

The effect of fast neutrino flavor conversions on Neutrino-driven wind and Supernova nucleosynthesis

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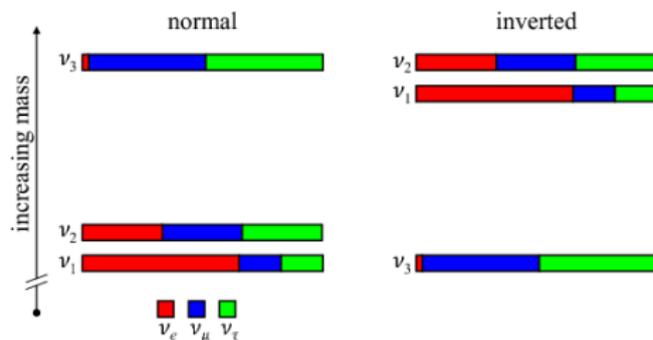
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Neutrino flavor transformations

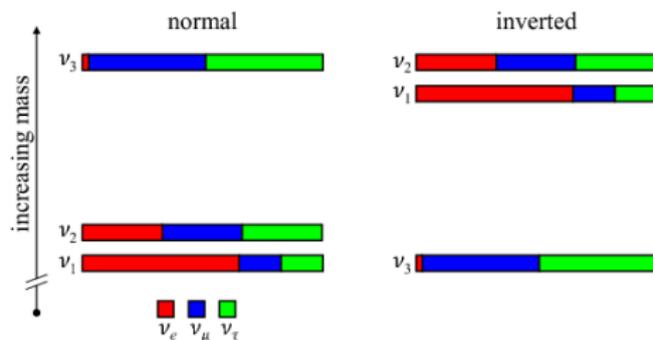
- Neutrino mass and weak interaction eigenstates are not the same
- Rich physics of neutrino flavor transformations ($\nu_e \leftrightarrow \nu_{\mu,\tau}$, $\bar{\nu}_e \leftrightarrow \bar{\nu}_{\mu,\tau}$)
- ν -matter and $\nu - \nu$ interactions can induce flavor instabilities



• Image credit: KM3NeT collaboration

Neutrino flavor transformations

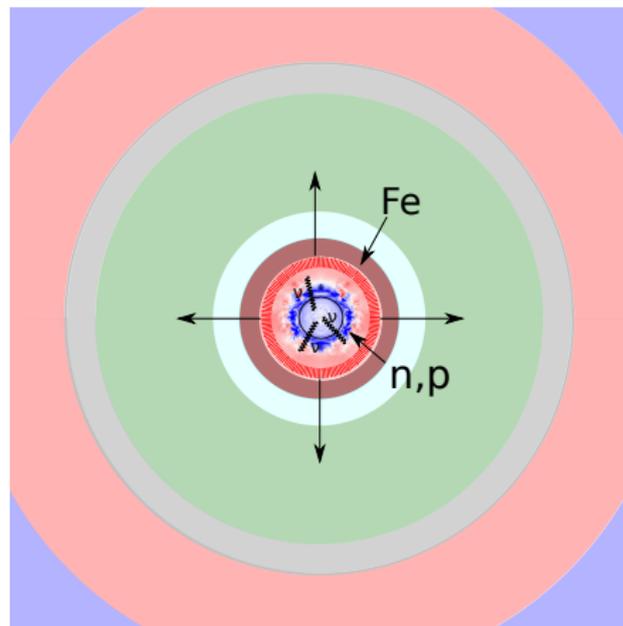
- Neutrino mass and weak interaction eigenstates are not the same
- Rich physics of neutrino flavor transformations ($\nu_e \leftrightarrow \nu_{\mu,\tau}$, $\bar{\nu}_e \leftrightarrow \bar{\nu}_{\mu,\tau}$)
- ν -matter and $\nu - \nu$ interactions can induce flavor instabilities
- **"Fast" neutrino flavor transformations may occur on very short scales and at high matter densities**



- *Image credit: KM3NeT collaboration*
- **Potential effect on the production of the elements in Supernovae**

