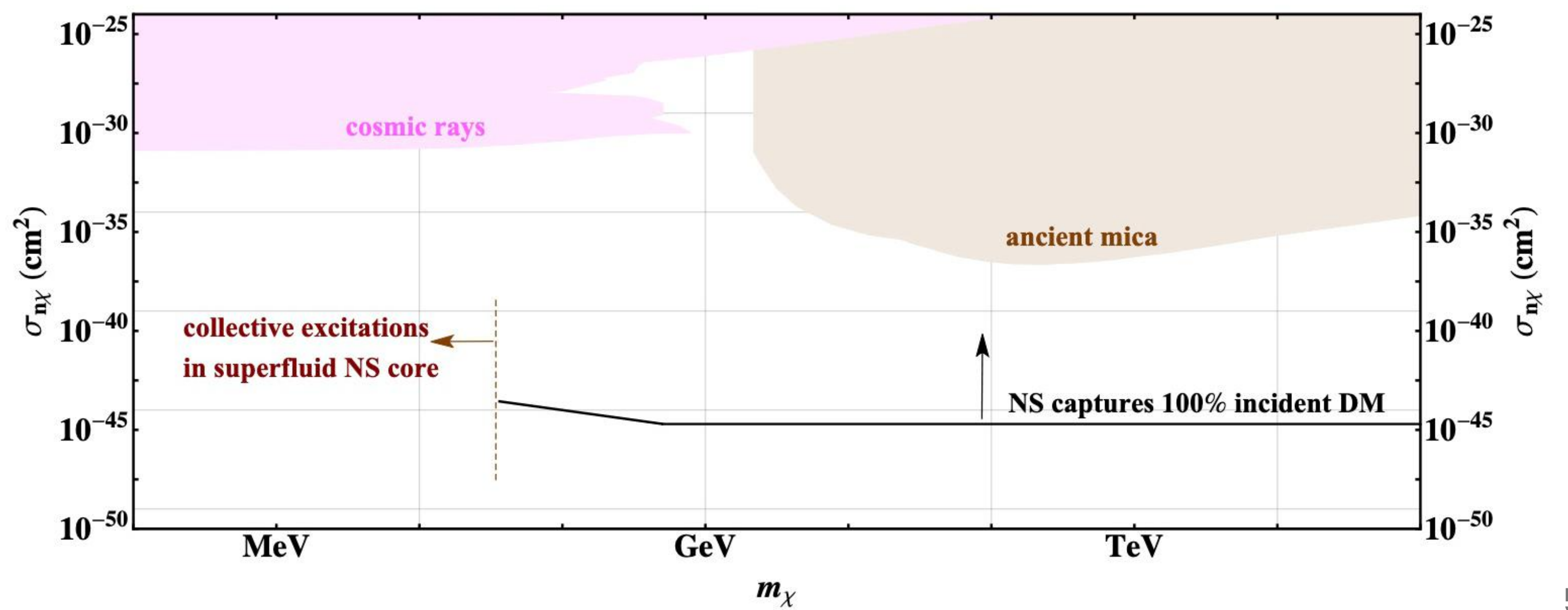
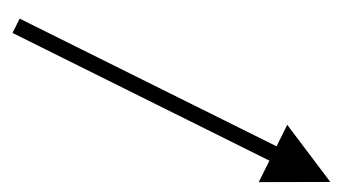
These bounds assume all DM is in subhalos



Subhalo mass/radius sensitivity regions in this plot



Correspond to the DM-nucleon cross-section sensitivity shown here

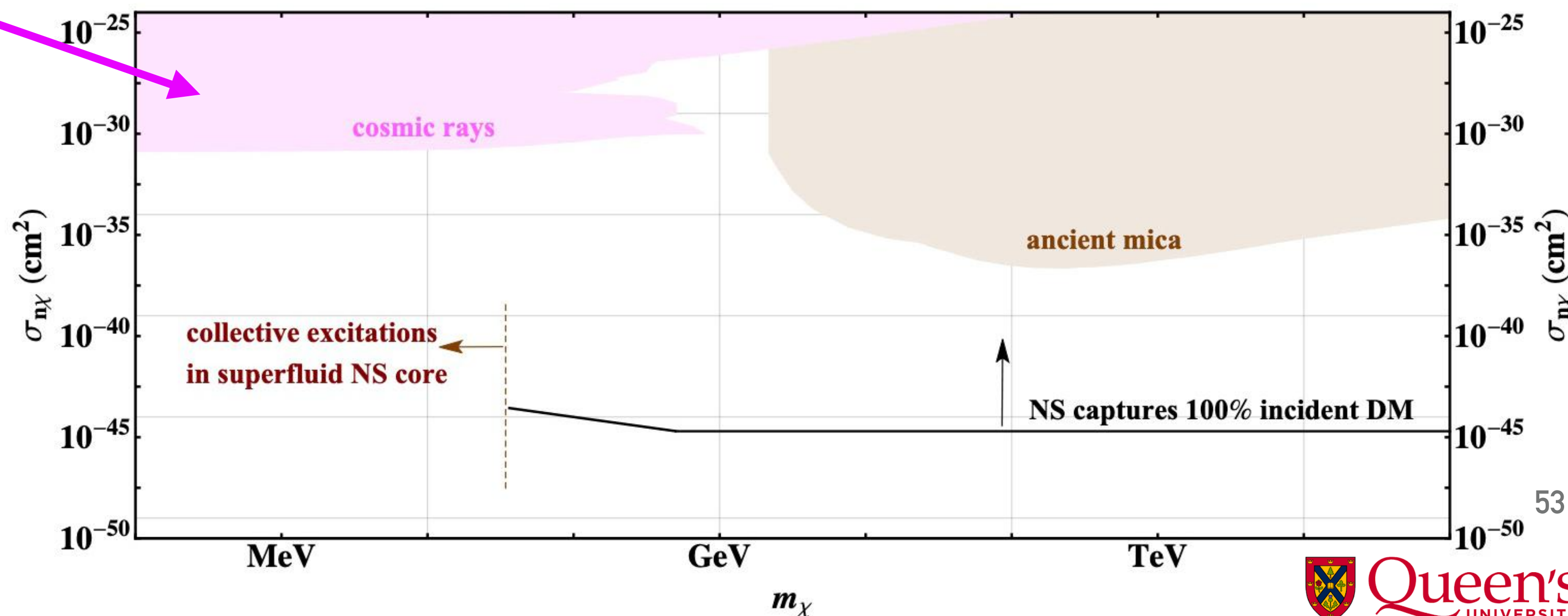
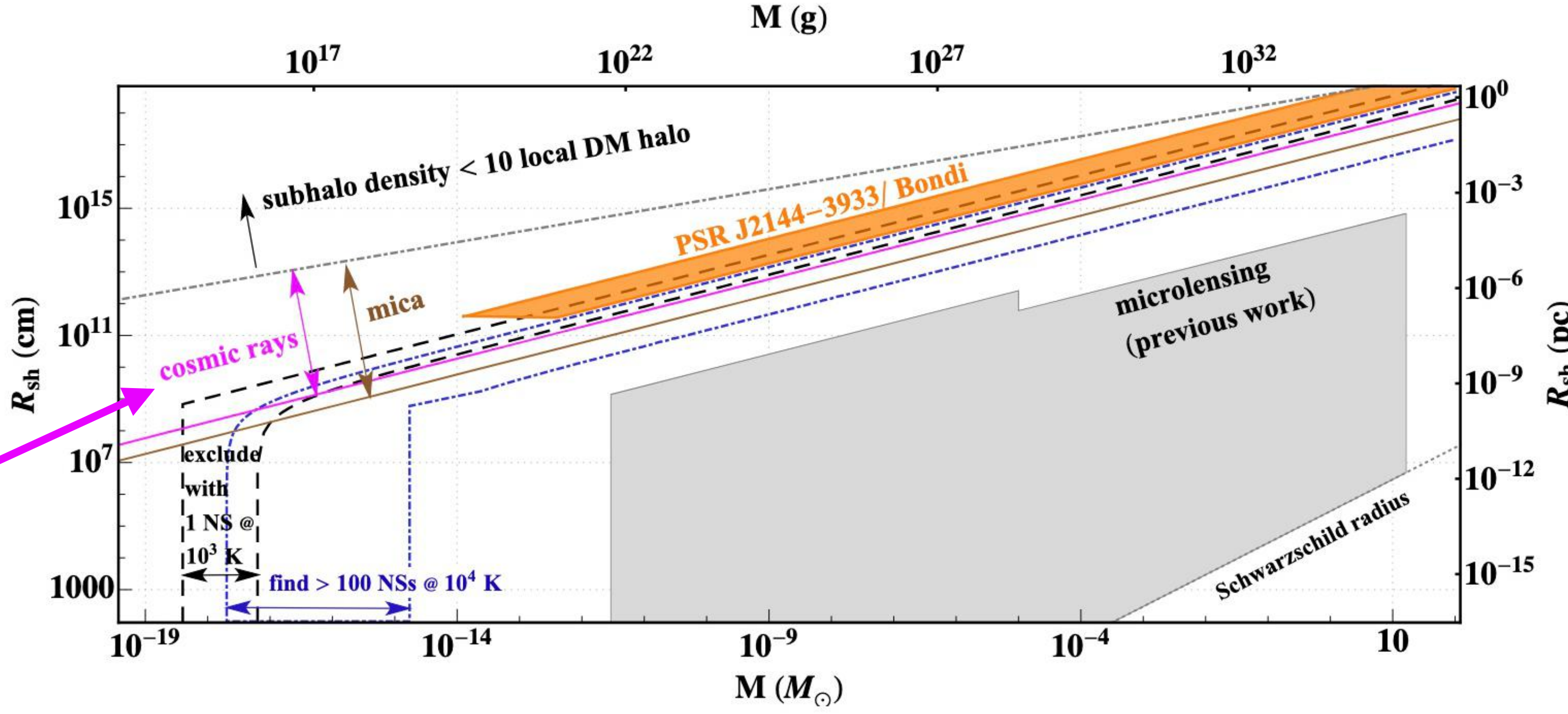


Cosmic rays will boost dark matter in subhalos so long as the interaction rate matches CR-diffuse DM interactions over ~8 kpc interaction lengths