Neutrinos and supernovae

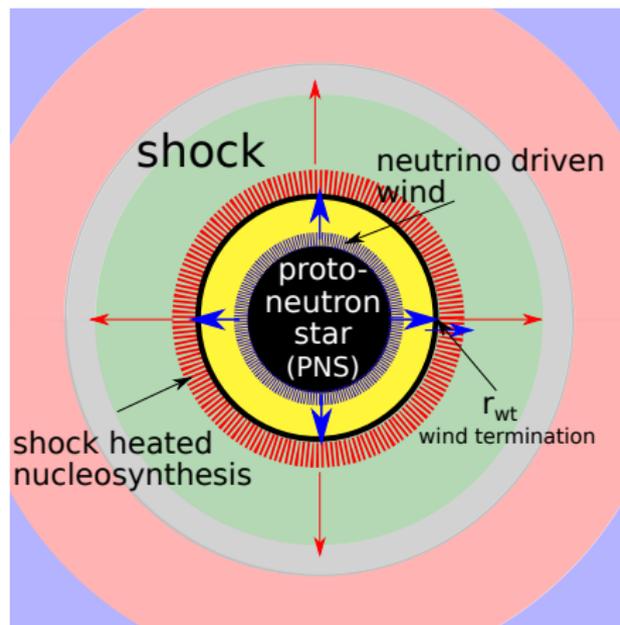
- Fe core of a massive star ($\geq 8 - 10 M_{\odot}$) collapses until nuclear densities are reached
- Strong neutrino emission from electron captures and cooling ($\sim 10^{56}/s$)
- Bounce shock stalls (ν losses, nuclear dissociation)
- Neutrino heating induces turbulence and revives the explosion



- *Inset from Pan et al. (2016)*

Neutrinos and supernovae

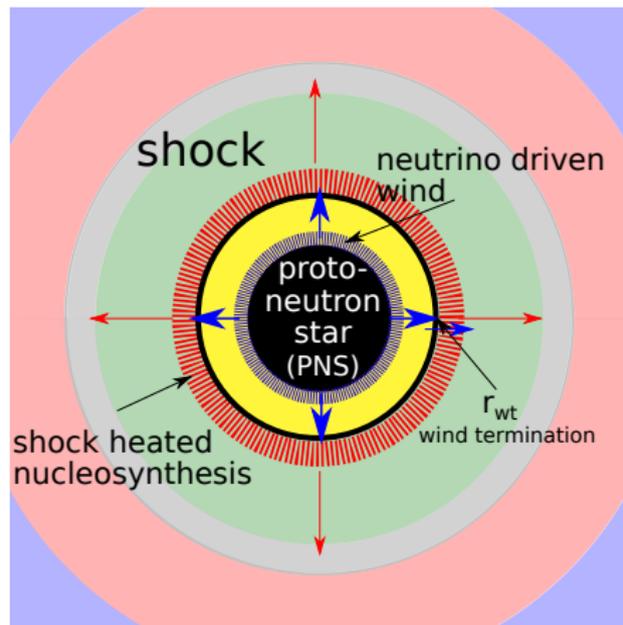
- Three sites of nucleosynthesis
 - 1 Inner, neutrino heated, shock-driven ejecta
 - 2 Outer shock-driven ejecta
 - 3 **Neutrino driven wind (NDW)**



Inset from Pan et al. (2016)

Neutrinos and supernovae

- Three sites of nucleosynthesis
 - 1 Inner, neutrino heated, shock-driven ejecta
 - 2 Outer shock-driven ejecta
 - 3 **Neutrino driven wind (NDW)**
- NDW: Fast, hot matter outflow from the PNS surface: few $10^{-3} M_{\odot}$
- NDW is determined by long-term neutrino cooling of the PNS, decoupled from the explosion mechanism
- Neutrinos determine Y_e of the ejecta

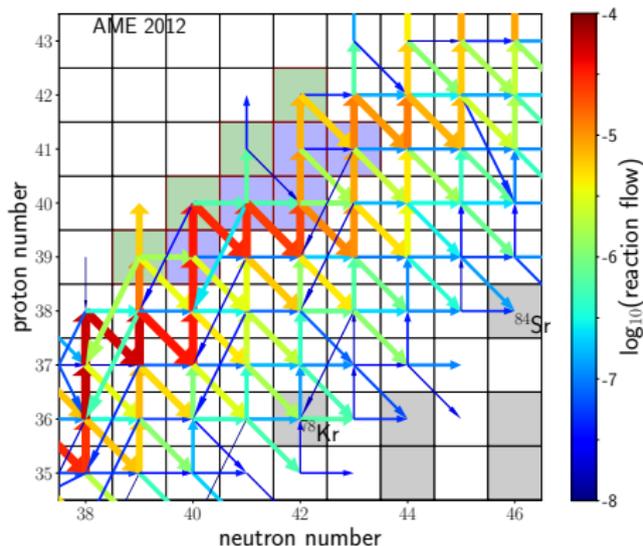


Inset from Pan et al. (2016)

- Wind slows down as it catches up with the shocked ejecta

Nucleosynthesis conditions: νp process

- NDW is expected to be proton-rich ($Y_e > 0.5$)
 - Proton capture nucleosynthesis inhibited by long β^+ decay lifetimes
 - If neutrons are present: (n, p) instead of β^+ decay
 - $\bar{\nu}_e + p \rightleftharpoons n + e^+$ as neutron source
-
- Important quantities:
 - ▶ Initial proton-richness (Y_e)
 - ▶ p-to-seed ratio (depends on expansion timescale, entropy)
 - ▶ $\bar{\nu}_e$ exposure while the temperature is right (3 – 1 GK)



Steady state model for NDW

- Time-independent eigenvalue problem for mass loss rate \dot{M}

$$\dot{M} = 4\pi r^2 \rho v, \quad (1a)$$

$$v \frac{dv}{dr} = -\frac{1}{\rho} \frac{dP}{dr} - \frac{GM_{\text{PNS}}}{r^2}, \quad (1b)$$

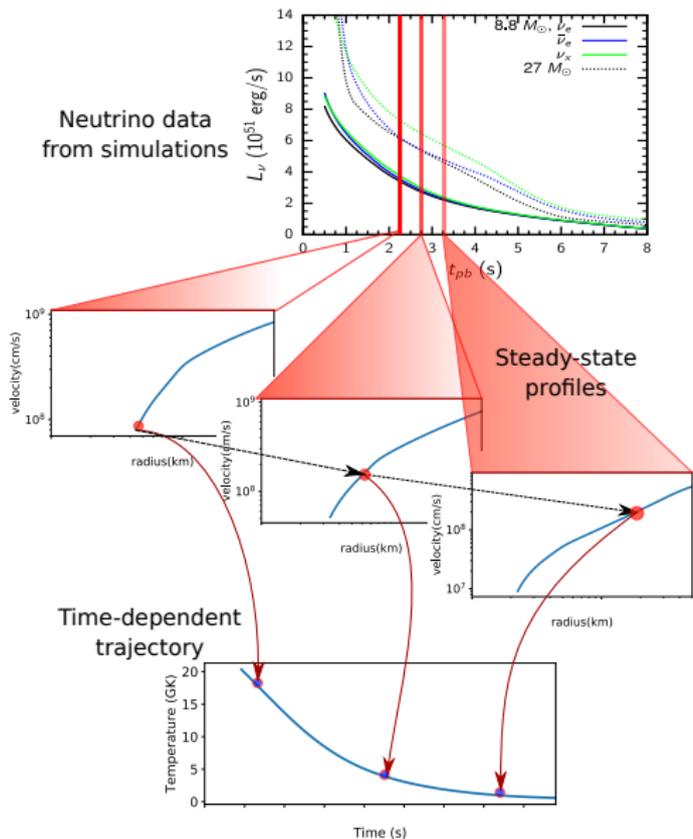
$$\frac{d\epsilon}{dr} = \frac{P}{\rho^2} \frac{d\rho}{dr} + \frac{\dot{q}}{v}, \quad (1c)$$

$$\frac{dY_e}{dr} = \frac{1}{v} [(\lambda_{\nu_e n} + \lambda_{e+n}) Y_n - (\lambda_{\bar{\nu}_e p} + \lambda_{e-p}) Y_p], \quad (1d)$$

- Inner boundary at $R_\nu(t)$: $\dot{q} = 0$, $dY_e/dr = 0$, and $T = T_{\nu_e}$
- Outer boundary : Transonic solution, i.e., $\frac{dv}{dr} \rightarrow 0$
- **Interaction and heating based on neutrino fluxes and spectra from simulations**
- Details in *Xiong et al. (2019) ApJ 880, 81*

Time-dependent model

- Time-dependent neutrino luminosity and spectra from simulations
- Sequence of steady state snapshots covering up to 1000 km
- Tracking "tracers" across snapshots gives time-dependent thermodynamic trajectories
- Allows to calculate total wind mass and integrated yields
- Two models: e8.8 and s27