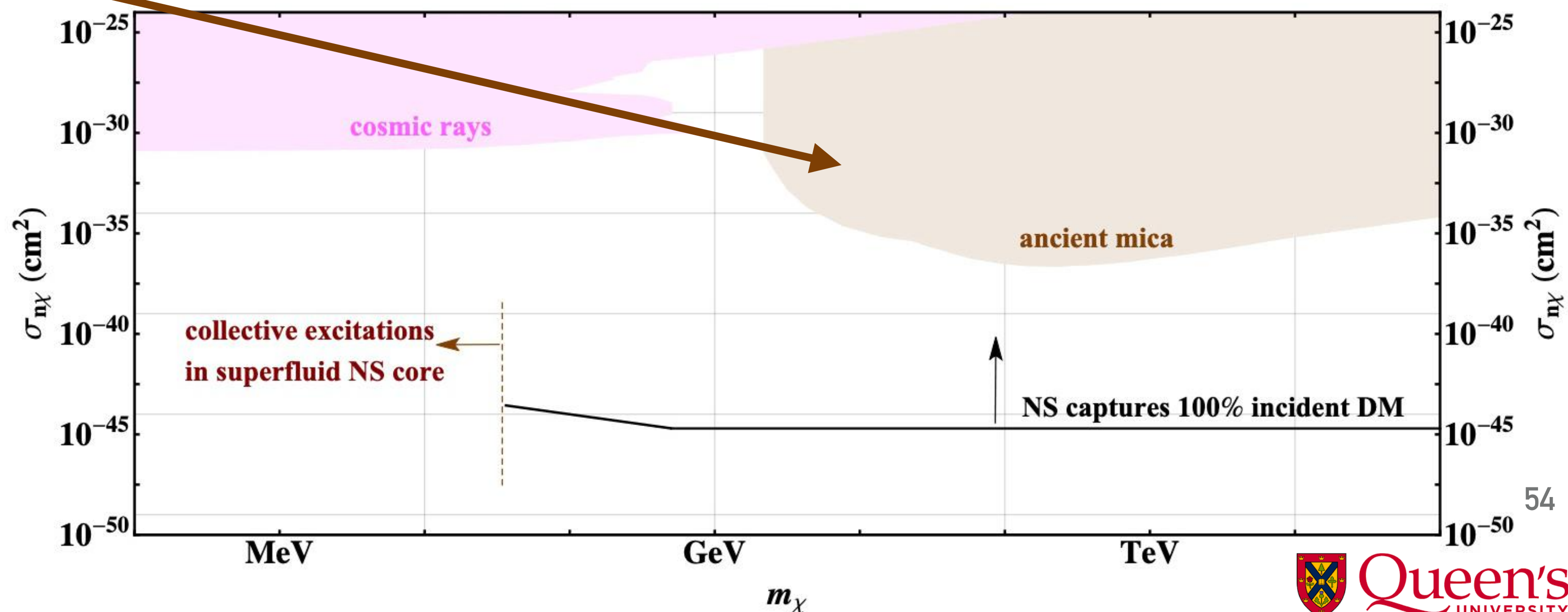
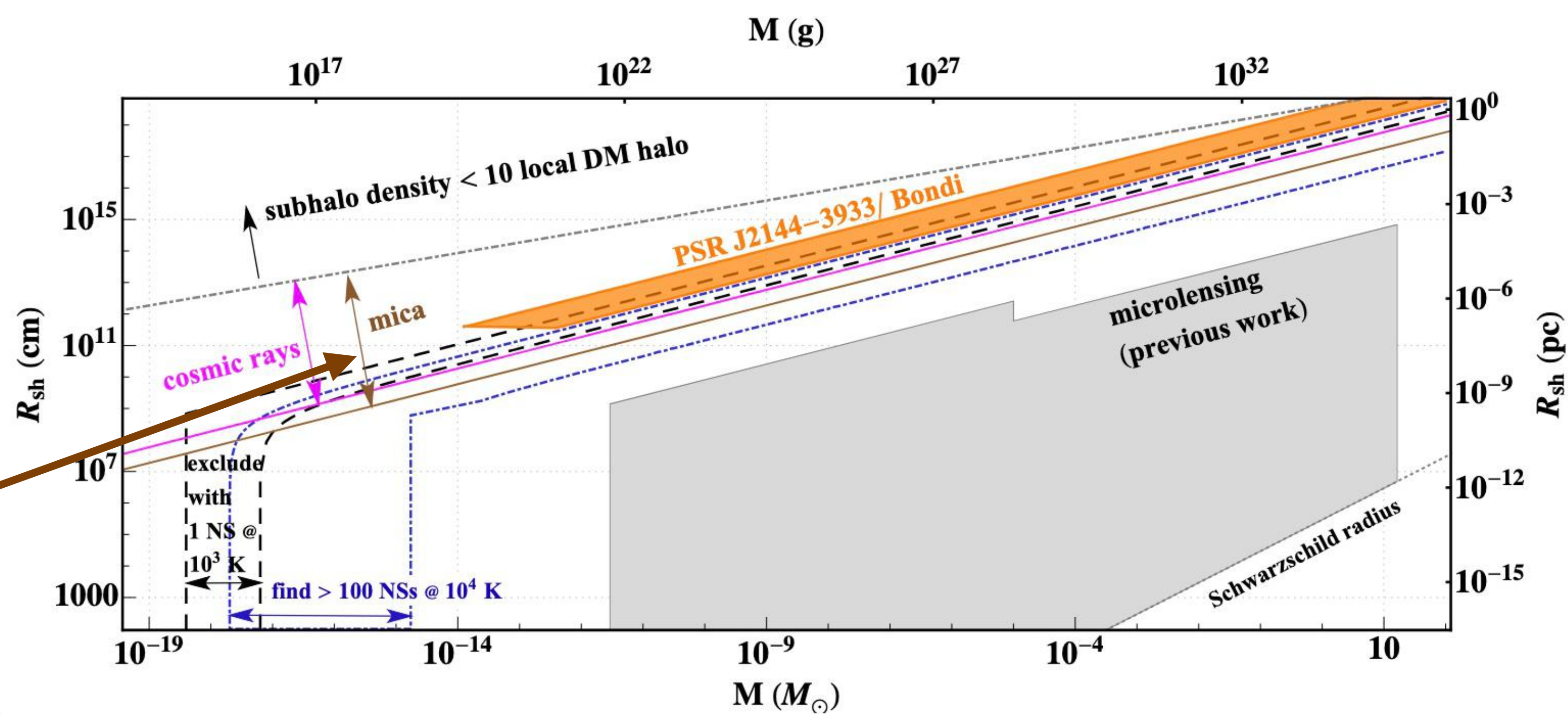
$$\tau_{\text{CR}} = \int_0^{L_{\text{dfs}}} ds n_{\chi}(s) \sigma_{\text{geo}} f_{\text{hit}} = \frac{\rho_{\odot}}{M} \sigma_{\text{geo}}$$

Low-mass DM boosted out of a subhalo yields similar CR-boosted DM bounds

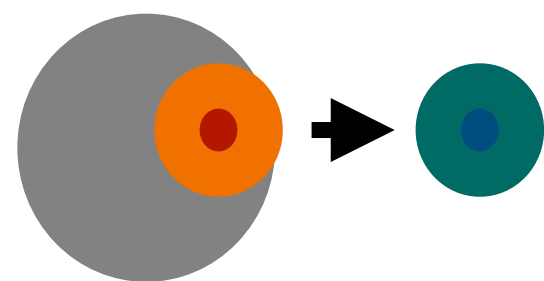
$\tau_{\text{CR}} > 1$



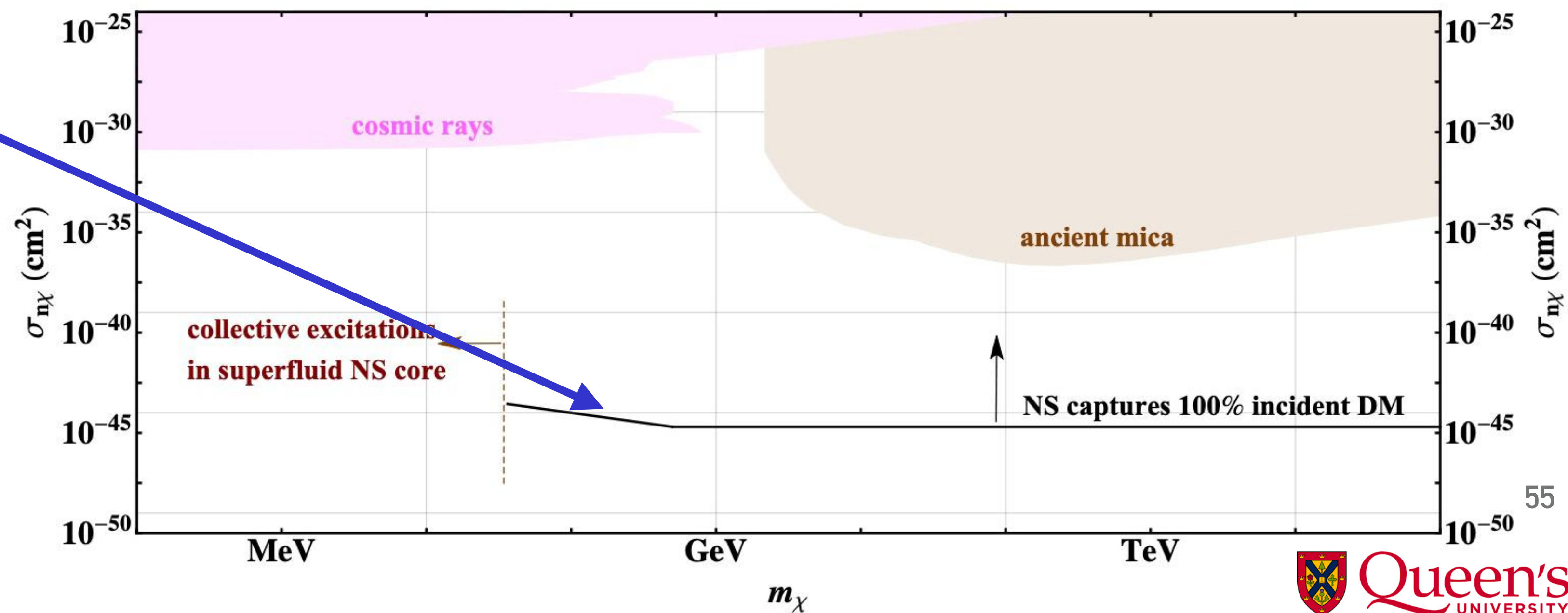
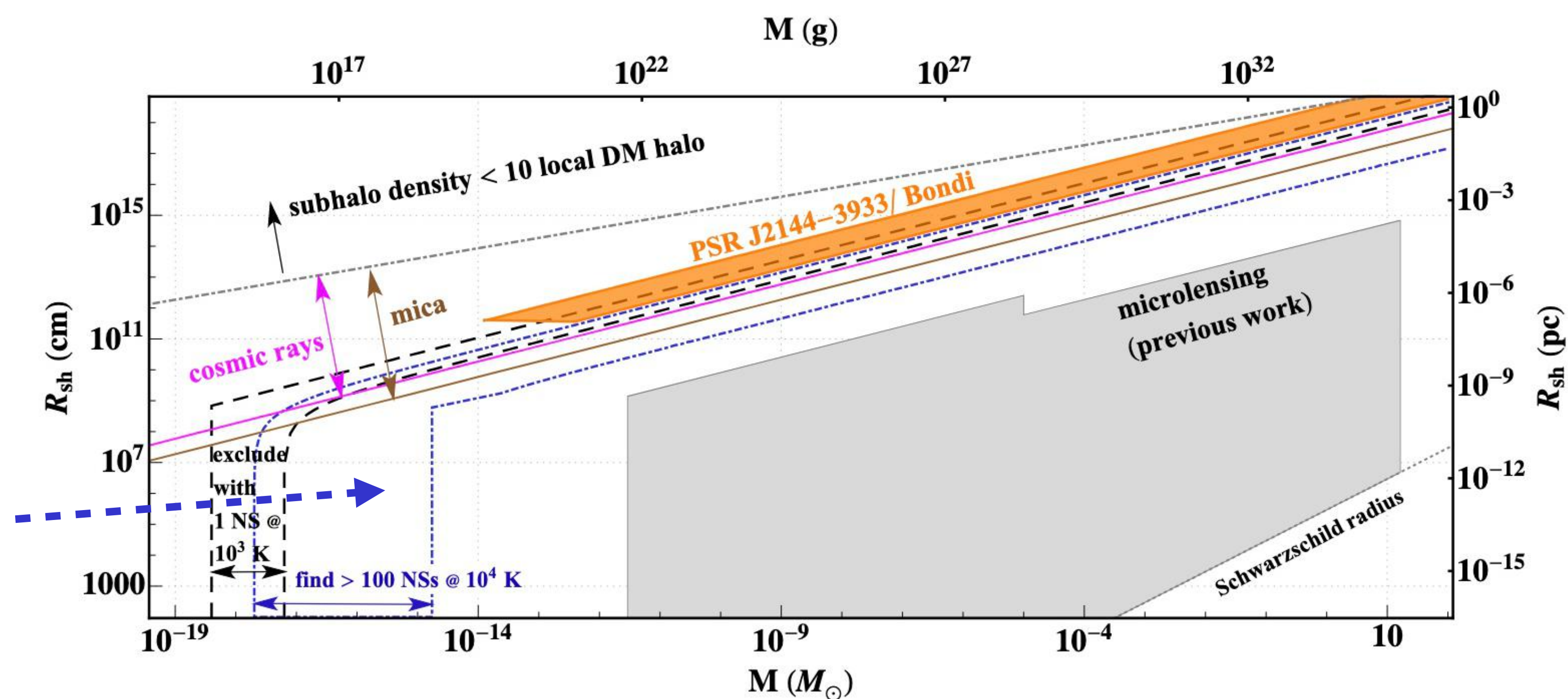
Requiring that subhalos interact with ancient mica in 500 Myr yields this bound