Neutrino flavor transformations

- In a supernova, $\langle E_{\nu_{\mu,\tau}} \rangle > \langle E_{\nu_e} \rangle$
- Flavor conversion can turn high energy $\nu_{\mu,\tau}$ into ν_e
- **Two cases:**

Flavor equilibrium

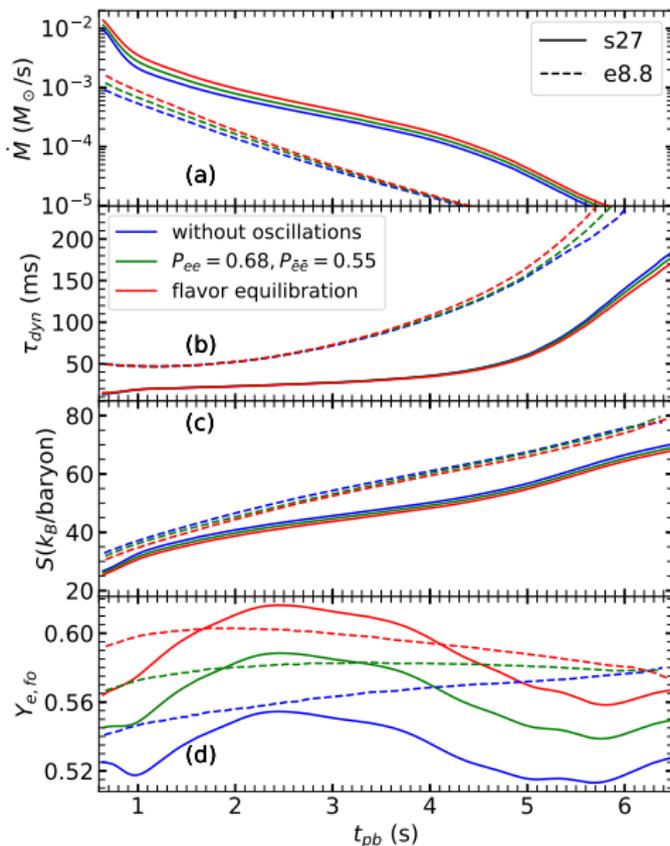
- $$F_{\nu_e}^{\text{osci}} = 1/3 \sum_{i=e,\mu,\tau} F_{\nu_i}^0$$

Intermediate mixing

- $$F_{\nu_e}^{\text{osci}} = P_{ee} F_{\nu_e}^0 + (1 - P_{ee}) F_{\nu_x}^0$$
- $$F_{\bar{\nu}_e}^{\text{osci}} = P_{\bar{e}\bar{e}} F_{\bar{\nu}_e}^0 + (1 - P_{\bar{e}\bar{e}}) F_{\bar{\nu}_x}^0$$
- $$P_{ee} = 0.68, P_{\bar{e}\bar{e}} = 0.55$$

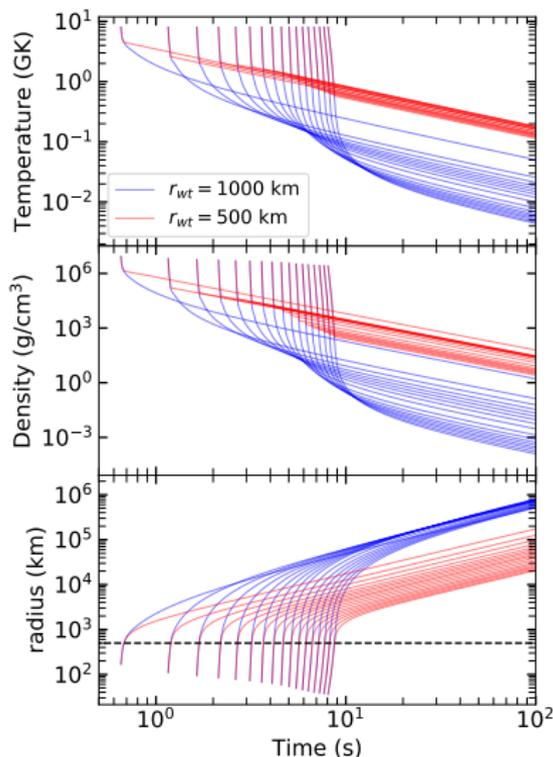
Effects on dynamics

- Flavor oscillations effectively make ν_e and $\bar{\nu}_e$ more energetic
- Charged current reactions are the main heating mechanism
- Increased wind outflow rate \dot{M} , increased total wind mass by up to 40%
- Otherwise, only minor changes of the dynamics
- More energetic ν_e lead to an increase of Y_e



Wind termination radius

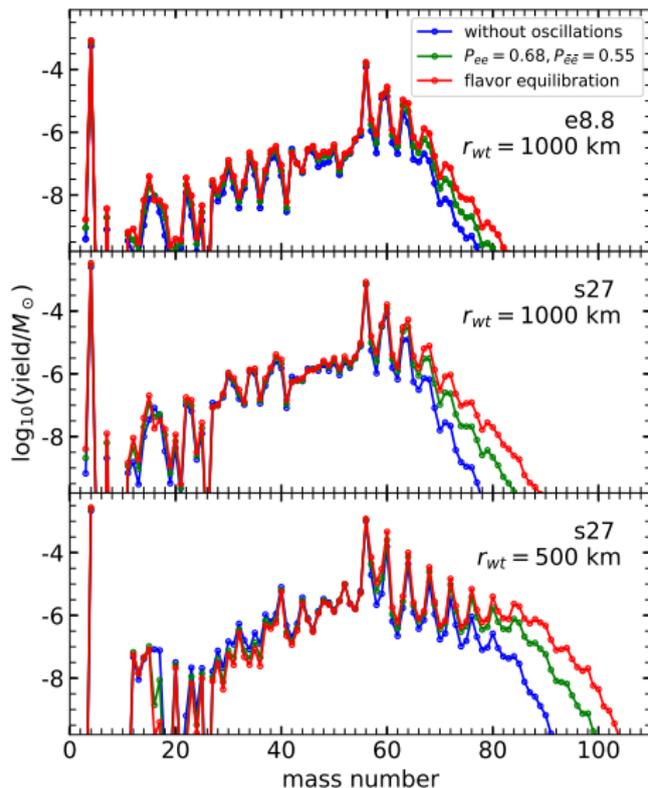
- Wind slows down when it catches up with the SN ejecta at r_{wt}
- Material spends more time in the relevant temperature range (*Wanajo et al. 2011*)
- Consistent model requires full simulation
- Simple model: switch to parametric expansion at smaller r_{wt}
- additional model: s27 with $r_{wt} = 500$ km



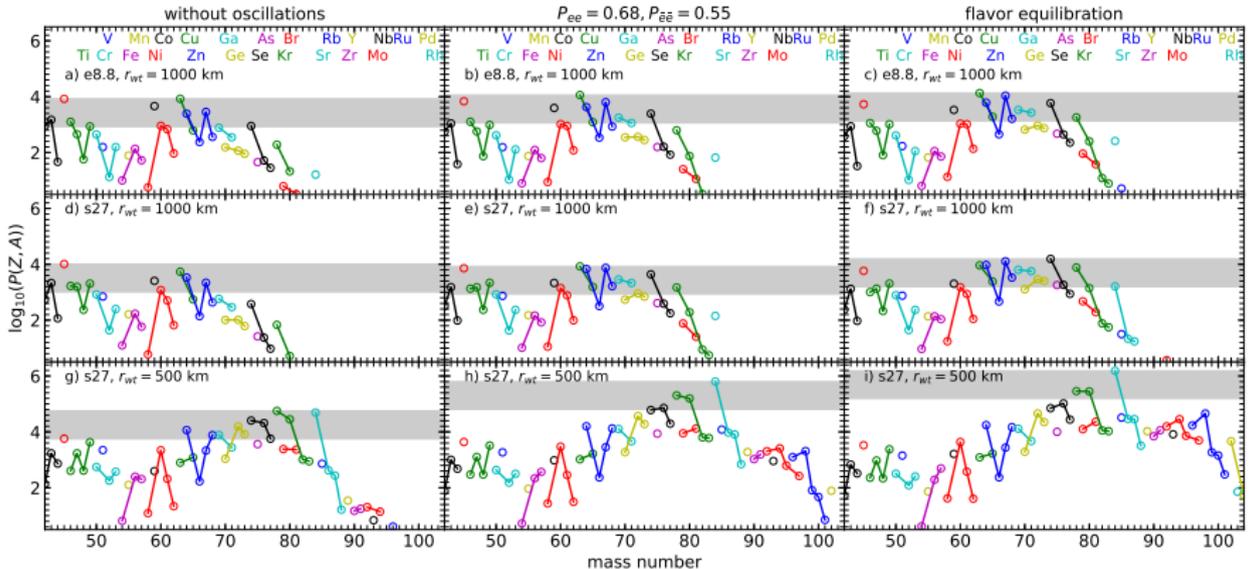
In total: 3 models x 3 oscillation cases