This region predicts a population of 100 NSs within a kpc of Earth heated to 10^4 K (flash-heated)



for this DM-nucleon cross-section



DETERMINING HOW DARK MATTER HEATS NEUTRON STARS

kinetic only

$$(\gamma - 1)\dot{m}_x$$

(very sensitive to escape velocity)

annihilation

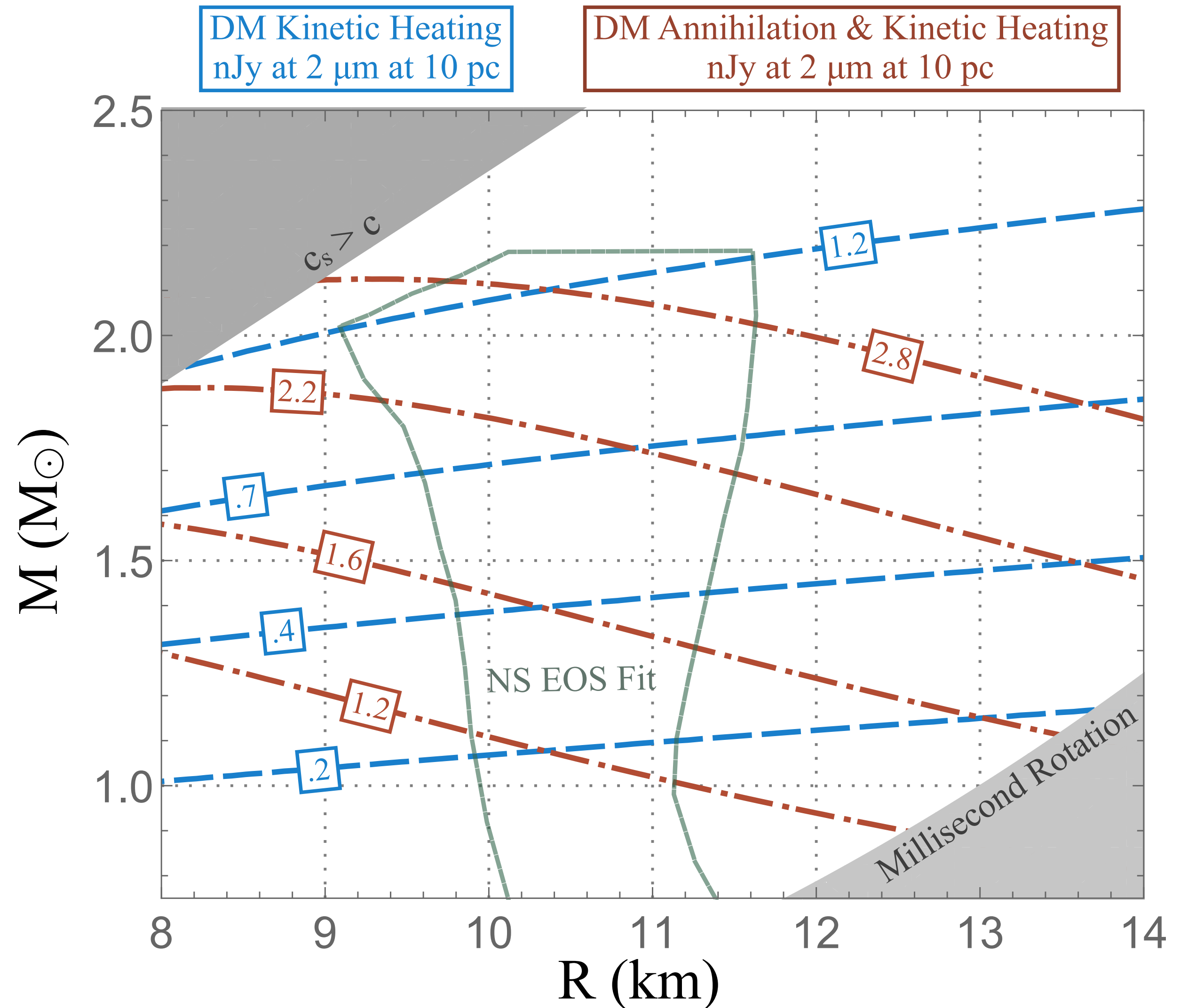
$$\gamma\dot{m}_x$$

(heating mostly from captured mass, scales with NS mass)

Test against ISM accretion:

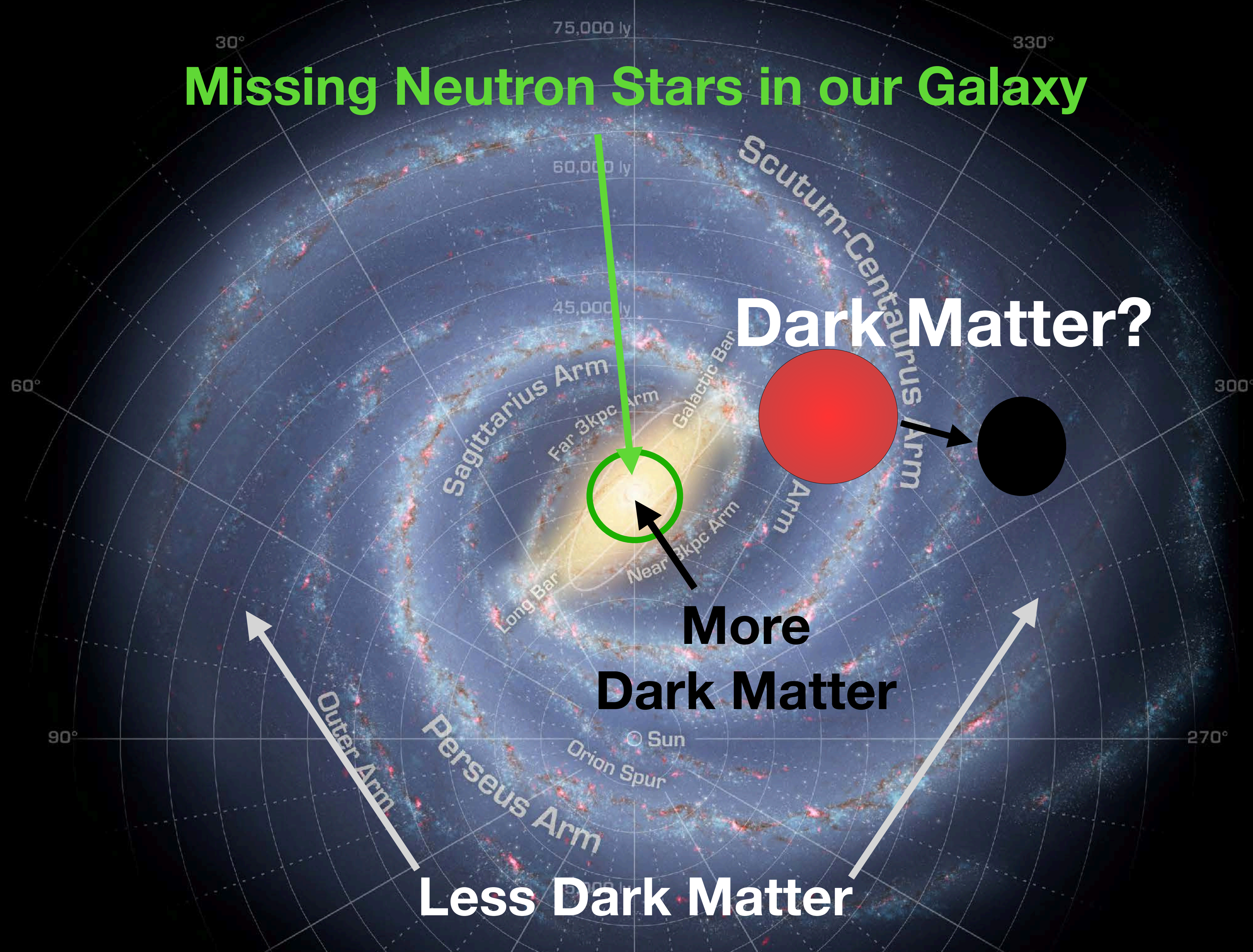
$$L_{dm} \propto \frac{1}{v_{dm-ns}}$$

$$L_{bar} \propto \frac{1}{v_{bar-ns}}$$



nanoJansky $\sim 10^{-30}$ GeV / (cm² s Hz)

Missing Neutron Stars in our Galaxy



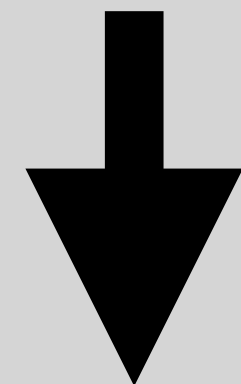
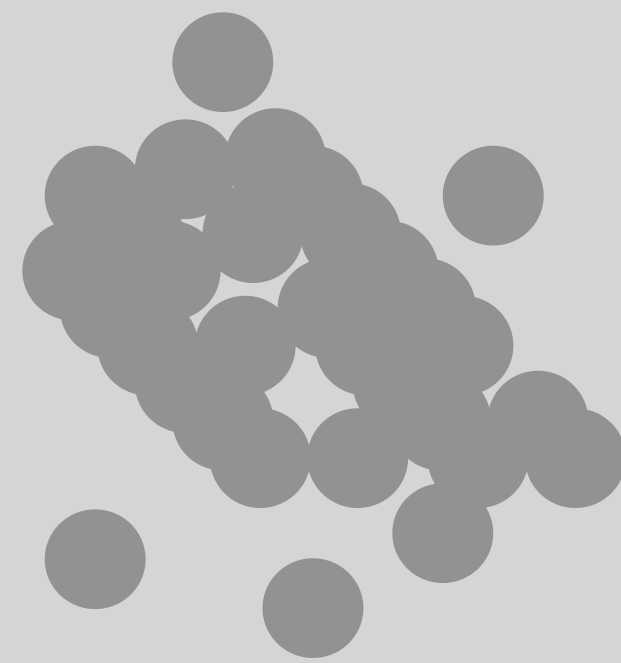
Dark Matter?

**More
Dark Matter**

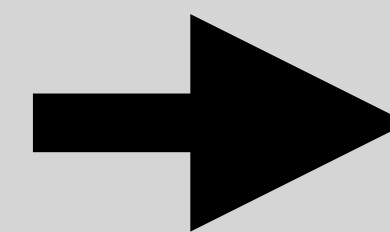
Less Dark Matter

What makes gold? (heavier elements)

Recipe: lots of neutrons, very hot (10^9 K)