Effects on nucleosynthesis

- Standard e8.8 and s27 models ($r_{wt} = 1000$ km) show little effect of the νp process
- With $r_{wt} = 500$ km, the production of heavier nuclei is enhanced
- More energetic $\bar{\nu}_e$ favor νp process by providing more neutrons
- More energetic ν_e favor lead to higher Y_e at freeze-out
 - ▶ higher Y_e at freeze-out
 - ▶ more neutrons

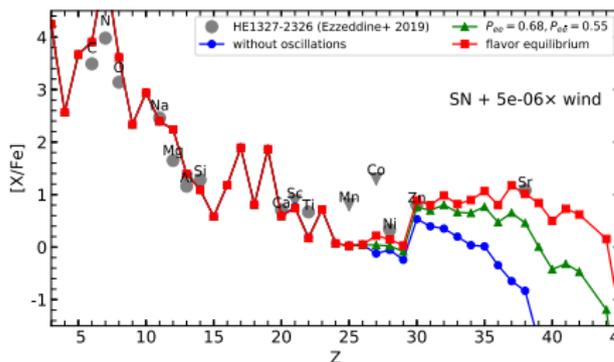
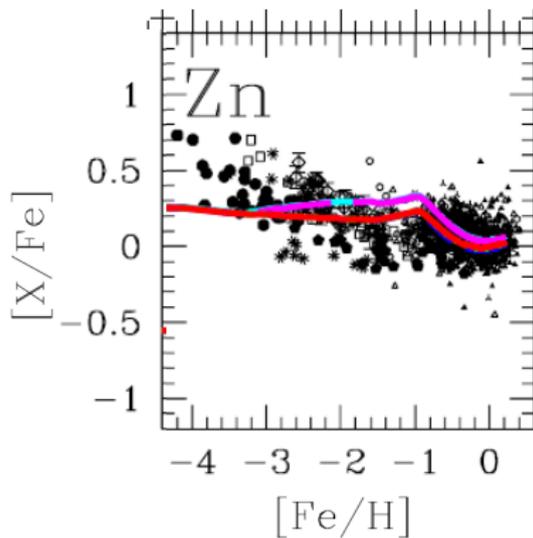


Effects on nucleosynthesis



- Significantly increased production of ^{64}Zn , ^{78}Kr , ^{84}Sr
- At most 3% of SN should produce conditions as in model s27, $r_{wt} = 500$ km and flavor equilibrium to avoid overproduction of ^{84}Sr

Implications



Model s27 with $r_{wt} = 500$ km, combined with SN yields from Heger et al. (2010)

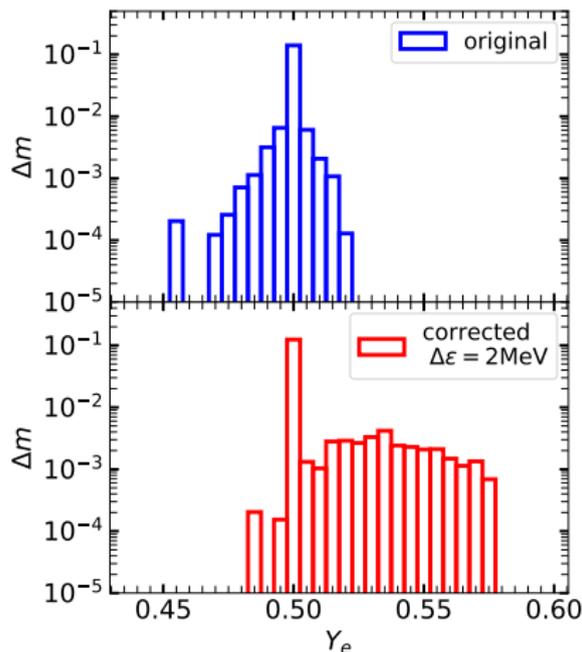
Kobayashi et al. (2020)