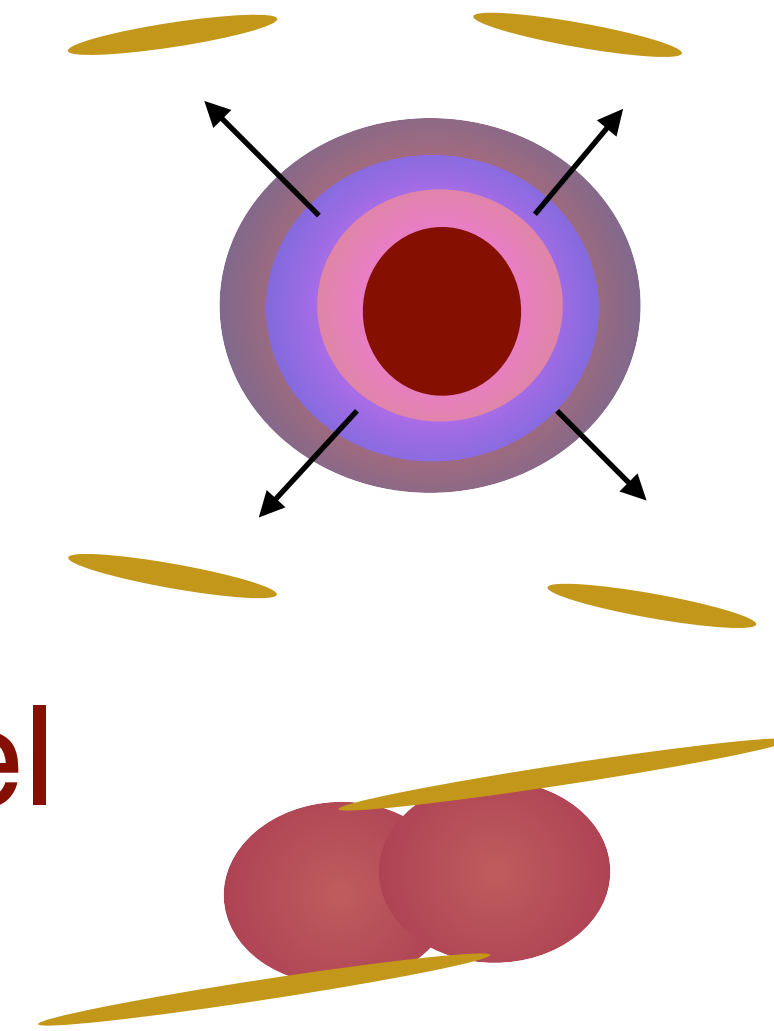
+ kilonova light



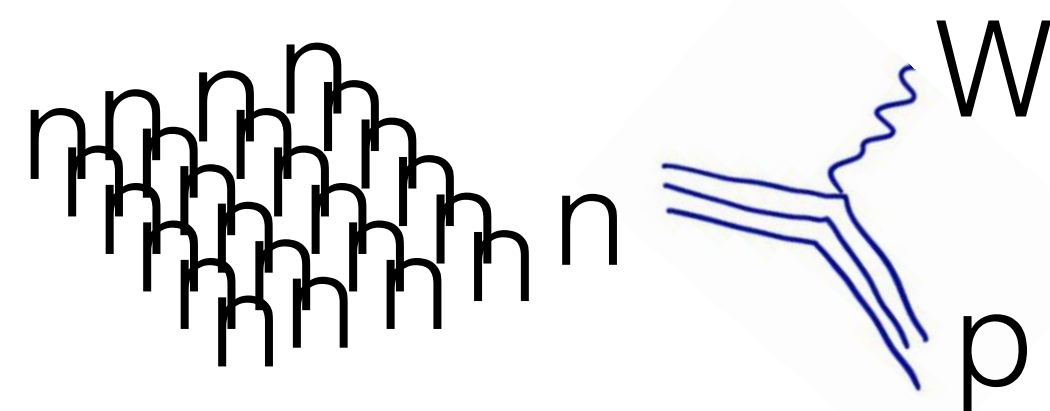
Possible r-process sites — total $10^4 M_{\odot}$ produced in Milky Way

-Neutrons ejected by neutrino wind during core collapse supernovae (frequent, $\sim 1/100$ years)

-Merging neutron star binaries, tidal forces expel dense neutron star fluid (rare, $\sim 1/10^4$ years)

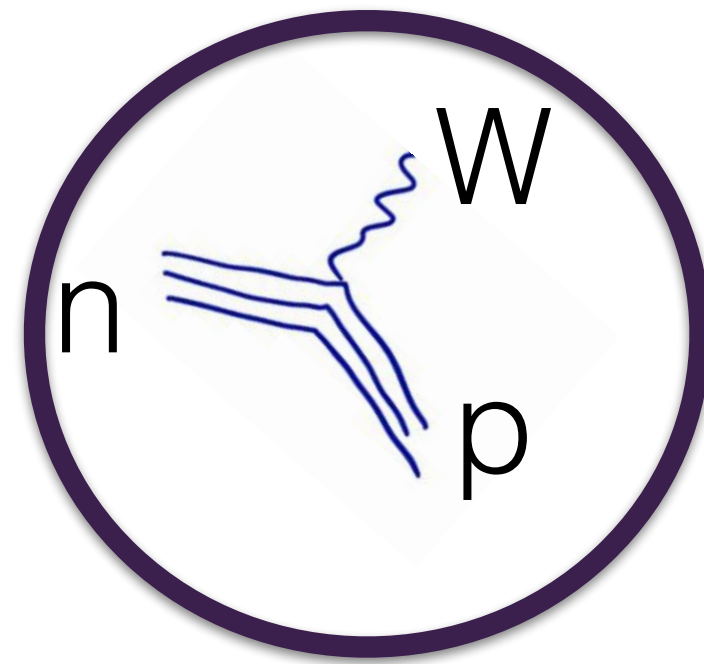


In each case, neutron rich fluid beta decays,
forms heavy neutron-rich elements.



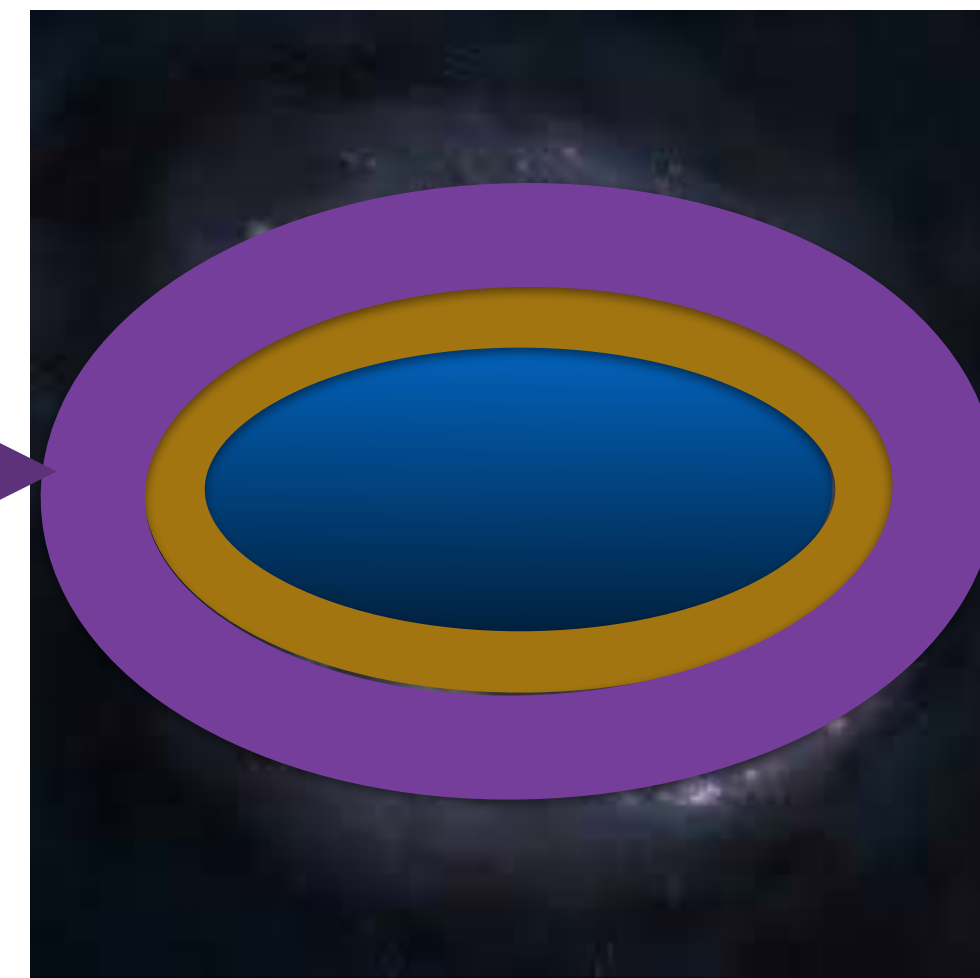
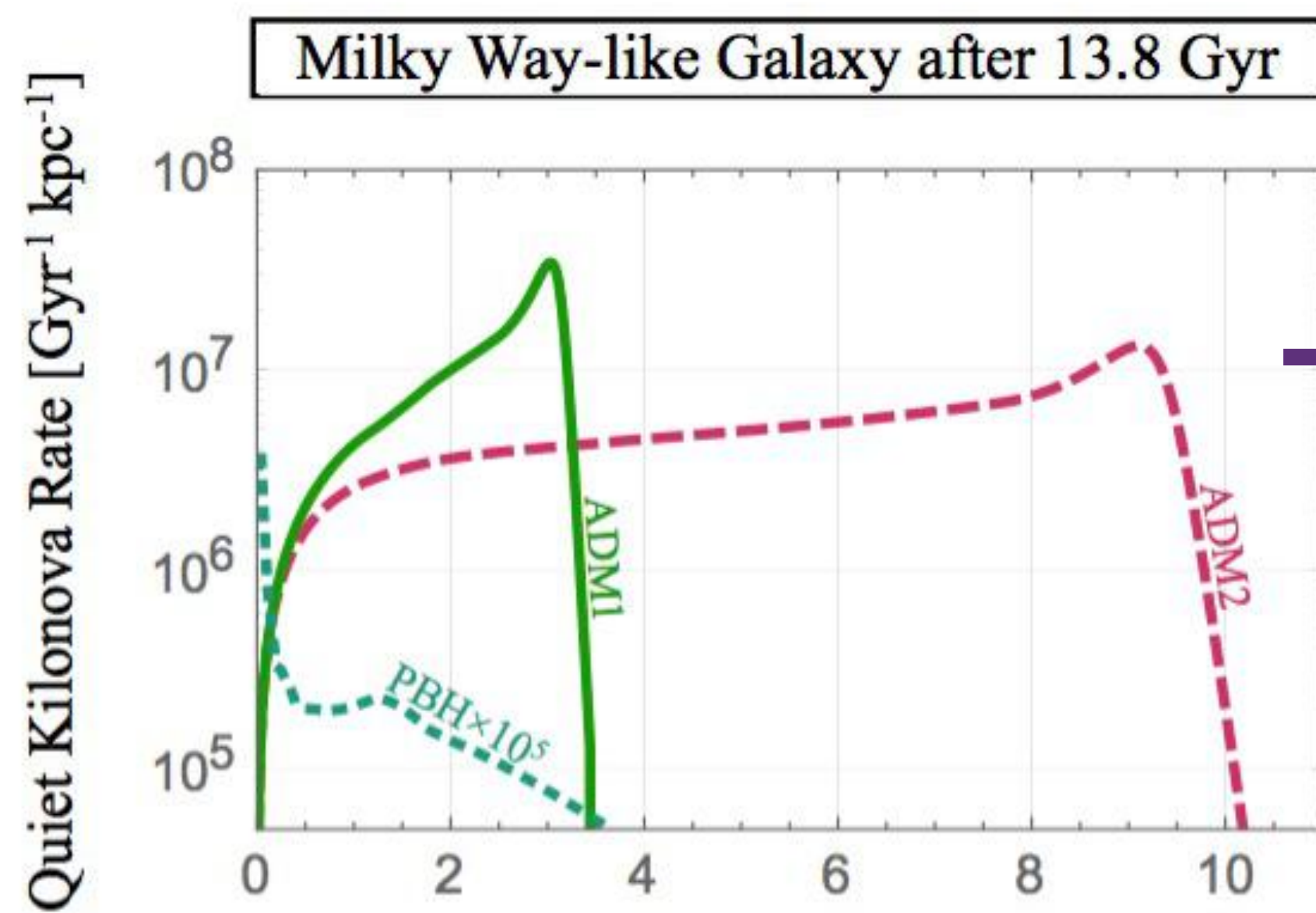
... Gold, Uranium,
Europium, Barium...

R-process donuts



Gravity Waves
From NS mergers

R-Process Donuts: Standard NS merger kilonovae occur in donut shaped external regions of disc galaxies



Asymmetric Dark Matter Using Kilonovae And NS Mergers