- High Zn abundances in metal poor stars are challenging for most supernova models
- Promising for Zn and Sr abundances in HE1327-2326

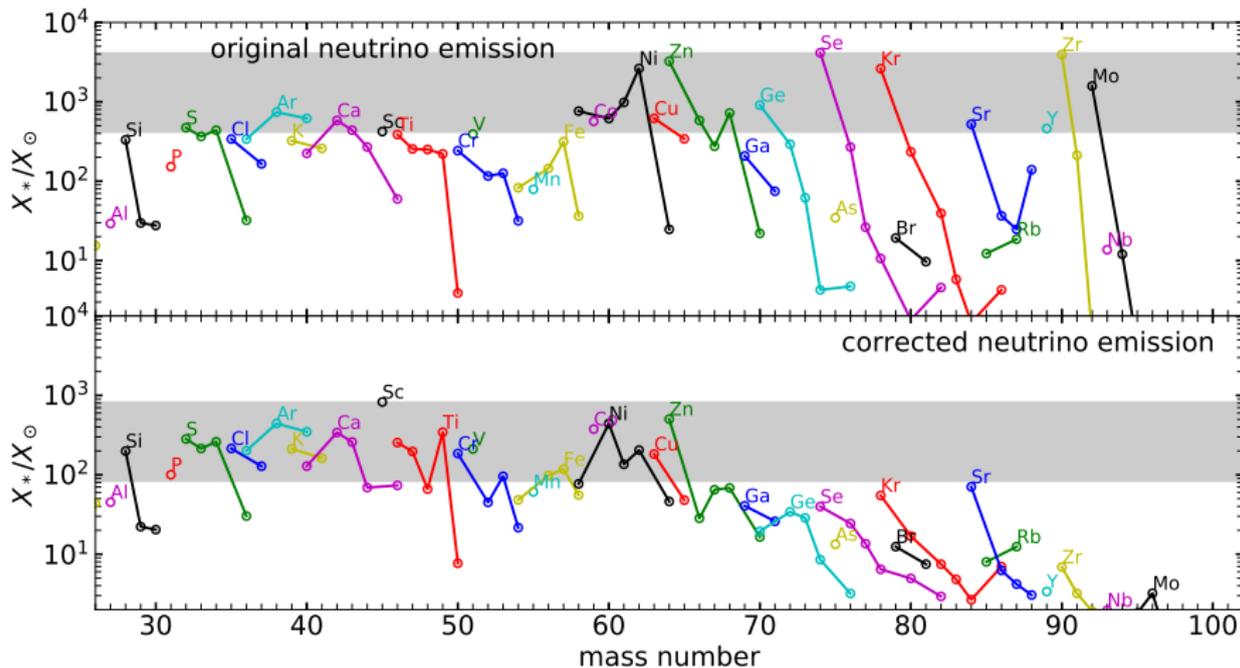
- 1 Introduction
- 2 Effects on the Neutrino Driven Wind
- 3 Impact on the innermost supernova ejecta

Supernova simulation

- Composition of the inner, neutrino heated supernova ejecta are also strongly affected by neutrino irradiation
- 3D Simulation for a $11.8 M_{\odot}$ star (Müller+ 2019, Janka+ 2015)
- Nucleosynthesis of the innermost $0.1 M_{\odot}$
- $\Delta\varepsilon = \langle \varepsilon_{\bar{\nu}_e} \rangle - \langle \varepsilon_{\nu_e} \rangle$
- Original: $\Delta\varepsilon \approx 4$ MeV
- Corrected: $\Delta\varepsilon \approx 2$ MeV (artificially increased $\langle \varepsilon_{\nu_e} \rangle$)
- Flavor conversions would have similar effects on the spectra



Nucleosynthesis results



- Flavor equilibrium would suppress production of p isotopes
- Enhanced yields of Sc
- Change pattern of Ti, Fe, Ni chains

Conclusions & Outlook

- Fast flavor transformation favor the νp process in neutrino driven winds
- Make NDWs more likely to be a major source of p -isotopes
- Overproduction is only problematic if the most favorable conditions are assumed
- Major effects also on the inner supernova ejecta, may suppress heavy element production
- (not in this talk) Affects the production of isotopes due to the ν -process (^{26}Al , ^{36}Cl , ^{138}La , ^{180}Ta)